# Flow equation, conformal symmetry and AdS geometries

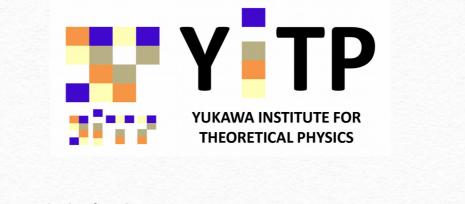
YITP Workshop

Holography, Quantum Entanglement and Higher Spin Gravity II

16. Mar. 2018 @ YITP

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Yukawa Institute for Theoretical Physics

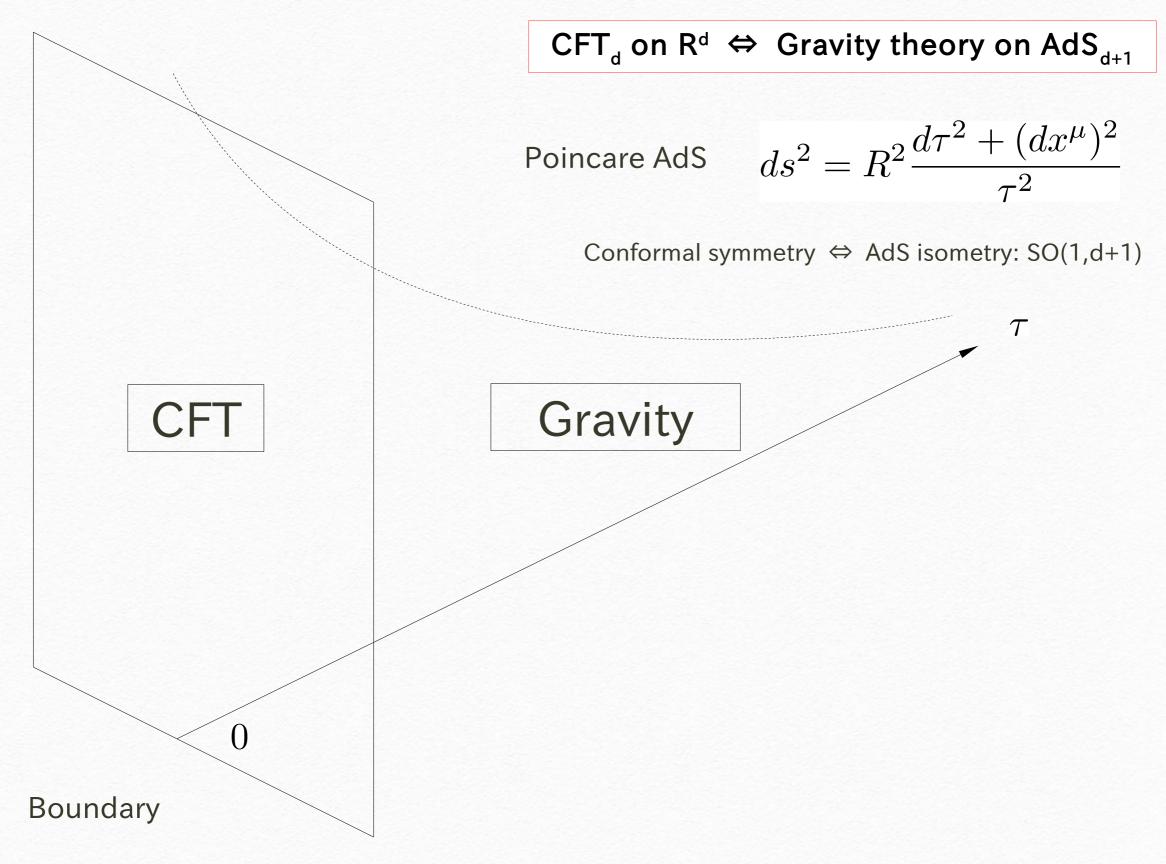


Ref.

S.Aoki-SY ArXiv:1707.03982 PTEP (2018) 031B01 S.Aoki-SY ArXiv:1709.07281

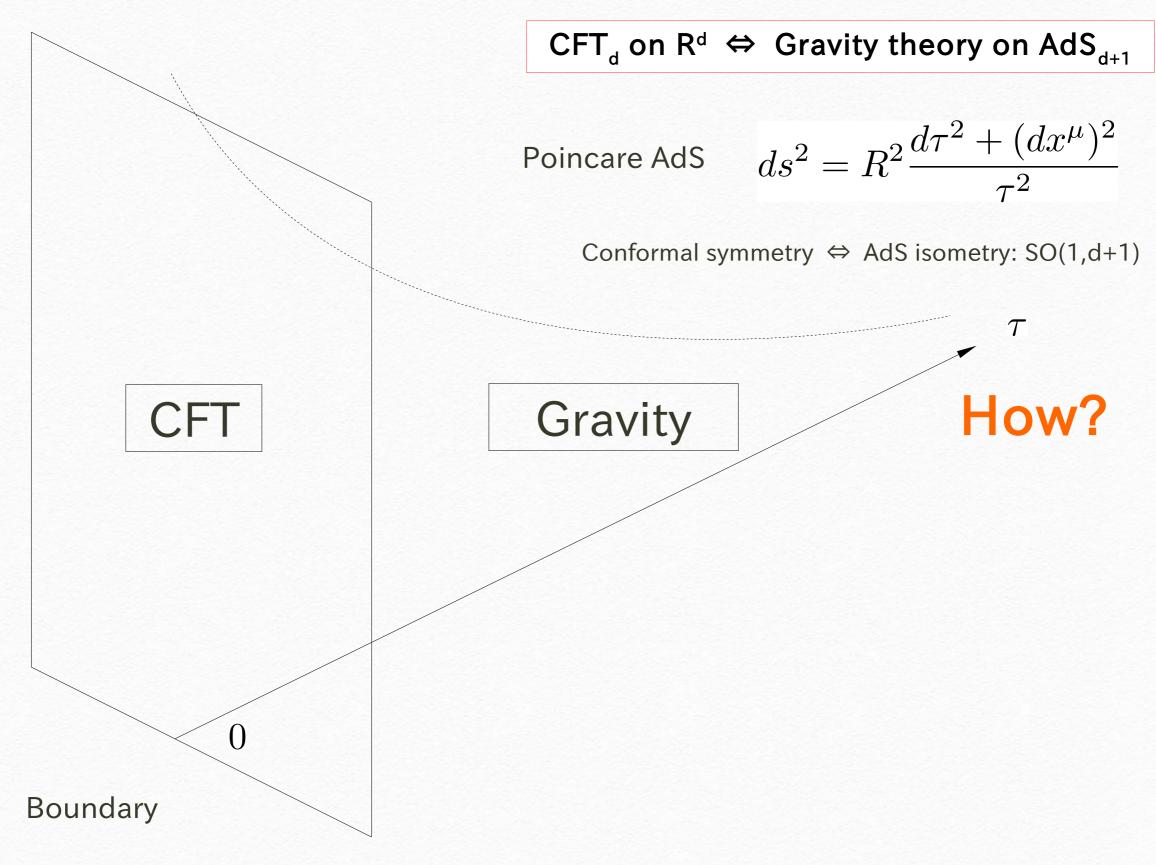
# Holography and AdS/CFT

[Maldacena '97]



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# AdS radial direction from CFT

#### 1. Relevant RG flow

Construction of gravity solutions corresponding to UV and IR CFTs in the asymptotic regions. [Girardello-Petrini-Porrati-Zaffaroni '98] [Distler-Zamora '98] [de-Boer-Verlinde-Verlinde '99] [Skenderis '00]

### 2. Wilsonian RG flow

The Wilsonian cut-off will correspond to sharp cut-off at the AdS radial direction [Heemskerk-Polchinski '10]

#### 3. Stochastic quantization

Euclidean path integral  $\equiv$  Equilibrium limit of statistical mechanical system coupled to a heat bath. [Lifshytz-Periwal '00]

#### 4. Entanglement enetropy

continuous multi-scale entanglement renormalization ansatz (cMERA) [Swingle '09]...

### 5. bilocal field

Relative coordinate of bi-local field in vector models [Das-Jevicki '03]

### 6. Flow equation

Smearing operators so as to resolve a UV singularity in the coincidence limit

[Aoki-Kikuchi-Onogi '15] [Aoki-Balog-Onogi-Weisz '16,'17]

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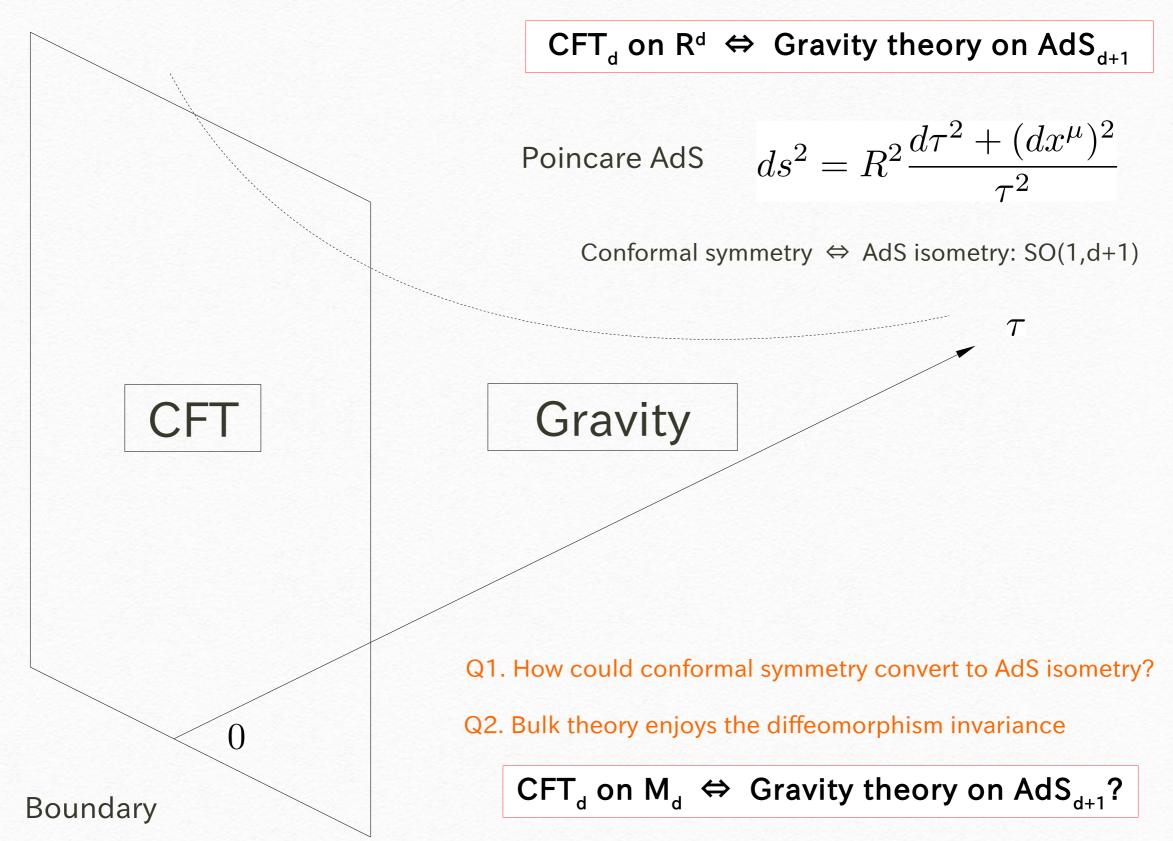
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Flow equation

# (Gradient) flow equation

smear operators to resolve a UV singularity in the coincidence limit

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"Regularization"

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"Regularization"

Consider CFT with a real scalar primary operator O in d dimensions

Free flow equation

 $\frac{\partial O(x;t)}{\partial t} = \partial^2 O(x;t) \qquad O(x;0) = O(x)$ 

t: flow parameter, O(x;t): flowed operator

# (Gradient) flow equation

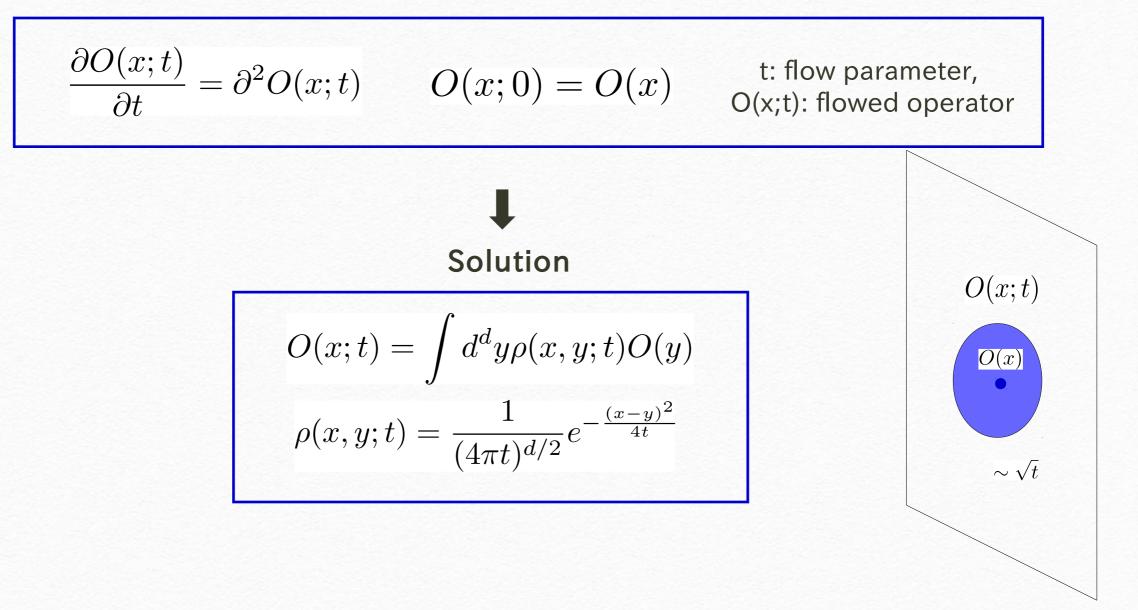
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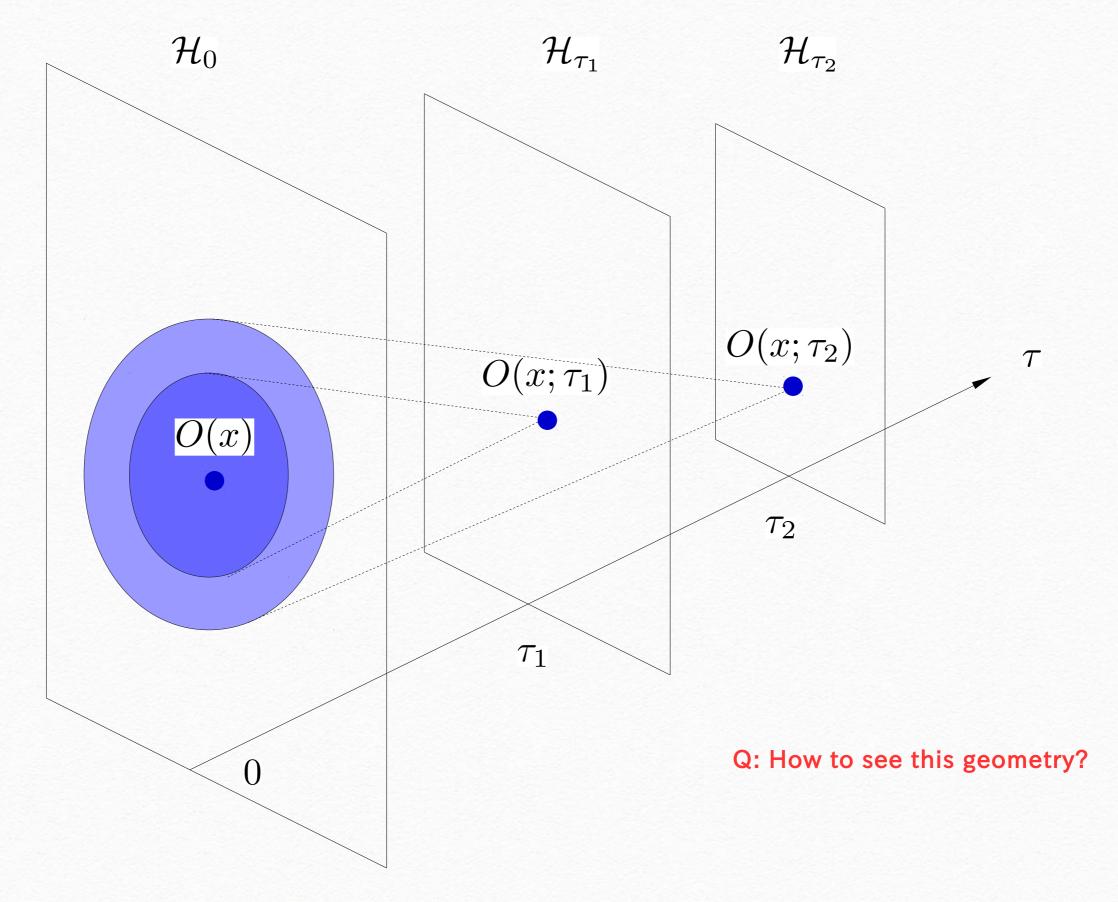
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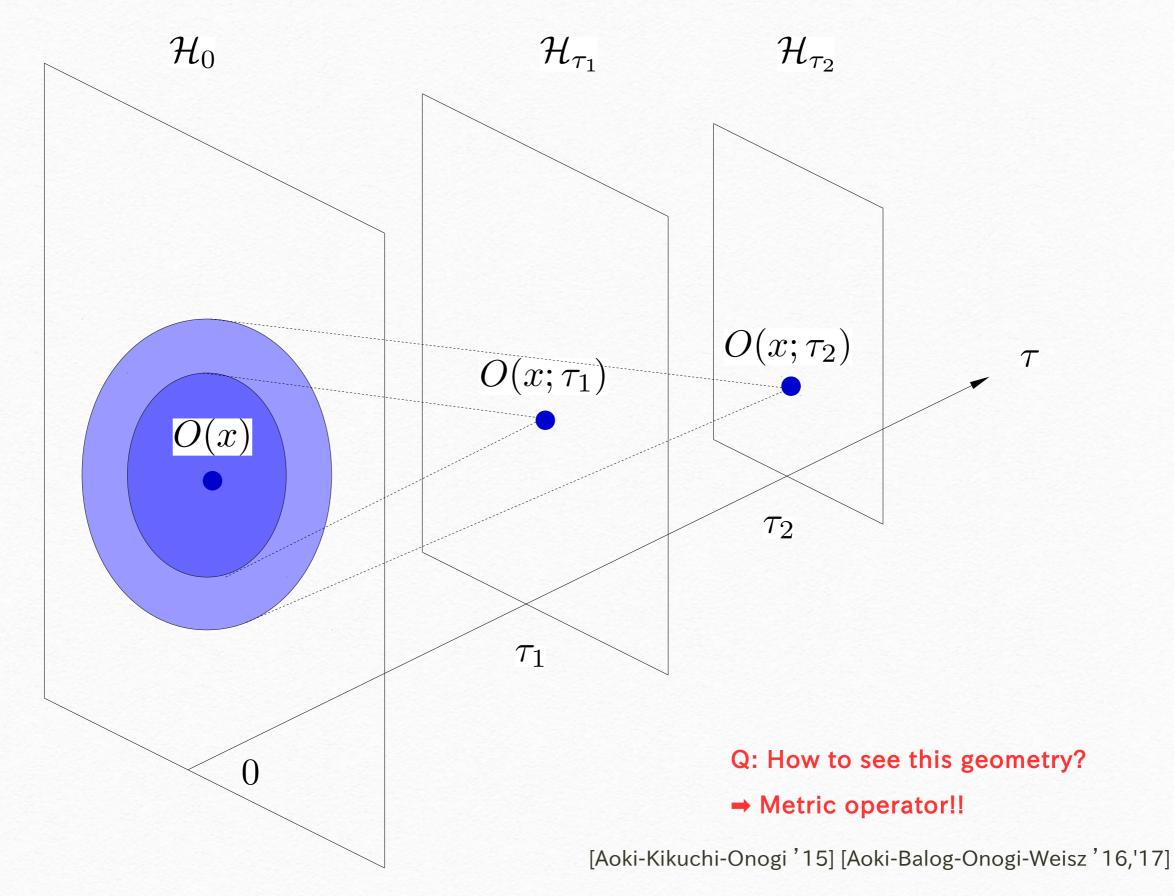
Free flow equation



# Sketch of smearing and extra direction



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# Metric operator & induced metric

[Aoki-Kikuchi-Onogi '15] [Aoki-Balog-Onogi-Weisz '16,'17]

### Metric operator and induced metric

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<u>Def.</u> (Dimensionless normalized operator)

$$\sigma(x;t) := \frac{O(x;t)}{\sqrt{\langle O(x;t)^2 \rangle_{CFT}}}$$

NOTE:  $\langle \sigma(x;t)\sigma(x;t)\rangle_{CFT}=1$ 

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<u>Def.</u> (Metric operator and induced metric)

$$\hat{g}_{MN}(x;t) := R^2 \frac{\partial \sigma(x;t)}{\partial z^M} \frac{\partial \sigma(x;t)}{\partial z^N} \qquad g_{MN}(z) := \langle \hat{g}_{MN}(x;t) \rangle_{CFT}$$

R: constant of length dimension  $z^M = (x^{\mu}, \tau)$  with  $\tau = \sqrt{2dt}$ 

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<u>Result 1</u> Induced metrics with the free flow equation become the AdS metric for a couple of vector models in the free or critical limit.

[Aoki-Kikuchi-Onogi '15] [Aoki-Balog-Onogi-Weisz '16,'17]

<u>Result 2</u> An induced metric becomes quantum information metric for a general QFT. [Aoki-SY '17]

<u>Result 3</u> An induced metric defined in this way becomes AdS for a general CFT. [Aoki-SY '17]

[S.Aoki-SY '17] PTEP (2018) 031B01

(Bure metric for a density matrix) Def.

$$D(\rho, \rho + d\rho)^2 = \frac{1}{2} \operatorname{tr}(d\rho \, G)$$

ho : density matrix G: hermitian 1 form operator satisfying  $ho \, G + G \, 
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ho$ 

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In CFT setup, the density matrix for a state localized at  $z = (x, \tau)$  is

 $\rho_z = |\sigma(x;t)\rangle \langle \sigma(x;t)|, \quad |\sigma(x;t)\rangle := \sigma(x;t)|0\rangle, \quad \langle \sigma(x;t)| := \langle 0|\sigma(x;t)| \\ z^M = (x^\mu, \tau) \text{ with } \tau = \sqrt{2dt}$ 

Def. (Inner product)

 $\langle \sigma(x;t) | \sigma(w;s) \rangle := \langle \sigma(x;t) \sigma(w;s) \rangle_{CFT}$ 

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$$D(\rho_z, (\rho + d\rho)_z)^2 = \frac{1}{2} \operatorname{tr}(d\rho_z d\rho_z) = g_{MN}(z) dz^M dz^N,$$

cf. Fischer information metric  $|\langle \Psi_{\lambda} | \Psi_{\lambda+\delta\lambda} \rangle|^2 = 1 - 2G_{\lambda\lambda}\delta\lambda^2 + \cdots$ 

 $|\Psi_{\lambda}\rangle$ : Vacuum state for H<sub>0</sub> +  $\lambda$  V

### [S.Aoki-SY '17] PTEP (2018) 031B01

cf. [Jevicki-Kazama-Yoneya '98]

[Das-Jevicki '03]

[Aoki-SY '17]

Conformal transformation:

$$\delta x^{\mu} = a^{\mu} + \omega^{\mu}{}_{\nu}x^{\nu} + \lambda x^{\mu} + b^{\mu}x^{2} - 2x^{\mu}(b_{\nu}x^{\nu}),$$
  
$$\delta^{\text{conf}}O(x) = -\delta x^{\mu}\partial_{\mu}O(x) - \frac{\Delta}{d}(\partial_{\mu}\delta x^{\mu})O(x)$$

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 $\delta^{\rm conf}\sigma(x;t) = -\left\{2\lambda - 4(b_{\mu}x^{\mu})\right\}t\partial_{t}\sigma(x;t) - \left\{\delta x^{\mu} + 2t(d-2-\Delta)b^{\mu}\right\}\partial_{\mu}\sigma(x;t) + 4t^{2}b^{\mu}\partial_{\mu}\partial_{t}\sigma(x;t)$ 

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 $\exists$  Higher derivative term  $\Rightarrow$  Naively does not work for special conformal transformation...

cf. [Das-Jevicki '03] or Das's talk in 'Strings and Fields 2017'

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∃ Higher derivative term → Naively does not work for special conformal transformation...

Let us decompose this into the following.

$$\begin{split} \delta^{\text{conf}}\sigma(x;t) &= \delta^{\text{diff}}\sigma(x;t) + \delta^{\text{extra}}\sigma(x;t), \\ \delta^{\text{diff}}\sigma(x;t) &= -\left(\bar{\delta}t\partial_t + \bar{\delta}x^{\mu}\partial_{\mu}\right)\sigma(x;t), \quad \delta^{\text{extra}}\sigma(x;t) = 4t^2b^{\nu}\partial_{\nu}(\partial_t + \frac{\Delta+2}{2t})\sigma(x;t), \\ \text{where} \quad \bar{\delta}x^{\mu} &= \delta x^{\mu} + 2dtb^{\mu}, \\ \bar{\delta}t &= (2\lambda - 4(b_{\mu}x^{\mu}))t \end{split}$$

[Aoki-SY '17]

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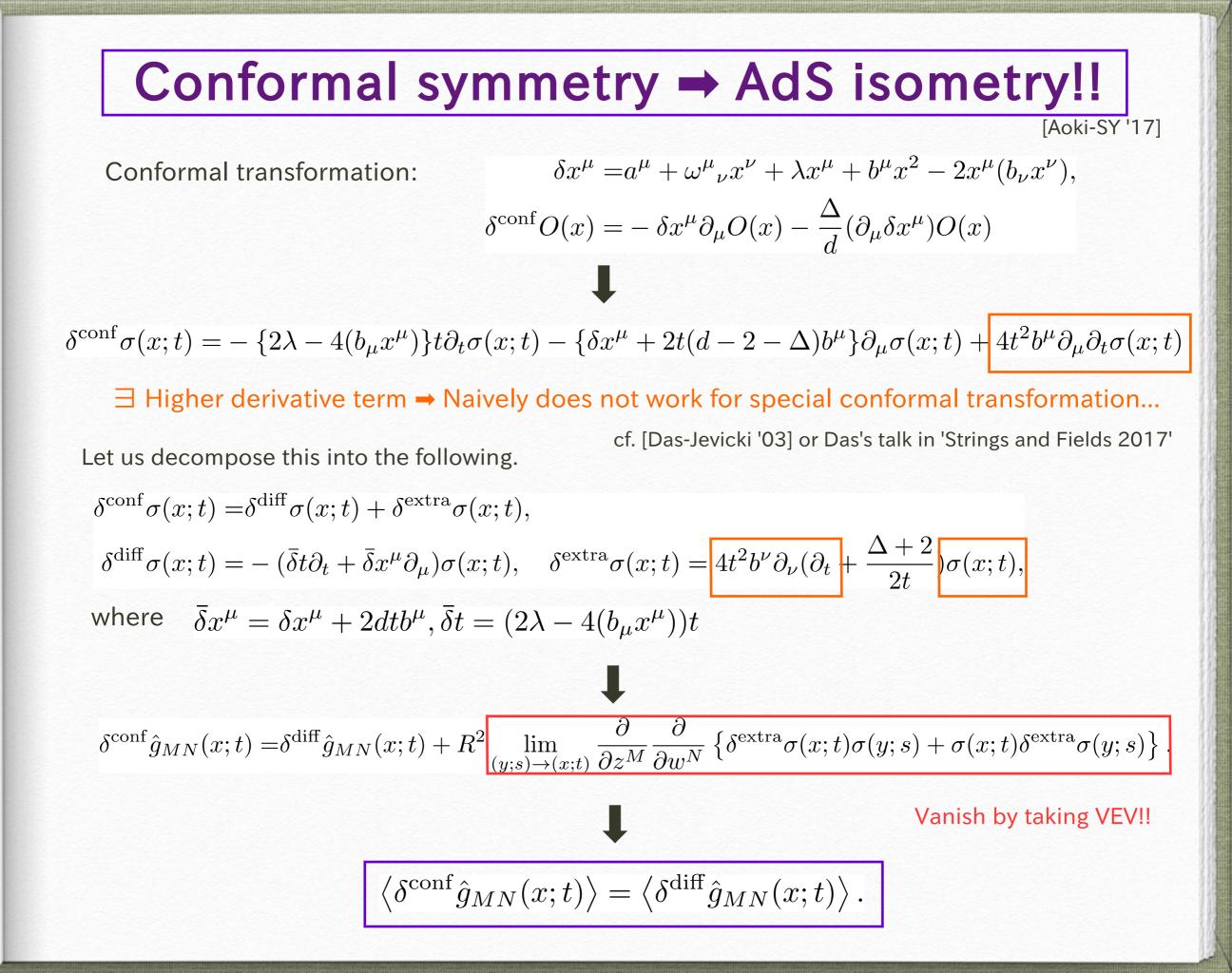
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# Q2: Generalization to non-trivial curved boundary?

[S.Aoki-SY '17]

A curved manifold needs to admit CFT to live. → Restrict ourselves to a conformally flat manifold. We need to construct a flow equation associated with the curved manifold.

$$\frac{\partial O(x;t)}{\partial t} = \partial^2 O(x;t) \quad \to \quad ???$$

A curved manifold needs to admit CFT to live. → Restrict ourselves to a conforamlly flat manifold. We need to construct a flow equation associated with the curved manifold.

### "Canonical" free flow equation (Primary flow equation)

$$\frac{\partial}{\partial \tilde{t}} O(\Omega_x; \tilde{t}) = g^{-\frac{1}{d}}(x) g^{-\frac{\Delta}{2d}}(x) \partial^2 \cdot g^{\frac{\Delta}{2d}}(x) O(\Omega_x; \tilde{t}), \qquad O(\Omega_x; 0) = O(\Omega_x)$$

 $g^{rac{1}{d}}(x)$  the conformal factor

$$\tilde{t} = g^{\frac{1}{d}}(x)t$$

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The induced metric:  $\tilde{g}_{\tilde{\tau}\tilde{\tau}}(z) = R^2 \frac{\Delta}{\tilde{\tau}^2}, \\ \tilde{g}_{\tilde{\tau}\mu}(z) = g_{\mu\tilde{\tau}}(z) = -R^2 \frac{\Delta}{\tilde{\tau}} \frac{\partial}{\partial x^{\mu}} \log\{g^{\frac{1}{2d}}(x)\}, \\ \tilde{g}_{\mu\nu}(z) = R^2 \Delta \left[\frac{\partial}{\partial x^{\mu}} \log\{g^{\frac{1}{2d}}(x)\} \frac{\partial}{\partial x^{\nu}} \log\{g^{\frac{1}{2d}}(x)\} + \frac{\delta_{\mu\nu}g^{\frac{1}{d}}(x)}{\tilde{\tau}^2}\right], \\ \text{This is the (local) AdS metric whose radius is Rv } \Delta \text{ with the boundary } M_{\mathtt{d}} \end{split}$ 

A curved manifold needs to admit CFT to live. → Restrict ourselves to a conformally flat manifold. We need to construct a flow equation associated with the curved manifold.

### "Canonical" free flow equation (Primary flow equation)

Summary

• An induced metric corresponds to the quantum information metric.

- For a general CFT, the induced metric for the free flow always becomes AdS.
- Conformal transf. converts to AdS isometry after quantum averaging.
- We have constructed a canonical flow equation for a primary scalar operator on a conformally flat manifold, called the primary flow equation.
- The induced metric associated with the primary flow equation becomes AdS with the conformally flat boundary.
- AdS with the conformally flat boundary is obtained from the usual Poincare AdS by a bulk finite diffeomorphism.

# **Future works**

How to encode dynamics beyond geometry? For excited states?

working in progress [Aoki-Balog-SY]

- How to reconstruct bulk operator?
- Canonical choice of a normalized field?

Spin 1,2 field? Fermion?

- de Sitter construction? Application to the real world?

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Thank you!!

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