

# OPE for Conformal Defects and Holography

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Based on the work (arXiv:1710.11165)  
with M. Fukuda (UT) and N. Kobayashi (Kavli-IPMU)

# Outline

- 1 Introduction
- 2 The structure of defect OPE blocks
- 3 Reconstruction of AdS scalar fields from conformal defects

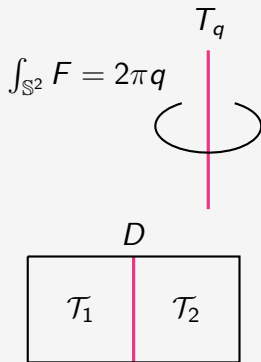
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# Introduction

**Defects** = Non-local objects in QFTs

- Defined by boundary conditions around them
- Many examples:
  - 1-dim : Line operators (Wilson-'t Hooft loops)
  - 2-dim : Surface operators
- Codim-1 : Domain walls and boundaries
- Codim-2 : Entangling surface for entanglement entropy



# Why defects?

- Characterize phases of QFTs:

$W$ : Wilson loop,  $T$ : 't Hooft loop

Confinement :  $\langle W \rangle \sim e^{-\text{Area}}$ ,  $\langle T \rangle \sim e^{-\text{Length}}$

Higgs :  $\langle W \rangle \sim e^{-\text{Length}}$ ,  $\langle T \rangle \sim e^{-\text{Area}}$

- Constrain bulk CFT data in defect CFT by conformal bootstrap

[Liendo-Rastelli-van Rees 12]

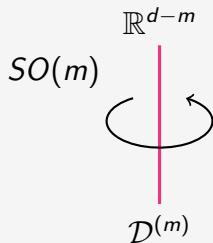
$$\sum_k \begin{array}{c} \mathcal{O}_1 \quad \mathcal{O}_2 \\ \diagdown \quad \diagup \\ \mathcal{O}_k \\ | \\ \hline \mathcal{D} \end{array} = \sum_l \begin{array}{c} \mathcal{O}_1 \quad \mathcal{O}_2 \\ | \quad | \\ \hat{\mathcal{O}}_l \\ | \\ \hline \mathcal{D} \end{array}$$

# Conformal defects

- In Euclidean  $CFT_d$ , the conformal group is  $SO(d+1, 1)$
- Codimension- $m$  conformal defects  $\mathcal{D}^{(m)}$  are either flat or spherical, preserving

$SO(d+1-m, 1)$  : conformal symmetry on defects

$SO(m)$  : rotation in the transverse direction



- Defects allow defect local operators  $\hat{\mathcal{O}}_n$

## OPE for conformal defects

There are two types of OPE in DCFT

- Bulk-to-defect OPE : [Cardy 84, McAvity-Osborn 95]

$$\begin{array}{c} \mathcal{D}^{(m)} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \circ \hat{\mathcal{O}} = \sum_n b_{\hat{\mathcal{O}} \hat{\mathcal{O}}_n}^{(m)} \begin{array}{c} \mathcal{D}^{(m)} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \hat{\mathcal{O}}_n + (\text{descendants})$$

- Defect OPE : [Berenstein-Corrado-Fischler-Maldacena 98, Gadde 16]

$$\begin{array}{c} \mathcal{D}^{(m)} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \\ \text{---} \end{array} \text{---} \hat{\mathcal{O}} = \sum_n c_{\hat{\mathcal{O}} \hat{\mathcal{O}}_n}^{(m)} \hat{\mathcal{O}}_n + (\text{descendants})$$

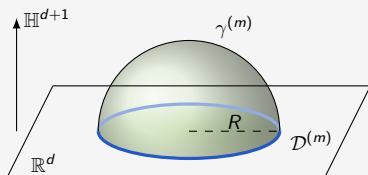
## Goal of the talk

The **defect OPE** has been less understood than the bulk-to-defect OPE

- 1 To what extent are we able to determine the structure of the defect OPE by conformal symmetry?
  - Decomposition by the irreducible representations

$$\mathcal{D}^{(m)} = \sum_{n \in \text{primaries}} \mathcal{B}^{(m)}[\mathcal{O}_n]$$

- 2 Can we probe the bulk AdS information by conformal defects on the boundary?





# Overview of our results

- 1 Give the integral representation of the **defect OPE blocks**

$$\mathcal{B}^{(m)}[\mathcal{O}_n] = \int d^d x \langle \mathcal{O}_n(x) \rangle_{\mathcal{D}^{(m)}} \tilde{\mathcal{O}}_n(x)$$

$\tilde{\mathcal{O}}$  : shadow operator with  $\tilde{\Delta} = d - \Delta$  for  $\mathcal{O}$  with  $\Delta$

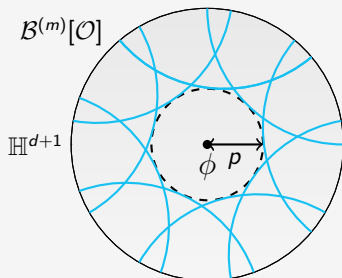
- 2 Reconstruct the AdS scalar field from the blocks

- $\hat{\phi} = \mathcal{B}^{(m)}[\mathcal{O}] :$

The Radon transform of the **AdS scalar field**  $\phi$  when  $\mathcal{O}$  scalar

- Reproduce the (Euclidean) **HKLL formula** :

$$\phi(Y) = \phi(\hat{\phi}) = \int d^d x K(Y|x) \mathcal{O}(x)$$



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## Embedding space formalism

- Uplift  $\mathbb{R}^d$  into the embedding space  $\mathbb{R}^{d+1,1}$

$$x^\mu \rightarrow X^A \quad (A = 1, \dots, d+2)$$

- Choose the  $d$ -dim hypersurface (the projective null cone)

$$X \cdot X \equiv X^A X_A = 0, \quad X^A \sim \lambda X^A \quad (\lambda \in \mathbb{R})$$

- Realize the conformal symmetry  $SO(d+1,1)$  linearly

$$X'^A = J^A_B X^B, \quad J^A_B \in so(d+1,1)$$

## Correlation functions in embedding space

Construct the invariants under  $SO(d+1, 1)$  in the embedding space

- Two-point function of scalar fields  $\Phi$  of dimension  $\Delta$ :

$$\langle \Phi(X_1) \Phi(X_2) \rangle = \frac{1}{(X_1 \cdot X_2)^\Delta}$$

For a spin  $l$  operator  $\mathcal{O}_{\Delta, l}(X)$ , define the encoding polynomial by contracting the indices with an auxiliary vector  $Z^A$ :

[Costa-Penedones-Poland-Rychkov 11]

$$\mathcal{O}_{\Delta}(X, Z) \equiv \mathcal{O}_{\Delta, A_1 \dots A_l}(X) Z^{A_1} \dots Z^{A_l}, \quad X \cdot Z = Z^2 = 0$$

- Two-point functions are uniquely determined by the invariants:

$$\langle \mathcal{O}_{\Delta, l}(X_1, Z_1) \mathcal{O}_{\Delta, l}(X_2, Z_2) \rangle = \frac{1}{(X_{12})^\Delta} \left[ \frac{(Z_1 \cdot Z_2)(X_1 \cdot X_2) - (Z_1 \cdot X_2)(Z_2 \cdot X_1)}{X_{12}} \right]^l.$$

# Conformal defects in embedding space

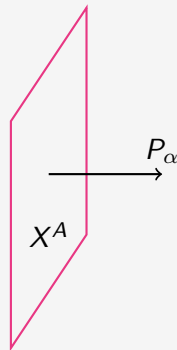
- Specify a codimension- $m$  defect by the frame vectors  $P_\alpha^A$ : [Gadde 16]

$$X \cdot P_\alpha = 0, \quad (\alpha = 1, \dots, m)$$

- Choosing an orthonormal basis,  $P_\alpha \cdot P_\beta = \delta_{\alpha\beta}$ , the defect preserves the symmetry group  $SO(d+1-m, 1) \times SO(m)$

- Correlation function should be invariant under the frame choice ( $P'_\alpha = g_\alpha^\beta P_\beta$ ,  $g \in SO(m)$ ):

$$\langle \mathcal{D}^{(m)}(P_\alpha) \mathcal{O}_\Delta(X) \rangle \propto \frac{1}{[(P_\alpha \cdot X)(P_\alpha \cdot X)]^{\Delta/2}}$$

 $\mathcal{D}^{(m)}(P_\alpha)$ 

 $X^2 = 0$

## Defect OPE blocks

We expect the defect OPE of the form

$$\mathcal{D}^{(m)}(P_\alpha) = \langle \mathcal{D}^{(m)}(P_\alpha) \rangle \left[ \sum_n c_{\mathcal{O}_n}^{(m)} R^{\Delta_n} \mathcal{O}_n(C) + (\text{descendants}) \right]$$

( $R$  : radius,  $C$  : center vector)

- The descendant terms are fixed by the primary  $\mathcal{O}_n$  and the conformal symmetry,

$$\mathcal{D}^{(m)}(P_\alpha) = \sum_n \mathcal{B}^{(m)}[P_\alpha, \mathcal{O}_n]$$

- The defect OPE block is in the irreducible representation of  $\mathcal{O}_n$ :

$$\langle \mathcal{B}^{(m)}[P_\alpha, \mathcal{O}_n] \mathcal{O}_n(X) \rangle = \langle \mathcal{D}^{(m)}(P_\alpha) \mathcal{O}_n(X) \rangle .$$

## Projectors and shadows

- Want to characterize the defect OPE blocks by their irreps
- Spectral decomposition by the irreps of the conformal group:

$$\mathbf{1} = \sum_n |\mathcal{O}_n|$$

- $|\mathcal{O}_n|$ : Projector onto the conformal multiplet of the primary  $\mathcal{O}_n$   
[Ferrara-Grillo-Parisi-Gatto 72, ..., Simmons-Duffin 12]

For a spin  $l$  operator,

$$|\mathcal{O}_{\Delta,l}| = \frac{1}{\mathcal{N}_{\Delta,l}} \int D^d X |\mathcal{O}_{\Delta,l}(X, D_Z)\rangle \langle \tilde{\mathcal{O}}_{d-\Delta,l}(X, Z)|$$

$\tilde{\mathcal{O}}_{d-\Delta,l}$ : the shadow operator of  $\mathcal{O}_{\Delta,l}$

$D_Z$ : Todorov's differential operator ( $\sim \partial_Z$ )

# Integral representation of defect OPE blocks

- Expand the defect by the projectors:

$$\begin{aligned}
 \langle \mathcal{D}^{(m)}(P_\alpha) \cdots \rangle &= \sum_{\Delta, l} \langle \mathcal{D}^{(m)}(P_\alpha) | \mathcal{O}_{\Delta, l} | \cdots \rangle + (\text{other irrep.}) \\
 &= \sum_{\Delta, l} \frac{1}{\mathcal{N}_{\Delta, l}} \int D^d X \langle \mathcal{D}^{(m)}(P_\alpha) \mathcal{O}_{\Delta, l}(X, D_Z) \rangle \langle \tilde{\mathcal{O}}_{d-\Delta, l}(X, Z) \cdots \rangle \\
 &\quad + (\text{other irrep.})
 \end{aligned}$$

- Can read off the block contribution:

The integral rep of the defect OPE block

$$\mathcal{B}^{(m)}[P_\alpha, \mathcal{O}_{\Delta, l}] = \frac{1}{\mathcal{N}_{\Delta, l}} \int D^d X \tilde{\mathcal{O}}_{d-\Delta, l}(X, D_Z) \langle \mathcal{D}^{(m)}(P_\alpha) \mathcal{O}_{\Delta, l}(X, Z) \rangle$$



## Alternative representation and monodromy condition

- The integral rep is **invariant under the exchange**  $\mathcal{O} \leftrightarrow \tilde{\mathcal{O}}$

$$\mathcal{B}^{(m)}[P_\alpha, \mathcal{O}_{\Delta,l}] = \frac{1}{\mathcal{N}_{d-\Delta,l}} \int D^d X \mathcal{O}_{\Delta,l}(X, D_Z) \langle \mathcal{D}^{(m)}(P_\alpha) \tilde{\mathcal{O}}_{d-\Delta,l}(X, Z) \rangle$$

- The block includes the **shadow contribution**  $g_{\tilde{\mathcal{O}}}$

$$\mathcal{B}^{(m)}[P_\alpha, \mathcal{O}_{\Delta,l}] = g_{\mathcal{O}} + K g_{\tilde{\mathcal{O}}}$$

- To extract the shadow contribution, impose the **monodromy condition**

$$\begin{aligned} g_{\mathcal{O}} &\rightarrow e^{2\pi i \Delta} g_{\mathcal{O}} \\ g_{\tilde{\mathcal{O}}} &\rightarrow e^{2\pi i (d-\Delta)} g_{\tilde{\mathcal{O}}} \end{aligned}$$

under the phase rotation  $P_\alpha \rightarrow e^{-2\pi i} P_\alpha$

## Constraint equations

There are two types of equations the defect OPE block satisfies

The conformal Casimir equation

$$(L^2(P_\alpha) + C_{\Delta,l}) \mathcal{B}^{(m)}[P_\alpha, \mathcal{O}_{\Delta,l}] = 0$$

- $L^2(P_\alpha) \equiv \frac{1}{2} L_{AB}(P_\alpha) L^{AB}(P_\alpha)$  : quadratic Casimir operator
- $C_{\Delta,l} = \Delta(\Delta - d) + l(l + d - 2)$  : the eigenvalue

“Trivial” equations for scalar primaries  $\mathcal{O}_\Delta$

$$C_{ABCD}(P_\alpha) \mathcal{B}^{(m)}[P_\alpha, \mathcal{O}_\Delta] = 0$$

- $C_{ABCD}(P_\alpha) \equiv \frac{1}{2} L_{A[B}(P_\alpha) L_{CD]}(P_\alpha)$ :  $\frac{d(d^2-1)(d+2)}{24}$  quadratic operators

## Moduli space of conformal defects

- The moduli space has a coset structure:

$$\mathcal{M}^{(d,m)} = \frac{SO(d+1,1)}{SO(m) \times SO(d+1-m,1)}$$

- The quadratic Casimir operator is the Laplacian on  $\mathcal{M}^{(d,m)}$

$$-L^2(P_\alpha) = \square_{\mathcal{M}^{(d,m)}}$$

- The defect OPE block is a **scalar field propagating on  $\mathcal{M}^{(d,m)}$**

Klein-Gordon equation on  $\mathcal{M}^{(d,m)}$

$$(\square_{\mathcal{M}^{(d,m)}} - M^2) \mathcal{B}^{(m)}[P_\alpha, \mathcal{O}_{\Delta,l}] = 0, \quad M^2 = \mathcal{C}_{\Delta,l}$$

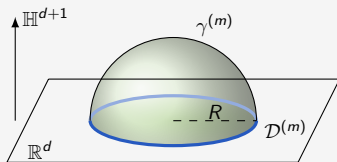
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# Conformal defects and submanifolds in AdS

- Associated to a given defect  $\mathcal{D}^{(m)}$  is a **unique** submanifold  $\gamma^{(m)}$  in AdS s.t.  $\partial\gamma^{(m)} = \mathcal{D}^{(m)}$
- Their moduli spaces are equivalent:

$$\mathcal{M}^{(d,m)} = \frac{\text{Isom}(\text{AdS}_{d+1})}{\text{Stab}(\gamma^{(m)} \in \text{AdS}_{d+1})}$$



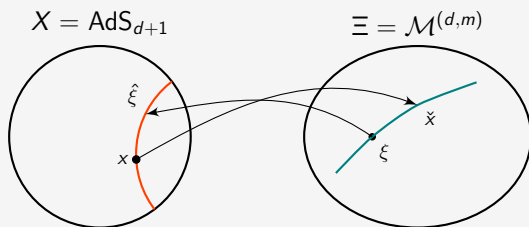
- What is the dual description of the defect OPE block in AdS?

## Moduli space and AdS

- The defect OPE block  $\mathcal{B}^{(m)}$  is a scalar field **not on  $\mathbb{H}^{d+1}$  (= Euclidean  $\text{AdS}_{d+1}$ )**, but on  $\mathcal{M}^{(d,m)}$
- The two spaces have similar coset structures:

$$\mathcal{M}^{(d,m)} = \frac{SO(d+1,1)}{SO(m) \times SO(d+1-m,1)}, \quad \mathbb{H}^{d+1} = \frac{SO(d+1,1)}{SO(d)}$$

- In this case, the Radon transform maps a function from one to the other [e.g. Helgason 10]



# Radon transform

From  $\mathbb{H}^{d+1}$  to  $\mathcal{M}^{(d,m)}$ :

$$\hat{\phi}(\xi) = \int_{x \in \xi} d\nu(x) \phi(x)$$

- $\xi$  : a codim- $m$  submanifold in  $\mathbb{H}^{d+1}$
- $\phi(x)$  : a function on  $\mathbb{H}^{d+1}$

From  $\mathcal{M}^{(d,m)}$  to  $\mathbb{H}^{d+1}$ :

$$\check{f}(x) = \int_{\xi \in \mathcal{M}^{(d,m)}} d\mu(\xi) f(\xi)$$

- $\xi$  : a codim- $m$  submanifold through  $x$
- $f(\xi)$  : a function on  $\mathcal{M}^{(d,m)}$

## Intertwining property

$$\left(\square_{\mathbb{H}^{d+1}} - M^2\right) \phi = 0 \quad \Leftrightarrow \quad \left(\square_{\mathcal{M}^{(d,m)}} - M^2\right) \hat{\phi} = 0$$

## Defect OPE blocks as Radon transformed fields

- We may regard the defect OPE block

$$\hat{\phi} = \mathcal{B}^{(m)}$$

as the Radon transform of **an AdS scalar field**  $\phi$

- The Radon transform maps  $\phi$  into the manifold with larger dimensions

$$\dim \mathcal{M}^{(d,m)} - \dim \mathbb{H}^{d+1} = (m-1)(d+1-m) \geq 0$$

Hence there should be constraints for  $\hat{\phi}$ , which are known as the range characterization [Ishikawa 97]

$$C_{ABCD} \hat{\phi} = 0$$

- This is what the block  $\mathcal{B}^{(m)}[\mathcal{O}_\Delta]$  satisfies for **a scalar primary**  $\mathcal{O}_\Delta$



## Dictionary for the Radon transform

In summary, the defect OPE block of a scalar primary is identified with an AdS scalar field

- $\mathcal{M}^{(d,m)}$

- $\mathbb{H}^{d+1}$

- scalar block:  $\mathcal{B}^{(m)}[\mathcal{O}_\Delta]$

- AdS scalar field:  $\phi$

quadratic Casimir equation

$$(\square_{\mathcal{M}^{(d,m)}} - M^2) \mathcal{B}^{(m)}[\mathcal{O}_\Delta] = 0$$

Klein-Gordon equation

$$(\square_{\mathbb{H}^{d+1}} - M^2) \phi = 0$$

constraint equations

$$C_{ABCD} \mathcal{B}^{(m)}[\mathcal{O}_\Delta] = 0$$

range characterization

$$C_{ABCD} \hat{\phi} = 0$$

## Construction of AdS scalar fields from defects

- Can construct an AdS scalar  $\phi$  from the scalar block  $\hat{\phi} = \mathcal{B}^{(m)}$  by the dual Radon transform  $f \rightarrow \check{f}$  ?

- The inversion formula allows such a construction (note  $(\hat{\phi})^\check{\phantom{\phi}} \neq \phi$ )

$$\phi = \mathcal{I} \circ \hat{\phi} = \mathcal{I} \circ \mathcal{B}^{(m)}[\mathcal{O}_\Delta]$$

$\mathcal{I}$ : an integral transformation [Helgason 10]

- Equivalent to the bulk reconstruction formula

$$\phi(Y) = \int D^d X K_\Delta(Y|X) \mathcal{O}_\Delta(X)$$

with the Euclidean version of the HKLL kernel

[Hamilton-Kabat-Lifschytz-Lowe 06]

$$(\square_{\mathbb{H}^{d+1}} - M^2) K_\Delta(Y|X) = 0, \quad M^2 = \Delta(d - \Delta)$$

# Summary

- 1 Give the integral representation of the **defect OPE blocks**

$$\mathcal{B}^{(m)}[\mathcal{O}_n] = \int d^d x \langle \mathcal{O}_n(x) \rangle_{\mathcal{D}^{(m)}} \tilde{\mathcal{O}}_n(x)$$

$\tilde{\mathcal{O}}$  : shadow operator with  $\tilde{\Delta} = d - \Delta$  for  $\mathcal{O}$  with  $\Delta$

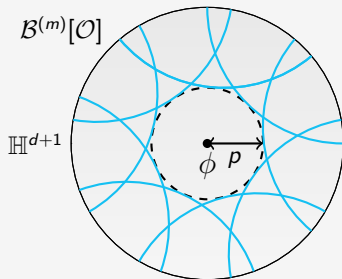
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The Radon transform of the **AdS scalar field**  $\phi$  when  $\mathcal{O}$  scalar

- Reproduce the (Euclidean) **HKLL formula** :

$$\phi(Y) = \phi(\hat{\phi}) = \int d^d x K(Y|x) \mathcal{O}(x)$$



## Future direction

- Constraints on the bulk OPE data from the block?
  - Need the relation to the bulk-to-defect OPE
- Can extend the construction to higher spin fields in AdS?
  - No known Radon transform beyond a scalar field in AdS
  - Spinning defects to incorporate spins  $\Rightarrow$  [Nozomu's talk]
- Generalized notion of entanglement entropy?
  - Codim-2 deficit angle  $\Rightarrow$  codim- $m$  deficit ??
  - Analytic continuation to Lorentzian signature?