

# Gravity Actions from Tensor Networks

[hep-th/1609.04645] with T. Takayanagi and K. Watanabe  
and ongoing work with  
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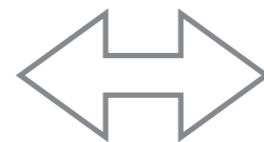
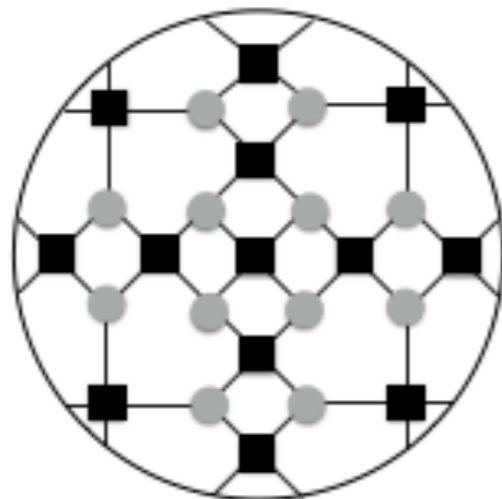
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# Motivation

- We do not know the precise mechanism of AdS/CFT.
- There is a significant similarity between Tensor network and gravity.

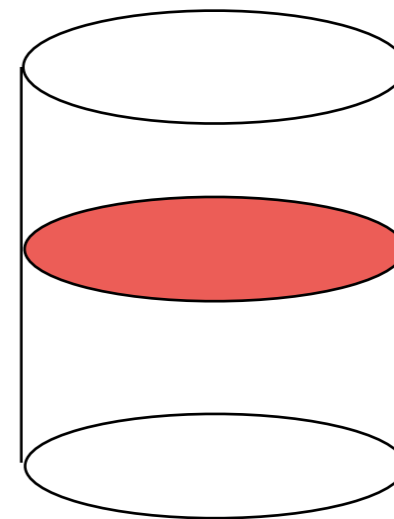
[Swingle]

Tensor Network



Bulk timeslice in AdS/CFT

Equivalent?



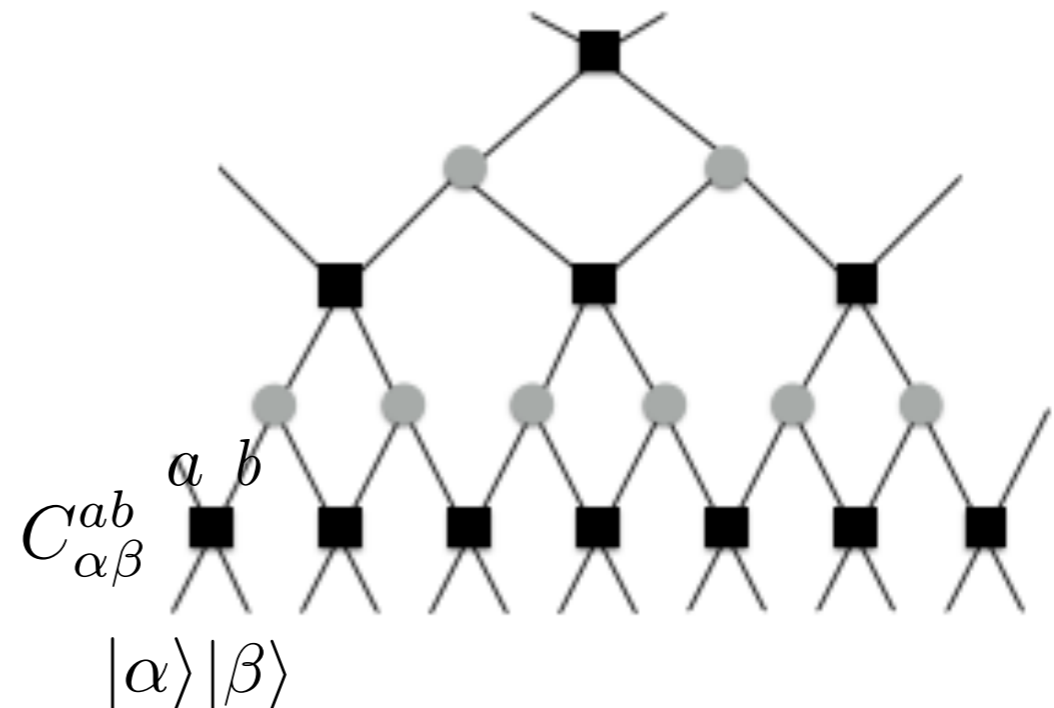
The aim is to find precise relation between Tensor Network and gravity.

# Tensor Network [White][Vidal]

Tensor network: Network of tensors. By contracting tensors, we can express complicated states.

- It can simulate ground state of field theory.

- RG flow of state.



- Grounds state gets simplified and loses short range entanglement gradually.

# Tensor Network

State without short range quantum entanglement

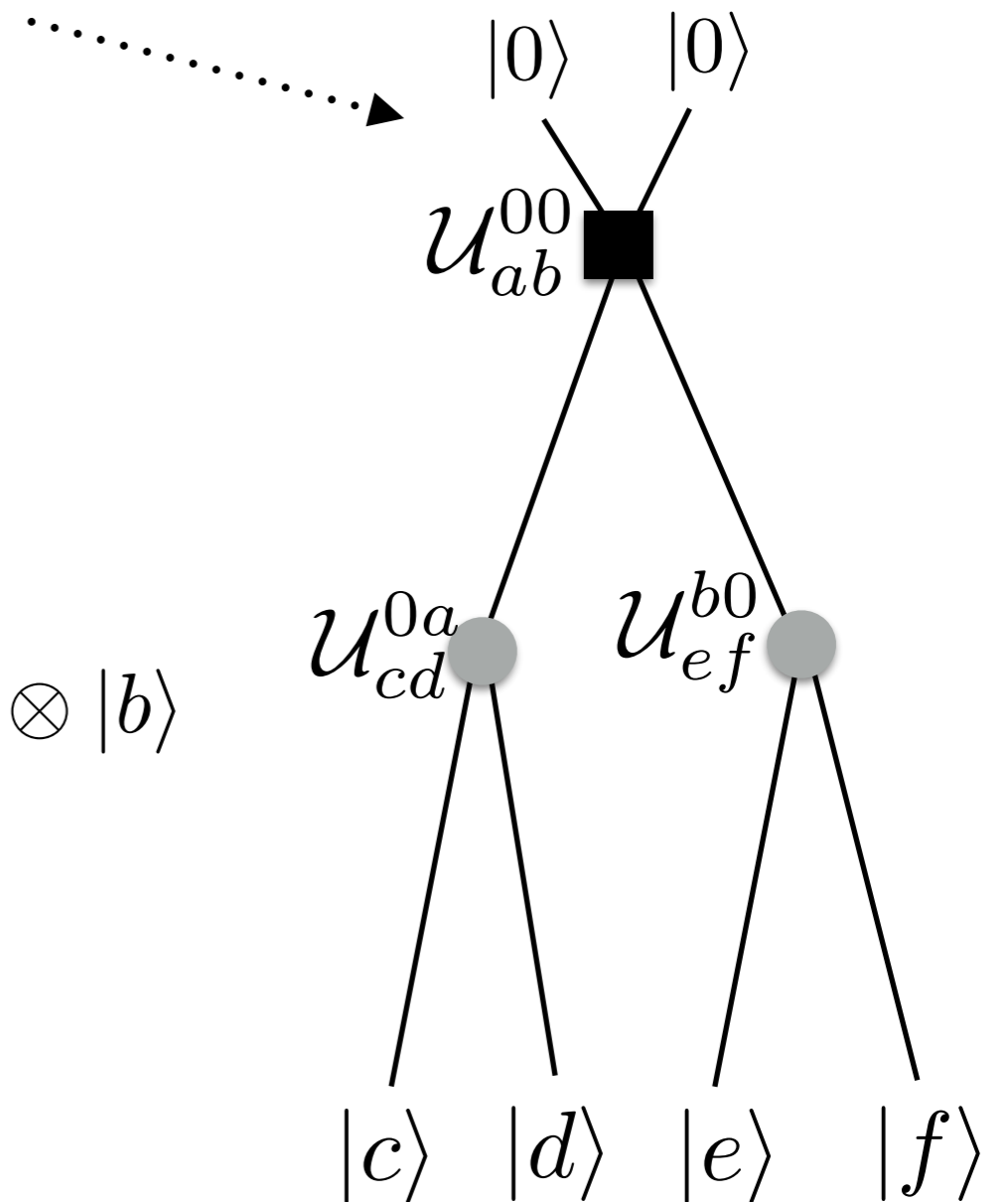
Reduce quantum entanglement

Coarse graining

Ground state

$$\sum_{c,d,e,f} \mathcal{U}_{0a}^{cd} \mathcal{U}_{ef}^{b0} \sum_{ab} \mathcal{U}_{ab}^{00} |c\rangle \otimes |d\rangle \otimes |e\rangle \otimes |f\rangle$$

$$\begin{aligned} & \dots \rightarrow |0\rangle \otimes |0\rangle \\ & \uparrow \\ & \sum_{a,b} \mathcal{U}_{ab}^{00} |a\rangle \otimes |b\rangle \\ & \uparrow \\ & \dots \end{aligned}$$



# Tensor Network and Ryu Takayanagi Formula

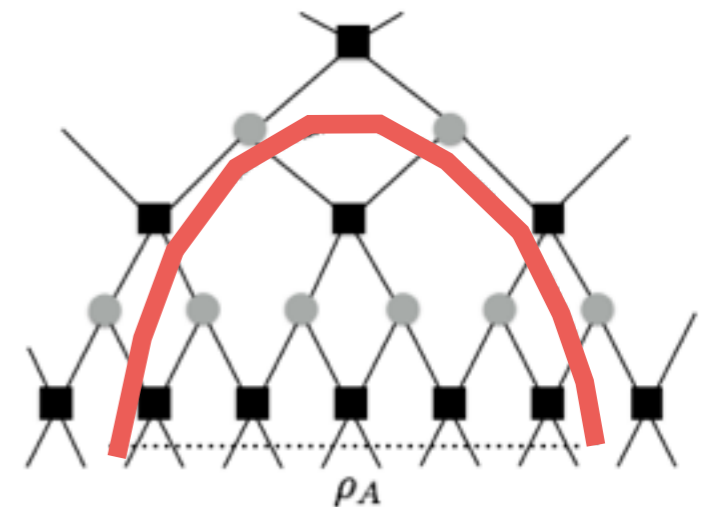
- For certain tensor networks, such as MERA, entanglement entropy of state can be expressed by number of bonds of minimal bond surface.

$$S[\rho_A] \sim \mathcal{O}(1) \text{ constant}$$

$$\times \#(\text{Bonds intersecting the minimal bonds surface})$$

$$\times \log(\text{dimensions of bond})$$

- We can identify number of intersecting bonds as area of surface.



$$S[\rho_A] \sim \mathcal{O}(1) \text{ constant} \quad \times \quad (\text{Area of the minimal surface})$$

$$\times \log(\text{dimensions of bond})$$

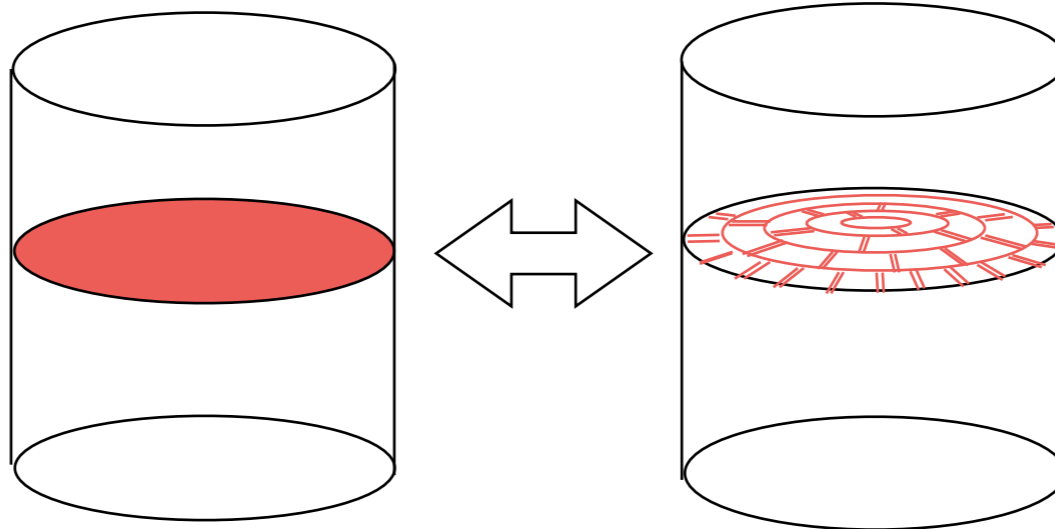
- Significant similarity to **Ryu-Takayanagi formula** in AdS/CFT. [Swingle]

# Tensor Network = Bulk timeslice?

Conjecture

[Swingle]

Tensor network of CFT ground state = Timeslice of bulk spacetime in AdS/CFT



Timeslice of  $AdS_{d+1}$

Tensor network

## Problems

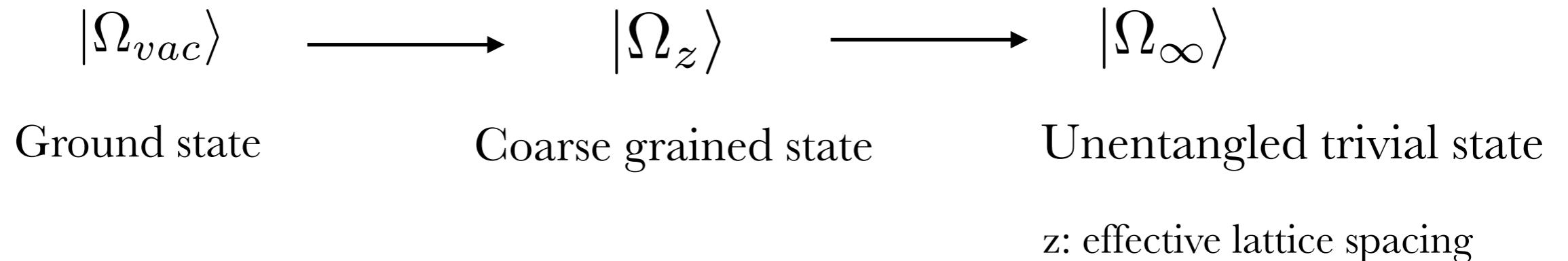
- How should we determine correct tensor network?
- Relation between gravity and tensor network is still ambiguous.

# Continuous Tensor Network

[Haegeman, Osborne, Verschelde and Verstraete]

[M.M, Ryu, Takayanagi, Wen]

- Interpolation from vacuum state to unentangled trivial state.



- Boundary states have no short range quantum entanglement, so we can use boundary states as  $|\Omega_\infty\rangle$ .

- In Euclidean path integral, wave function is

$$\langle \phi_0 | \Omega_{vac} \rangle = \int_{\tau < 0} \mathcal{D}\phi(x, \tau) e^{-\int_{-\infty}^0 d\tau \int dx \mathcal{L}_E[\phi(\tau, x)]} \Bigg/ \sqrt{\int \mathcal{D}\phi(x, \tau) e^{-S_E[\phi]}}$$

$$\begin{array}{l}
 \phi(x, 0) = \phi_0(x) \\
 \phi(x, \tau = -\infty) = \phi_B(x)
 \end{array}
 \qquad
 \begin{array}{l}
 \phi(x, \infty) = \phi(x, -\infty) \\
 = \phi_B(x)
 \end{array}$$

# Deforming action

[M.M, Takayanagi, Watanabe]

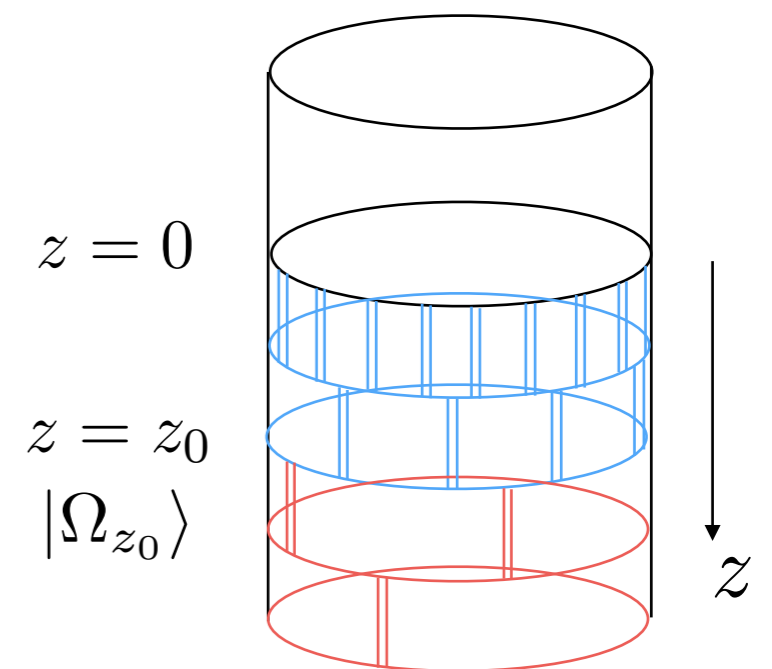
- By deforming this Euclidean path integral, we can simulate tensor network.
- In the deformation, we identify  $z = -\tau$  as **effective lattice spacing**.  
In other word, UV cut off is **z dependent**, and modes with spatial momentum larger than  $1/z$  are discarded.

$$\langle \phi_0 | \Omega_{vac} \rangle \propto \int_{z>0} \mathcal{D}\phi(x, z) e^{-\int_{\epsilon}^{\infty} dz \int dk d\omega \theta(z \sqrt{(k^2 + \omega^2)/2}) \mathcal{L}_E[\phi]}$$

$$\phi(x, \epsilon) = \phi_0(x)$$

where  $\theta(x) = 0 (x > 1)$

- By this deformation, we can discard redundant path integral, without changing wavefunction.





# Most efficient path integral proposal

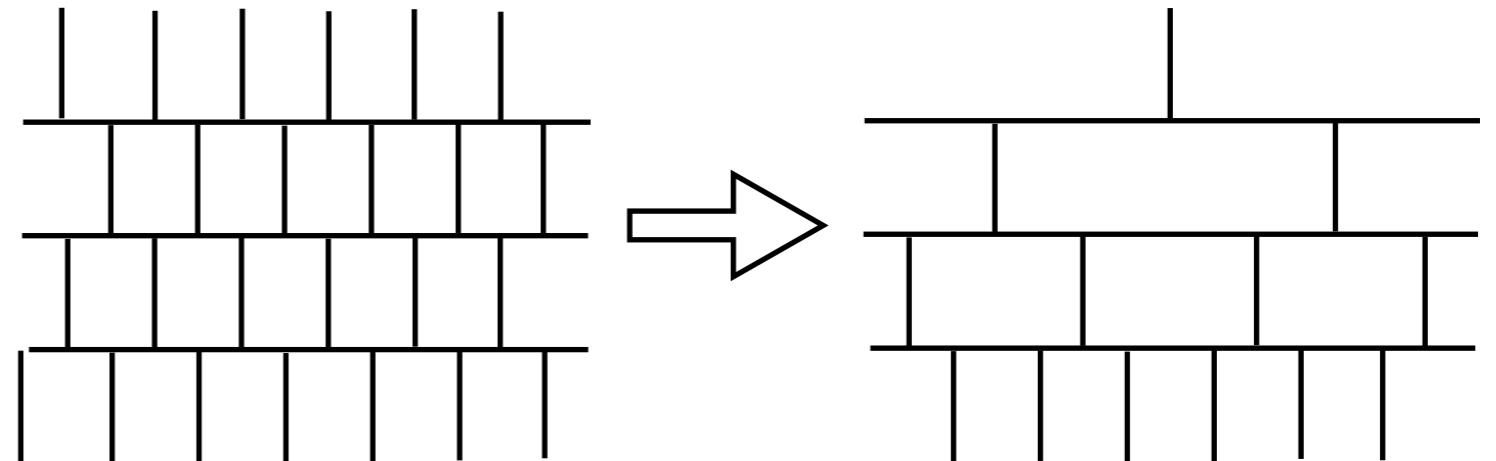
Work in progress [Pawel, Kundu, M.M, Takayanagi, Watanabe]

## Questions

- The deformation of the path integral is not unique.
- No criterion relate it to gravity.

## Proposal to fix the deformation

We need to eliminate redundant tensors from the path integral, in order to relate tensor network with AdS/CFT and Ryu-Takayanagi formula.



Proposal

The deformation with minimum redundant tensors corresponds to bulk time slice in AdS/CFT.

# Most efficient path integral proposal

- We deform the path integral by Weyl transformation to the CFT, without changing the metric at  $z = 0$ .

$$ds^2 = dz^2 + dx^2 \quad \Rightarrow \quad ds^2 = e^{\gamma\phi(x,z)}(dz^2 + dx^2)$$

## A definition of redundancy

- There is a variety of ways to define redundancy of tensor network.
- Let's consider path integral representation of the vacuum state.

$$\langle \psi_0 | \Omega_{vac} \rangle = \int_{\psi(x,\epsilon)=\psi_0(x)} \mathcal{D}\psi \, e^{-S[\psi]} = \boxed{e^{-S_L[\phi]}} \int_{\psi(x,\epsilon)=\psi_0(x)} \mathcal{D}\psi \, e^{-S[\psi,\phi]}$$

$\boxed{e^{-S_L[\phi]}}$  corresponds to amount of redundancy of the original path integral compared to deformed one.

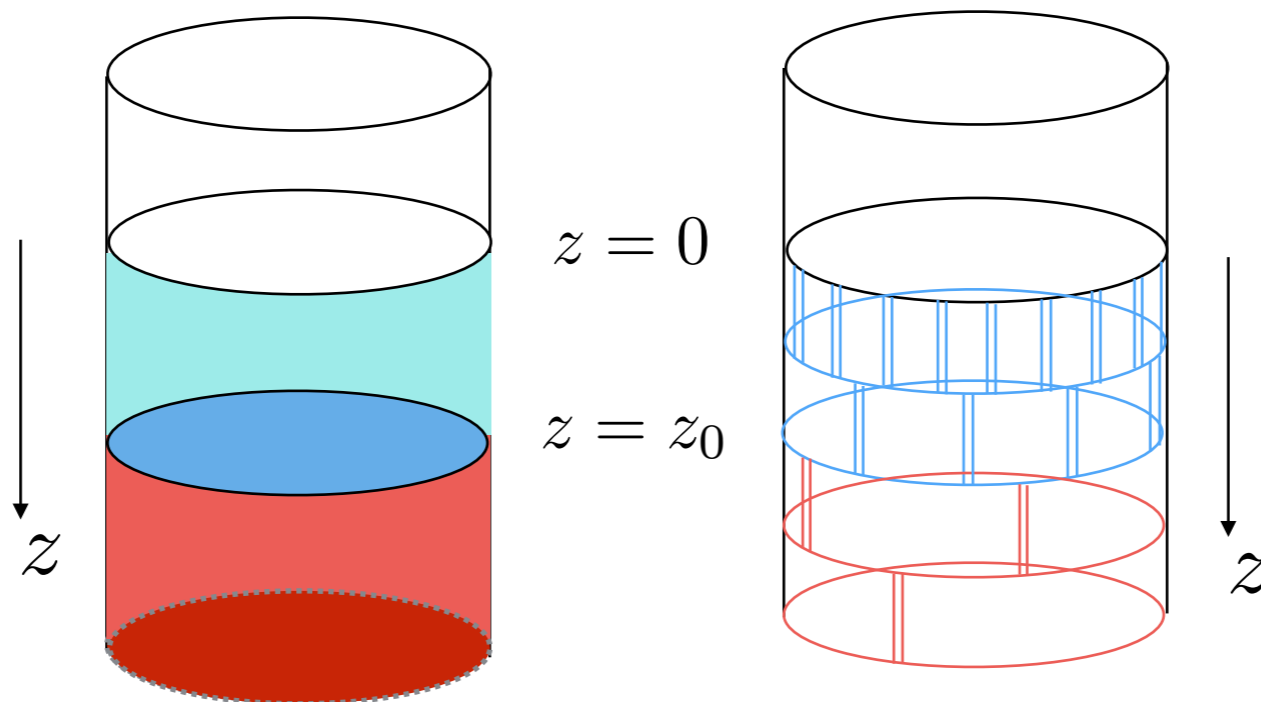
# Most efficient path integral proposal

## Emergent metric

“The most efficient path integral”  $\longleftrightarrow$   $e^{-S_L[\phi]}$  takes maximal value.

We identify the metric  $ds^2 = e^{\gamma\phi(x,z)}(dz^2 + dx^2)$  as metric on the bulk timeslice.

$\Rightarrow$   $\phi(x, z)$  is determined by minima of  $S_L[\phi]$



# Example: 2d CFT

For Weyl transformation,

$$[\mathcal{D}\phi_{ds^2=e^{\gamma\phi}(dz^2+dx^2)}] = e^{S_L} [\mathcal{D}\phi_{ds^2=dz^2+dx^2}]$$

where  $S_L$  is the Liouville action

$$S_L = \frac{c}{24\pi} \int_{\epsilon}^{\infty} dz \int_{-\infty}^{\infty} dx \left[ (\partial_x \phi)^2 + (\partial_z \phi)^2 + \frac{\mu}{\gamma^2} (e^{\gamma\phi} - 1) \right] \quad \mu > 0$$

is minimized when

$$e^{\gamma\phi} = \frac{4}{\mu} z^{-2} \quad ds^2 = \frac{4}{\mu} \frac{dz^2 + dx^2}{z^2}$$

So we get AdS solution as true minimum.

In contrast to Einstein-Hilbert action, Liouville action is **bounded from below**.

# Example: 2d CFT

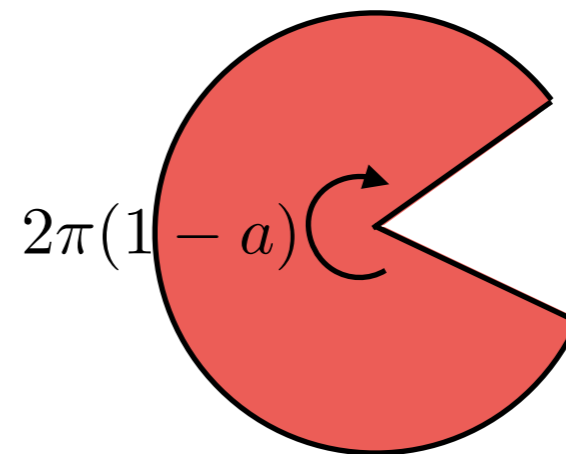
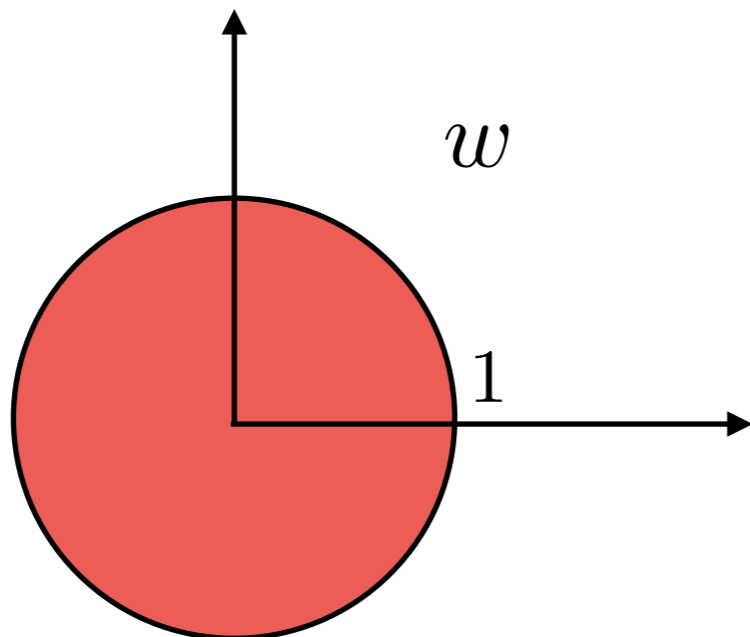
## Excited state

For state excited by primary operator with dimension  $h$  at  $w = 0$  ( $|w| < 1$ )  
assuming large  $c$ , corresponding solution is

$$ds^2 = \frac{4}{\mu} \cdot \frac{a^2}{|w|^{2(1-a)} (1 - |w|^{2a})^2} dw d\bar{w} \quad \left( a = 1 - \frac{12h}{c} \right)$$

deficit angle geometry with angle  $2\pi(1 - a)$ .

In  $\text{AdS}_3/\text{CFT}_2$ , the relation is  $a = \sqrt{1 - \frac{24h}{c}}$ , consistent when  $h \ll c$ .



# Example: 2d CFT

## BTZ black hole

Similarly, by considering wave function of Thermo field double state,

$$|\Psi_{TFD}\rangle = \frac{1}{\sqrt{Z(\beta)}} \sum_i e^{-\beta E_i} |E_i\rangle \otimes |E_i\rangle$$

$$\begin{aligned} \langle \tilde{\psi}_1, \tilde{\psi}_2 | \Psi_{TFD} \rangle &= \int_{-\beta/4 < z < \beta/4} \mathcal{D}\psi(x, z) e^{-S_{CFT}} \\ \psi(x, -\beta/4) &= \tilde{\psi}_1(x) \\ \psi(x, \beta/4) &= \tilde{\psi}_2(x) \end{aligned}$$

we get Liouville action.

As the minimum of the action, we obtain BTZ solution.

$$ds^2 = \frac{4}{\mu} \cdot \frac{4\pi^2}{\beta^2} \cdot \frac{dz^2 + dx^2}{\cos^2\left(\frac{2\pi z}{\beta}\right)}$$

# Summary

- We expressed flow of tensor network states by path integral with deformed action.
- We fixed one tensor network flow by assuming “most efficient path integral”.
- We identified bulk metric with background metric, which is consistent with several examples.

## Questions

- Entanglement entropy? Quantum corrections? Higher dimensions?