



Center for Gravitational Physics
Yukawa Institute for Theoretical Physics



Conformal blocks from Wilson lines with loop corrections

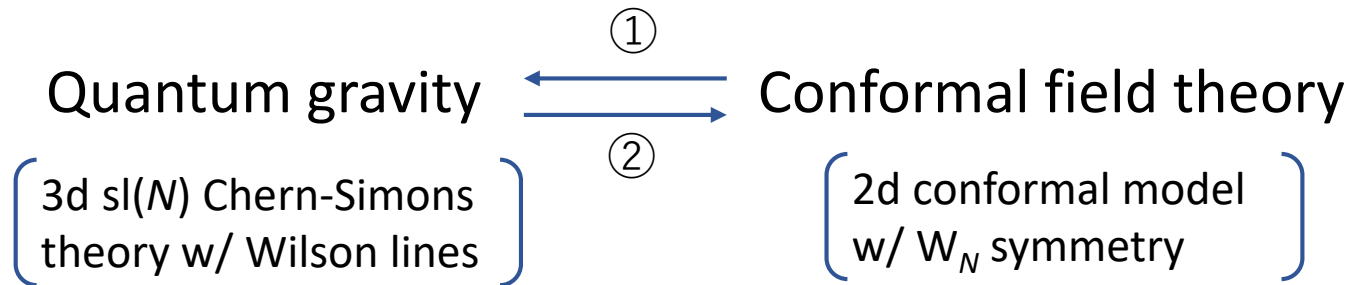
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Refs: YH-Uetoko, PTEP(2017)113B03; 1801.08549

Mar. 16th, 2018

YITP Workshop “Holography, Quantum Entanglement and Higher Spin Gravity II”

Holography and its use



- ① **Quantum** gravity from well-defined CFT
 - Fix the renormalization prescription for quantum gravity corrections
- ② CFT data of CFT from gravity
 - New way to obtain the expressions of conformal blocks

Previous works

□ Formal argument (Verlinde '90)

- CS gravity + Wilson line → Liouville conformal blocks
- A modification of Witten's work '89

□ A practical computation (Fitzpatrick-Kaplan-Li-Wang '16)

- Wilson line + CFT → $1/c$ expansion of conformal blocks
- Useful for real computations, not purely bulk picture

□ A bulk interpretation (Besken-Hedge-Kraus '17)

- Wilson line → Conformal weight at $1/c$ order
- **NOT** succeeded at $1/c^2$ order, no prescription of regularization

Plan of this talk

- Introduction
- General methods
- Conformal blocks from Wilson lines
- Conclusion

General methods

sl(2) CS gravity

□ Solution to EOM (in a gauge)

$$A = e^{-\rho L_0} a(z) e^{\rho L_0} dz + L_0 d\rho$$

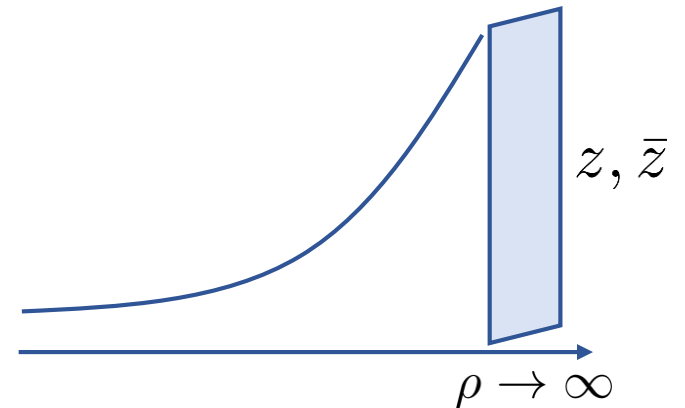
- Boundary DOM are in $a(z)$
 \Leftrightarrow sl(2) WZW model (affine sl(2) sym.)

□ Asymptotic AdS condition

$$(A - A_{\text{AdS}})|_{\rho \rightarrow \infty} = \mathcal{O}((e^\rho)^0) \quad \longrightarrow \quad a(z) = L_1 + \frac{1}{\hat{k}} T(z) L_{-1}$$

- AdS background $\Leftrightarrow a(z) = L_1$
- DS reduction of affine sl(2)

$$\Leftrightarrow \text{Virasoro symmetry by } T(z) \text{ with } c = 6\hat{k} = \frac{3\ell_{\text{AdS}}}{2G_N}$$



$$\left(\begin{aligned} [L_m, L_n] &= (m - n) L_{m+n} \\ m, n &= 0, \pm 1 \end{aligned} \right)$$

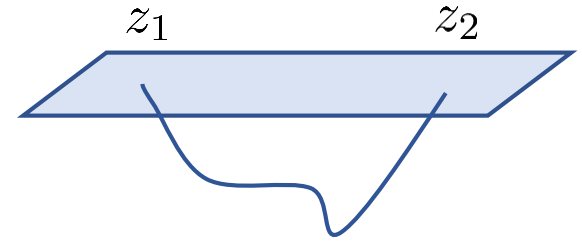
Wilson line in $sl(2)$ CS theory

□ Wilson line operator

$$W(z_2, z_1) = P \exp \left[\int_{z_1}^{z_2} A_\mu dx^\mu \right]$$

$\rightarrow P \exp \left[\int_{z_1}^{z_2} a(z) dz \right] = P \exp \left[\int_{z_1}^{z_2} \left(L_1 + \frac{6}{c} T(z) L_{-1} \right) dz \right]$

A gauge transformation 1/c-expansion



□ Path integral over gauge fields

- We use correlators of $T(z)$ in a CFT with center c

$$\langle T(z_1) T(z_2) \rangle = \frac{c/2}{(z_1 - z_2)^4}$$

Correlators from Wilson lines

- Natural extension of leading order conjecture

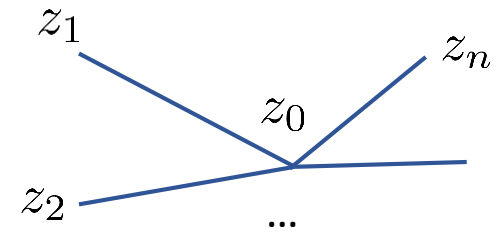
(Bhatta-Raman-Suryanarayana, Besken-Hegde-Hijano-Kraus '16, '17)

$$\langle \mathcal{O}_{h_1}(z_1) \cdots \mathcal{O}_{h_n}(z_n) \rangle = \langle S | \prod_{i=1}^n W_{j_i}(z_i, z_0) | j_i, j_i \rangle$$

$$(L_1 | j, m \rangle = 0 \Rightarrow m = j)$$

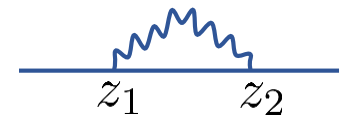
- $\langle S |$ is a singlet in $\mathcal{R}_{j_1} \otimes \cdots \otimes \mathcal{R}_{j_n}$
- A singlet \Leftrightarrow A conformal block
- 2pt function ($z_0 \rightarrow z_2$)

$$\langle \mathcal{O}_h(z_1) \mathcal{O}_h(z_2) \rangle = \langle j, -j | W_j(z_1, z_2) | j, j \rangle$$



Renormalization prescription

□ Divergences would arise at the coincident points of $T(z)$ in the integral

$$\langle T(z_1)T(z_2) \rangle = \frac{c/2}{(z_1 - z_2)^4}$$


□ Our proposal for removing divergences

1. Introduce a regulator ϵ

$$\langle T(z_1)T(z_2) \rangle = \frac{c/2}{(z_1 - z_2)^{4-2\epsilon}}$$

Scaling invariance
is preserved

2. Remove divergences $1/\epsilon^n$ by redefining \mathcal{N}_2, c_2

$$W(z_2, z_1) = \mathcal{N}_2 P \exp \left[\int_{z_1}^{z_2} (L_1 + \frac{6}{c} c_2 T(z) L_{-1}) dz \right]$$

3. ϵ^0 term in c_2 is fixed by Ward identity at the boundary

$$\langle \mathcal{O}_h(z) \mathcal{O}_h(0) T(\infty) \rangle = h z^2 \langle \mathcal{O}_h(z) \mathcal{O}_h(0) \rangle$$

Conformal blocks from
Wilson lines

CFT results

□ \mathcal{O}_{h_j} : Operator corresponding to $|j,j\rangle$

- Obtained by DS reduction of spin j operator in $\mathfrak{sl}(2)$ WZW model
- Conformal weight in $1/c$ expansion


$$\begin{aligned} h_j &= -j - \frac{6j(j+1)}{c} - \frac{78j(j+1)}{c^2} + \dots \\ &= h_0 + \frac{h_1}{c} + \frac{h_2}{c^2} + \dots \end{aligned}$$

□ 2pt function

$$\begin{aligned} \langle \mathcal{O}_{h_j}(z) \mathcal{O}_{h_j}(0) \rangle &= \frac{1}{z^{2h_j}} \\ &= \frac{1}{z^{2h_0}} \left[1 - \frac{2h_1 \log z}{c} + \frac{2h_1^2 (\log z)^2 - 2h_2 \log z}{c^2} \right] + \dots \end{aligned}$$

Wilson line computations

- $\langle \mathcal{O}(z)\mathcal{O}(0) \rangle$ at $1/c$ order

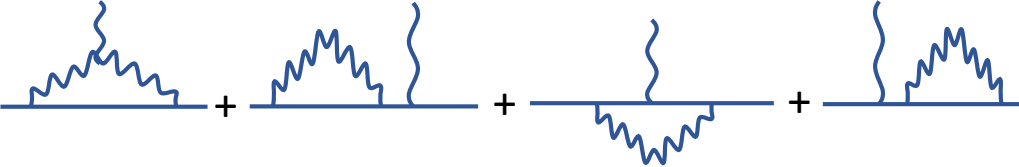


$$= \frac{1}{cz^{2h_0}} \left[\frac{6(h_0 - 1)h_0}{\epsilon} + 12h_0(h_0 - 1) \log z + 2h_0(5h_0 - 2) \right]$$

h_1 reproduced

Absorbed by redefining N_2

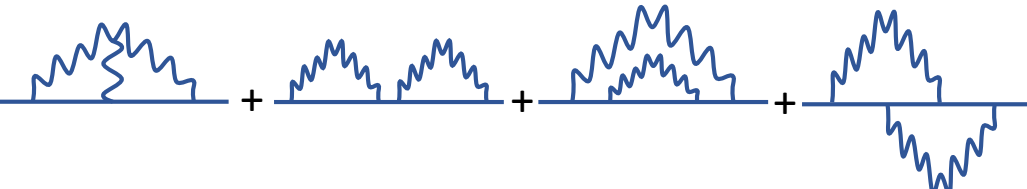
- $\langle \mathcal{O}(z)\mathcal{O}(0)\mathcal{T}(\infty) \rangle$ at $1/c$ order



$$\Rightarrow \left(h_0 + \frac{h_1}{c} \right) z^2 \langle \mathcal{O}(z)\mathcal{O}(0) \rangle$$

$$c_2 = 1 + \frac{1}{c} \left(\frac{6}{\epsilon} + \mathbf{3} \right) + \dots$$

- $\langle \mathcal{O}(z)\mathcal{O}(0) \rangle$ at $1/c^2$ order



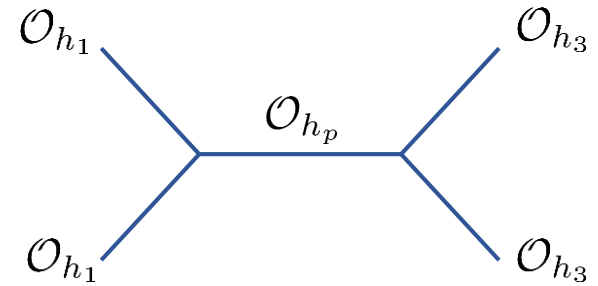
h_2 reproduced with redefined N_2, c_2

Conformal blocks

□ The expansion of 4pt function

$$\begin{aligned} & \langle \mathcal{O}_{h_1}(z) \mathcal{O}_{h_1}(0) \mathcal{O}_{h_3}(\infty) \mathcal{O}_{h_3}(1) \rangle \\ &= \sum_p C_{11}^p C_{33}^p \underbrace{\mathcal{F}(h_1, h_3; h_p; z)}_{\text{Conformal block}} \end{aligned}$$

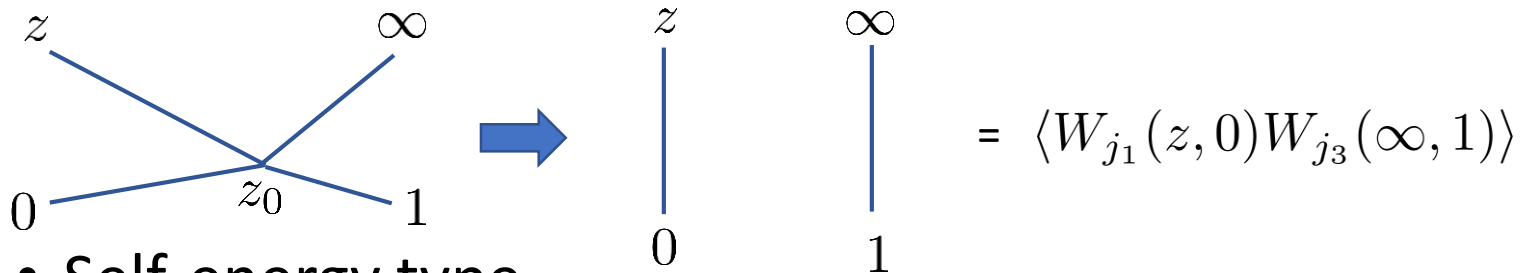
\nearrow
3pt functions



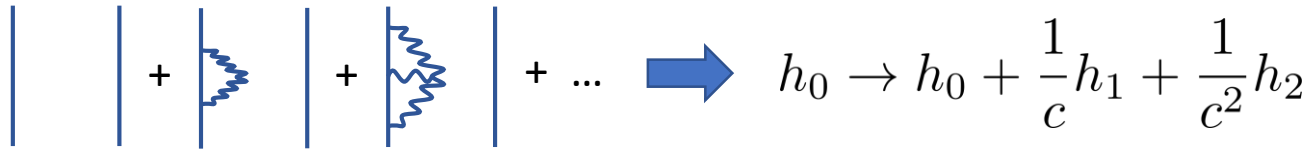
□ Conformal blocks of \mathcal{O}_{h_p}

- $p=0$: Identity conformal block
- $p \neq 0$: General conformal block

Identity block up to $1/c^2$ order



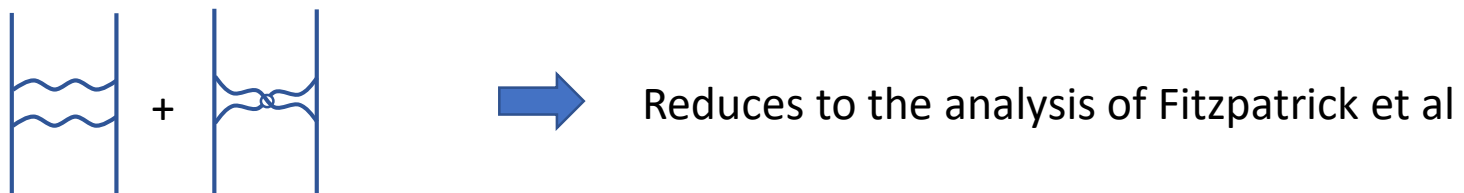
- Self-energy type



- T exchange



- TT exchange



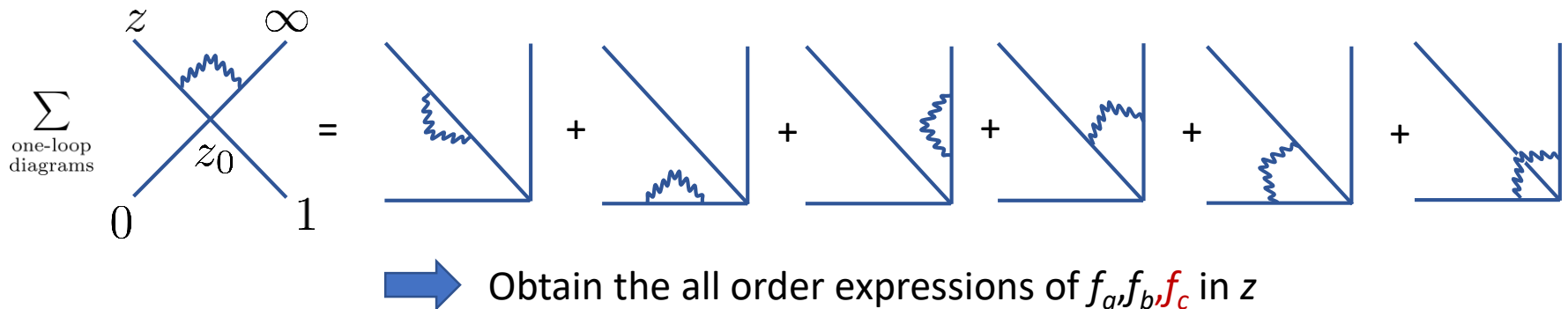
General blocks at $1/c$ order

Known results (e.g. Fitzpatrick et al '16)

$$\mathcal{F}(h_1, h_3; h_p; z) = z^{-2h_3+h_p} {}_2F_1(h_p, h_p; 2h_p; z) + \frac{1}{c} [h_1 h_3 f_a(h_p, z) + (h_1 + h_3) f_b(h_p, z) + f_c(h_p, z)] + \dots$$

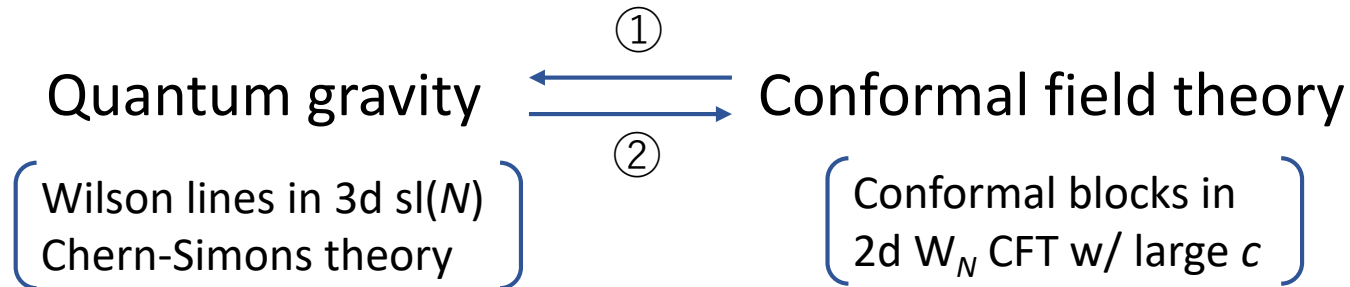
Known in all orders of z
Known in **first few orders** of z

Our computations



Conclusion

Summary



① Quantum gravity from well-defined CFT

- Fix the **renormalization prescription** from the boundary theory
→ Reproduced conformal weight at $1/c^2$ order

② CFT data of CFT from gravity

- A new systematic way for $1/c$ expansion of conformal blocks
→ Obtained **new** expressions of conformal blocks

(The analysis has been generalized from $N=2$ to $N=3$)

Future directions

□ Treatment of heavy operators

Operators with $h_{1,s} \propto \mathcal{O}(c) \iff$ Conical spaces

- Quantization of conical spaces (cf. Raeymaekers '14)
- Application to heavy-light correlators (YH-Uetoko '17)
- Extension to W_3 case is an open problem

□ Introduce supersymmetry (YH-Uetoko in progress)

- Important to see relation to superstring theory
 - $N=3$ SUSY \Leftrightarrow $\text{AdS}_3 \times \text{SU}(3)/\text{U}(1) (\times \text{SO}(5)/\text{SO}(3))$ (Creutzig-YH-Rønne'13,'14)
 - $N=4$ SUSY \Leftrightarrow $\text{AdS}_3 \times \text{S}^3 \times \text{T}^4 (\times \text{S}^3 \times \text{S}^1)$ (Gaberdiel-Gopakumar '13,'14)