

Ising model on random networks and canonical tensor model

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N. Sasakura, Y.S., PTEP 2014 (2014) 5, 053B03.

N. Sasakura, Y.S., 1402.0740 [hep-th] (accepted in SIGMA).

What happens when combining Quantum Mechanics and General Relativity?

[Quantum mechanics (\hbar)] \rightarrow

$$(\Delta X) \cdot (M \Delta V) \geq \hbar/2$$

[Special Relativity (c)] \rightarrow

$$\Delta V \leq c$$

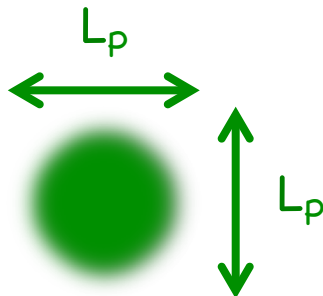
$$(\Delta X) \geq \hbar/(2Mc)$$

Compton wavelength

[General Relativity (G_N)] \rightarrow $M \leq \Delta X c^2 / (2G_N)$

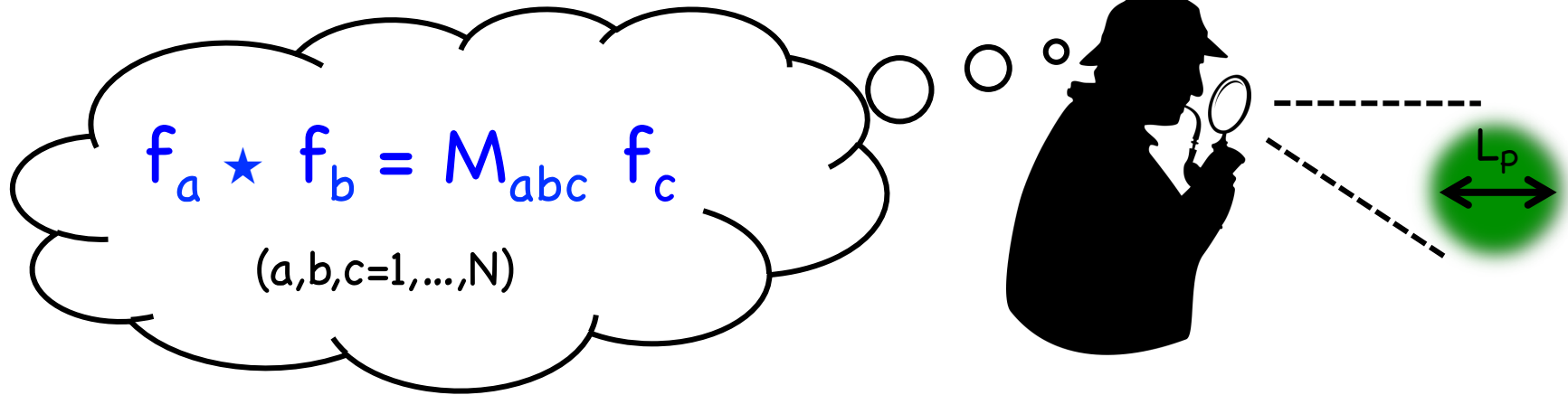
$$(\Delta X) \geq [\hbar G_N / c^3]^{1/2} = L_p$$

Planck length



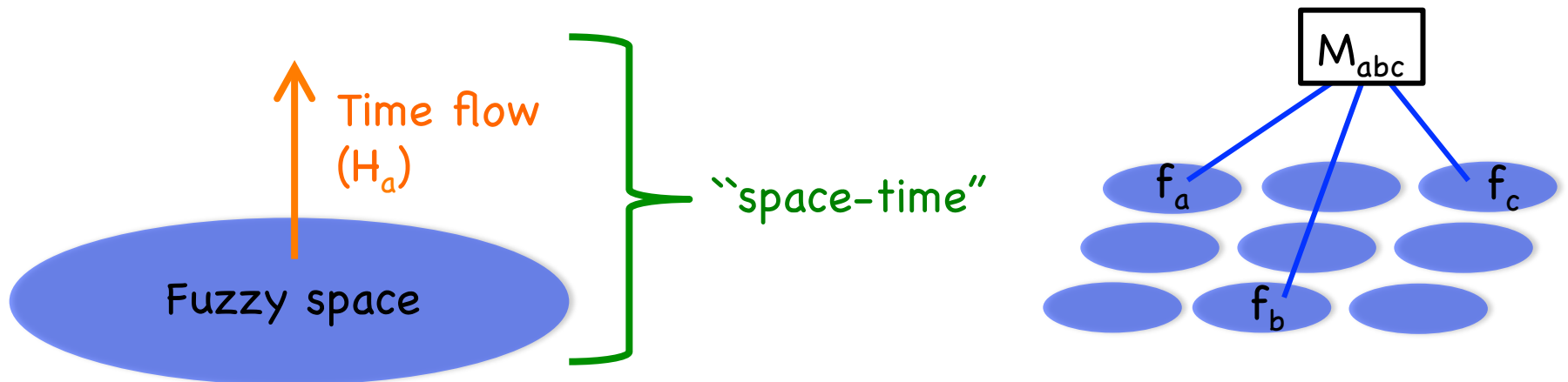
We are ignorant of a length shorter than L_p

Canonical Tensor Model = Theory of dynamical **fuzzy space** [N. Sasakura, 2011]



Fuzzy space is defined by functions $\{f_1, f_2, \dots, f_N\}$ and the product, $f_a \star f_b$.

$N = \#[\text{functions}] = \#[\text{"points" in space}]$



Canonical Tensor Model [N. Sasakura, 2011]

Our knowledge is quite limited...

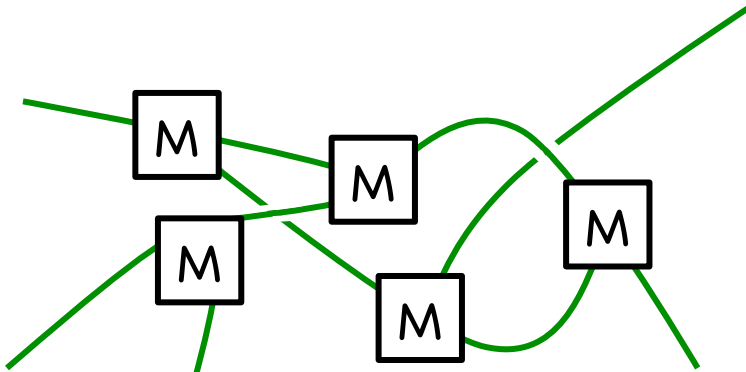
N=1: General Relativity in minisuperspace [N. Sasakura and Y.S., 2014]

N=2: Locality is favoured [N. Sasakura, 2013]

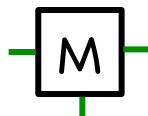
Difficulty is tied to much DOF of M_{abc} ($a=1,\dots,N$).

Overcome

Statistical system on random networks [N. Sasakura and Y.S., 2014]

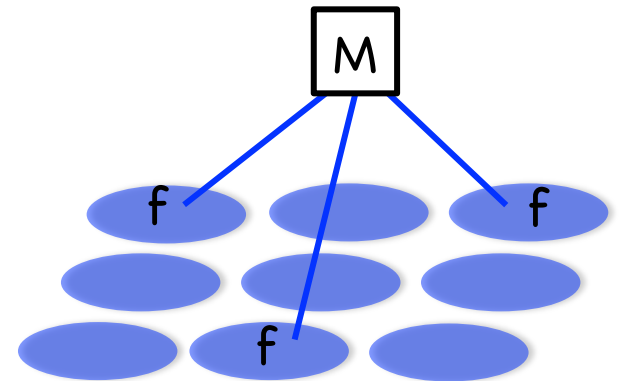


Statistical system on random network

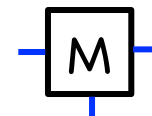


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Canonical tensor model



Same symmetry

Canonical Tensor Model [N. Sasakura, 2011]

Our knowledge is quite limited...

N=1: General Relativity in minisuperspace [N. Sasakura and Y.S., 2014]

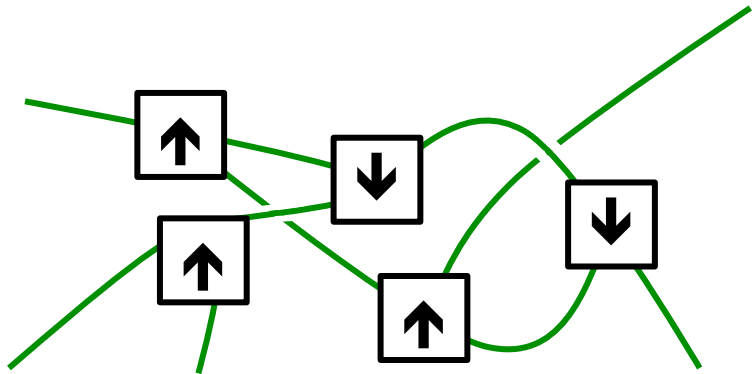
N=2: Locality is favoured [N. Sasakura, 2013]

Difficulty is tied to much DOF of M_{abc} ($a=1,\dots,N$).



Overcome

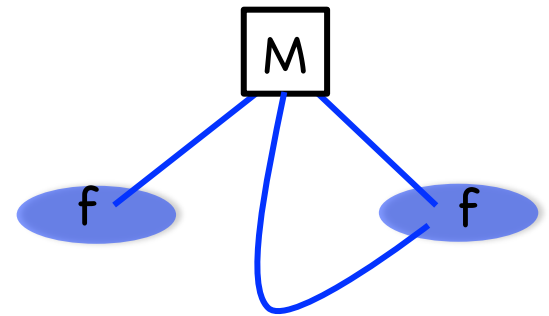
Statistical system on random networks [N. Sasakura and Y.S., 2014]



Ising model on random network

Phase transition line

"="



N=2 canonical tensor model

Boundary of Hamiltonian flow

match



Canonical Tensor Model [N. Sasakura, 2011]

Hamiltonian:

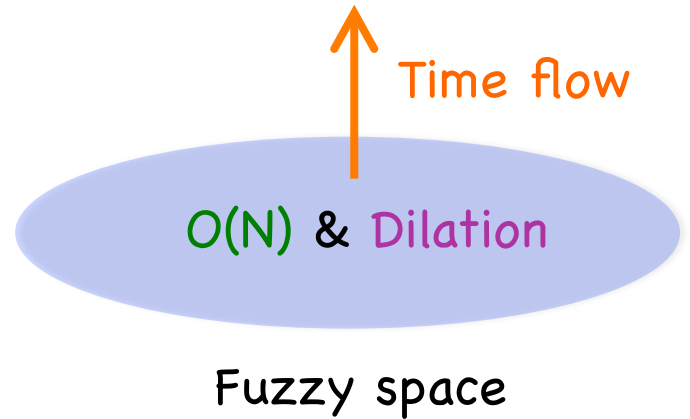
$$H = N_a H_a + N_{[ab]} J_{[ab]} + ND$$

where

$$H_a = (1/2) P_{abc} P_{bde} M_{cde} \quad \text{Generator for time flow}$$

$$J_{[ab]} = (P_{acd} M_{bcd} - P_{bcd} M_{acd})/4 \quad o(N) \text{ generator}$$

$$D = (1/6) M_{abc} P_{abc} \quad \text{Dilation generator}$$



Poisson's bracket:

$$\{M_{abc}, P_{def}\} = \delta_{ad} \delta_{be} \delta_{cf} + \text{perm. of (def)}$$

Hamiltonian vector flow of N=2 model [N. Sasakura and Y.S., 2014]

Gauge fixing:

$$P_{111} = 1, \quad P_{112} = 0, \quad P_{122} = x, \quad P_{222} = y$$

[(M,P) are real & symmetric]

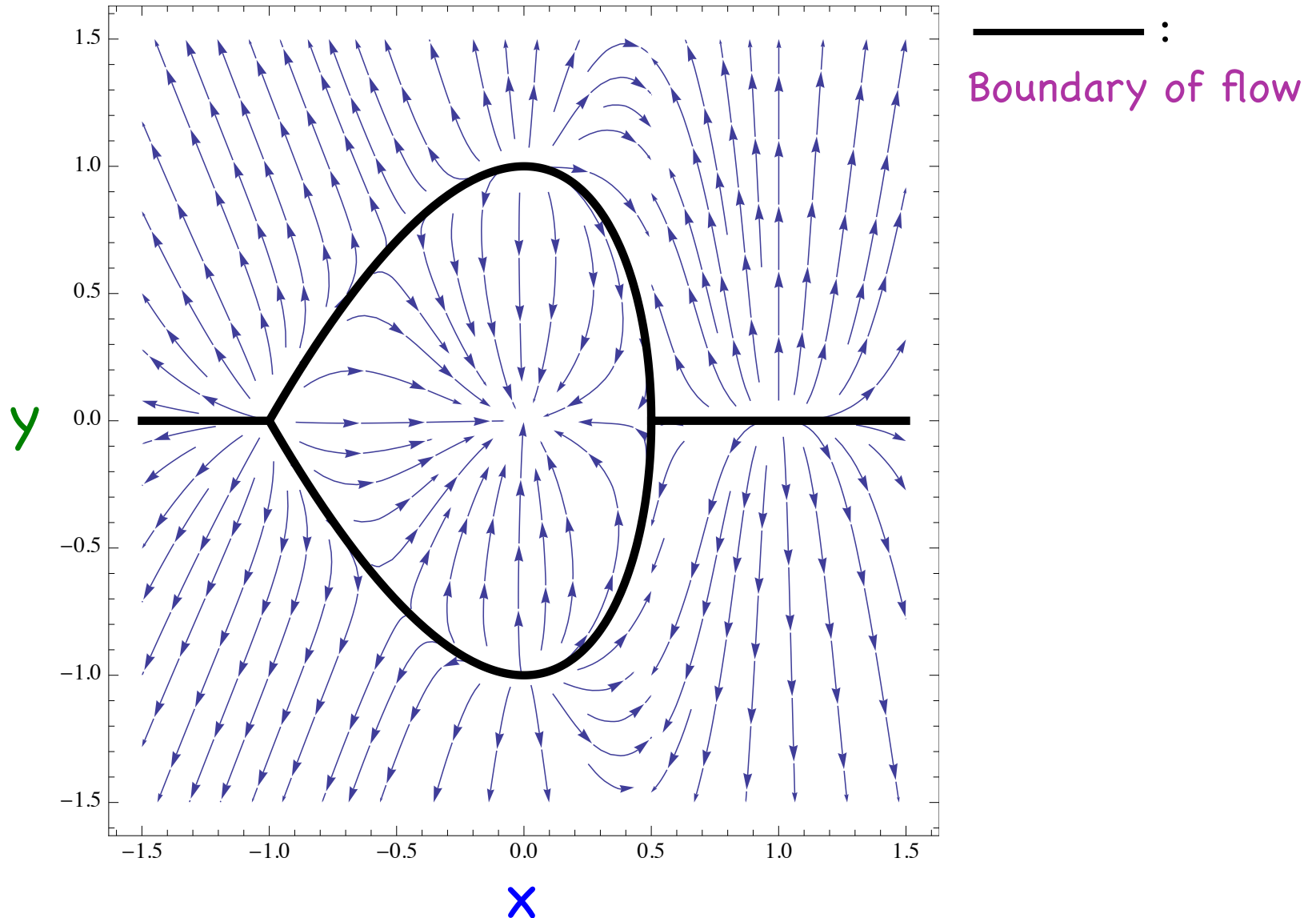
Hamiltonian vector flow:

$$H \propto A(x,y) \frac{\partial}{\partial x} + B(x,y) \frac{\partial}{\partial y}$$

[A(x,y) & B(x,y) are functions]

Hamiltonian vector flow of N=2 model [N. Sasakura and Y.S., 2014]

$$H \propto A(x,y) \frac{\partial}{\partial x} + B(x,y) \frac{\partial}{\partial y}$$



Statistical systems on random network [N. Sasakura and Y.S., 2014]

Grand-type partition function:

$$Z(M, t) = \int \prod_{d=1}^N d\Phi_d \exp[-\Phi_a \Phi_a + t M_{abc} \Phi_a \Phi_b \Phi_c]$$

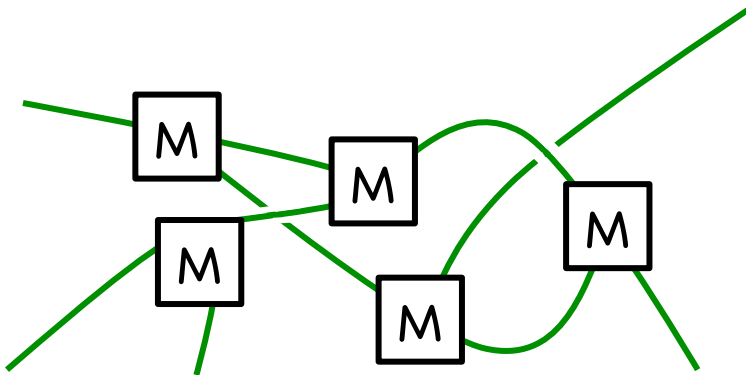
$O(N)$ and dilation invariance:

$$Z(L(M), t) = Z(M, t)$$

$$L(M)_{ace} = L_{a\acute{a}} L_{c\acute{c}} L_{e\acute{e}} M_{\acute{a}\acute{c}\acute{e}}$$

$$Z(e^\sigma M, e^{-\sigma} t) = Z(M, t)$$

Generate statistical system (M) on random network (Feynman graph):



Statistical systems on random network [N. Sasakura and Y.S., 2014]

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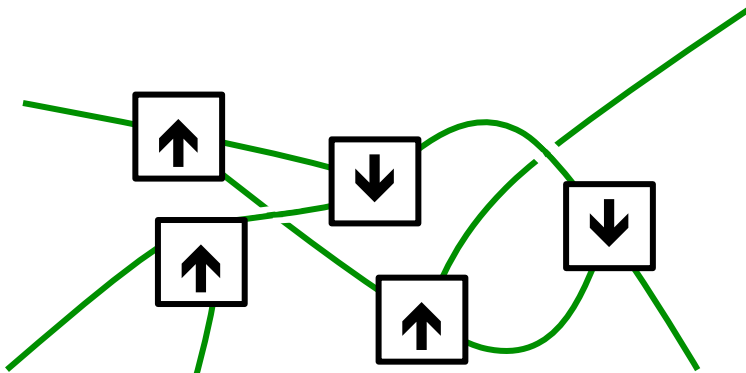
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Generate statistical system (M) on random network (Feynman graph):



Ex.) $N=2$ (Ising model)

$$M_{abc} = \sum_i R_{ai} R_{bi} R_{ci} e^{H\sigma_i}$$

$$(R^T R)_{ij} = e^{J\sigma_i \sigma_j} \quad (J \geq 0)$$

Statistical systems on random network [N. Sasakura and Y.S., 2014]

To check thermodynamic properties, introduce **partition function**:

$$\begin{aligned} Z(M, t) &\cong \sum_{n=0}^{\infty} t^n Z_n(M) \\ &= \sum_{n=0}^{\infty} t^n \left[(1/n!) \int \prod_{d=1}^N d\Phi_d (M_{abc} \Phi_a \Phi_b \Phi_c)^n e^{-\Phi_e \Phi_e} \right] \end{aligned}$$

Define the free energy per vertex (n):

$$f_{2n} = -(1/2n) \log[Z_{2n}(M)] \quad [Z_n \text{ vanishes for } n \text{ odd}]$$

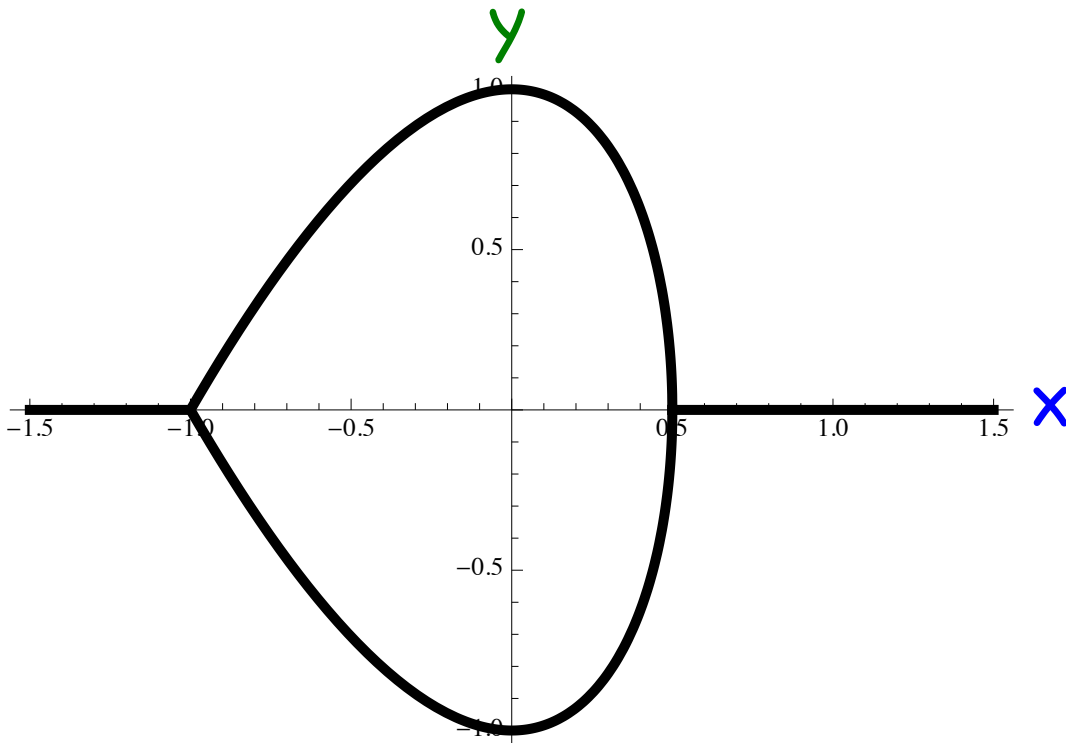
For $N=2$ (Ising model), check the behaviour of f_{2n} in the thermodynamic limit, $n \rightarrow \infty$.

$$M_{111} = 1, \quad M_{112} = 0, \quad M_{122} = x, \quad M_{222} = y \quad (\text{Gauge fixing})$$

Statistical systems on random network [N. Sasakura and Y.S., 2014]

$$f_{2n}(x, y) = -(1/2n) \log[Z_{2n}(x, y)]$$

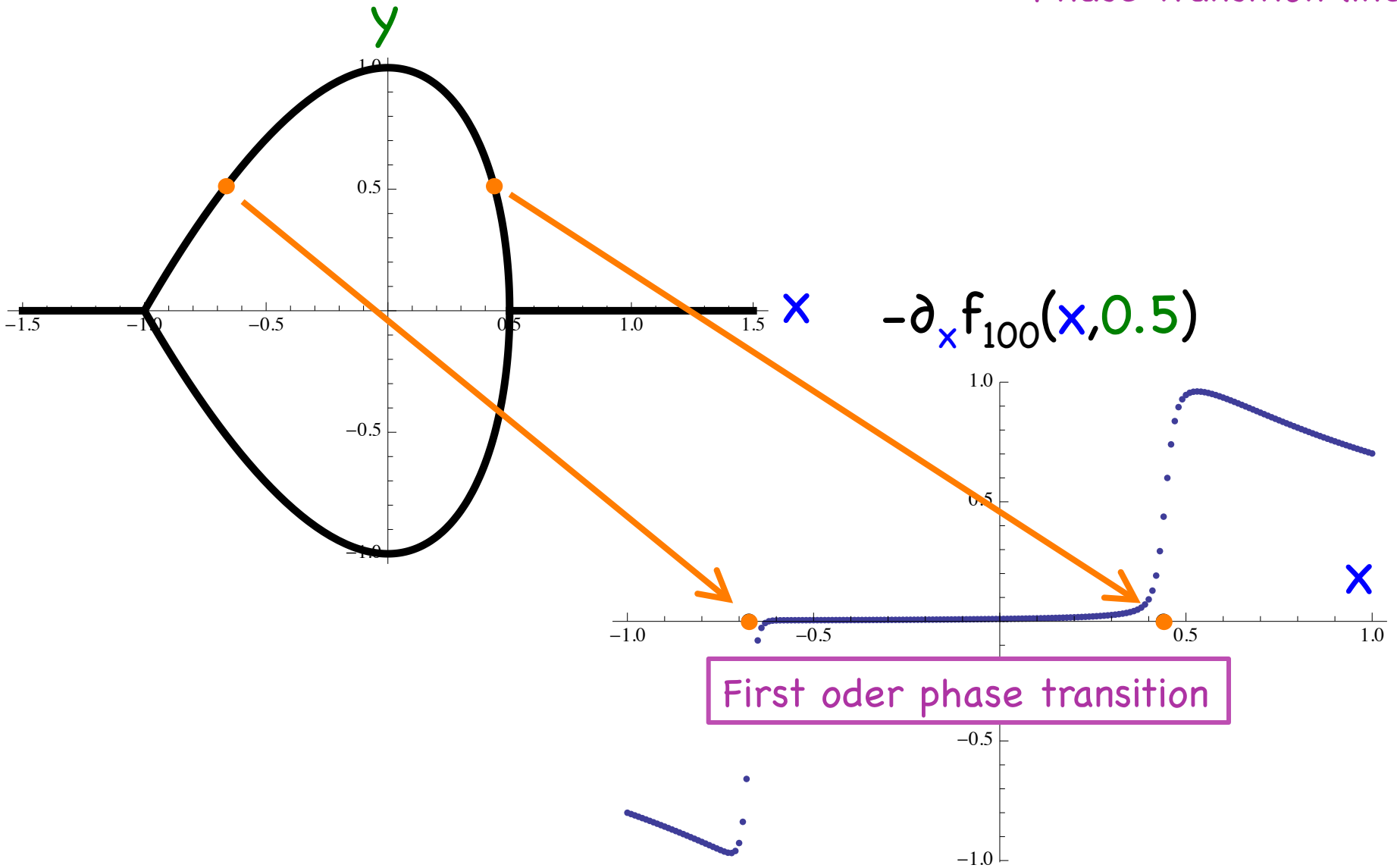
————— :
Phase transition line



Statistical systems on random network [N. Sasakura and Y.S., 2014]

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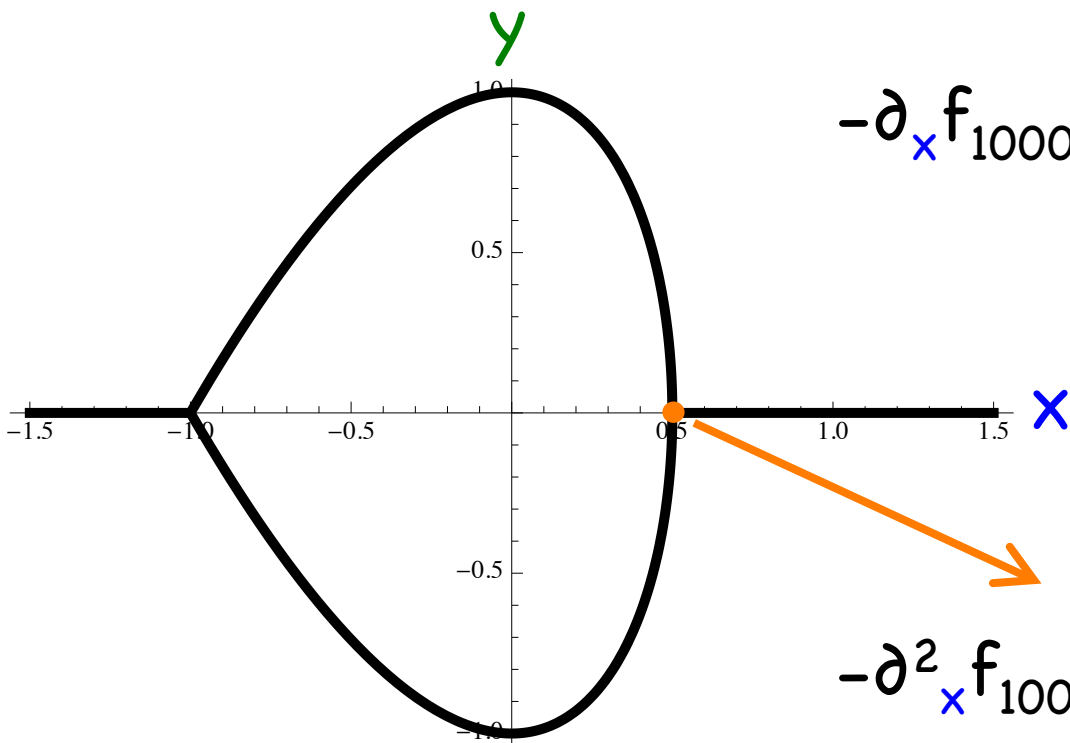
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Phase transition line



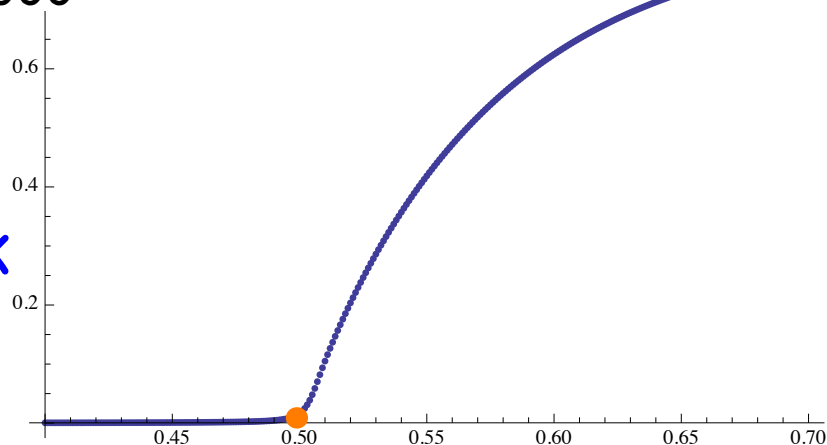
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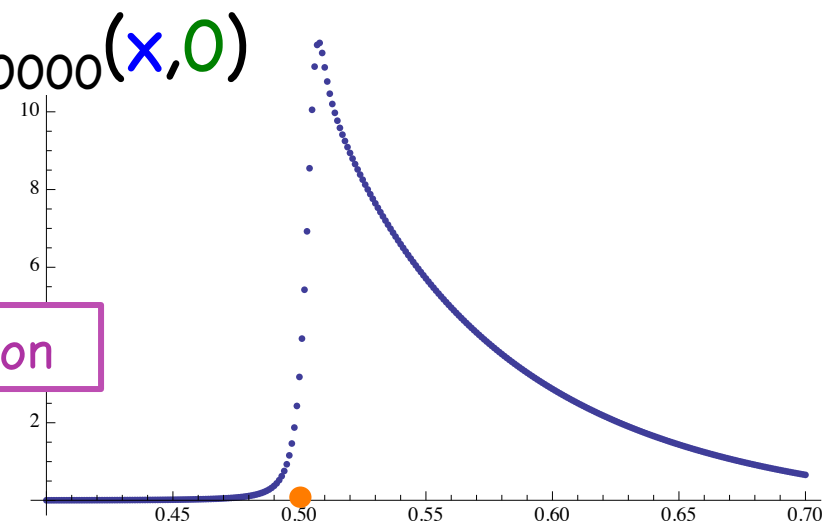
— :
Phase transition line



$$-\partial_x f_{100000}(x, 0)$$



$$-\partial_x^2 f_{100000}(x, 0)$$



Second order phase transition

Conclusion

[N=2 canonical tensor model]

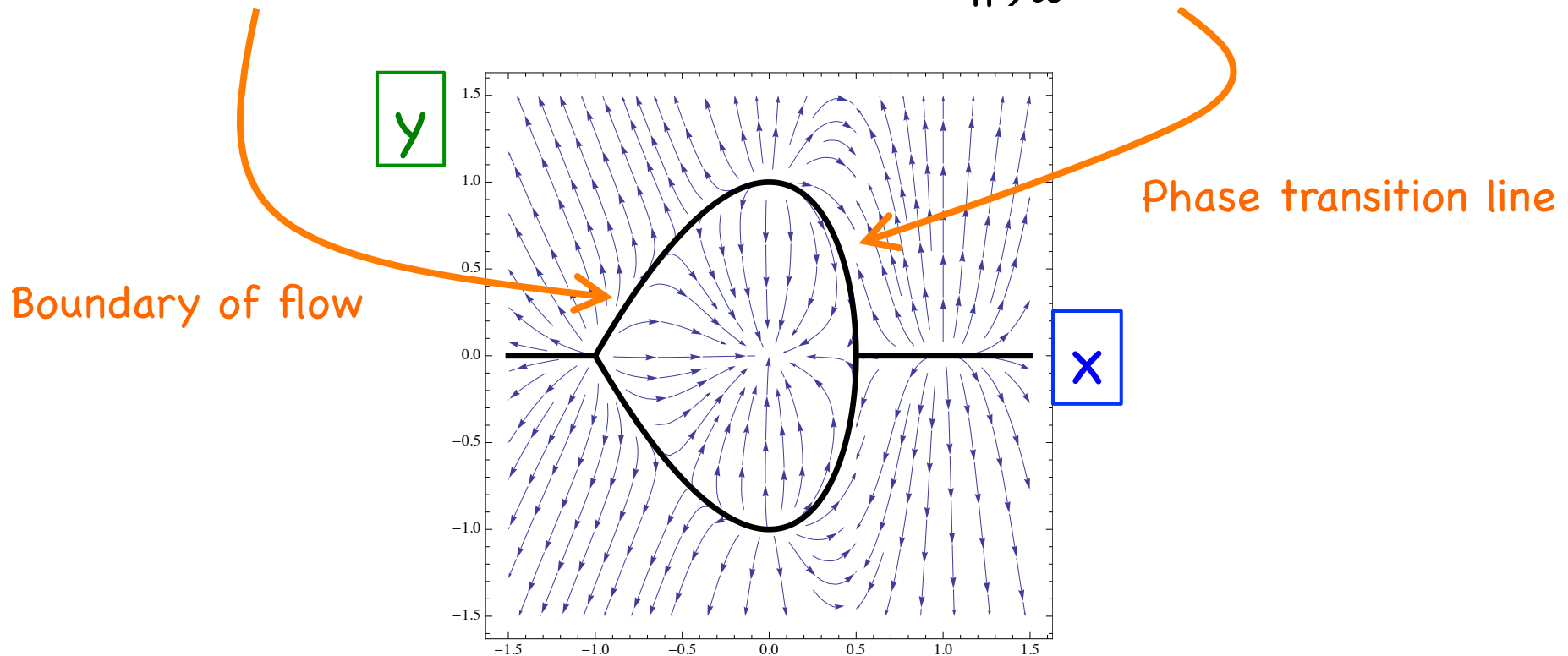
Hamiltonian flow:

$$H \propto A(x,y) \partial/\partial x + B(x,y) \partial/\partial y$$

[Ising model on random networks]

Free energy:

$$f(x,y) = \lim_{n \rightarrow \infty} -(1/2n) \log[Z_{2n}(x,y)]$$



Implication:

RG-like procedure may exist in Ising model on random networks

Large-N (#[points] are large) canonical tensor model
can be described by statistical systems on Random networks?

Works!



Exact solutions to the Wheeler-DeWitt equation
can be found in the language of random networks
[work in progress with G. Narain and N. Sasakura]



$$H_a \psi = J_{[ab]} \psi = D \psi = 0$$

