

Workshop “Holography, Quantum Entanglement
and Higher Spin Gravity”
Feb. 6–7, 2017, at YITP, Kyoto U.

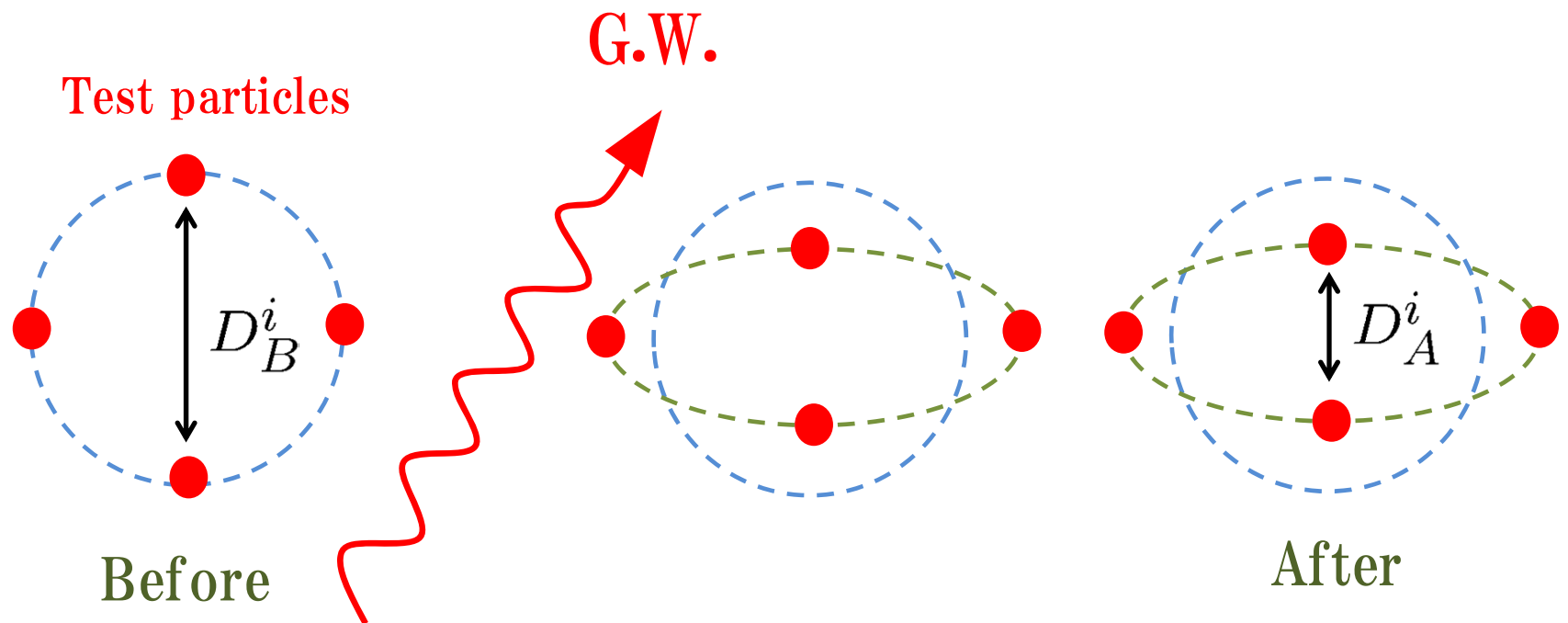
BMS Supertranslations and Gravitational Memory

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1612.03290 w/ S. Hollands and R.M. Wald
1702.00095 D. Garfinkle and A. Tolish

Memory Effect

Permanent displacement of test particles produced by radiation burst



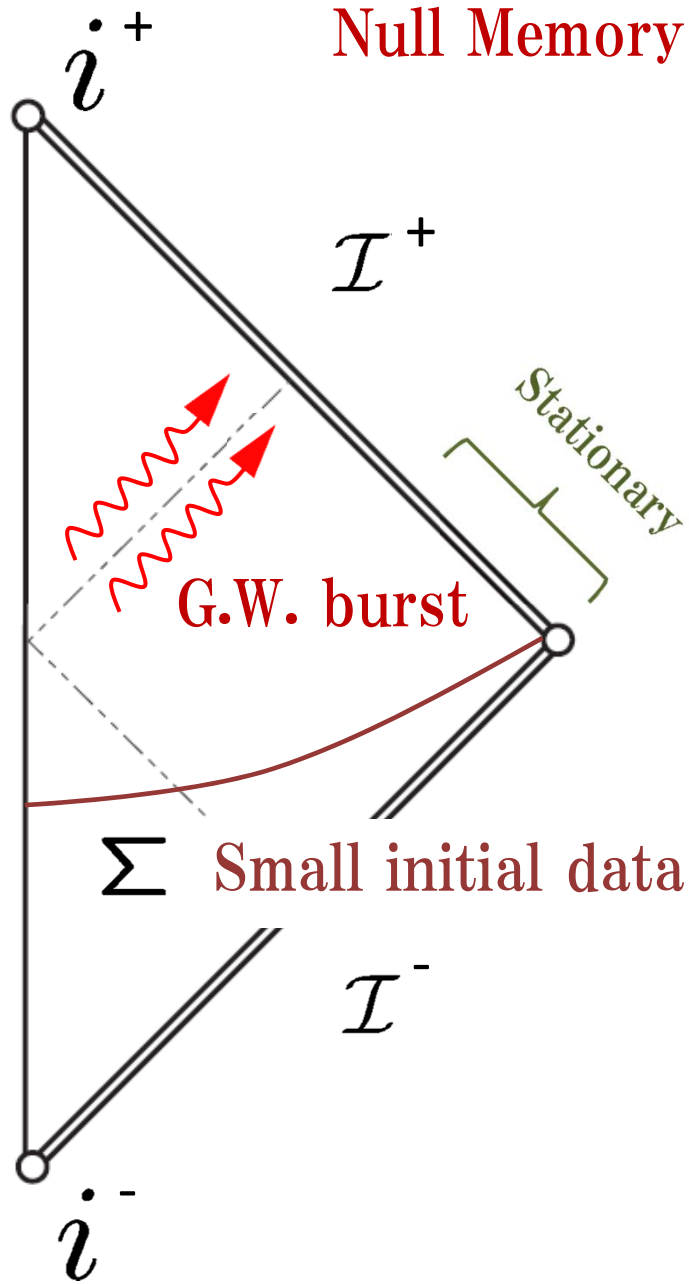
$$\Delta h_{ij}^{\text{TT}} = \frac{1}{r} \Delta \sum_a \frac{4M_a}{\sqrt{1 - v_a^2}} \left(\frac{v_a^i v_a^j}{1 - v_a \cos \theta} \right)^{\text{TT}}$$



Two types of memory effects

Null Memory (c.f. Non-linear memory)

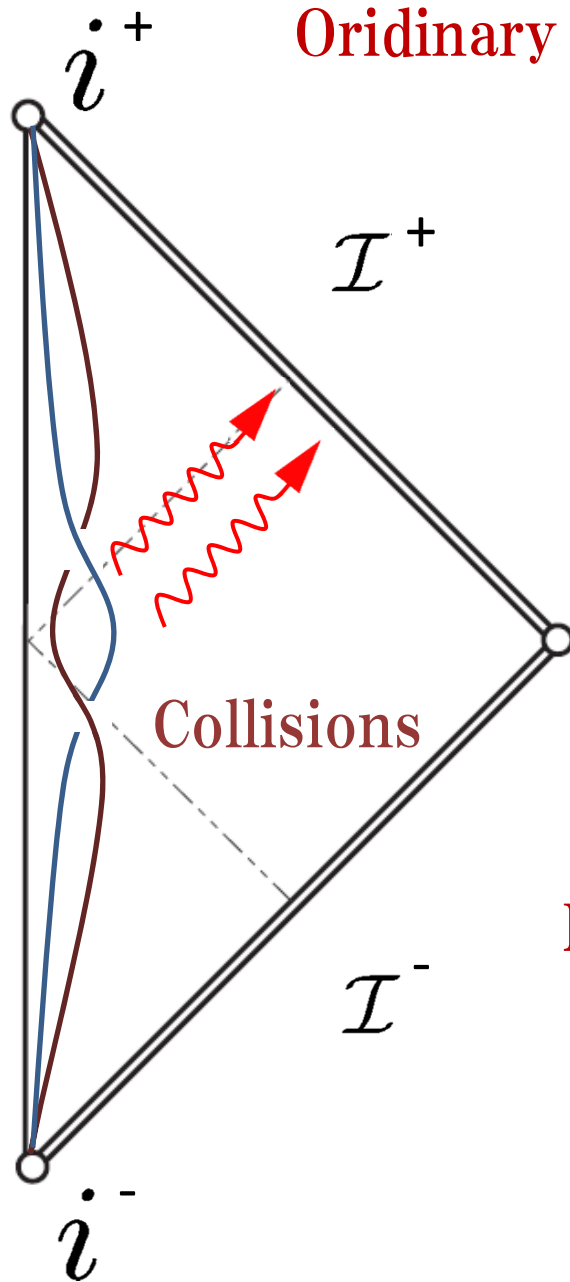
Christodoulou 91



Energy-Momentum is carried mainly by **radiation**

Ordinary Memory (c.f. Linear memory)

Zel'dovich-Polnarev 74



Most Energy-Momentum is carried by the matter

**Particles/matter come in from i^-
go out to i^+**

Basic ideas

Perturbations $g_{ab} = \eta_{ab} + h_{ab}$

Einstein equation + gauge condition

$$\nabla^c \nabla_c h_{ab} = 16\pi \underline{S_{ab}} \quad \text{Source}$$

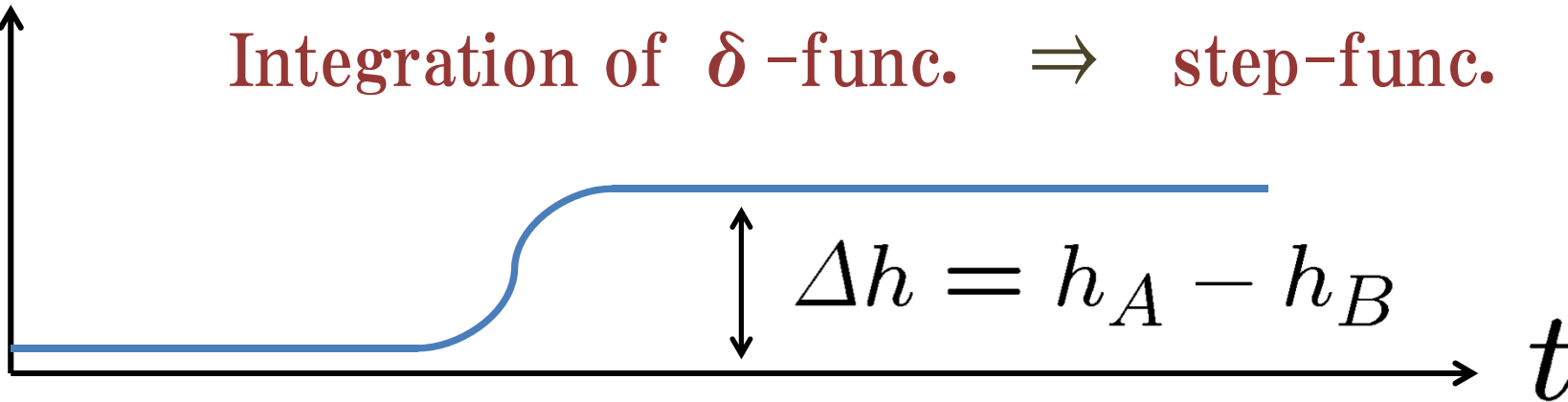
$$h_{ab}(\mathbf{x}) = 16\pi \int d^D \mathbf{x}' \underline{G(\mathbf{x} - \mathbf{x}') S_{ab}(\mathbf{x}')}$$

Green's function

$$D = 4 \quad \text{Green's function} \quad G \sim \frac{1}{4\pi} \frac{\delta(t - r)}{r}$$

h

Integration of δ -func. \Rightarrow step-func.



Geodesic deviation $\frac{d^2}{dt^2} D^i = \underline{R_{titj}} D^j$
 curvature wrt h_{ab}

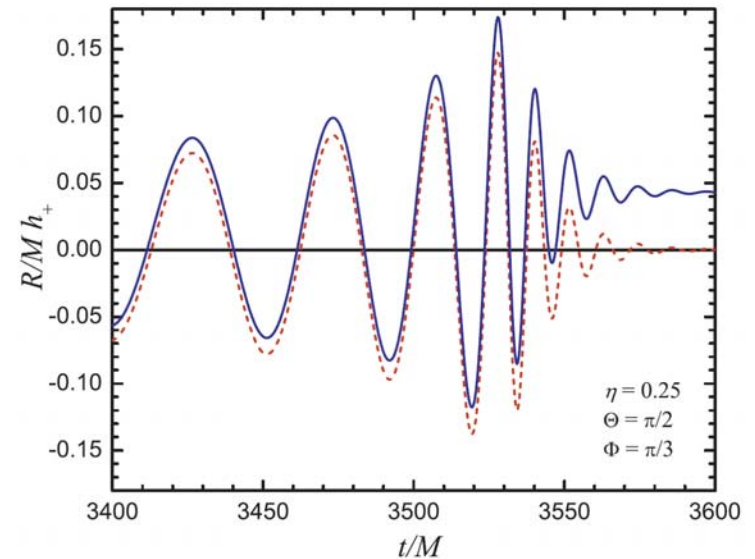
\rightarrow Displacement $\Delta D^i = \frac{1}{2} \Delta h_{ij}^{\text{TT}} D^j$

Why interesting?

I. Memory effects \Rightarrow detectable

e.g. LISA

Supermassive BH
mergers out to $z \leq 2$



e.g. Favata '10

Why interesting?

II. BMS symmetry and Soft graviton

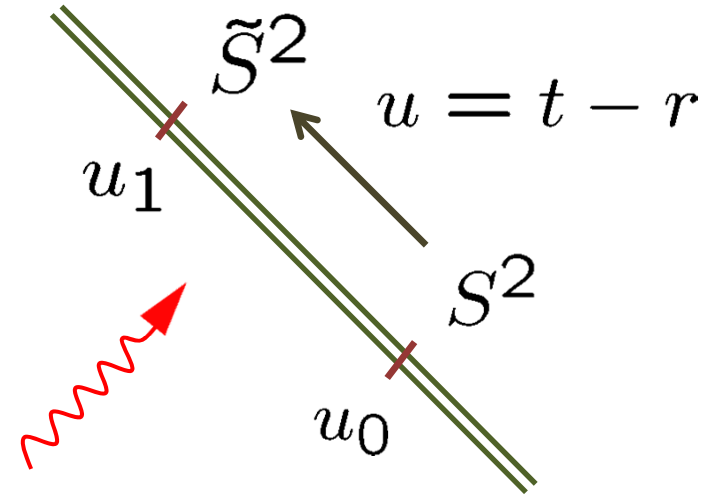
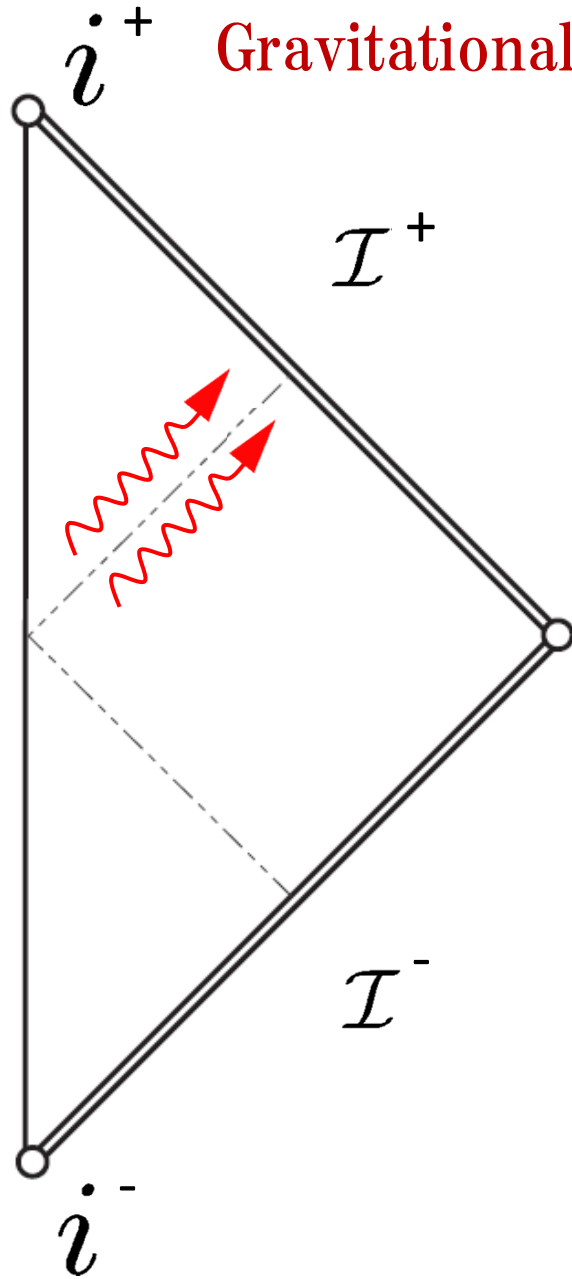
Memory,
BMS supertranslations, and
Soft theorems

Strominger '14, Strominger-Zhiboedov '15

BMS supertranslations, and
Black hole's soft hair

Hawking-Perry-Strominger '16

Gravitational Wave Memory and BMS Symmetry



G.W. burst

$$\mathcal{I}^+ \simeq \mathbf{R} \times S^2$$

$$\tilde{g}_{ab} = \Omega^2 g_{ab}$$

$$\Omega \sim \frac{1}{r} \rightarrow 0 \quad \text{at null infinity } \mathcal{I}$$

Asymptotic Symmetry

e.g. Minkowski space $\tilde{g}_{ab} = \Omega^2 \eta_{ab}$

In physical spacetime

Poincare group

10 parameters

- Translations

4

- Lorentz transf.

6

At null infinity

BMS group

Infinite dimensions

- Super translations

$$\mathbf{X} = T(\theta, \phi) \cdot \frac{\partial}{\partial u} + \dots$$

map generator to generator

- \sim BMS/super translations

BMS Supertranslations : angle dependent translations

- Generator : $\mathbf{X} = T(z) \cdot \left(\frac{\partial}{\partial u} - \frac{\partial}{\partial r} \right) - \frac{1}{r} \cdot D^A T(z) \cdot \frac{\partial}{\partial z^A}$

$z^A = (\theta, \phi)$: angle coordinates on S^2

$T(z)$: spherical harmonics on S^2 w/ $\ell \geq 2$

(c.f. $\ell = 1 \Rightarrow \mathbf{X}$: Lorentz boost)

- Displacement tensor : $\Delta_A^B = - \left(D_A D^B - \frac{1}{2} \delta_A^B D^c D_c \right) T(z)$
- Displacement : $\Delta D^B = \frac{1}{r} \Delta_A^B D^A$

Memory effects in Higher Dimensions?

Our claim

No Memory in HD General Relativity

1. Explicit calculations in a simple model 1702.00095
2. General analysis 1612.03290

1. Explicit calculations in a simple model

1702.00095

Garfinkle-Hollands-AI-Tolish-Wald

$$D = 4 \quad \text{Green's function} \quad G \sim \frac{1}{4\pi} \frac{\delta(t-r)}{r}$$

Integration of δ -func. \Rightarrow step-func.

$D > 4$ (even dimensions)

Leading term is given by derivatives of δ -func.

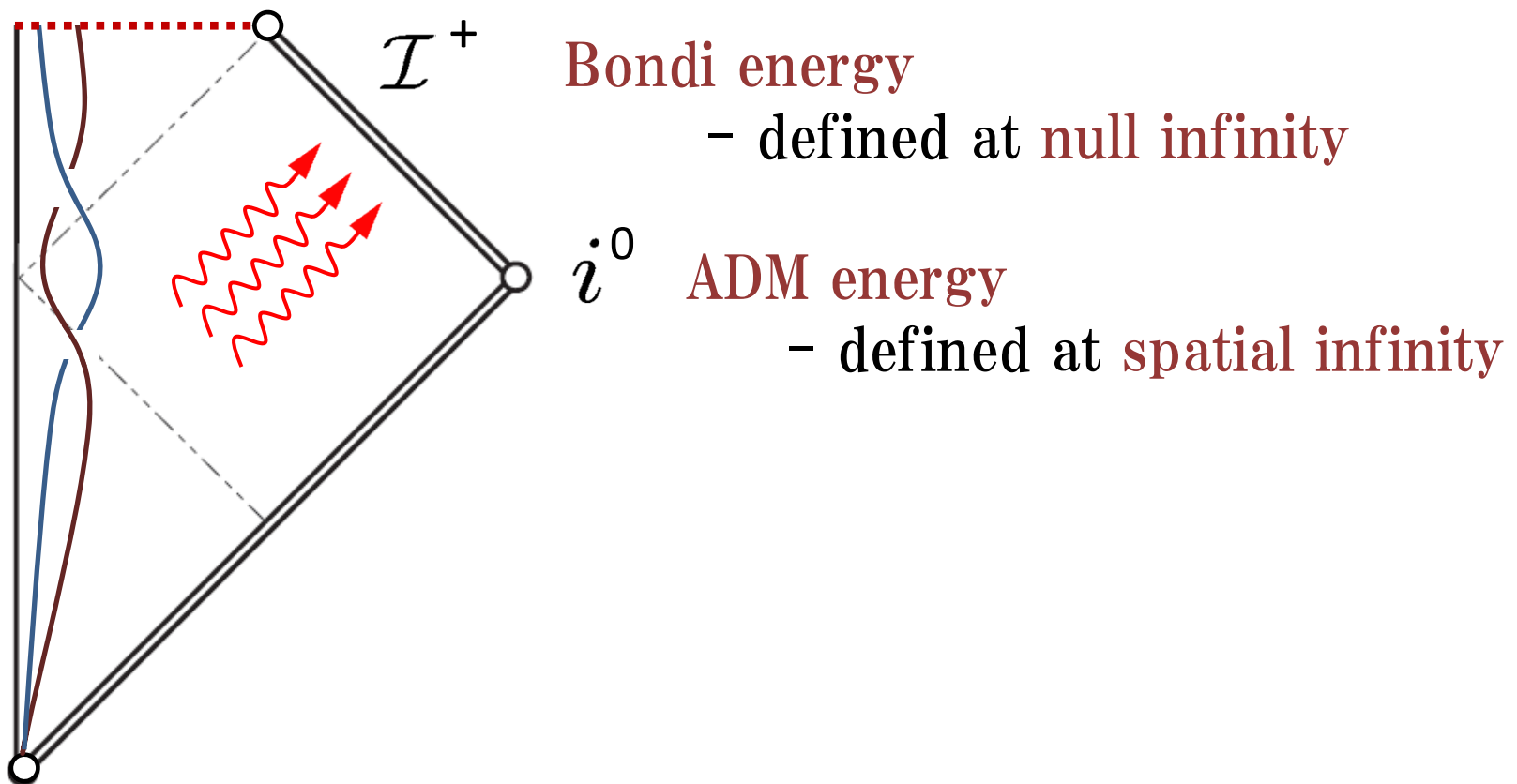
$$\text{e.g. } D = 6 \quad G \sim \frac{\delta'(u)}{r^2} + \dots$$

2. General analysis

1612.03290

Hollands-AI-Wald

Notion of asymptotic flatness and energy



How to define asymptotic flatness?

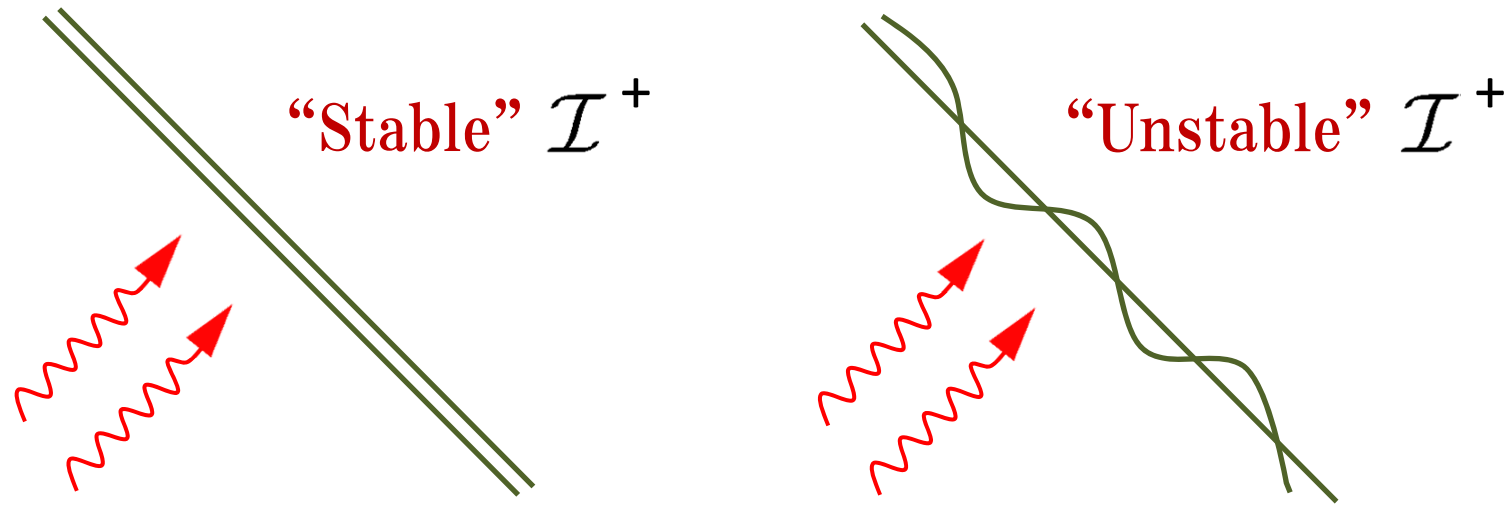
I. The fall off conditions should be **general enough**

→ to include interesting phenomena
e.g. radiating spacetimes

II. _____ should be **strength enough**

→ to allow one to derive physically meaningful expression for energy

Stability of conformal null infinity



Perturbations

- Null infinity \mathcal{I} should remain a **regular null surface**

Perturbed unphysical metric $\delta\tilde{g}_{ab} = \Omega^2\delta g_{ab}$

\Rightarrow should be extended to \mathcal{I}

$D = 4$ case \Rightarrow stable : Geroch - Xanthopoulos '78

Stability of conformal null infinity

In even dimensions $D > 4$

Desirable fall-off conditions

$$\delta\tilde{g}_{ab} = O(\Omega^{(d-2)/2})$$

$$\delta\tilde{g}_{ab}\tilde{\nabla}^a\Omega = O(\Omega^{d/2})$$

$$\delta\tilde{g}_{ab}\tilde{\nabla}^a\Omega\tilde{\nabla}^b\Omega = O(\Omega^{(d+2)/2}) \leftarrow \text{remain null}$$

$$\tilde{g}^{ab}\delta\tilde{g}_{ab} = O(\Omega^{d/2}) \leftarrow \text{particular in } D > 4$$

— Linearized version of our general analysis of 1612.03290

Proof :

Hollands-AI 05

{ TT gauge works !
Choice of variables ϕ_α :

$$\tau_{ab} := \Omega^{-(D-2)/2} \delta \tilde{g}_{ab} \quad \tau_a := \Omega^{-1} \tau_{ab} \tilde{\nabla}^b \Omega$$

Einstein's equations reduce to a hyperbolic system:

$$\tilde{\nabla}^c \tilde{\nabla}_c \phi_\alpha = A_\alpha^{\beta a} \tilde{\nabla}_a \phi_\beta + B_\alpha^\beta \phi_\beta$$

- can show this possesses a well-posed initial value problem under our fall-off conditions



$\Rightarrow \phi_\alpha$ has smooth extension on \mathcal{I}^+





- Difference between $4D$ and $D > 4$

radiations drop off $\sim \frac{1}{r^{(D-2)/2}}$

static Coulomb part drops off $\sim \frac{1}{r^{D-3}}$

$D = 4$  both parts behave $\sim \frac{1}{r}$  Memory

$D > 4$  $1 < \frac{D-2}{2} < D - 3$  No Memory

Our fall-off conditions

⇒ guarantee **stability** of conformal null infinity

⇒ **general enough**

to allow for **radiating spacetimes**

⇒ **stringent enough**

to allow for physically meaningful **energy**

(i.e. well-defined symplectic current)

..... but NOT allow Supertranslations

Generator of asymptotic symmetry

$$\mathbf{X} = T(\text{angle}) \cdot \frac{\partial}{\partial u} + \dots$$

$$\left\{ \begin{array}{l} D = 4 \quad T(z) : \text{arbitrary func. on } S^2 \\ D > 4 \quad T(z) : \text{spherical harmonics on } S^{D-2} \end{array} \right.$$

$$\text{w/ } \ell = 0, 1$$

 $\mathbf{X} : D$ -dimensional translations

How to define Bondi energy?

— follow Noether charge method Wald-Zoupas '00

$$\left\{ \begin{array}{l} L = \frac{1}{16\pi} R \cdot \epsilon \\ \theta : (D-1)\text{-form} \quad \delta L = \text{EOM} \cdot \delta g + d\theta \\ \omega : \text{symplectic current} \\ \omega(\delta_1 g, \delta_2 g) = \delta_1 \theta(\delta_2 g) - \delta_2 \theta(\delta_1 g) \end{array} \right.$$

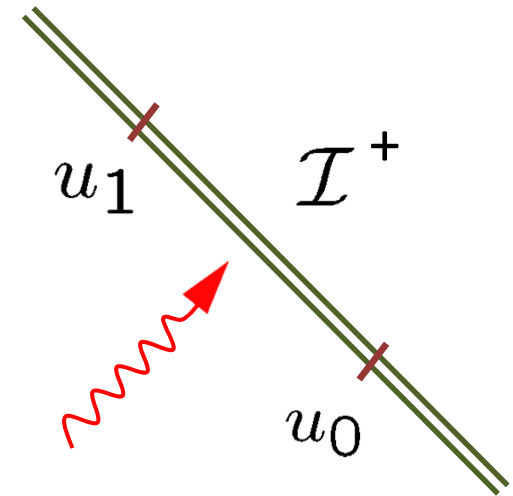
Under **our fall-off conditions**,
the symplectic current ω is **well-defined**

News tensor :

$$N_{ab} := \Omega^{-(D-4)/2} \times (\text{pull back of Schouten tensor})$$

Flux formula :

$$\mathcal{H}_X(u) = -\frac{1}{32\pi} \int_{S^{D-2}(u)} T N^{ab} N_{ab}$$



Grav. Waves

All these are well-defined on our fall-off conditions

Summary :

- $D = 4$
 - BMS Supertranslations must be included
 - \Rightarrow parametrized by $\forall T(z)$ on S^2
- $T(z)$ acts as “potential” of Memory
- $D > 4 \Rightarrow$ need to impose reasonable fall-off conditions in order to guarantee
 - stability of null-infinity
 - well-defined of symplectic current \Rightarrow No Supertranslations \Rightarrow No memory