


*Photoproduction of  $K^*$   
for the study of the  
structure of  $\Lambda(1405)$*



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## Motivations : Two poles?

There are two poles of the scattering amplitude around nominal  $\Lambda(1405)$  energy region.

- Cloudy bag model  
(1990)

J. Fink *et al.* PRC41, 2720

- Chiral unitary model  
(2001~)

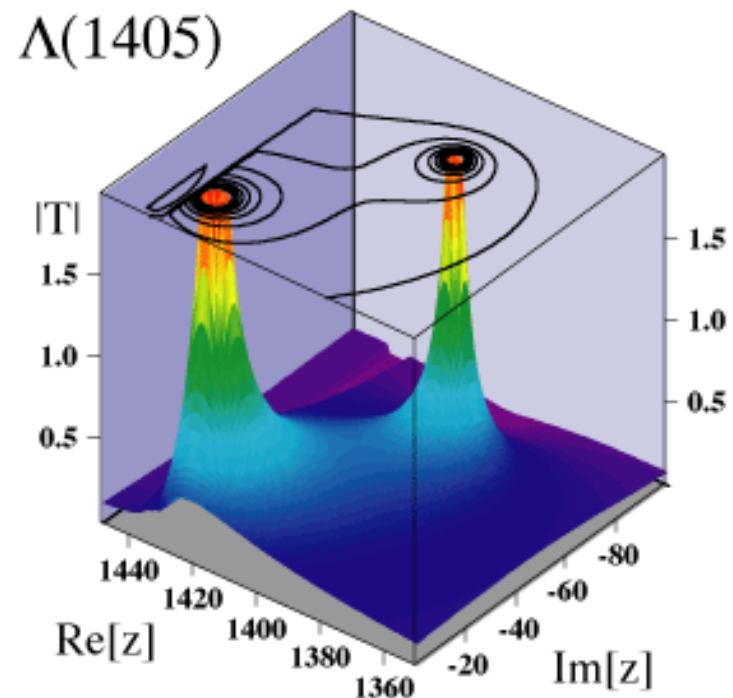
J. A. Oller *et al.* PLB500, 263

E. Oset *et al.* PLB527, 99

D. Jido *et al.* PRC66, 025203

T. Hyodo *et al.* PRC68, 018201

$\Lambda(1405) : J^P=1/2^-, I=0$



ChU model, T. Hyodo

# Chiral unitary model

Flavor SU(3) meson-baryon scatterings (s-wave)

**Chiral symmetry**

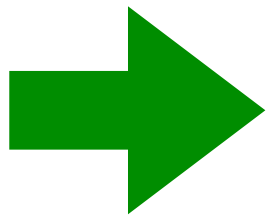
**Low energy  
behavior**



**Unitarity of S-matrix**

**Non-perturbative  
resummation**

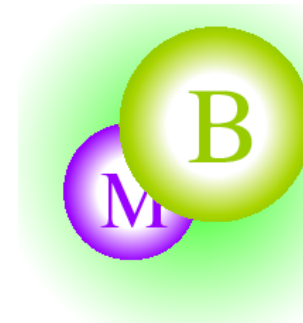
**Dynamical  
generation**



$$J^P = 1/2^-$$

**Resonances**

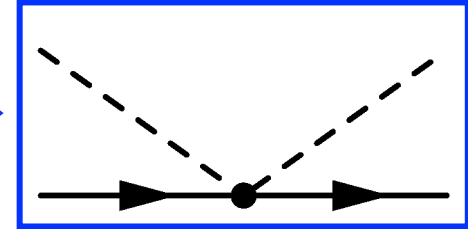
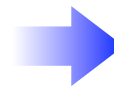
$\Lambda(1405)$ ,  $\Lambda(1670)$ ,  $N(1535)$ ,  
 $\Sigma(1620)$ ,  $\Xi(1620)$



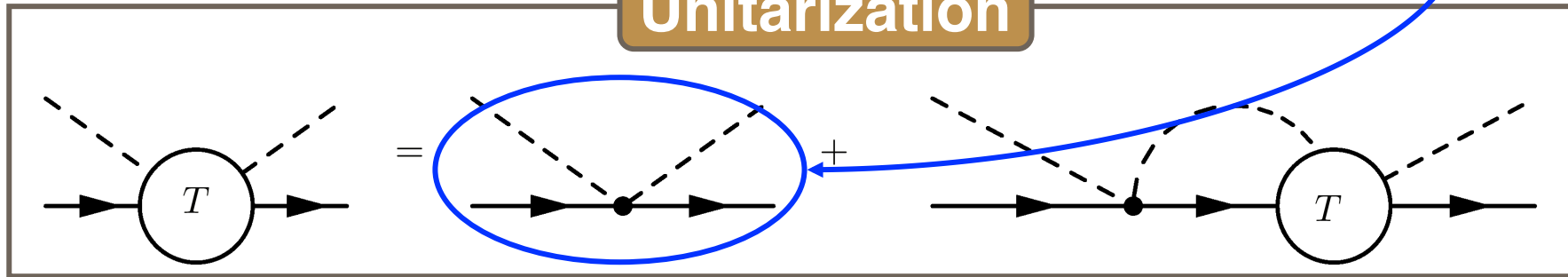
# Framework of the chiral unitary model

## Chiral perturbation theory

$$\mathcal{L}_{WT} = \frac{1}{4f^2} \text{Tr}(\bar{B}i\gamma^\mu[(\Phi\partial_\mu\Phi - \partial_\mu\Phi\Phi), B])$$

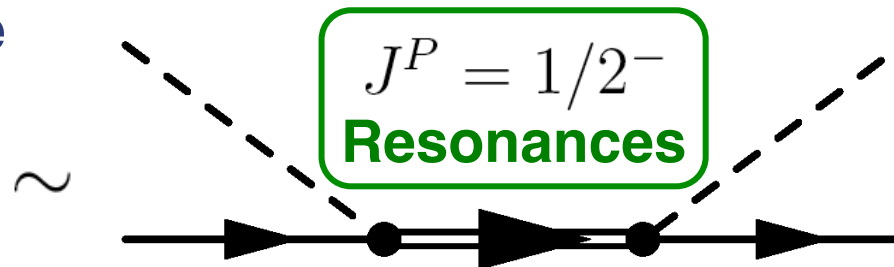


## Unitarization



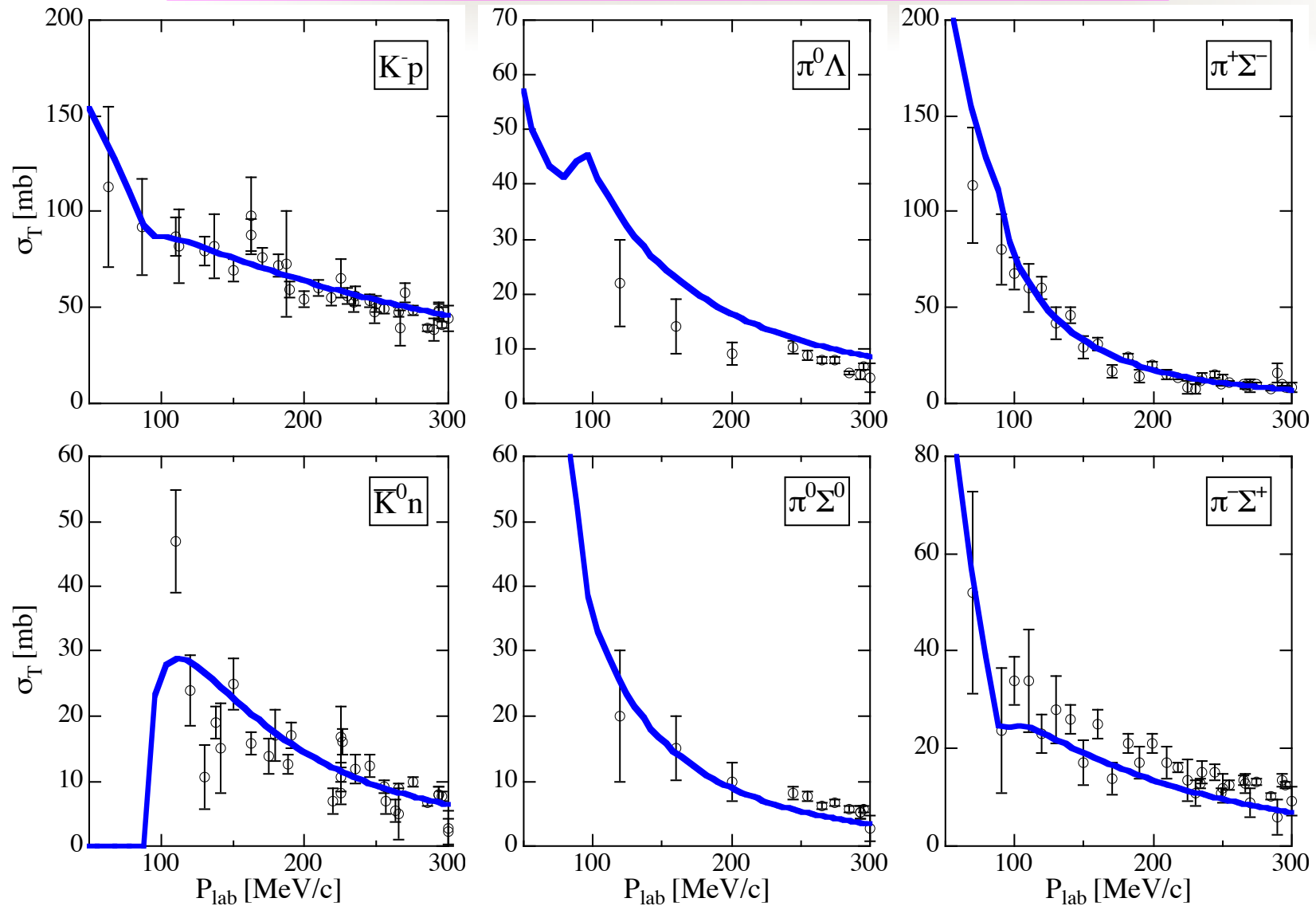
$$T_{ij}(\sqrt{s}) \sim \frac{g_i g_j}{\sqrt{s} - M_R + i\Gamma_R/2} + T_{ij}^{BG}$$

Generated resonances are expressed as poles of the scattering amplitude.





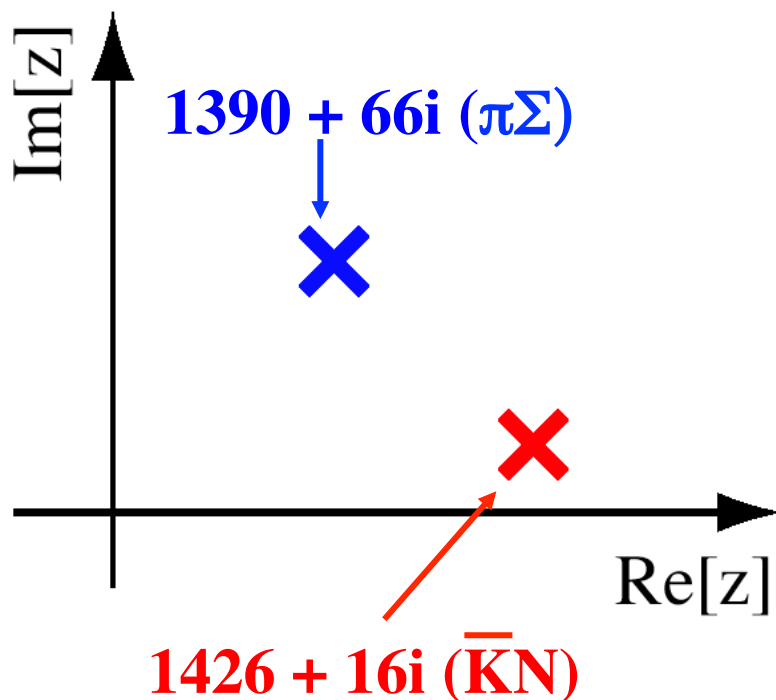
# Total cross sections of $K^-p$ scatterings



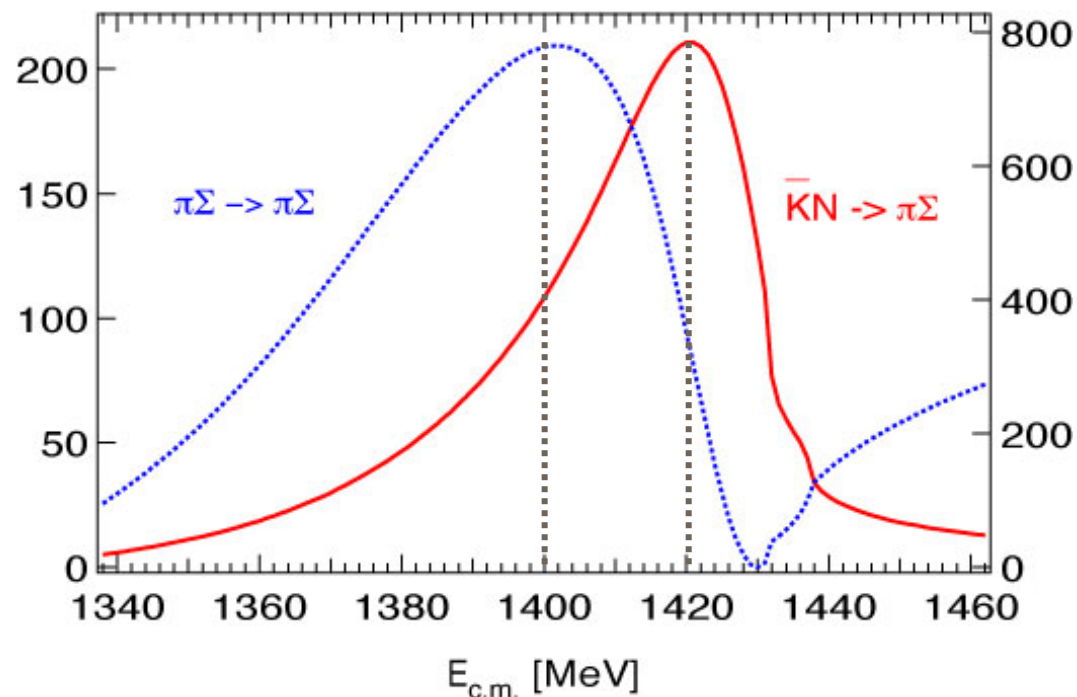
**T. Hyodo, et al., Phys. Rev. C 68, 018201 (2003)**

# $\Lambda(1405)$ in the chiral unitary model

## position of poles



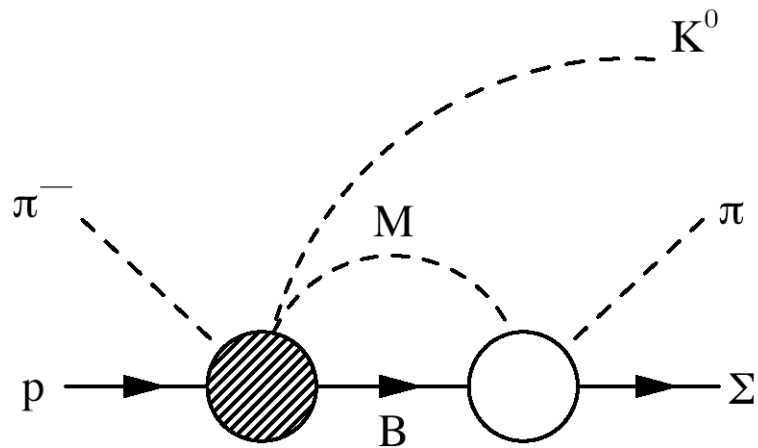
## $\pi\Sigma$ mass distribution



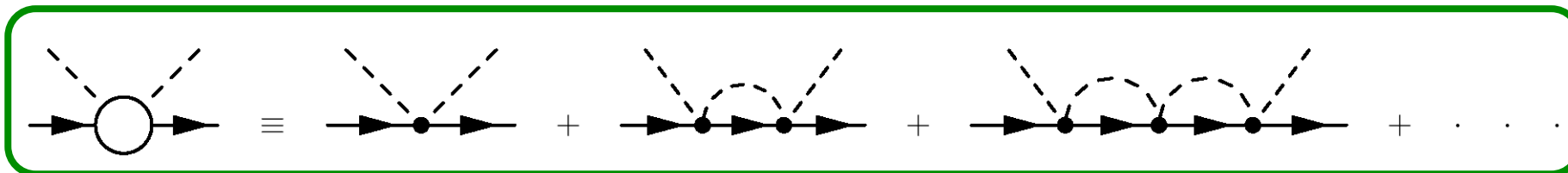
$$\frac{d\sigma}{dM_I} = C |t_{\pi\Sigma \rightarrow \pi\Sigma}|^2 p_{CM} \quad \longrightarrow \quad \frac{d\sigma}{dM_I} = \left| \sum_i C_i t_{i \rightarrow \pi\Sigma} \right|^2 p_{CM}$$

D. Jido, et al., Nucl. Phys. A 723, 205 (2003)

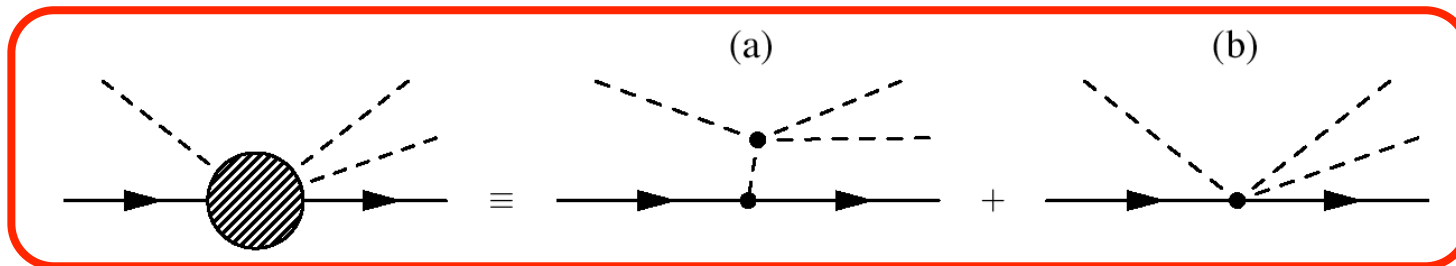
# Example : the $\pi^- p \rightarrow K^0 \pi \Sigma$ reaction



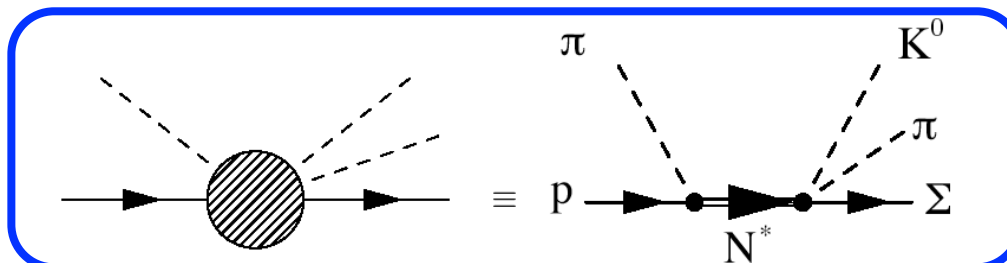
**Chiral unitary model**



**Chiral term**



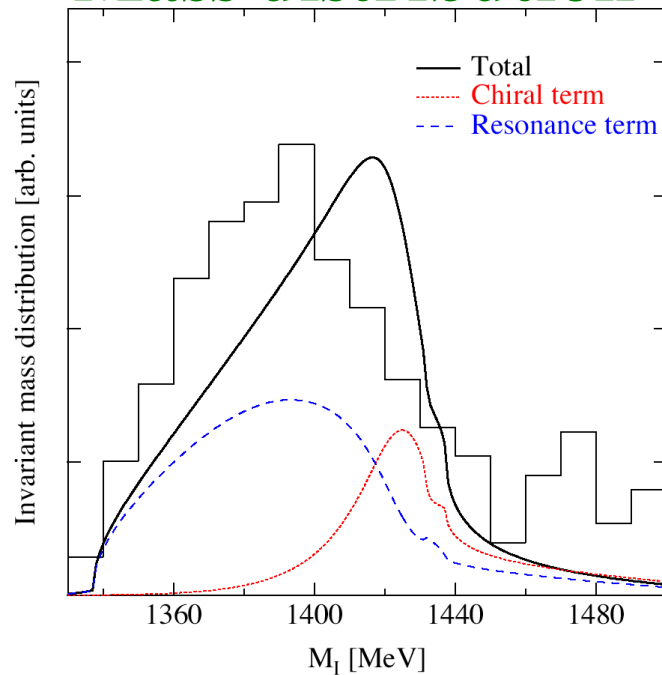
**N(1710)**





# Results for $\pi^- p \rightarrow K^0 \pi \Sigma$

## Mass distribution



## Total cross sections [mb]

final state	$K^0 K^- p$	$K^0 \bar{K}^0 n$	$K^0 \pi^0 \Lambda$	$K^0 \pi^+ \Sigma^-$	$K^0 \pi^- \Sigma^+$
Exp.	2.9	8.3	104.0	25.1	20.2
total	3.75	5.98	6.02	21.32	20.01
chiral	2.36	2.84	3.14	3.04	6.78
resonance	0.70	0.67	10.85	16.18	5.43

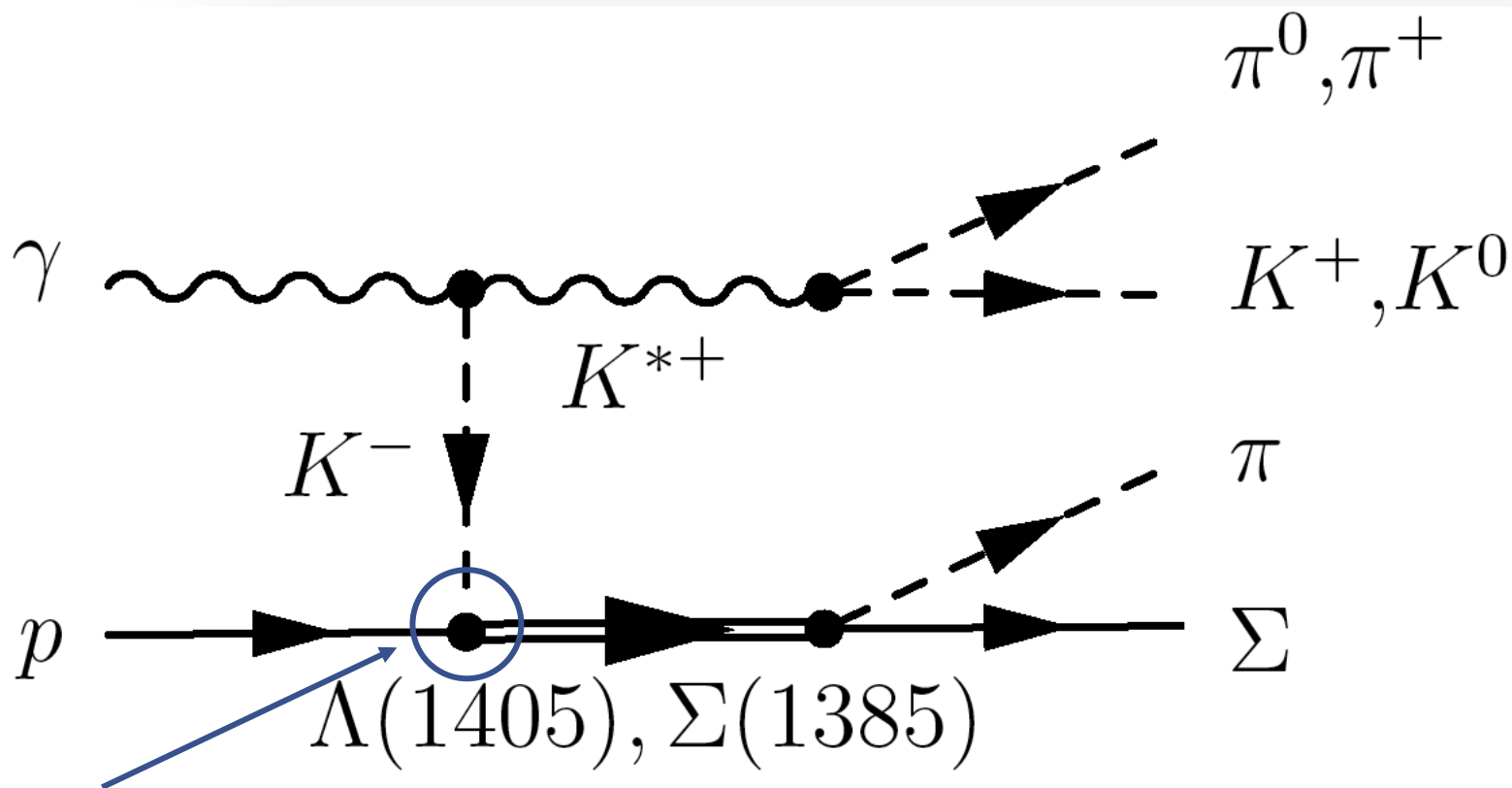
$\Sigma(1385)$  effect

Good agreement

❖ There are **two mechanisms** in the initial stage interaction, which **filter each one of the resonances**.

T. Hyodo, *et al.*, nucl-th/0307005, Phys. Rev. C, in press

## Photoproduction of $K^* \Lambda(1405)$



**Only  $K^-p$  channel appears at the initial stage**

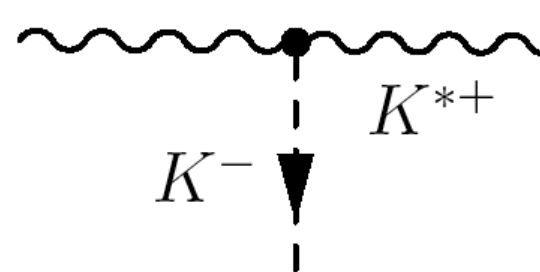
**Higher pole ??**

# Effective interactions for meson part

## 1. $\gamma VP$ coupling

$$-it = ig_{\gamma K^* K} \epsilon^{\mu\nu\alpha\beta} P_\mu \epsilon_\nu(K^{*+}) k_\alpha \epsilon_\beta(\gamma), \quad \gamma$$

$$|g_{\gamma K^{*\pm} K^\pm}| = 0.252 \text{ [GeV}^{-1}\text{]},$$

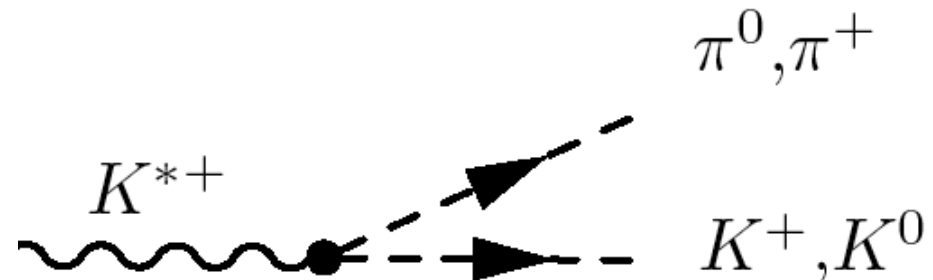
$$|g_{\gamma K^{*0} K^0}| = 0.385 \text{ [GeV}^{-1}\text{]}.$$


## 2. VPP coupling

$$-it(K^{*+} \rightarrow K^+ \pi^0) = i \frac{g_{VPP}}{\sqrt{2}} \frac{1}{\sqrt{2}} [p_\mu(K^+) - p_\mu(\pi^0)] \epsilon^\mu(K^{*+}),$$

$$-it(K^{*+} \rightarrow K^0 \pi^+) = i \frac{g_{VPP}}{\sqrt{2}} [p_\mu(K^0) - p_\mu(\pi^+)] \epsilon^\mu(K^{*+}),$$

$$g_{VPP} = -6.05$$



## Effective interaction for $\Sigma(1385)$

### 3. $\Sigma(1385)$ MB coupling

$$-it_{\Sigma^*i} = c_i \frac{12D + F}{5} \frac{1}{2f} \mathbf{S} \cdot \mathbf{k}_i$$

**SU(6) symmetry**



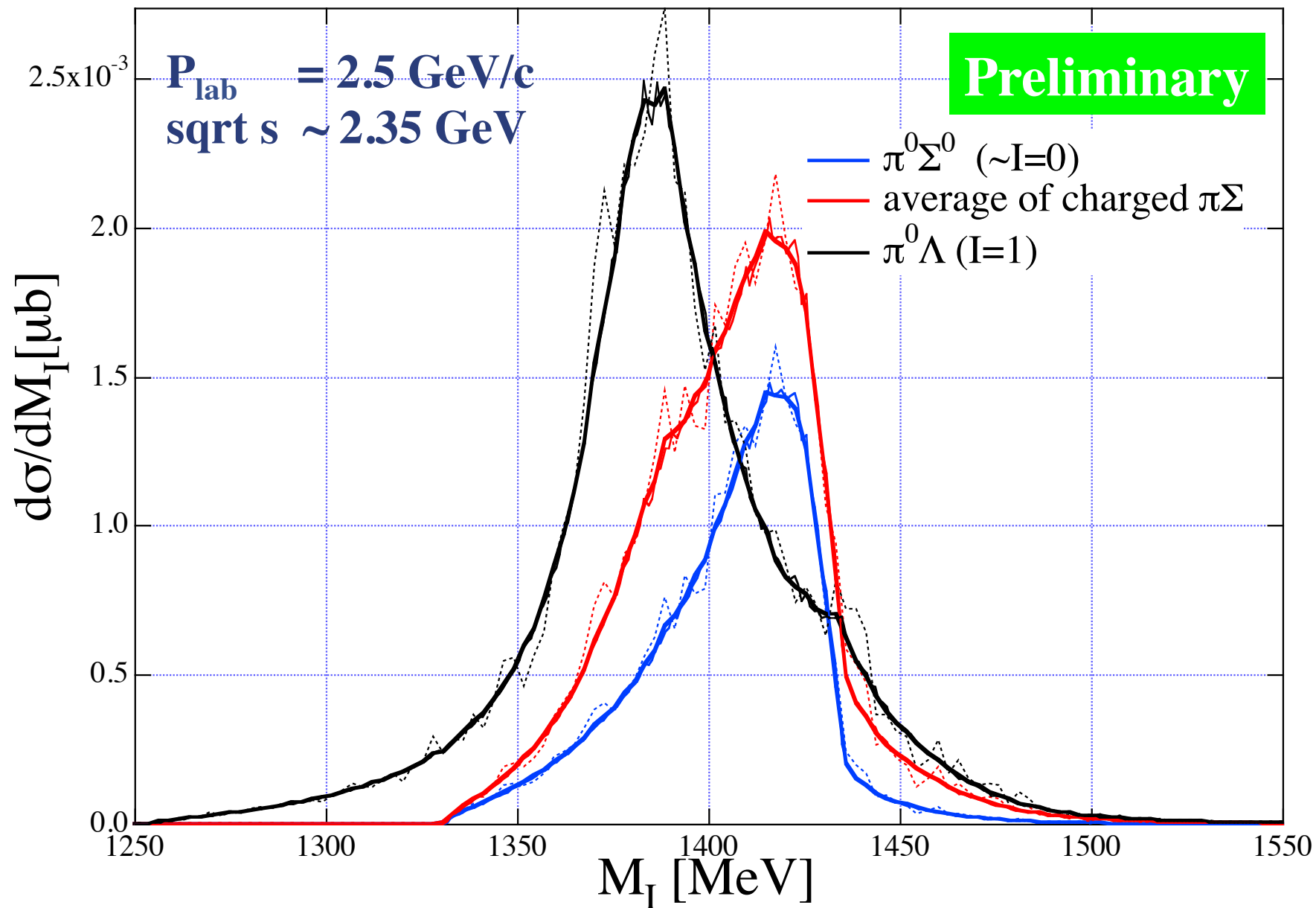
channel $i$	$K^-p$	$\bar{K}^0n$	$\pi^0\Lambda$	$\pi^0\Sigma^0$	$\eta\Lambda$	$\eta\Sigma^0$	$\pi^+\Sigma^-$	$\pi^-\Sigma^+$	$K^+\Xi^-$	$K^0\Xi^0$
$c_i$	$-\sqrt{\frac{1}{12}}$	$\sqrt{\frac{1}{12}}$	$\sqrt{\frac{1}{4}}$	0	0	$-\sqrt{\frac{1}{4}}$	$-\sqrt{\frac{1}{12}}$	$\sqrt{\frac{1}{12}}$	$\sqrt{\frac{1}{12}}$	$-\sqrt{\frac{1}{12}}$

### 4. $K\text{-}P \rightarrow \Sigma(1385) \rightarrow \text{MB}$ amplitude

$$\begin{aligned}
 -it_{1i} &= c_1 c_i \left( \frac{12D + F}{5} \frac{1}{2f} \right)^2 \mathbf{S} \cdot \mathbf{k}_1 \mathbf{S}^\dagger \cdot \mathbf{k}_i \frac{i}{M_I^{(b)} - M_{\Sigma^*} + i\Gamma_{\Sigma^*}/2} F_f(k_1) \\
 &= c_1 c_i \left( \frac{12D + F}{5} \frac{1}{2f} \right)^2 (k_1)_l (k_i)_m \left( \frac{2}{3} \delta_{lm} - \frac{i}{3} \epsilon_{lmn} \sigma_n \right) \frac{i}{M_I^{(b)} - M_{\Sigma^*} + i\Gamma_{\Sigma^*}/2} F_f(k_1)
 \end{aligned}$$

$$F_f(k_1) = \frac{\Lambda^2 - m_K^2}{\Lambda^2 - (k_1)^2}$$

# Experiments : $\pi\Sigma$ mass distribution



## Summary and conclusions

We study the **structure of  $\Lambda(1405)$**  using the chiral unitary model.

🍏 There are **two poles** of the scattering amplitude around nominal  $\Lambda(1405)$ .

**Pole 1 (1426+16i) : strongly couples to KN state**

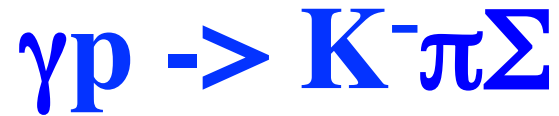
**Pole 2 (1390+66i) : strongly couples to  $\pi\Sigma$  state**

🍏 By observing the  **$\pi\Sigma$  mass distribution** in the  **$\gamma p \rightarrow K^* \Lambda(1405)$**  reaction, it could be possible to isolate **higher energy pole**.

[http://www.rcnp.osaka-u.ac.jp/~hyodo/index\\_e.html](http://www.rcnp.osaka-u.ac.jp/~hyodo/index_e.html)



## Appendix : other processes



J.C. Nacher, *et al.*, PLB445, 55(1999)

Spring-8



J.C. Nacher, *et al.*, PLB461, 299(1999)

J-PARC?

