Determination of the $\pi\Sigma$ scattering length from Λ_c decay

Tetsuo Hyodo,

and Makoto Oka

Tokyo Institute of Technology

supported by Global Center of Excellence Program
"Nanoscience and Quantum Physics"

2011, Mar. 28th

Scattering length

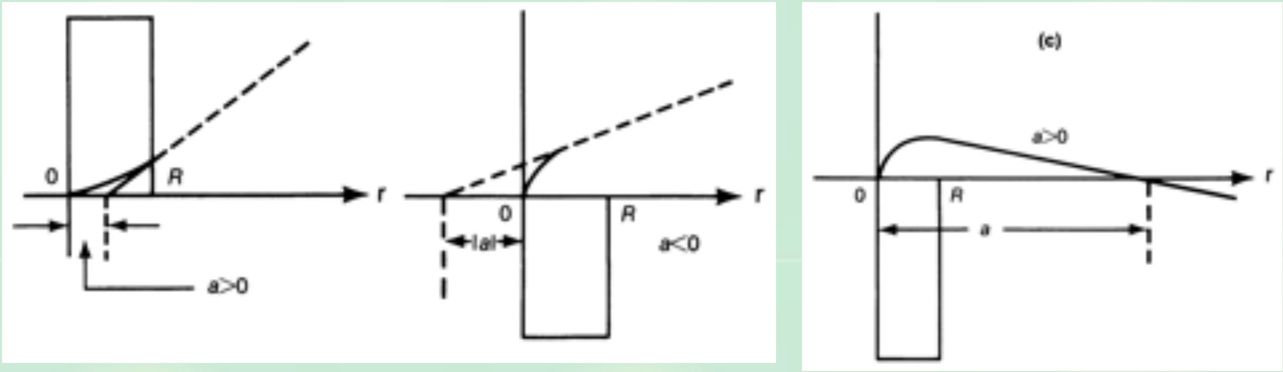
Scattering length: amplitude at threshold

- **characterizes the low energy scattering**
- **changes the sign if there is a bound state**

$$a = -f(k,\theta)|_{k\to 0}$$

$$\frac{d\sigma}{d\Omega} = |f(k,\theta)|^2, \quad \lim_{k\to 0}\sigma(k) = 4\pi a^2$$

$$f_{l=0}(k) = \frac{1}{k\cot\delta_0 - ik}, \quad k\cot\delta_0 = -\frac{1}{a} + r_e\frac{k^2}{2} + \dots$$



J.J. Sakurai, Modern Quantum Mechanics, p. 415

(In hadron physics we usually adopt the opposite sign
--> positive for attraction, negative for repulsion)

Importance of the πΣ scattering length

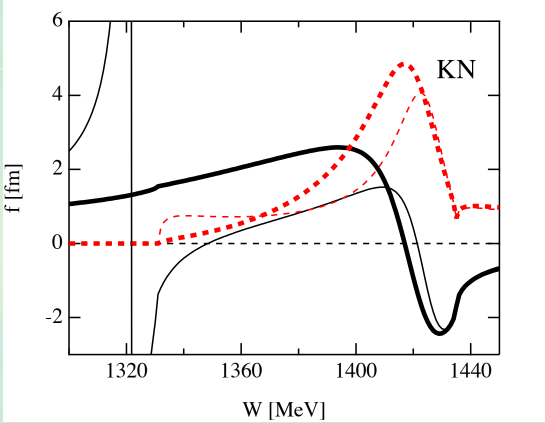
Structure of Λ(1405) and threshold behavior of πΣ scattering

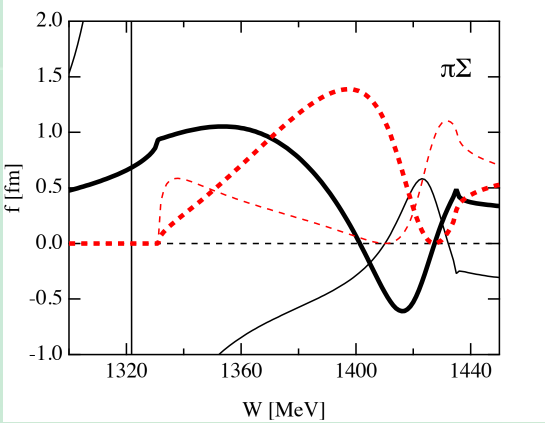
Y. Ikeda, T. Hyodo, D. Jido, H. Kamano, T. Sato, K. Yazaki, arXiv:1101.5190 [nucl-th]

Simple model extrapolation with $\bar{K}N$(I=0) being fixed
--> large uncertainty at low energy

Model	A1	A2	B E-dep	B E-indep
parameter ($\pi\Sigma$)	$d_{\pi\Sigma} = -1.67$	$d_{\pi\Sigma} = -2.85$	$\Lambda_{\pi\Sigma} = 1005$ MeV	$\Lambda_{\pi\Sigma} = 1465$ MeV
parameter ($\bar{K}N$)	$d_{\bar{K}N} = -1.79$	$d_{\bar{K}N} = -2.05$	$\Lambda_{\bar{K}N} = 1188$ MeV	$\Lambda_{\bar{K}N} = 1086$ MeV
pole 1 [MeV]	$1422 - 16i$	$1425 - 11i$	$1422 - 22i$	$1423 - 29i$
pole 2 [MeV]	$1375 - 72i$ (R)	1321 (B)	$1349 - 54i$ (R)	1325 (V)
$a_{\pi\Sigma}$ [fm]	0.934	-2.30	1.44	5.50
r_e [fm]	5.02	5.89	3.96	0.458
$a_{\bar{K}N}$ [fm] (input)	$-1.70 + 0.68i$	$-1.70 + 0.68i$	$-1.70 + 0.68i$	$-1.70 + 0.68i$

- **πΣ threshold behavior**
  --> pole structure
  --> K (Λ*) nuclei



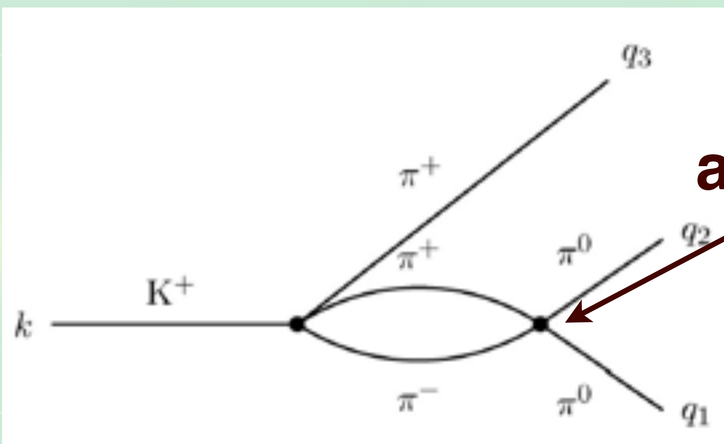


Determination of the scattering length

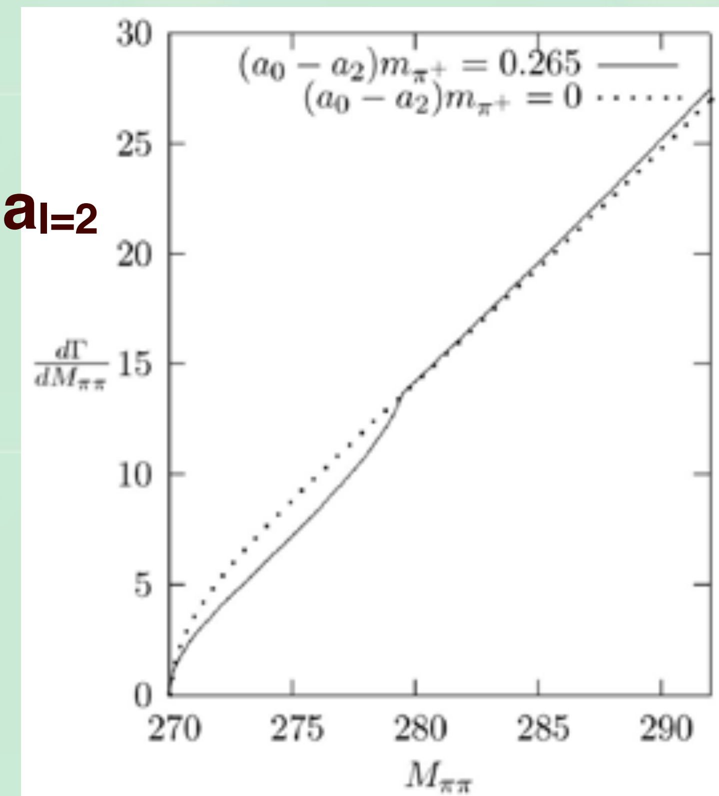
Extraction of hadron scattering length
- shift and width of atomic state (c.f. Kaonic hydrogen)
- extrapolation of low energy phase shift
- final state interaction from heavy particle's decay

Cabibbo's method for π-π scattering length

N. Cabibbo, Phys. Rev. Lett. 93, 121801 (2004)



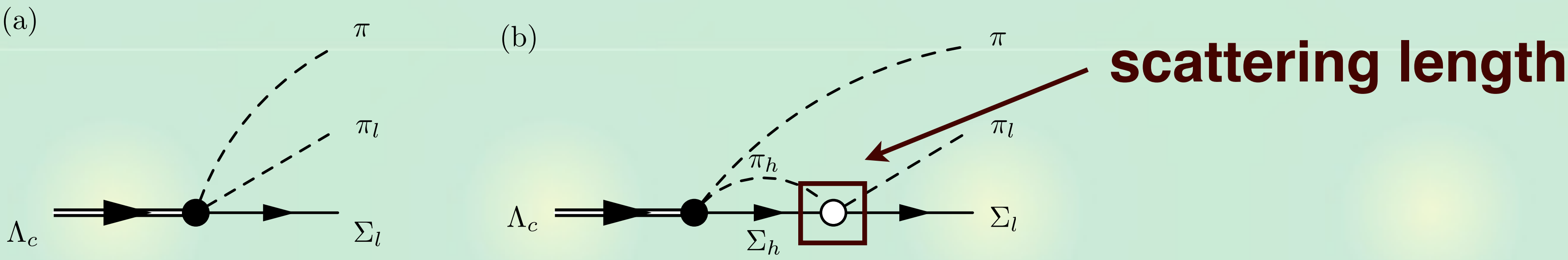
$a_{I=0} - a_{I=2}$



isospin violation
+ threshold cusp
+ amplitude interference
--> extraction of $a_{I=0} - a_{I=2}$

Decay process and loop function

Decay diagrams for Λ_c --> $\pi\ \pi_l\ \Sigma_l$ process



Spectral representation of the loop function

$$G(W) = \frac{1}{2\pi}\int_{W_{th}}^{\infty} dW' \frac{\rho(W')}{W - W' + i\epsilon} + (\text{subtractions}) \qquad \rho(W) = 2M_h \frac{q(W)}{4\pi W}$$

ρ: phase space, q: three-momentum

Imaginary part of the loop function (on-shell part):

$$\mathrm{Im}\ G(W) = -\frac{\rho(W)}{2}\Theta(W - W_{th})$$

amplitude (a) : real
amplitude (b) : real ($W < W_{th}$), complex ($W > W_{th}$)

Threshold cusp in the spectrum

Real and imaginary part of the amplitude:

$$\mathcal{M}(\Lambda_c \to \pi\pi_l\Sigma_l) = M_0(W) + \boxed{M_1(W)m_h\delta} \longleftarrow$$ **on-shell part**

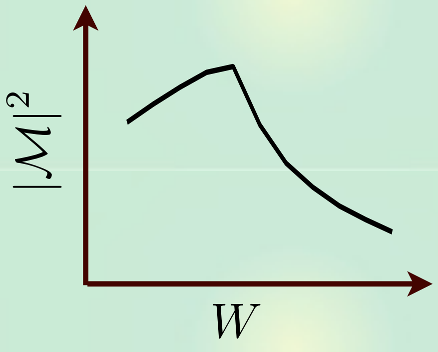
δ ~ velocity, vanish at threshold

The $\pi_l\,\Sigma_l$ invariant mass spectrum

$$|\mathcal{M}|^2 = \begin{cases} (M_0)^2 - (M_1 m_h)^2|\delta|^2 & \text{for} \quad W > W_{th} \\ (M_0)^2 + \underline{2M_0M_1m_h\delta} + (M_1m_h)^2\delta^2 & \text{for} \quad W < W_{th} \end{cases}$$

- **Spectrum is continuous (δ vanishes at threshold)**
- **Derivative of the spectrum is discontinuous**

$$\left.\frac{d|\mathcal{M}|^2}{dW}\right|_{W\to W_{th}-0} - \left.\frac{d|\mathcal{M}|^2}{dW}\right|_{W\to W_{th}+0} \propto 2M_0M_1m_h\frac{d\delta}{dW}$$



--> threshold cusp

It is purely kinematical effect. General phenomena.

M_1 amplitude: proportional to the scattering length

Determination of πΣ scattering length

Expansion of the decay spectrum --> scattering length

$$|\mathcal{M}|^2 = \begin{cases} A + B|\delta| + C|\delta|^2 + \mathcal{O}(|\delta|^3) & \text{for} \quad W > W_{th} \\ A' + B'\delta + C'\delta^2 + \mathcal{O}(\delta^3) & \text{for} \quad W < W_{th} \end{cases}$$

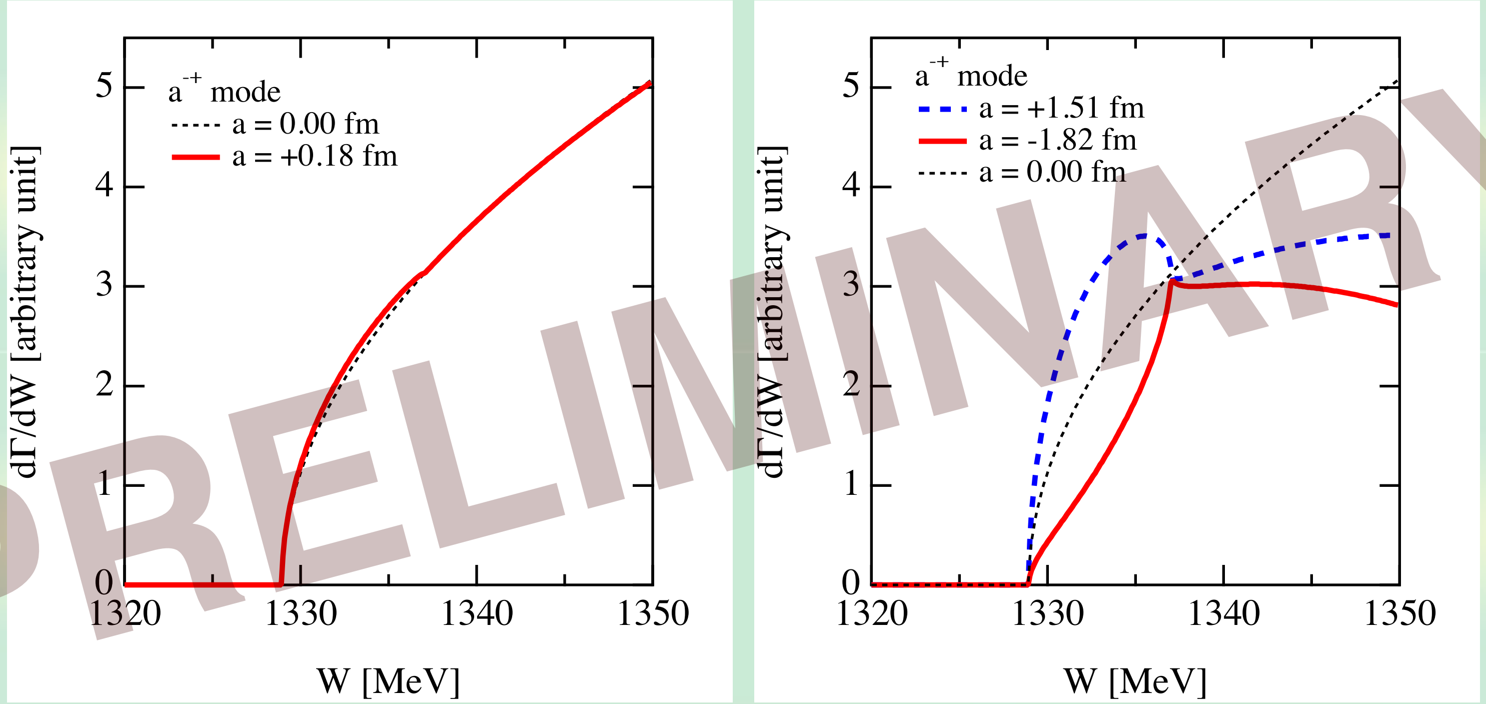
$$|a| = \frac{|B'|}{2m_h\sqrt{A}|M_0^{h(0)}|}, \qquad \frac{a}{|a|} = \frac{M_0^{(0)}}{|M_0^{(0)}|} \cdot \frac{M_0^{h(0)}}{|M_0^{h(0)}|} \cdot \frac{B'}{|B'|}$$

weak process

$$\boxed{M_0^{h(0)}, M_0^{(0)}}$$

spectrum

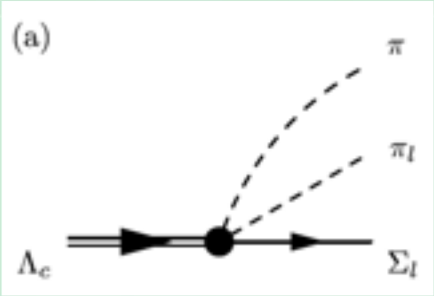
$$\boxed{A, B, C, A', B', C'}$$

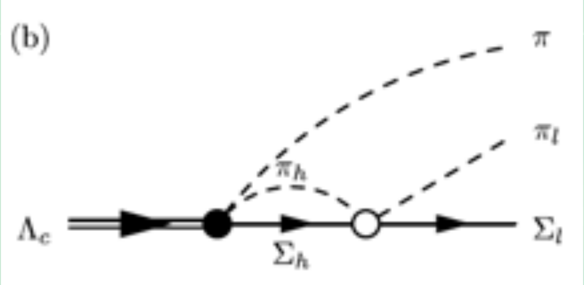


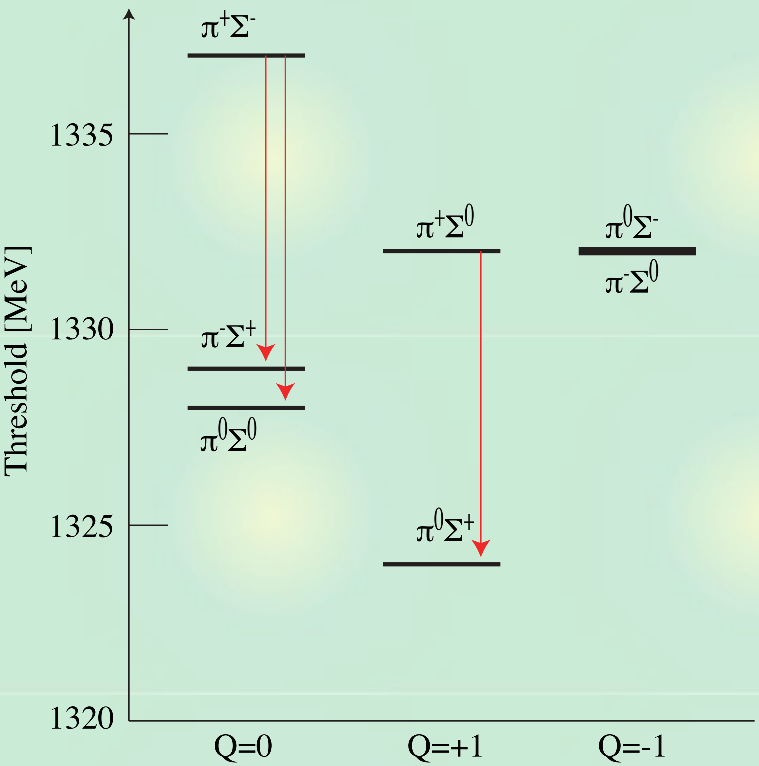
possible decay modes

Determination of πΣ scattering length

Similar approach to πΣ spectrum in $\Lambda_c \to \pi\,(\pi\Sigma)$







Σ^+(~uus) < Σ^0(~uds) < Σ^-(~dds)
--> complicated spectrum

To utilize threshold cusp, appreciable mass difference between $(\pi\Sigma)_h$ and $(\pi\Sigma)_l$ is necessary.

$$\pi^+\Sigma^- \to \pi^-\Sigma^+, \quad \pi^+\Sigma^- \to \pi^0\Sigma^0, \quad \pi^+\Sigma^0 \to \pi^0\Sigma^+,$$

Determination of πΣ scattering length

Three decay channels

$$\langle \pi^-\Sigma^+ |T| \pi^+\Sigma^- \rangle\Big|_{\text{threshold}} = \frac{1}{3}a^0 - \frac{1}{2}a^1 + \frac{1}{6}a^2 \equiv a^{-+}$$

$$\langle \pi^0\Sigma^0 |T| \pi^+\Sigma^- \rangle\Big|_{\text{threshold}} = \frac{1}{3}a^0 - \frac{1}{3}a^2 \equiv a^{00}$$

$$\langle \pi^0\Sigma^+ |T| \pi^+\Sigma^0 \rangle\Big|_{\text{threshold}} = -\frac{1}{2}a^1 + \frac{1}{2}a^2 \equiv a^{0+}$$

mode	$\Lambda_c \to \pi(\pi\Sigma)_h$	$\Lambda_c \to \pi(\pi\Sigma)_l$
a^{-+}	1.7 ± 0.5 %	3.6 ± 1.0 %
a^{00}	1.7 ± 0.5 %	1.8 ± 0.8 %
a^{0+}	1.8 ± 0.8 %	not known

A lot of Λ_c in B decay (Belle, Babar) --> feasible?

Structure around the cusp in $(\pi\Sigma)_l$ + spectrum in $(\pi\Sigma)_h$
--> extraction of the scattering length

Three unknown scattering lengths, two constraints

$$a^{-+} - a^{00} = a^{0+}$$

I=2 scattering length: lattice QCD --> Y. Ikeda et al

Summary

$\pi\Sigma$ scattering length from Λ_c decay

- **$\pi\Sigma$ scattering length is important for low energy KN-$\pi\Sigma$ amplitude.**
  --> K nuclei and Λ(1405) physics

  Y. Ikeda, T. Hyodo, D. Jido, H. Kamano, T. Sato, K. Yazaki, arXiv:1101.5190, to appear in Prog. Theor. Phys.

- **Threshold cusp in Λ_c --> $\pi\pi\Sigma$ decay is related with the $\pi\Sigma$ scattering length.**

  T. Hyodo, M. Oka, work in progress

- **3 isospin states v.s. 2 decay modes: Lattice QCD can help to complete.**

  Y. Ikeda, HAL QCD collaboration, work in progress