

Θ^+ production in the $K^+ p \rightarrow \pi^+ KN$ reaction



Tetsuo Hyodo^a,

F. J. Llanes-Estrada^b, E. Oset^c, A. Hosaka^a,

J. R. Pelaez^b and M. J. Vicente Vacas^c

RCNP, Osaka^a Madrid^b IFIC, Valencia^c 2004, July 21st 1

Motivations

Results of (π^-, K^-) reaction at KEK

Upper limit of cross section $\sim 2 \mu\text{b}$

K. Miwa, talk given at PENTAQUARK04

Why so small? What about (K^+, π^+) ?

Two meson coupling effects for Θ^+

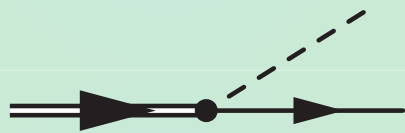
Possibility of $K\pi N$ bound state

P. Bicudo, *et al.*, Phys. Rev. C69, 011503 (2004)

T. Kishimoto, *et al.*, hep-ex/0312003

F. J. Llanes-Estrada, *et al.*, Phys. Rev. C69, 055203 (2004)

Two-meson coupling

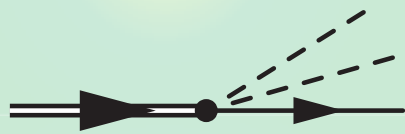


$$\Theta^+ \rightarrow KN$$

Very narrow

$$N(1710) \rightarrow \pi N$$

10–20 %

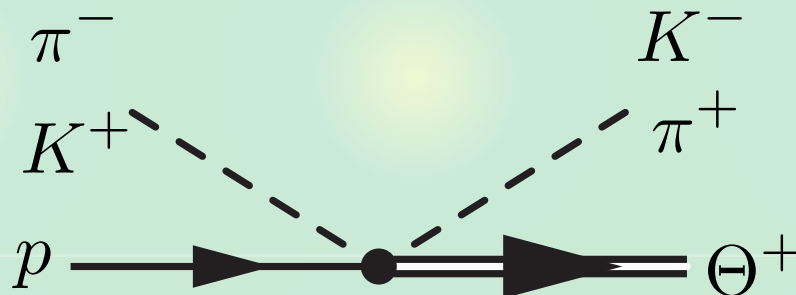


$$\Theta^+ \rightarrow K\pi N$$

Forbidden

$$N(1710) \rightarrow \pi\pi N$$

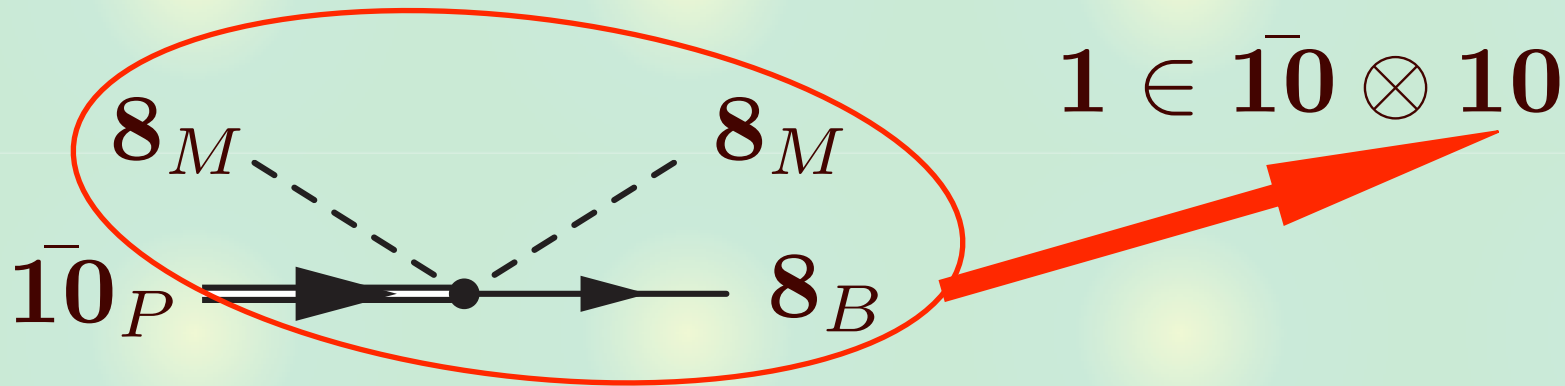
40–90 %



Large??

Effective interactions which account for the $N(1710) \rightarrow \pi\pi N$ decay

SU(3) structure of effective Lagrangian



$$\mathbf{8}_M \otimes \mathbf{8}_M \otimes \mathbf{8}_B = (\mathbf{1} \oplus \mathbf{8}^s \oplus \mathbf{8}^a \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27})_{MM} \otimes \mathbf{8}_B$$

$$= \mathbf{8} \quad \leftarrow \text{from } \mathbf{1}_{MM} \otimes \mathbf{8}_B$$

$$\oplus (\mathbf{1} \oplus \mathbf{8} \oplus \mathbf{8} \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27}) \quad \leftarrow \text{from } \mathbf{8}_{MM}^s \otimes \mathbf{8}_B$$

$$\oplus (\mathbf{1} \oplus \mathbf{8} \oplus \mathbf{8} \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27}) \quad \leftarrow \text{from } \mathbf{8}_{MM}^a \otimes \mathbf{8}_B$$

$$\oplus (\mathbf{8} \oplus \mathbf{10} \oplus \mathbf{27} \oplus \mathbf{35}) \quad \leftarrow \text{from } \mathbf{10}_{MM} \otimes \mathbf{8}_B$$

$$\oplus (\mathbf{8} \oplus \bar{\mathbf{10}} \oplus \mathbf{27} \oplus \mathbf{35}') \quad \leftarrow \text{from } \bar{\mathbf{10}}_{MM} \otimes \mathbf{8}_B$$

$$\oplus (\mathbf{8} \oplus \mathbf{10} \oplus \bar{\mathbf{10}} \oplus \mathbf{27} \oplus \mathbf{27} \oplus \mathbf{35} \oplus \mathbf{35}'' \oplus \mathbf{64}) \quad \leftarrow \text{from } \mathbf{27}_{MM} \otimes \mathbf{8}_B$$

Interaction Lagrangians 1

Antidecuplet field

$$P^{333} = \sqrt{6}\Theta_{10}^+$$

$$P^{133} = \sqrt{2}N_{10}^0 \quad P^{233} = -\sqrt{2}N_{10}^+$$

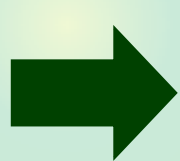
$$P^{113} = \sqrt{2}\Sigma_{10}^- \quad P^{123} = -\Sigma_{10}^0 \quad P^{223} = -\sqrt{2}\Sigma_{10}^+$$

$$P^{111} = \sqrt{6}\Xi_{10}^{--} \quad P^{112} = -\sqrt{2}\Xi_{10}^- \quad P^{122} = \sqrt{2}\Xi_{10}^0 \quad P^{222} = -\sqrt{6}\Xi_{10}^+$$

Construction of 8s Lagrangian

$$D^i_j[\mathbf{8}_{MM}^s] = 2\phi^i_k\phi^k_j - \frac{2}{3}\delta^i_j\phi^a_b\phi^b_a$$

$$T^{ijk}[\mathbf{10}_{MM(8s)B}] = \epsilon^{abk}D^i_a B^j_b + (\text{symmetrized})$$



$$\mathcal{L}^{8s} = \frac{g^{8s}}{2f}\bar{P}_{ijk}\epsilon^{abk}\phi^i_c\phi^c_a B^j_b + h.c.$$

Interaction Lagrangians 2

$$\mathcal{L}^{8s} = \frac{g^{8s}}{2f} \bar{P}_{ijk} \epsilon^{abk} \phi_c^i \phi_a^c B_b^j + h.c.$$

$$\mathcal{L}^{8a} = 0$$

$$\mathcal{L}^{10} = 0$$

← symmetry under exchange of mesons

$$\mathcal{L}^{27} = \frac{g^{27}}{2f} \left[4\bar{P}_{ijk} \epsilon^{ibc} \phi_b^j \phi_a^k B_c^a - \frac{4}{5} \bar{P}_{ijk} \epsilon^{abk} \phi_c^i \phi_a^c B_b^j \right] + h.c.$$

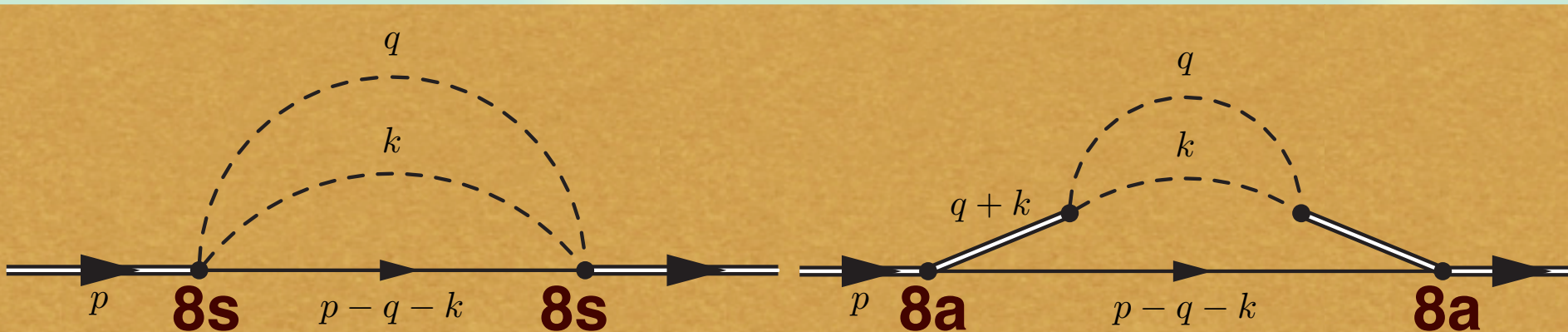
Experimental information

$$N(1710) \rightarrow \pi\pi (s\text{-wave}, I = 0) N$$

$$N(1710) \rightarrow \pi\pi (p\text{-wave}, I = 1) N$$

$$\mathcal{L}^{8a} = i \frac{g^{8a}}{4f^2} \bar{P}_{ijk} \epsilon^{abk} \gamma^\mu (\Phi \partial_\mu \Phi - \partial_\mu \Phi \Phi)^i{}_a B_b^j + h.c.$$

Diagrams for self-energy



Real part : mass shift

Imaginary part : decay width

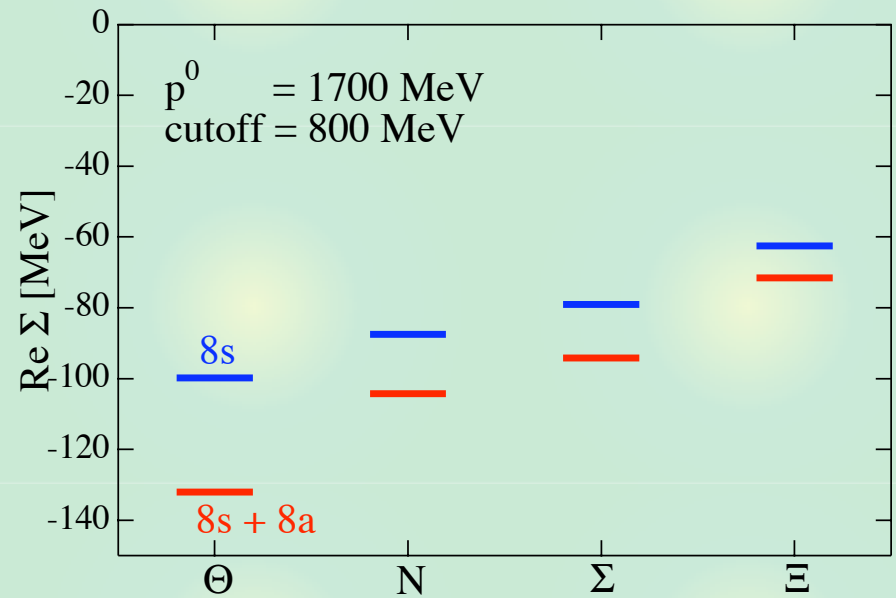
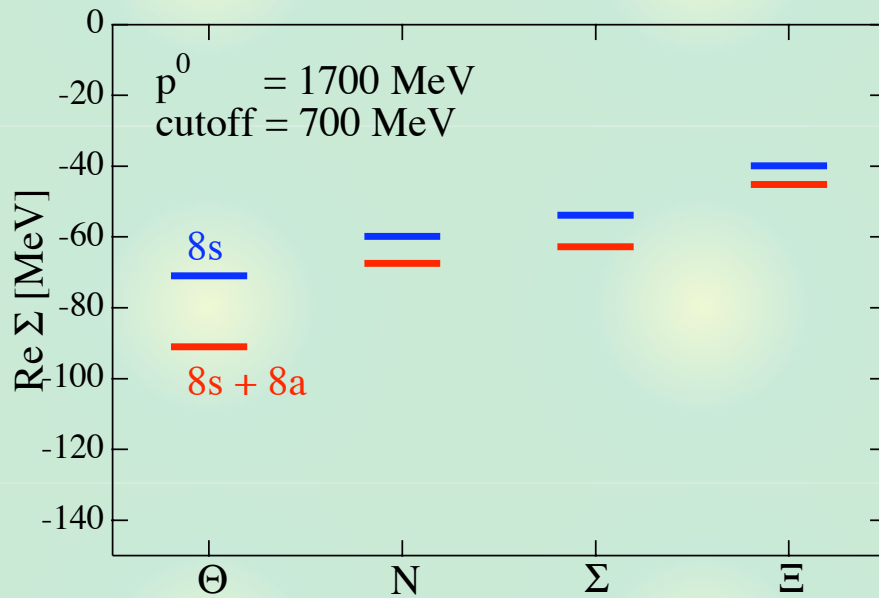
SU(3) breaking : masses of particles

$$N(1710) \rightarrow \pi\pi(s\text{-wave}, I = 0)N \quad \mathbf{25 \text{ MeV}}$$

$$N(1710) \rightarrow \pi\pi(p\text{-wave}, I = 1)N \quad \mathbf{15 \text{ MeV}}$$

➔ $g^{8s} = 1.88$, $g^{8a} = 0.315$

Results of self-energy : Real part (mass shift)



All mass shifts are attractive.

More bound for larger strangeness.

Mass difference between Θ and Ξ

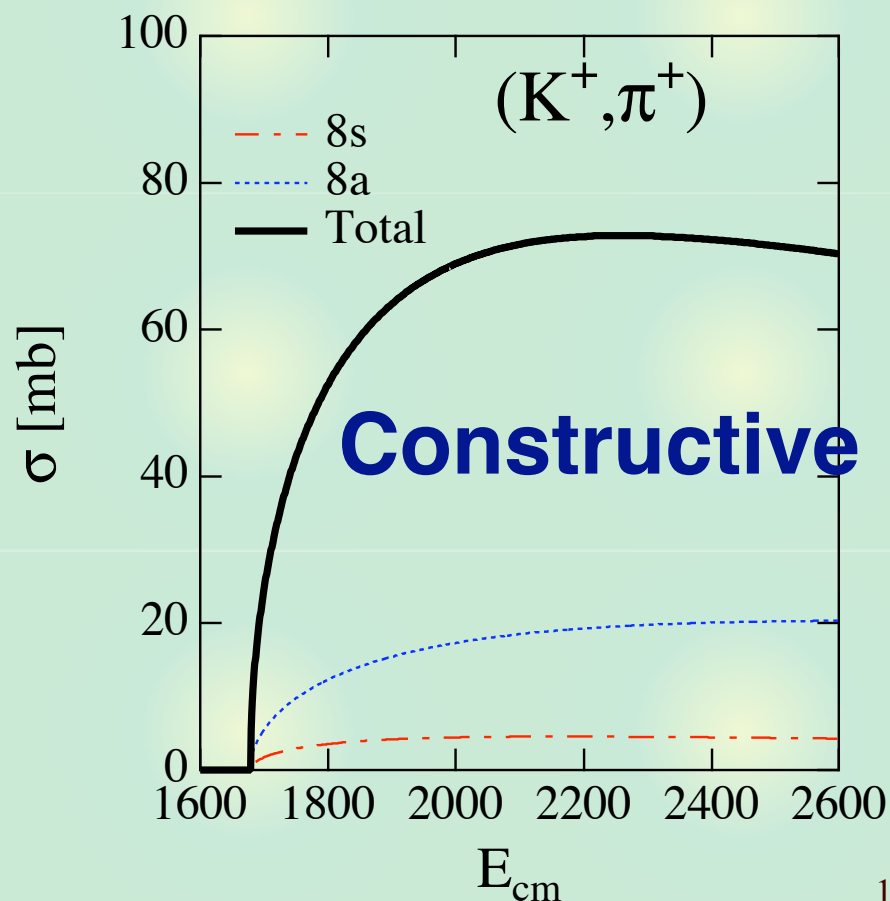
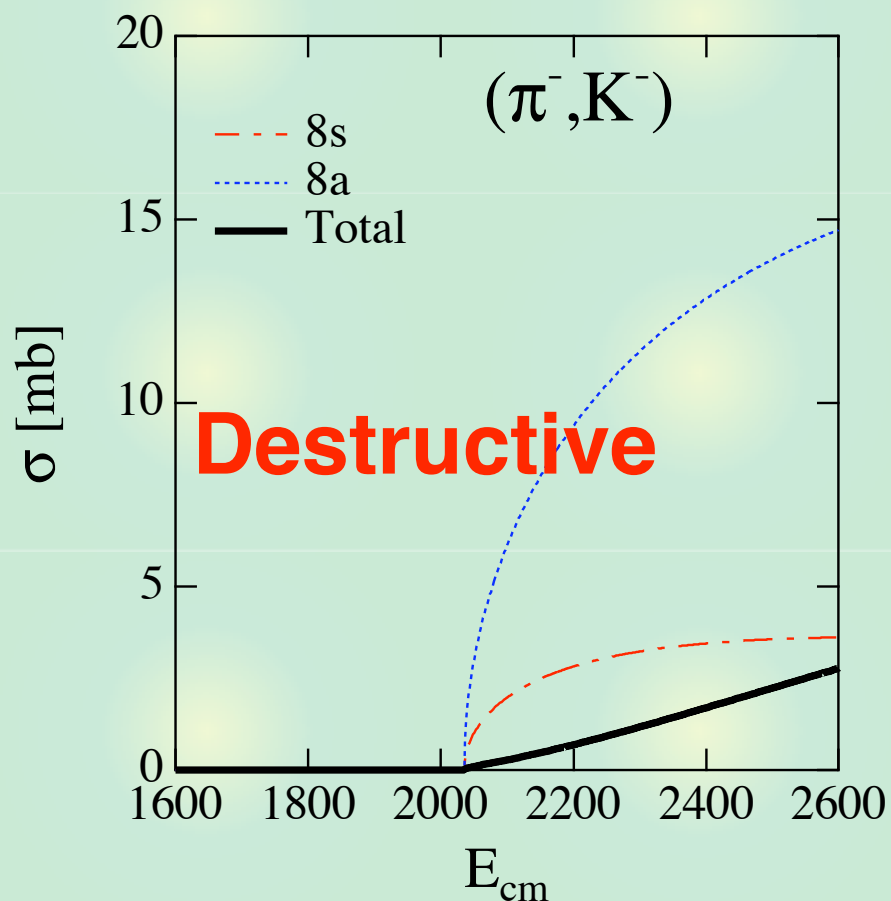
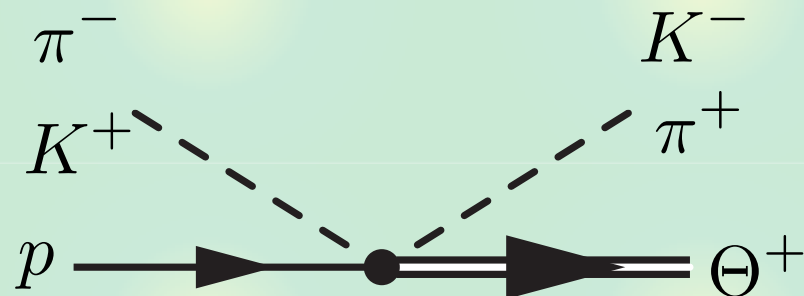
-> 60 MeV : ~20 % of 320 = 1860–1540

Results of self-energy : Imaginary part (decay width)

Decay [MeV]	$\Gamma^{(8s)}$	$\Gamma^{(8a)}$	$\Gamma_{BMM}^{(tot)}$
$N(1710) \rightarrow N\pi\pi$ (inputs)	25	15	40
$N(1710) \rightarrow N\eta\pi$	0.58	-	
$\Sigma(1770) \rightarrow N\bar{K}\pi$	4.7	6.0	24
$\Sigma(1770) \rightarrow \Sigma\pi\pi$	10	0.62	
$\Sigma(1770) \rightarrow \Lambda\pi\pi$	-	2.9	
$\Xi(1860) \rightarrow \Sigma\bar{K}\pi$	0.57	0.46	2.1
$\Xi(1860) \rightarrow \Xi\pi\pi$	-	1.1	


Results of reaction : cross sections

Total cross section of



Conclusion 1 : self-energy

We study the two-meson virtual cloud effect to the self-energy of baryon antidecuplet.

 Two-meson cloud effects are always **attractive**, and contribute to the antidecuplet **mass splitting**, of the order of **20%**.

 Antidecuplet members have relatively **small decay widths to MMB channel**.

T. H., F. J. Llanes-Estrada, E. Oset, A. Hosaka, J. R. Pelaez, M. J. Vicente Vacas, *in preparation*

Conclusion 2 : reactions

We investigate the Θ production in (π^-, K^-) and (K^+, π^+) reactions, with the vertices obtained from the self-energy study.



The small cross section in (π^-, K^-) reaction could be qualitatively explained by **the interference of two amplitudes.**

Future works



Self-energy

Possible mixing with the other flavor multiplets (8, 27, ...), other types of interactions (chiral?), ...



Reaction

Quantitative analysis (From factor), background cross section