

125GeV Higgs Boson

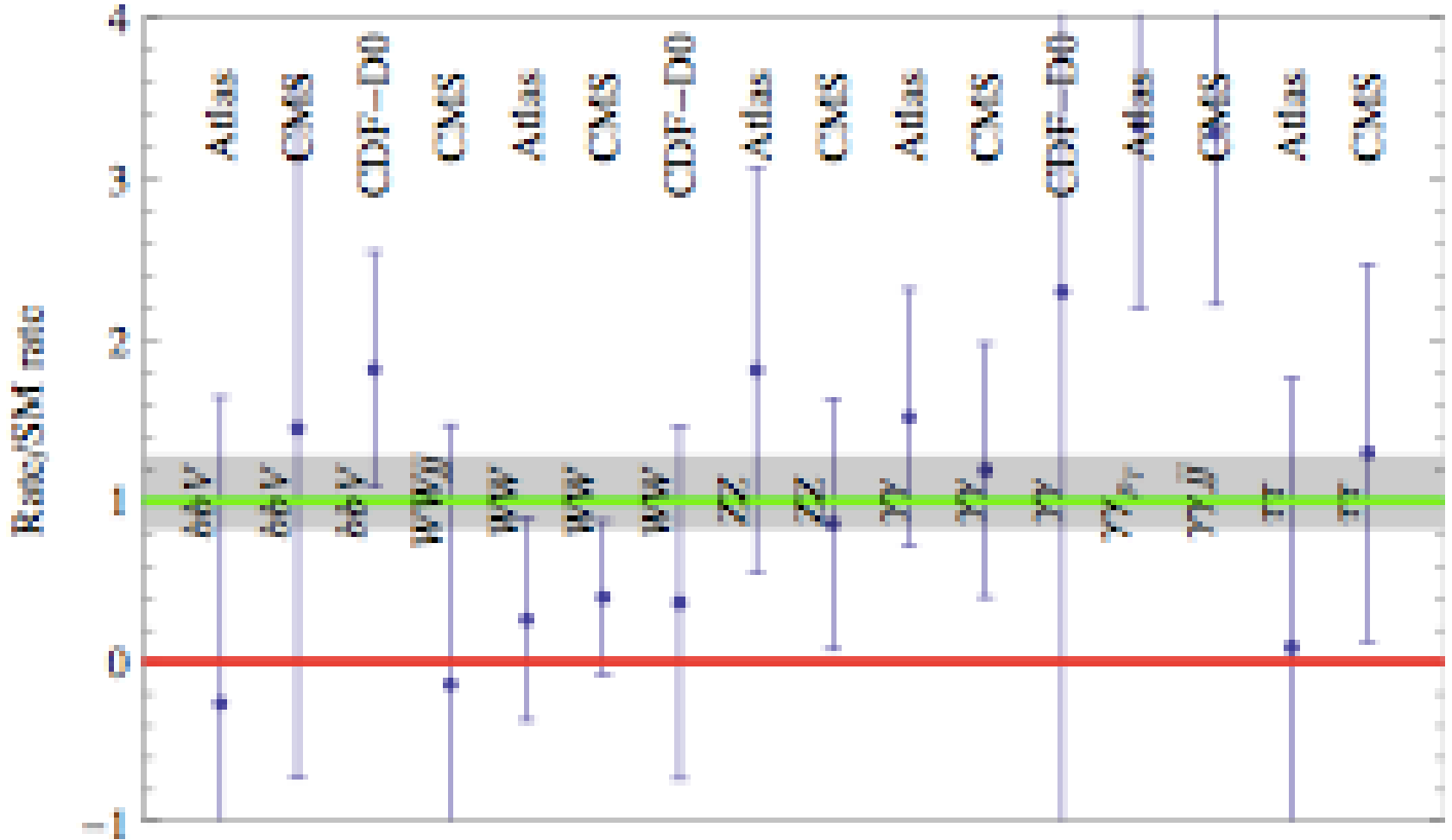
--Cosmological viewpoint--

Harigaya, Nagata, Hisano, Kikukawa,
Takayama, Tanabashi, Yamanaka

ルール

- 125GeVヒッグスを宇宙論の観点で考える
「????」という問題を真剣に考えると次に見えるのは「????」実験で「????」である。
- Higgs Portal Dark Matter (Nagata)
- Higgs Portal(?) Dark Radiation (Harigaya)
- Electroweak Baryogenesis (Hisano)
- Higgs Inflation (Yamanaka)

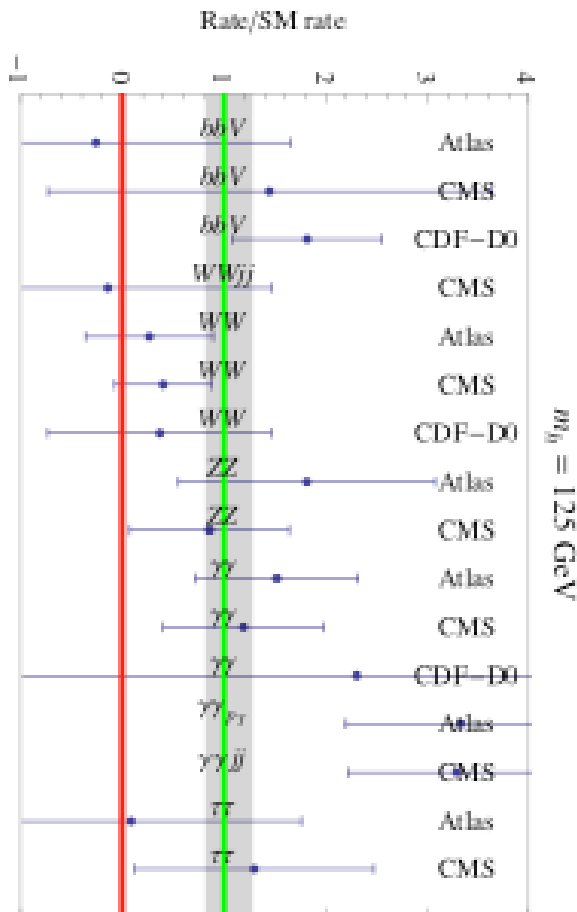
$m_h = 125 \text{ GeV}$



Invisible decay width

$$\mathcal{L}_S = \mathcal{L}_{\text{SM}} + \frac{1}{2}(\partial\phi)^2 - \frac{M_S^2}{2}\phi^2 - \frac{c_S}{2}|H|^2\phi^2 - \frac{d_S}{4!}\phi^4$$

ϕ : Scalar DM



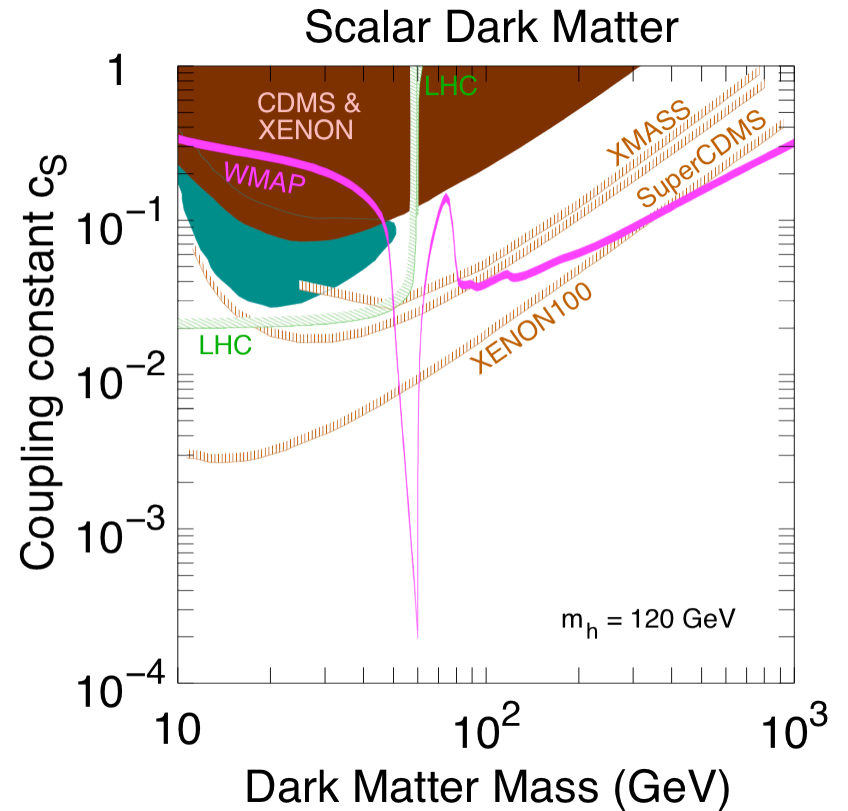
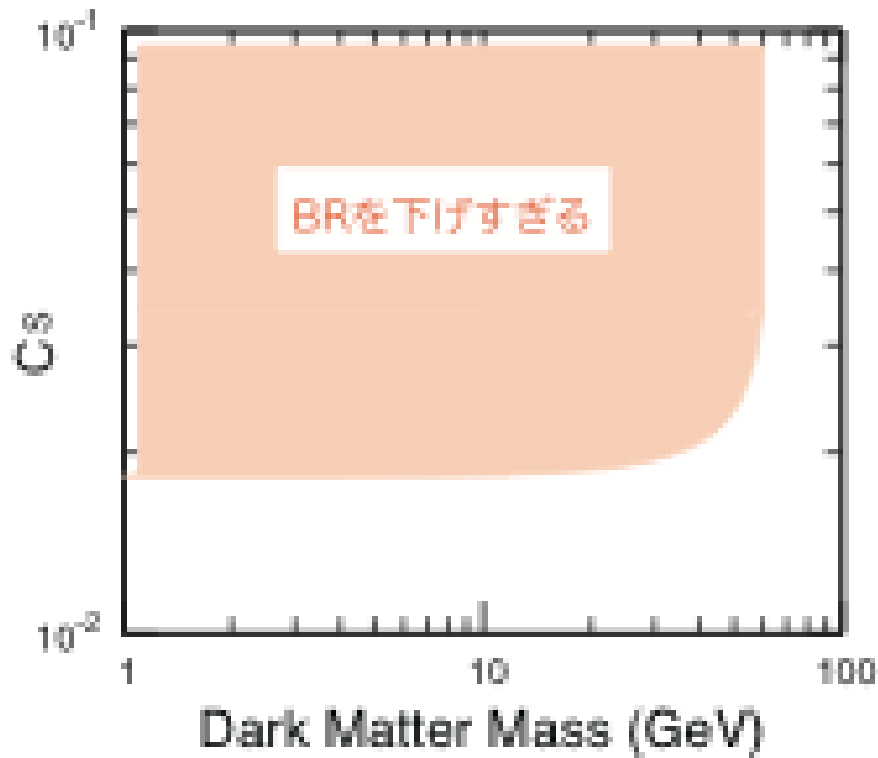
Partial Decay Width

$$\Gamma(H \rightarrow \phi\phi) = \frac{c_S^2 v^2}{32\pi m_h} \sqrt{1 - \frac{4m_S^2}{m_h^2}}$$

SMのGlobal fitから 2σ までずれてもよいとする

$$\Gamma(H \rightarrow \phi\phi) < 2 \times 0.2 \times 4 \text{ MeV}$$

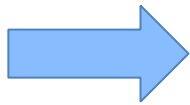
Couplingに対する制限



Dark Radiation

Dark Radiation

WMAP7 + ACBAR + ACT + SPR + SDSS-DR7



M. Archidiacono, E. Calabrese and A. Melchiorri [hep-ph 1109.2767](#)

背後にあるエネルギースケールを知りたい

とりあえず MeV scale での effective operator を書いてみる

scalar

$$\frac{1}{\Lambda^2} \phi^2 F_{\mu\nu} F^{\mu\nu} \quad \frac{m_f}{\Lambda^2} \phi^2 \bar{f} f \quad \frac{m_f}{\Lambda} \phi \bar{f} f$$
$$\frac{1}{\Lambda^4} (\partial\phi)^2 F_{\mu\nu} F^{\mu\nu} \quad \frac{1}{\Lambda} \phi F_{\mu\nu} F^{\mu\nu}$$

fermion

$$\frac{1}{\Lambda^2} \bar{\psi} \gamma_\mu \gamma_5 \psi \bar{f} \gamma^\mu f \quad \frac{m_\psi}{\Lambda^2} \bar{\psi} \gamma^{\mu\nu} \psi F_{\mu\nu} \quad \frac{m_f m_\phi}{\Lambda^4} \bar{\psi} \psi \bar{f} f$$

etc.

Thermal abundance を仮定して、cut off はいくらか？

エネルギースケール

Decoupling temperature $\sim 10\text{MeV}$, $f = e, m_\psi = 0.1\text{eV}$

Operators

Λ

$$\frac{1}{\Lambda^2} \phi^2 F_{\mu\nu} F^{\mu\nu}$$

TeV

$$\frac{m_f}{\Lambda^2} \phi^2 \bar{f} f$$

TeV

$$\frac{m_f}{\Lambda} \phi \bar{f} f$$

10^{15} GeV

$$\frac{1}{\Lambda^4} (\partial\phi)^2 F_{\mu\nu} F^{\mu\nu}$$

10 GeV

$$\frac{1}{\Lambda} \phi F_{\mu\nu} F^{\mu\nu}$$

10^8 GeV

$$\frac{1}{\Lambda^2} \bar{\psi} \gamma_\mu \gamma_5 \psi \bar{f} \gamma^\mu f$$

TeV

$$\frac{m_f}{\Lambda^3} \bar{\psi} \psi \bar{f} f$$

10 GeV

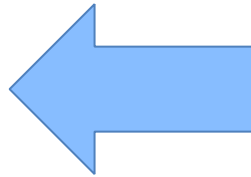
$$\frac{m_\psi}{\Lambda^2} \bar{\psi} \gamma^{\mu\nu} \psi F_{\mu\nu}$$

eV

TeV -> Higgs ?

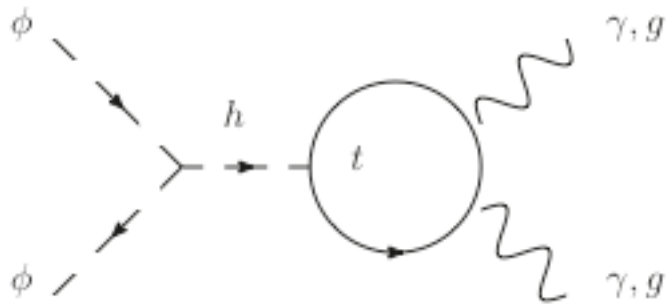
$$\frac{1}{\Lambda^2} \phi^2 F_{\mu\nu} F^{\mu\nu}$$

$$\frac{m_f}{\Lambda^2} \phi^2 \bar{f} f$$

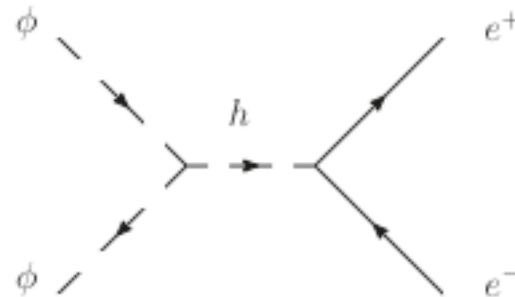


$$\lambda_{\phi H} \phi^2 (|H|^2 - v^2)$$

質量が出ない様に全オーダーでtune



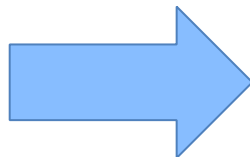
$$\lambda_{\phi H} \gtrsim 1000$$



$$\lambda_{\phi H} \gtrsim 0.01$$

Higgs invisible decayからの制限

$$\lambda_{\phi H} \gtrsim 0.01$$

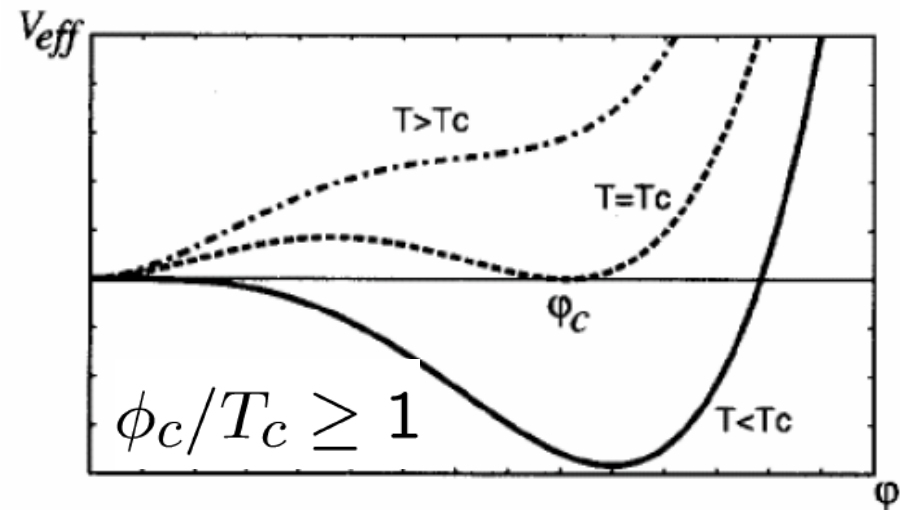
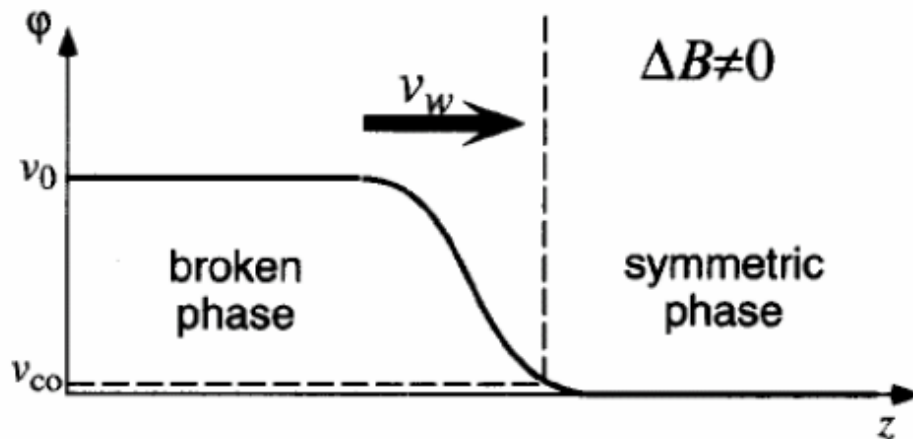

$$\Gamma(h \rightarrow \psi\psi) \gtrsim 0.001(\text{GeV})$$

ずれが見えるかも!?

Electroweak baryogenesis

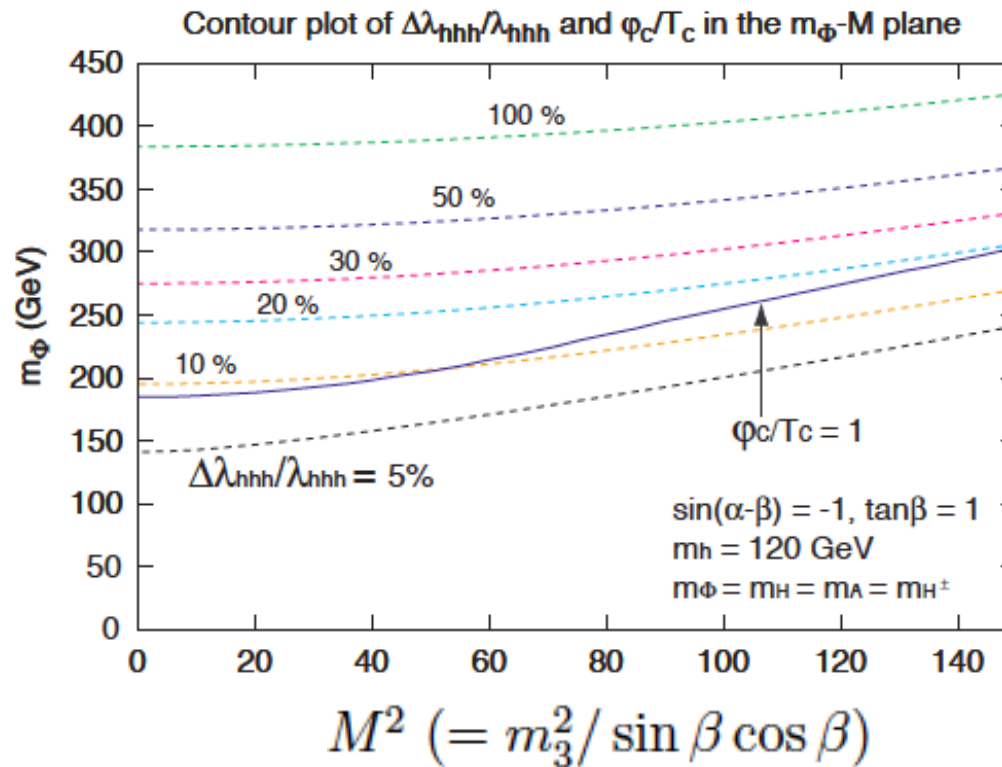
サハロフの3条件

- 1, バリオン数の破れ ← 電弱量子異常
- 2, CPの破れ ← とりあえず新物理を仮定
- 3, 熱平衡からのズレ ← 一次の電弱相転移



Triple vertex in 2HDM

$$V_{\text{tree}} = m_1^2 |\Phi_1|^2 + m_2^2 |\Phi_2|^2 - (m_3^2 \Phi_1^\dagger \Phi_2 + \text{h.c.}) \\ + \frac{\lambda_1}{2} |\Phi_1|^4 + \frac{\lambda_2}{2} |\Phi_2|^4 + \lambda_3 |\Phi_1|^2 |\Phi_2|^2 + \lambda_4 |\Phi_1^\dagger \Phi_2|^2 + \left[\frac{\lambda_5}{2} (\Phi_1^\dagger \Phi_2)^2 + \text{h.c.} \right],$$



Kanemura et al,
 Phys.Lett. B606
 (2005) 361-366

FIG. 1: The straight line stands for the critical line which satisfied the condition, $\varphi_c/T_c = 1$. The dashed lines are the deviation of hhh coupling from the SM value, where $\Delta\lambda_{hhh}^{\text{THDM}} \equiv \lambda_{hhh}^{\text{eff}}(\text{THDM}) - \lambda_{hhh}^{\text{eff}}(\text{SM})$.

Color assisted EWBG

Cohen et al
1203.2924 [hep-ph]

$$-\mathcal{L} \supset M_X^2 |X|^2 + \frac{K}{6} |X|^4 + Q |X|^2 |H|^2,$$

$N_{\text{scalars}} = 1$

$q_X = 2/3, N_{\text{scalars}} = 1$

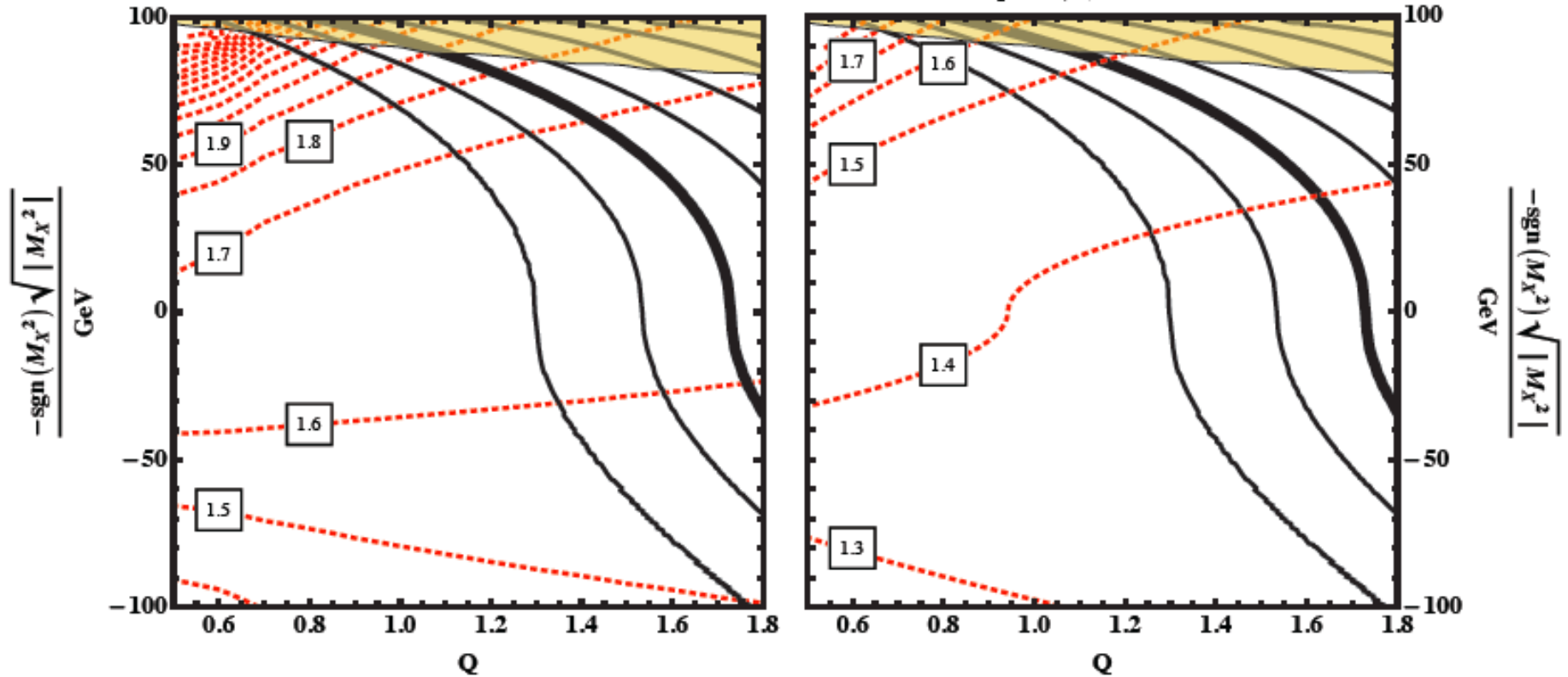


Figure 1: Contours of ϕ_C/T_C [black, solid lines] in the $-\text{sgn}(M_X^2)\sqrt{|M_X^2|}$ vs. Q plane for one new color-triplet scalar. (The most negative mass squared values are at the top of the plot.) The bolded line corresponds to $\phi_C/T_C = 0.9$ and the adjacent solid lines delineate steps of $\Delta(\phi_C/T_C) = 0.2$. The yellow shaded region is excluded because for these parameters the Universe would have evolved to a charge-color breaking minimum. In the *left* plot we also show contours of the ratio of the gluon fusion cross section to the SM value [red, dotted lines]. In the *right* plot we show contours of the ratio of the gluon fusion cross section times the branching ratio to di-photons to the SM value [red, dotted lines] when the charge of the colored scalar is taken to be $q_X = 2/3$.

Higgs inflationとは？

「Higgs = inflaton」としたslow-roll型inflation scenario

$$\text{Action} \quad S = \int d^4x \sqrt{-g} \left[\frac{1}{2} m_{\text{Pl}}^2 f(\phi) \mathcal{R} - \frac{1}{2} k(\phi) (\partial\phi)^2 - V(\phi) \right]$$

\mathcal{R} : Ricci scalar

重力との大きな値を持つnon-minimal couplingを導入することにより、

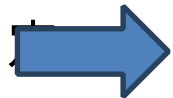
slow-rollingと初期揺らぎの生成を実現 [F. Bezrukhin and M. Shaposhnikov, arXiv:0710.3755]

メリット1: SM particle contentsのみで説明可能

メリット2: Higgsの物理を明らかにすることで地上実験にて検証可能

125GeV Higgs massとの相性は？

Tree level Higgs inflationでは、 n_s は標準模型パラメーターに依存せず—



Higgs massに予言与えず、地上実験での検証も困難

SM couplingのrunningを含めて、effective Planck massやPotentialのrunning



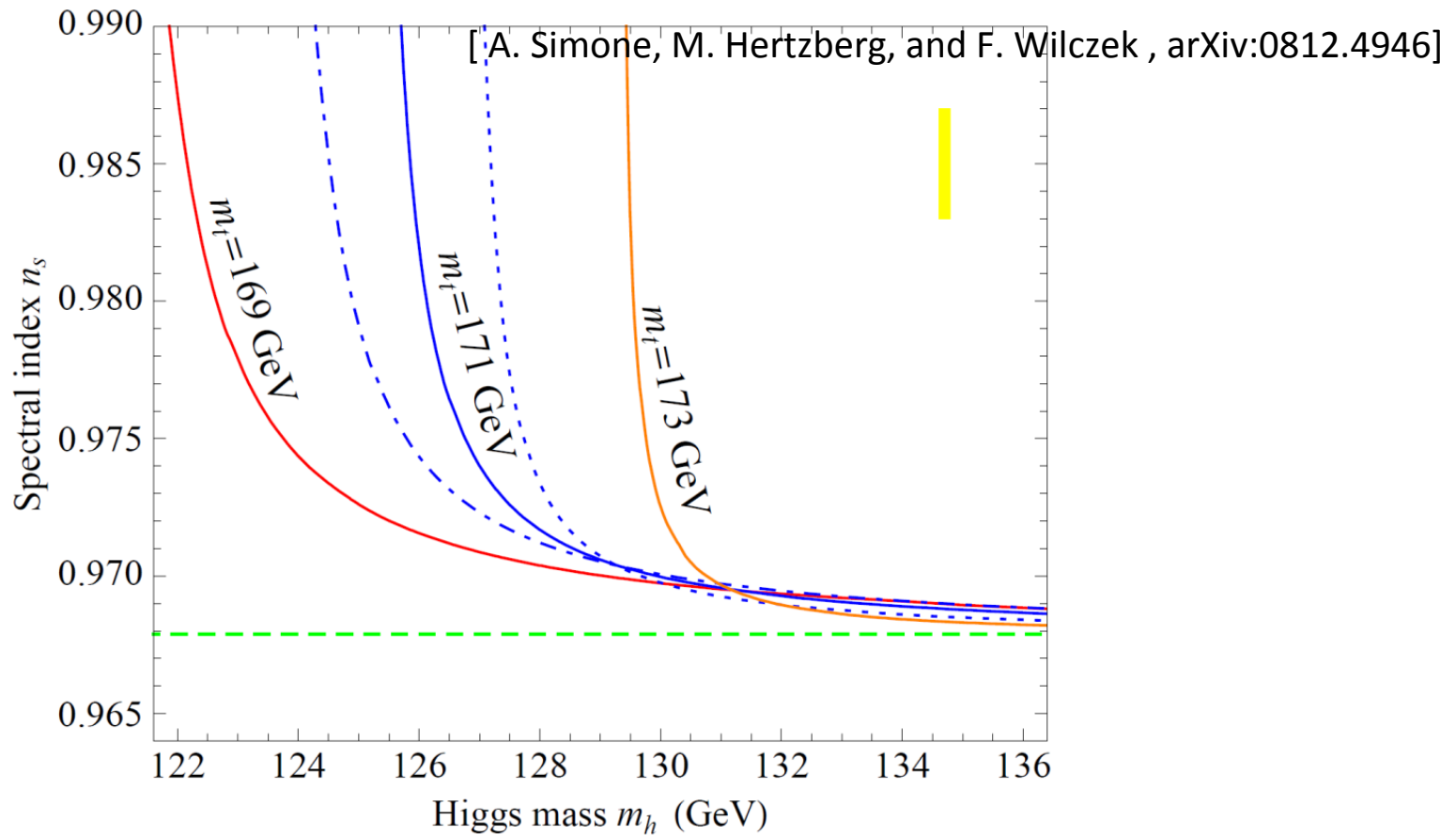
$n_s < 0.99$ を出せるHiggs mass (実際にはI) の領域に予言

$$m_h > 125.7 \text{ GeV} + 3.8 \text{ GeV} \left(\frac{m_t - 171 \text{ GeV}}{2 \text{ GeV}} \right) - 1.4 \text{ GeV} \left(\frac{\alpha_s(m_Z) - 0.1176}{0.0020} \right) \pm \delta$$

$\delta \sim 2 \text{ GeV}$: theoretical uncertainty from higher order corrections

125GeV HiggsはギリギリOK

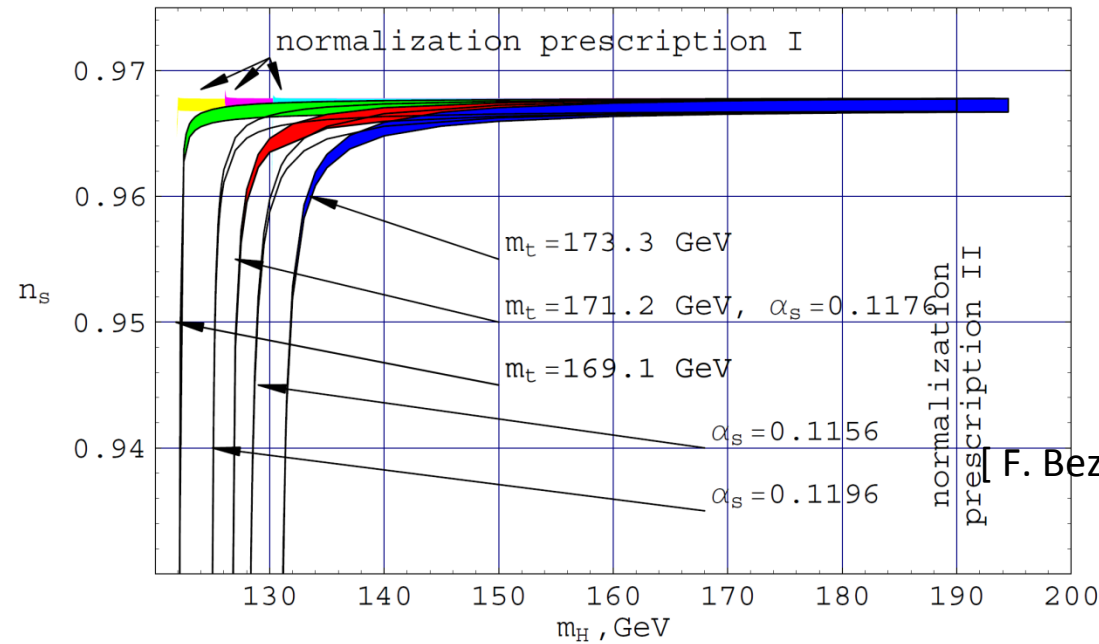
125GeV Higgs massとの相性は？



宇宙観測から課されている制限

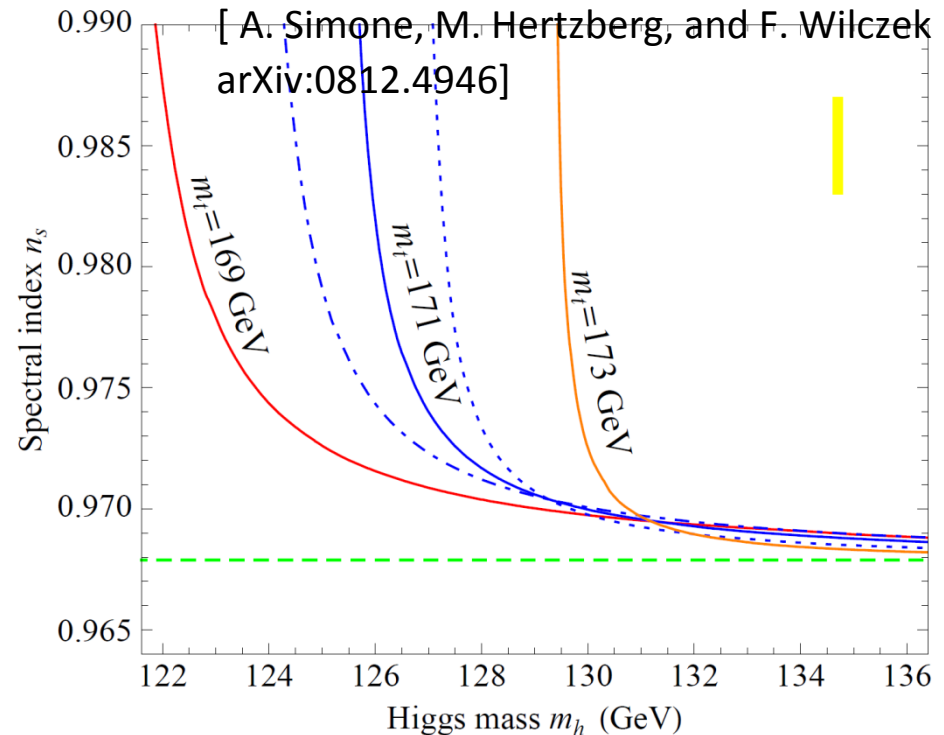
$0.93 < n_s < 0.99$ and $r < 0.22$ (at 95% confidence level)

Spectral index と Higgs mass



[F. Bezrukov and M. Shaposhnikov arXiv:0904.1537]

- ☑ Running orderや繰り込みの手法に大きく依存し、結果が全く異なる
- ☑ 125GeV Higgsがconfirmされ、Higgs inflationを信じるなら、計算方法を詰めるべき



補足(n_s 計算の両者の違い)

Comment from A. Simone, M. Hertzberg, and F. Wilczek arXiv:0812.4946

In our analysis, we computed the full RG improved effective potential. We did this including (i) 2-loop beta functions, (ii) the effect of curvature in the RG equations (through the function s), (iii) wave-function renormalization, and (iv) accurate specification of the initial conditions through proper pole matching. On the other hand, [32] did not compute the full effective potential or include any of the items (i)–(iv).⁹ Instead Ref. [32] approximated the potential at leading log order with couplings evaluated at an inflationary scale after running them at 1-loop (this is one step beyond [18] where couplings were not run).

Ref. [32] = [F. Bezrukov and M. Shaposhnikov arXiv:0904.1537]