Warped Black Holes in Lower-Spin Gravity

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[T. Azeyanagi, S. Detournay, M.R; 1801.07263]



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Introduction WAdS₃, WCFT and Lower-Spin Gravity





AdS₃ deformation:

$$ds^2_{\mathrm{WAdS}_3} = ds^2_{\mathrm{AdS}_3} - 2H\xi \otimes \xi$$

Introduction WAdS₃, WCFT and Lower-Spin Gravity





 $\mathsf{NHEK} \to \mathsf{global}\ \mathfrak{sl}(2,\mathbb{R}) \oplus \mathfrak{u}(1)$ isometry.



NHEK \rightarrow global $\mathfrak{sl}(2,\mathbb{R}) \oplus \mathfrak{u}(1)$ isometry.

At fixed polar angle \rightarrow (selfdual spacelike) WAdS_3.



 $\mathsf{NHEK} \to \mathsf{global}\ \mathfrak{sl}(2,\mathbb{R}) \oplus \mathfrak{u}(1)$ isometry.

At fixed polar angle \rightarrow (selfdual spacelike) WAdS_3.

 \rightarrow study NHEK.

Introduction WAdS₃, WCFT and Lower-Spin Gravity





QFT with global $\mathfrak{sl}(2,\mathbb{R})\oplus\mathfrak{u}(1)$ symmetries.



QFT with global $\mathfrak{sl}(2,\mathbb{R})\oplus\mathfrak{u}(1)$ symmetries.

Has ∞ conserved charges but not Lorentz invariant.





Possible WAdS₃ dual.



Possible WAdS₃ dual.

 \rightarrow better understanding of holography in a realistic setup!



What?

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 $\mathfrak{sl}(2,\mathbb{R})\oplus\mathfrak{u}(1)$ Chern-Simons theory.





Very simple model of WAdS₃ without local DOF.



Black Holes?



- Black Holes?
- ► Higher-Spin WAdS₃ black holes?



 \blacktriangleright Boundary conditions \rightarrow asymptotic symmetries.

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Determine thermal entropy.

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- Consistency checks:

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 - Compare with WCFT results.

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 - Holographic entanglement entropy.

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- Consistency checks:
 - Compare with WCFT results.
 - Holographic entanglement entropy.
 - Metric interpretation.

WAdS₃ BHs in Lower-Spin Gravity Boundary Conditions and Asymptotic Symmetries

$$I_{\rm CS} = \frac{k}{4\pi} \int_{\mathcal{M}} \langle \mathcal{A} \wedge \mathsf{d}\mathcal{A} + \frac{2}{3}\mathcal{A} \wedge \mathcal{A} \wedge \mathcal{A} \rangle + \frac{\kappa}{8\pi} \int_{\mathcal{M}} \langle \mathcal{C} \wedge \mathsf{d}\mathcal{C} \rangle$$

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Boundary Conditions

$$\mathcal{A}(\rho, t, \varphi) = b^{-1}(\rho) \left[a(t, \varphi) + d \right] b(\rho)$$

$$\mathcal{C}(\rho, t, \varphi) = c(t, \varphi)$$

$$a_{\varphi} = \mathrm{L}_{1} - \left(rac{2\pi}{k}\left(\mathcal{L} - rac{2\pi}{\kappa}\mathcal{K}^{2}
ight)
ight)\mathrm{L}_{-1} \qquad a_{t} = 0$$
 $c_{\varphi} = rac{4\pi}{\kappa}\mathcal{K}\mathrm{S} \qquad c_{t} = \mathrm{S}$

WAdS₃ BHs in Lower-Spin Gravity Boundary Conditions and Asymptotic Symmetries

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Asymptotic Symmetries

$$\begin{split} [L_{n}, L_{m}] &= (n-m) L_{n+m} + \frac{c}{12} n (n^{2}-1) \delta_{n+m,0} \\ [L_{n}, K_{m}] &= -m K_{n+m} \\ [K_{n}, K_{m}] &= \frac{\kappa}{2} n \delta_{n+m,0} \end{split}$$

[Hofman, Rollier '14]

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▶ Determine mass *M* and angular momentum *J*.

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- Relate β , $\Omega \Leftrightarrow M$, J.

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•
$$\delta S_{\mathrm{Th}} = \beta \left(\delta M - \Omega \delta J \right) \Rightarrow S_{\mathrm{Th}}$$
.

ULB

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Assumption

$$\varepsilon = \xi^{\mu} \mathcal{A}_{\mu}, \qquad \bar{\varepsilon} = \xi^{\mu} \mathcal{C}_{\mu},$$

Determine mass M and angular momentum J.

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Assumption

$$\varepsilon = \xi^{\mu} \mathcal{A}_{\mu}, \qquad \bar{\varepsilon} = \xi^{\mu} \mathcal{C}_{\mu},$$

$$\begin{split} \delta \boldsymbol{M} &:= \delta \boldsymbol{Q}[\boldsymbol{\varepsilon}\big|_{\partial_t}] + \delta \boldsymbol{Q}[\boldsymbol{\bar{\varepsilon}}\big|_{\partial_t}] = 2\pi\delta\mathcal{K} \\ \delta \boldsymbol{J} &:= \delta \boldsymbol{Q}[\boldsymbol{\varepsilon}\big|_{-\partial_{\boldsymbol{\omega}}}] + \delta \boldsymbol{Q}[\boldsymbol{\bar{\varepsilon}}\big|_{-\partial_{\boldsymbol{\omega}}}] = -2\pi\delta\mathcal{L} \end{split}$$

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$$\delta S_{\rm Th} = \beta \left(\delta M - \Omega \delta J \right) = k \langle h \, \delta a_{\varphi} \rangle + \frac{\kappa}{2} \langle \bar{h} \, \delta c_{\varphi} \rangle$$

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where

$$h = \frac{\beta}{2\pi} \oint d\varphi \left(a_t + \Omega a_{\varphi} \right)$$
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Assumption

ULB

 $\operatorname{Eigen}[h] = \operatorname{Eigen}[2\pi L_0] \qquad \operatorname{Eigen}[\bar{h}] = \operatorname{Eigen}[2\pi\gamma S]$

$$\delta S_{
m Th} = eta \left(\delta M - \Omega \delta J
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Assumption
$$Eigen[h] = Eigen[2\pi L_0]$$
 $Eigen[\bar{h}] = Eigen[2\pi\gamma S]$

$$\beta = 2\pi \left(\gamma - \frac{M}{\kappa \sqrt{\frac{1}{k} \left(-J - \frac{M^2}{\kappa} \right)}} \right) \qquad \Omega = \frac{1}{2\gamma \sqrt{\frac{1}{k} \left(-J - \frac{M^2}{\kappa} \right)} - \frac{2M}{\kappa}}$$

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$$\delta S_{\mathrm{Th}} = \beta \left(\delta M - \Omega \delta J \right) \Rightarrow S_{\mathrm{Th}}.$$

Thermal Entropy

$$S_{\mathrm{Th}} = 2\pi \left(M\gamma + \sqrt{rac{c}{6} \left(-J - rac{M^2}{\kappa}
ight)}
ight)$$

WAdS₃ BHs in Lower-Spin Gravity

$$S_{\mathrm{Th}} = -rac{4\pi i M M^{\mathrm{v}}}{\kappa} + 4\pi \sqrt{-\left(-J^{\mathrm{v}}-rac{\left(M^{\mathrm{v}}
ight)^{2}}{\kappa}
ight)\left(-J-rac{M^{2}}{\kappa}
ight)}$$

ULB

[Detournay, Hartman, Hofman '12]

WAdS₃ BHs in Lower-Spin Gravity

$$S_{\rm Th} = -\frac{4\pi i M M^{\rm v}}{\kappa} + 4\pi \sqrt{-\left(-J^{\rm v} - \frac{(M^{\rm v})^2}{\kappa}\right)\left(-J - \frac{M^2}{\kappa}\right)}$$
[Detournay, Hartman, Hofman '12]
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(Warped) vacuum:
$$e^{\oint d\varphi \, a_{\varphi}} = -1 \qquad e^{\oint d\varphi \, c_{\varphi}} = e^{2\pi i \gamma}$$

WAdS₃ BHs in Lower-Spin Gravity

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ULB

[Detournay, Hartman, Hofman '12]

Assumption

(Warped) vacuum:

$$\mathrm{e}^{\oint \mathsf{d}arphi \, a_arphi} = - \mathbb{1} \qquad \mathrm{e}^{\oint \mathsf{d}arphi \, c_arphi} = \mathrm{e}^{2\pi i \gamma}$$

$$J^{\mathrm{v}}=rac{c}{24}+rac{\gamma^{2}\kappa}{4}\qquad M^{\mathrm{v}}=rac{i\kappa\gamma}{2}$$



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6



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[Castro, Hofman, Iqbal '15]

$$\mathcal{S}_{\mathrm{EE}} = -\log\left[\mathcal{W}^{\mathfrak{sl}(2,\mathbb{R})}_{\mathcal{R}}(\mathit{C};\mathcal{A})\right] - \log\left[\mathcal{W}^{\mathfrak{u}(1)}_{\mathcal{R}}(\mathit{C};\mathcal{C})\right]$$

ULB

[Castro, Hofman, Iqbal '15]

Holographic Entanglement Entropy

$$S_{\rm EE} = \gamma \left(\frac{\kappa}{2}\Delta t + M\Delta\varphi\right) + \frac{c}{6}\log\left[\frac{\beta_{\varphi}}{\pi\epsilon}\sinh\left[\frac{\pi\Delta\varphi}{\beta_{\varphi}}\right]\right]$$

where

$$\beta_{\varphi} = \frac{\pi}{\sqrt{\frac{6}{c}\left(-J - \frac{M^2}{\kappa}\right)}}$$

Holographic Entanglement Entropy
$$S_{\rm EE} = \gamma \left(\frac{\kappa}{2}\Delta t + M\Delta\varphi\right) + \frac{c}{6}\log\left[\frac{\beta_{\varphi}}{\pi\epsilon}\sinh\left[\frac{\pi\Delta\varphi}{\beta_{\varphi}}\right]\right]$$

6

$$S_{\rm EE} = iM^{\rm v} \left(-\Delta t + \frac{\beta - \delta}{\beta_{\varphi}} \Delta \varphi \right) + \left(i\frac{\delta}{\pi}M^{\rm v} - 4J^{\rm v} \right) \log \left[\frac{\beta_{\varphi}}{\pi\epsilon} \sinh \left[\frac{\pi \Delta \varphi}{\beta_{\varphi}} \right] \right]$$

[Song, Wen, Xu '16]

Holographic Entanglement Entropy
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6

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[Song, Wen, Xu '16]

$$\Rightarrow \delta = 2\pi\gamma$$



$$S_{
m EE} \stackrel{rac{\Delta arphi}{eta arphi} \gg 1}{pprox} rac{S_{
m Th}}{2\pi} \Delta arphi$$



• Warped geometry \Leftrightarrow lower-spin gravity.

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- $\blacktriangleright Warped geometry \Leftrightarrow \text{lower-spin gravity.}$
- Determine geometric variables:

$$\mathcal{B}' = \mathcal{B}'(\mathcal{A}, \mathcal{C}; \alpha, \mathfrak{b}, \mathfrak{c})$$



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Metric:

$$\mathrm{d} s^2 = g_{\mu
u} \mathrm{d} x^\mu \mathrm{d} x^
u = \langle \mathcal{B}, \mathcal{B}
angle$$

- ► Warped geometry ⇔ lower-spin gravity.
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ULB

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u = \langle {\cal B}, {\cal B}
angle$$

	WAdS ₃ BH Parameters			
$\mathfrak{b}^2 =$	$\frac{\nu^2}{2\ell^2}$ c	$=\frac{\nu^2+3}{2\ell^2}$	$\alpha = -$	$\frac{16}{k^2k\ell^2}$

- ► Warped geometry ⇔ lower-spin gravity.
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ULB

Metric:

$${\sf d} {\it s}^2 = {\it g}_{\mu
u} {\it d} {\it x}^\mu {\it d} {\it x}^
u = \langle {\cal B}, {\cal B}
angle$$



 \Rightarrow spacelike (stretched) warped AdS₃ black hole!



Killing vectors:

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- Killing vectors:
 - Globally well defined Killing vectors \checkmark

Initial Assumption

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(Warped) vacuum:

$$e^{\oint \mathrm{d} arphi \, a_arphi} = -1 \qquad e^{\oint \mathrm{d} arphi \, c_arphi} = e^{2\pi i \gamma}$$

Killing vectors:

► Globally well defined Killing vectors √

•
$$J^{\mathrm{v}} = \frac{c}{24} + \frac{\gamma^2 \kappa}{4}$$
 $M^{\mathrm{v}} = \frac{i\kappa\gamma}{2} \checkmark$

Initial Assumption

ULB

 $\operatorname{Eigen}[h] = \operatorname{Eigen}[2\pi L_0] \qquad \operatorname{Eigen}[\bar{h}] = \operatorname{Eigen}[2\pi\gamma S]$

- Killing vectors:
 - Globally well defined Killing vectors

•
$$J^{\mathrm{v}} = \frac{c}{24} + \frac{\gamma^2 \kappa}{4}$$
 $M^{\mathrm{v}} = \frac{i\kappa\gamma}{2} \checkmark$

• Thermodynamics: $\beta(M, J), \Omega(M, J) \checkmark$

- Killing vectors:
 - Globally well defined Killing vectors \checkmark

ULB

•
$$J^{\mathrm{v}} = \frac{c}{24} + \frac{\gamma^2 \kappa}{4}$$
 $M^{\mathrm{v}} = \frac{i\kappa\gamma}{2} \checkmark$

• Thermodynamics: $\beta(M, J), \Omega(M, J) \checkmark$

$$\blacktriangleright \ \gamma = \frac{\mathfrak{b}}{2}\sqrt{\frac{k\alpha}{2}} = \frac{2\nu}{\nu^2 + 3}$$

Conclusion

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 \blacktriangleright Warped BHs and lower-spin gravity . \checkmark

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Conclusion

- \blacktriangleright Warped BHs and lower-spin gravity . \checkmark
- Easiest toy model for WCFTS with T > 0.

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- Straightforward extension to higher-spins.

Conclusion

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- Warped BHs and lower-spin gravity .
- Easiest toy model for WCFTS with T > 0.
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Outlook

Higher-spin warped BHs.

Conclusion

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- \blacktriangleright Warped BHs and lower-spin gravity . \checkmark
- Easiest toy model for WCFTS with T > 0.
- Straightforward extension to higher-spins.

Outlook

- Higher-spin warped BHs.
- Full partition function as in [lizuka, Tanaka, Terashima '15].

Conclusion

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- ► Warped BHs and lower-spin gravity . √
- ► Easiest toy model for WCFTS with *T* > 0.
- Straightforward extension to higher-spins.

Outlook

- Higher-spin warped BHs.
- Full partition function as in [lizuka, Tanaka, Terashima '15].
- Near horizon soft hair?

Thank you for your attention!