

Entanglement Behavior of 2D Quantum Models

Shu Tanaka (YITP, Kyoto University)

Collaborators:

Hosho Katsura (Univ. of Tokyo, Japan)

Anatol N. Kirillov (RIMS, Kyoto Univ., Japan)

Vladimir E. Korepin (YITP, Stony Brook, USA)

Naoki Kawashima (ISSP, Univ. of Tokyo, Japan)

Lou Jie (Fudan Univ., China)

Ryo Tamura (NIMS, Japan)



VBS on symmetric graphs, J. Phys. A, 43, 255303 (2010)

“VBS/CFT correspondence”, Phys. Rev. B, 84, 245128 (2011)

Quantum hard-square model, Phys. Rev. A, 86, 032326 (2012)

Nested entanglement entropy, Interdisciplinary Information Sciences, 19, 101 (2013)



京都大学
KYOTO UNIVERSITY



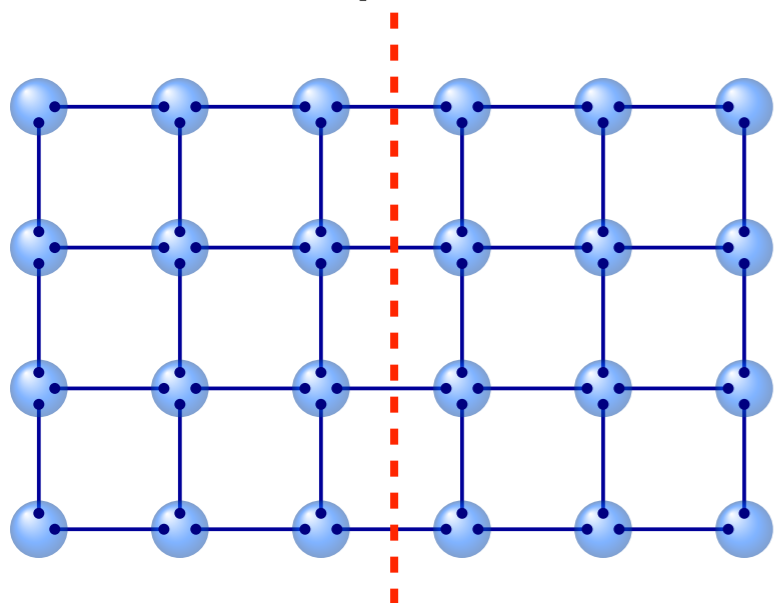
Digest

Entanglement properties of 2D quantum systems

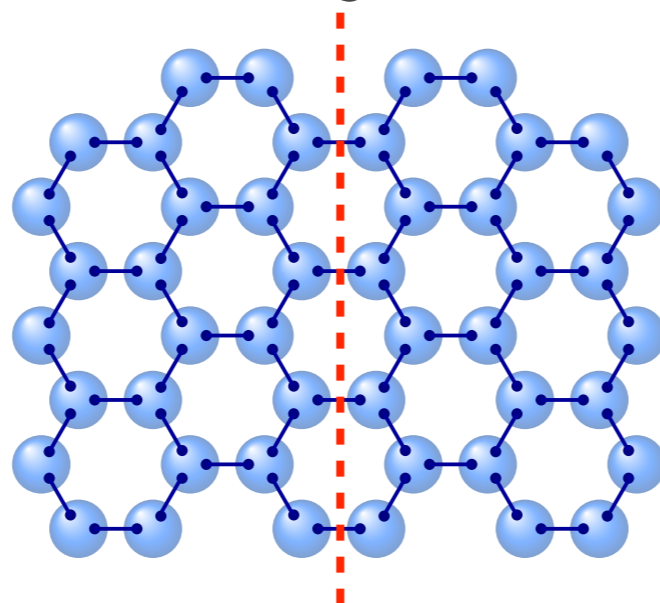
Physical properties of 1D quantum systems



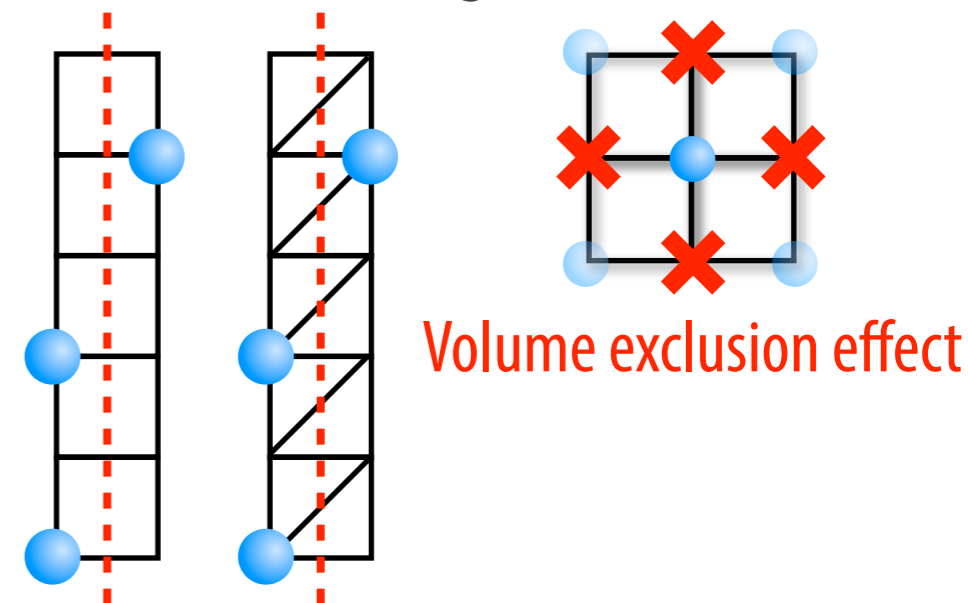
VBS on square lattice



VBS on hexagonal lattice



Quantum lattice gas on ladder



VBS state on 2D lattice

Total system	Entanglement Hamiltonian
Square lattice	1D AF Heisenberg
Hexagonal lattice	1D F Heisenberg

Quantum lattice gas on ladder

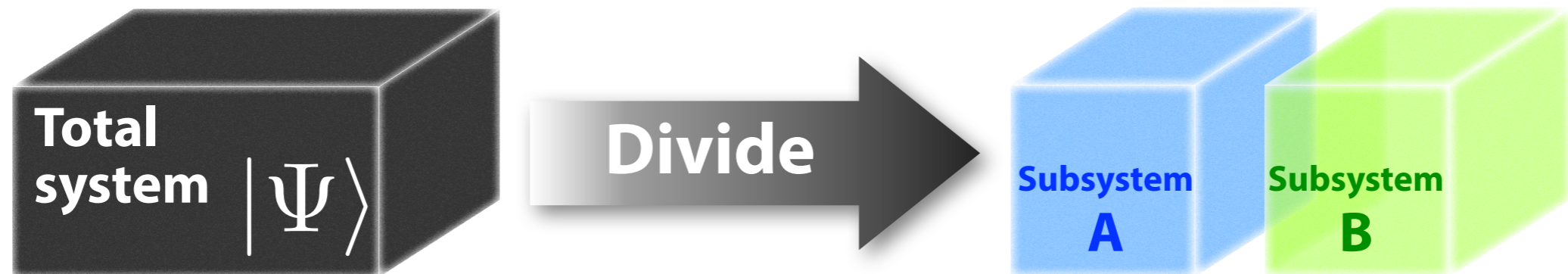
Total system	Entanglement Hamiltonian
Square ladder	2D Ising
Triangle ladder	2D 3-state Potts

Introduction

- ***Entanglement***
- ***Motivation***
- ***Preliminaries***

Introduction

EE is a measure to quantify entanglement.



Schmidt decomposition

$$|\Psi\rangle = \sum_{\alpha} \lambda_{\alpha} |\phi_{\alpha}^{[A]}\rangle \otimes |\phi_{\alpha}^{[B]}\rangle$$



Normalized GS

$$\phi_{\alpha}^{[A]} \in \mathcal{H}_A, \phi_{\alpha}^{[B]} \in \mathcal{H}_B$$

$\{|\phi_{\alpha}^{[A]}\rangle\}, \{|\phi_{\alpha}^{[B]}\rangle\}$: Orthonormal basis

Reduced density matrix

$$\rho_A = \text{Tr}_B |\Psi\rangle\langle\Psi| = \sum_{\alpha} \lambda_{\alpha}^2 |\phi_{\alpha}^{[A]}\rangle\langle\phi_{\alpha}^{[A]}|$$

von Neumann entanglement entropy

$$\mathcal{S} = \text{Tr} \rho_A \ln \rho_A = - \sum_{\alpha} \lambda_{\alpha}^2 \ln \lambda_{\alpha}^2$$

Introduction

Entanglement properties in **1D** quantum systems!!

1D **gapped** systems: EE **converges** to some value.

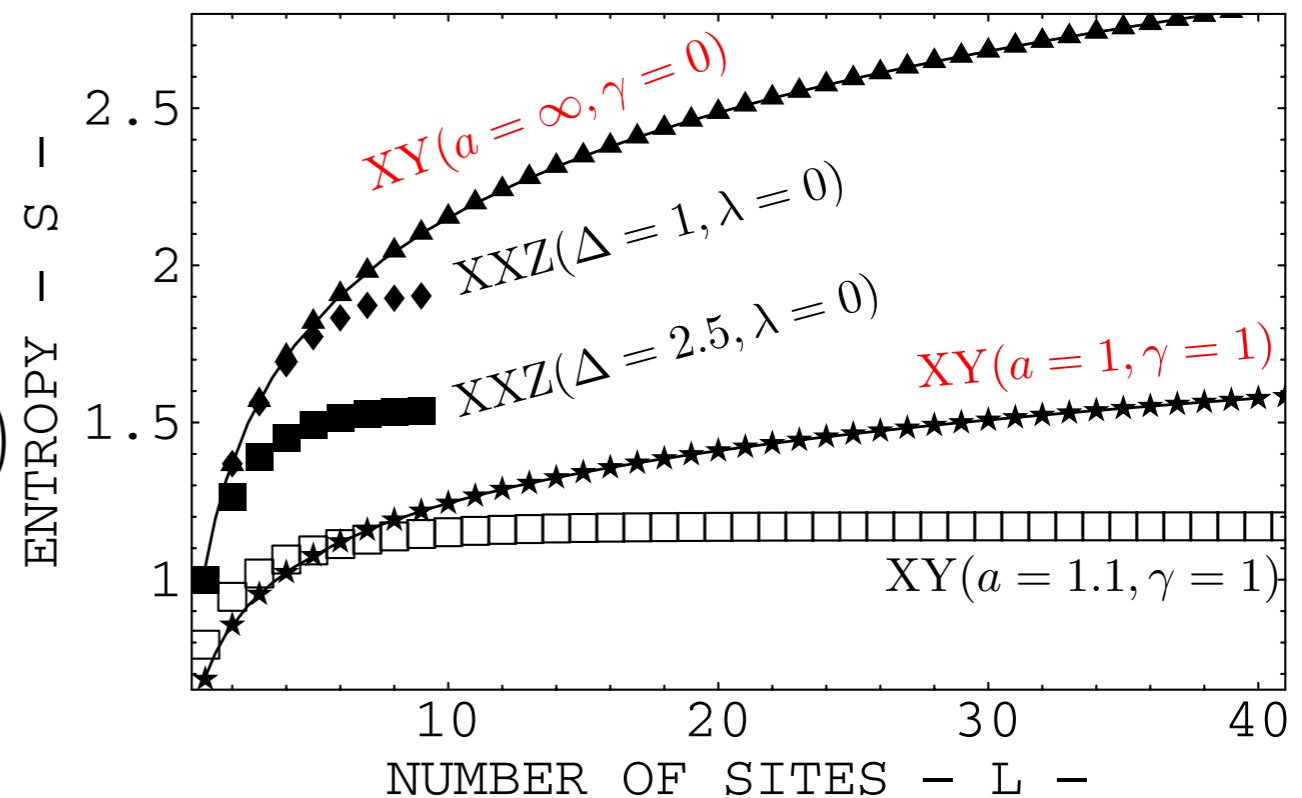
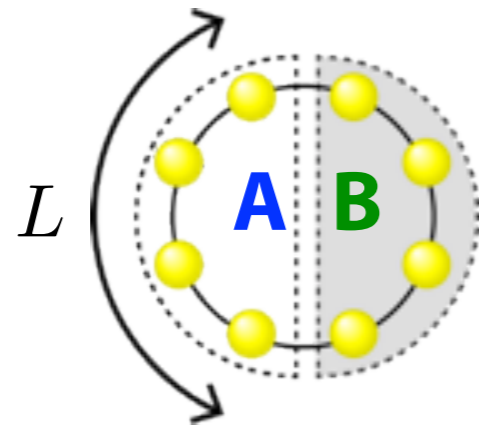
1D **critical** systems: EE **diverges** logarithmically with L .
coefficient is related to the central charge.

XXZ model under magnetic field

$$\mathcal{H}_{\text{XXZ}} = \sum_i (\sigma_i^x \sigma_{i+1}^x + \sigma_i^y \sigma_{i+1}^y + \Delta \sigma_i^z \sigma_{i+1}^z - \lambda \sigma_i^z)$$

XY model under magnetic field

$$\mathcal{H}_{\text{XY}} = - \sum_{i=0}^{N-1} \left(\frac{a}{2} [(1 + \gamma) \sigma_i^x \sigma_{i+1}^x + (1 - \gamma) \sigma_i^y \sigma_{i+1}^y] + \sigma_i^z \right)$$



G. Vidal et al. PRL 90, 227902 (2003)

Entanglement properties in **2D** quantum systems??

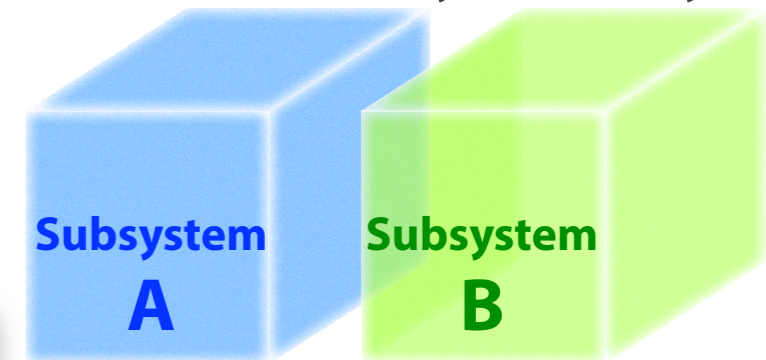
Preliminaries: reflection symmetric case

Pre-Schmidt decomposition

$$|\Psi\rangle = \sum_{\alpha} |\phi_{\alpha}^{[A]}\rangle \otimes |\phi_{\alpha}^{[B]}\rangle \quad \{|\phi_{\alpha}^{[A]}\rangle\}, \{|\phi_{\alpha}^{[B]}\rangle\}$$

Linearly independent
(but not orthonormal)

Reflection symmetry



Overlap matrix

$$(M^{[A]})_{\alpha\beta} := \langle \phi_{\alpha}^{[A]} | \phi_{\beta}^{[A]} \rangle, \quad (M^{[B]})_{\alpha\beta} := \langle \phi_{\alpha}^{[B]} | \phi_{\beta}^{[B]} \rangle$$

Reflection symmetry $\longrightarrow M^{[A]} = M^{[B]} = M$

Useful fact

If $M^{[A]} = M^{[B]} = M$ and M is real symmetric matrix,

$$\mathcal{S} = - \sum_{\alpha} p_{\alpha} \ln p_{\alpha}, \quad p_{\alpha} = \frac{d_{\alpha}^2}{\sum_{\alpha} d_{\alpha}^2}$$

where d_{α} are the eigenvalues of M .

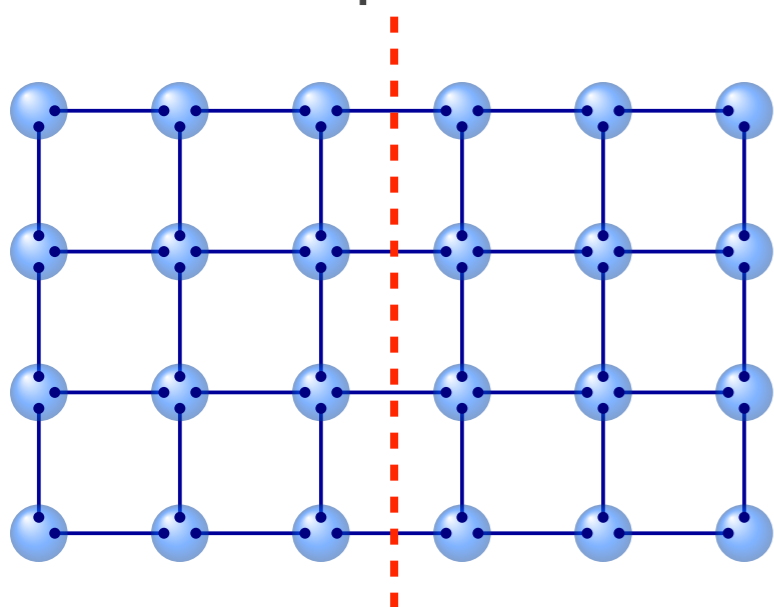
Digest

Entanglement properties of 2D quantum systems

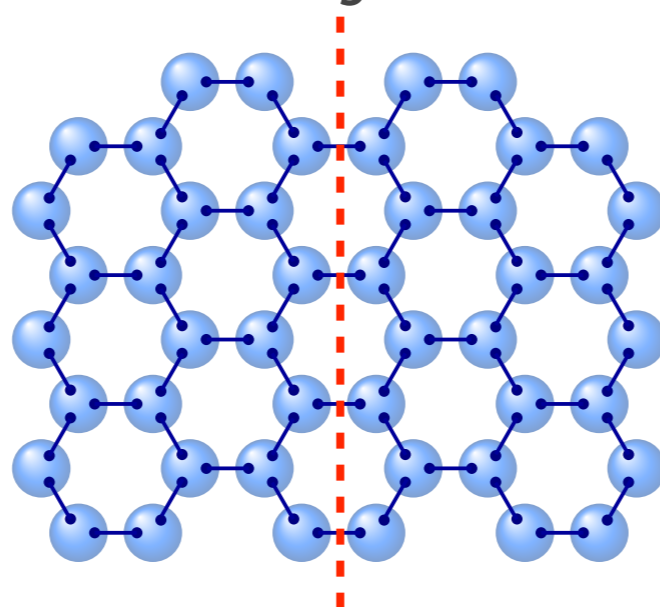


Physical properties of 1D quantum systems

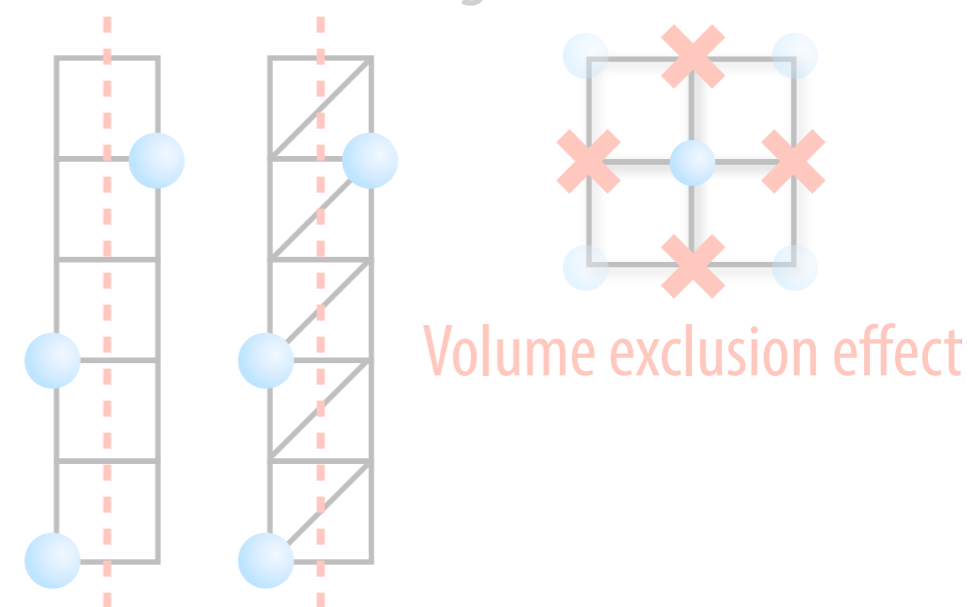
VBS on square lattice



VBS on hexagonal lattice



Quantum lattice gas on ladder



VBS state on 2D lattice

Total system	Entanglement Hamiltonian
Square lattice	1D AF Heisenberg
Hexagonal lattice	1D F Heisenberg

Quantum lattice gas on ladder

Total system	Entanglement Hamiltonian
Square ladder	2D Ising
Triangle ladder	2D 3-state Potts

VBS (Valence-Bond-Solid) state

Valence bond = Singlet pair

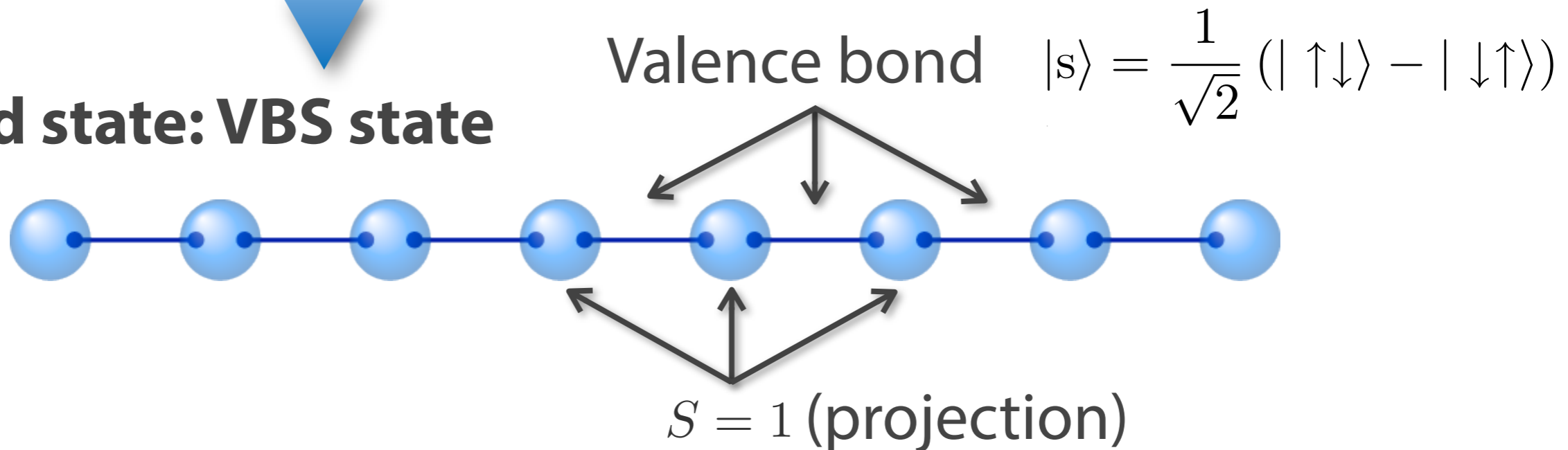
$$|s\rangle = \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle)$$

AKLT (Affleck-Kennedy-Lieb-Tasaki) model

*I. Affleck, T. Kennedy, E. Lieb, and H. Tasaki, PRL **59**, 799 (1987).*

$$\mathcal{H} = \sum_i \left[\vec{S}_i \cdot \vec{S}_{i+1} + \frac{1}{3} \left(\vec{S}_i \cdot \vec{S}_{i+1} \right)^2 \right] \quad (S = 1)$$

Ground state: VBS state

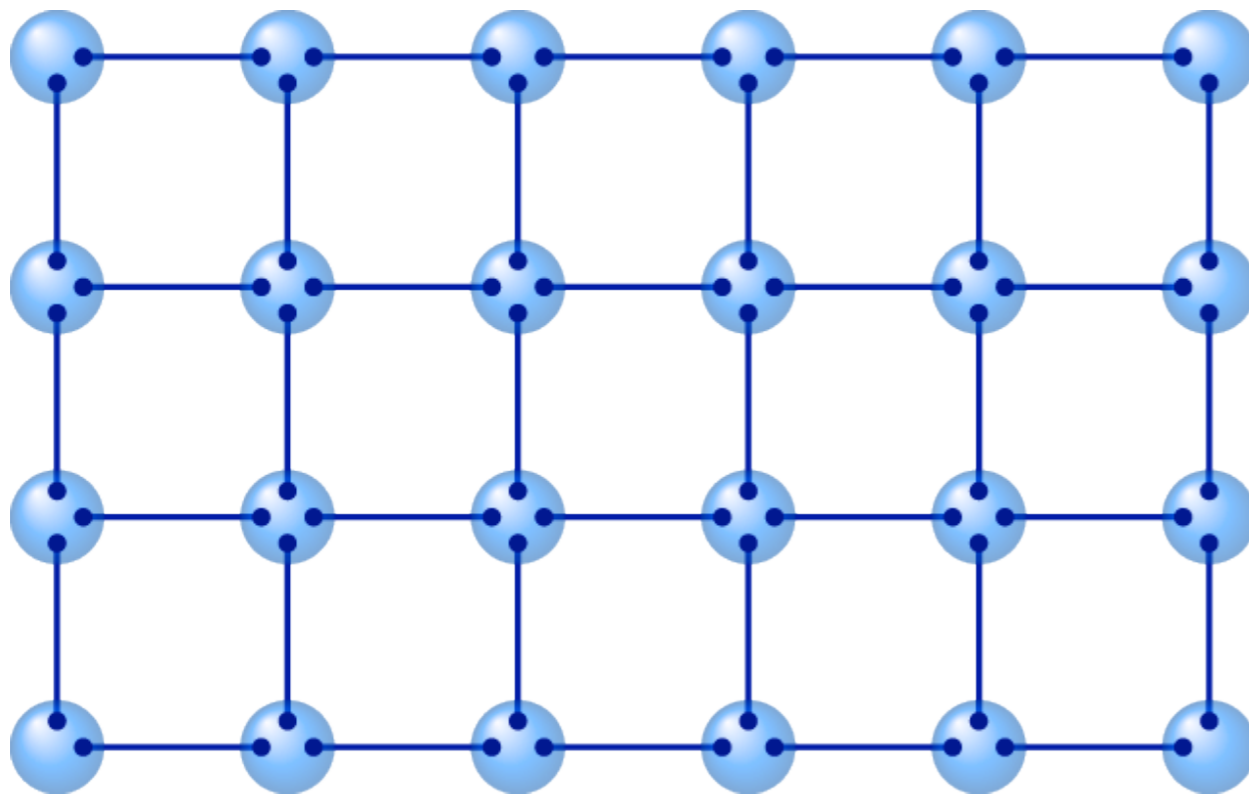


- Exact unique ground state; $S=1$ VBS state
- Rigorous proof of the "Haldane gap"
- AFM correlation decays fast exponentially

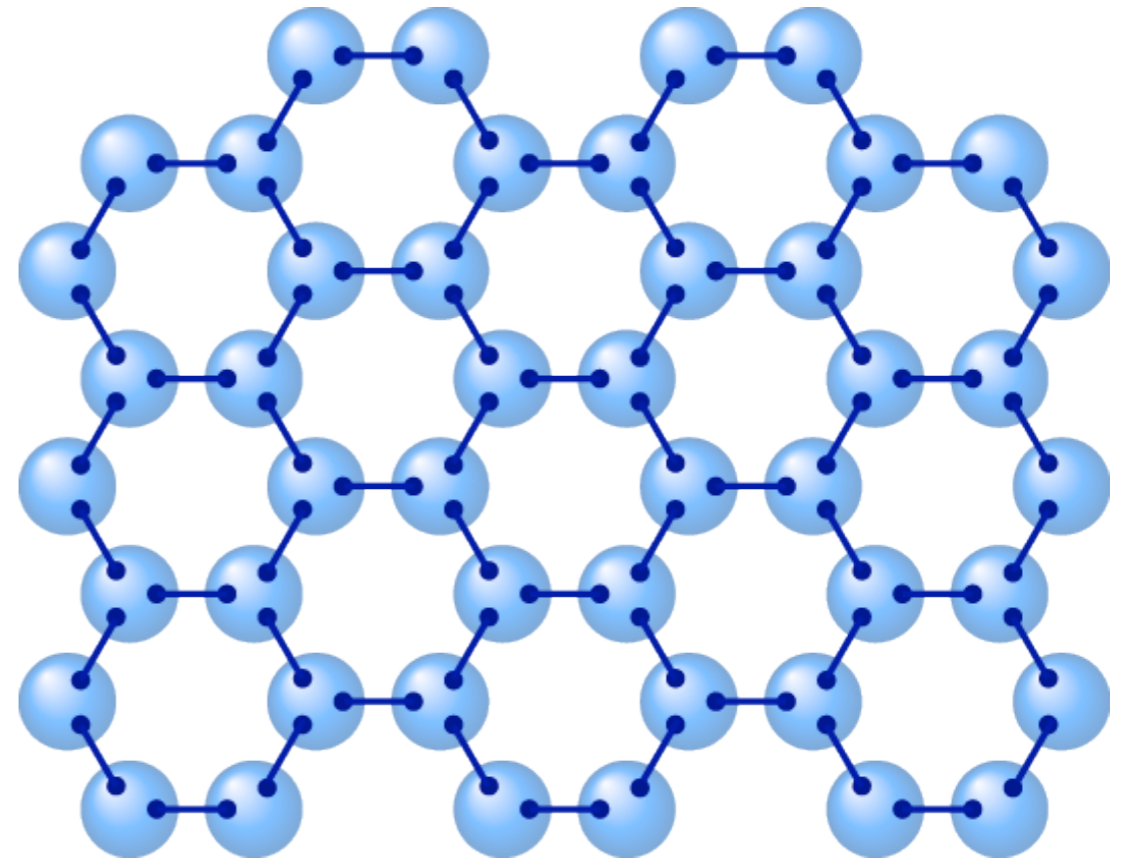
VBS (Valence-Bond-Solid) state

VBS state = Singlet-covering state

2D square lattice



2D hexagonal lattice



MBQC using VBS state

*T-C. Wei, I. Affleck, and R. Raussendorf, Phys. Rev. Lett. **106**, 070501 (2011).*

*A. Miyake, Ann. Phys. **326**, 1656 (2011).*

VBS (Valence-Bond-Solid) state

VBS state = Singlet-covering state

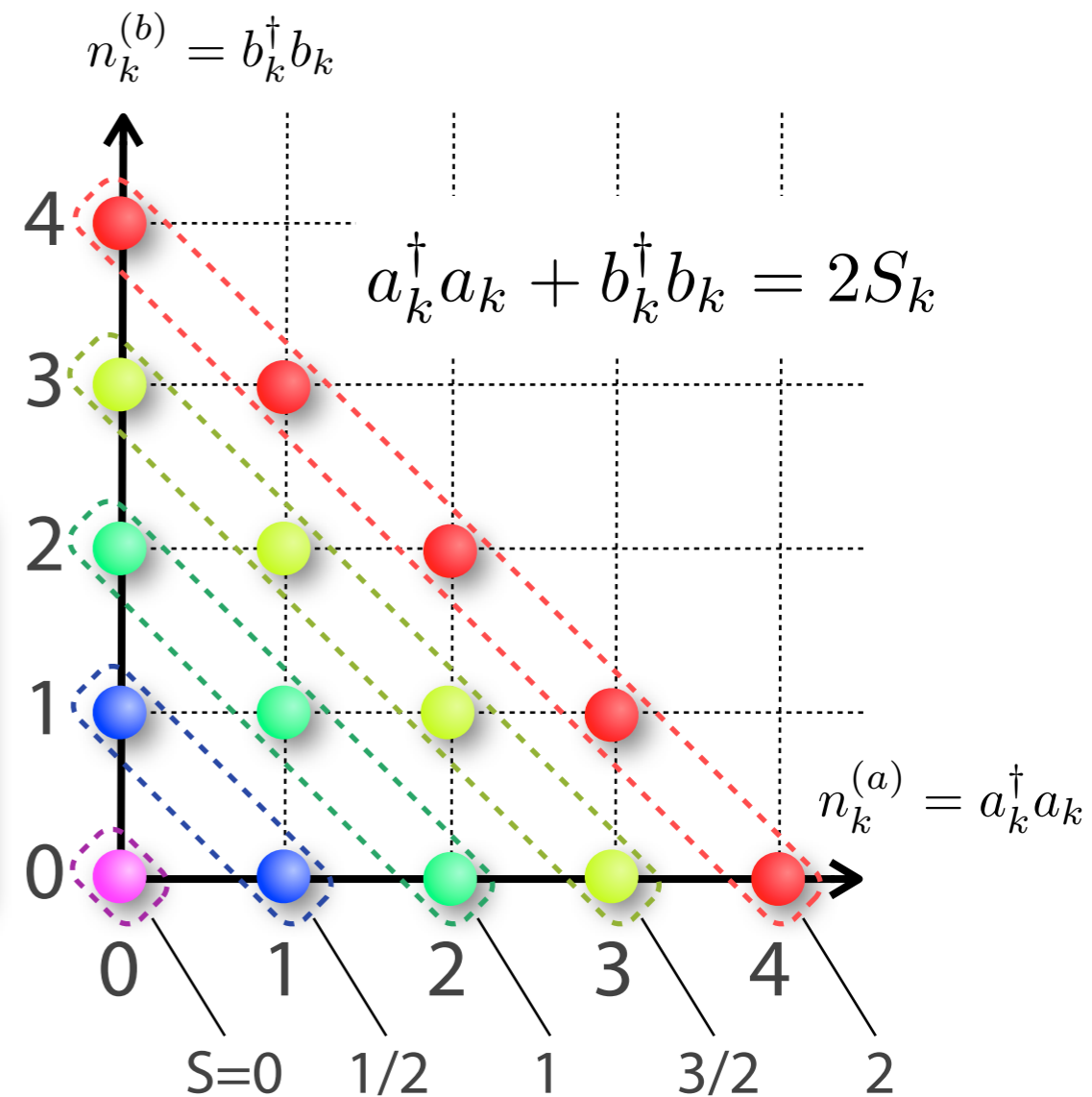
Schwinger boson representation

$$|\uparrow\rangle = a^\dagger |\text{vac}\rangle, \quad |\downarrow\rangle = b^\dagger |\text{vac}\rangle$$



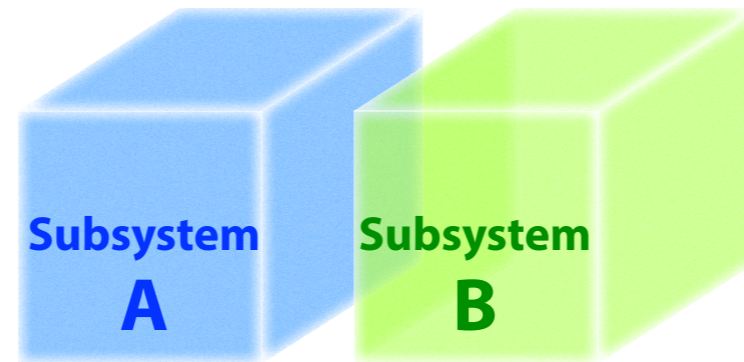
Valence bond solid (VBS) state

$$|\text{VBS}\rangle = \prod_{\langle k,l \rangle} (a_k^\dagger b_l^\dagger - b_k^\dagger a_l^\dagger) |\text{vac}\rangle$$

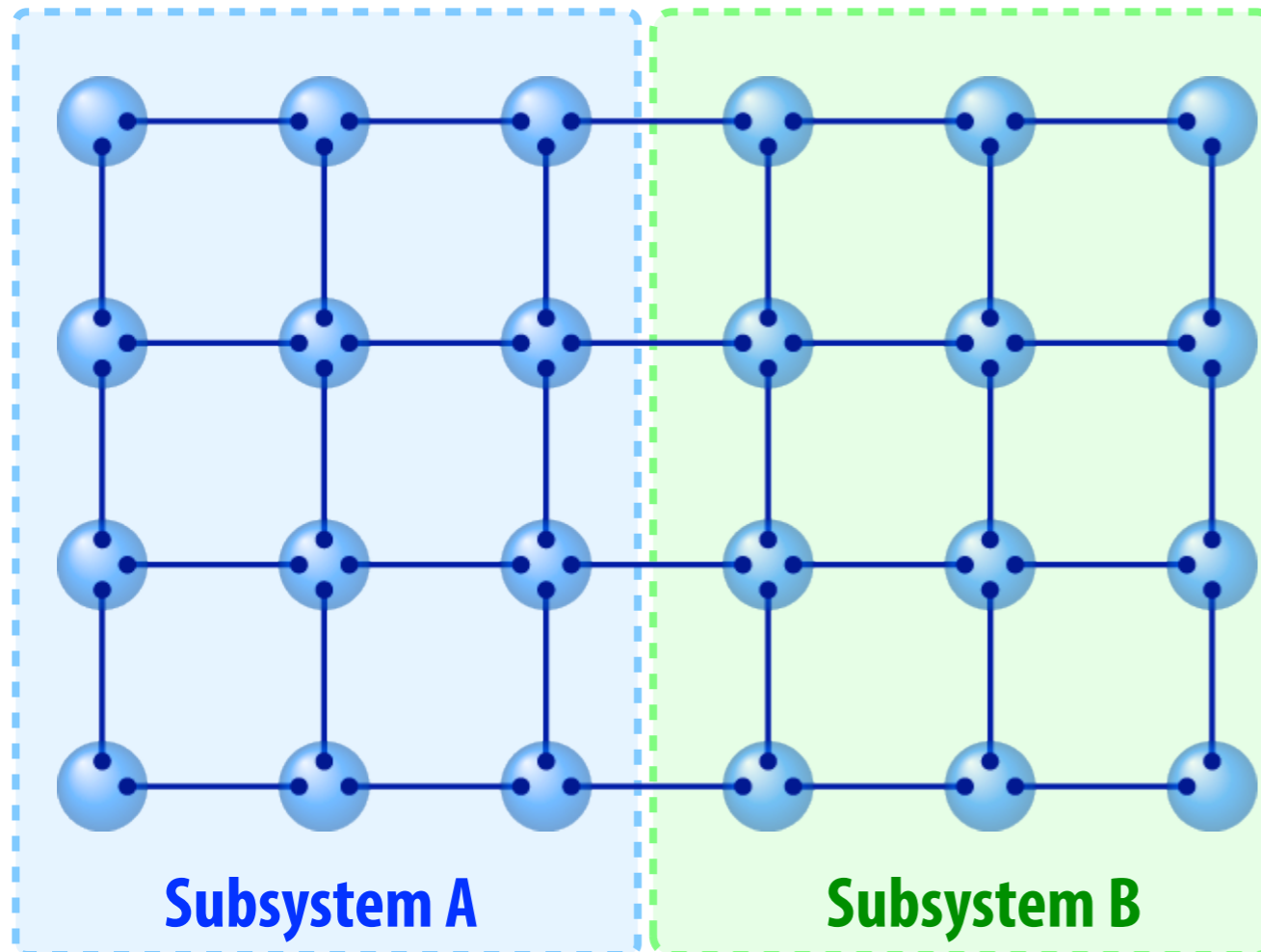


VBS (Valence-Bond-Solid) state

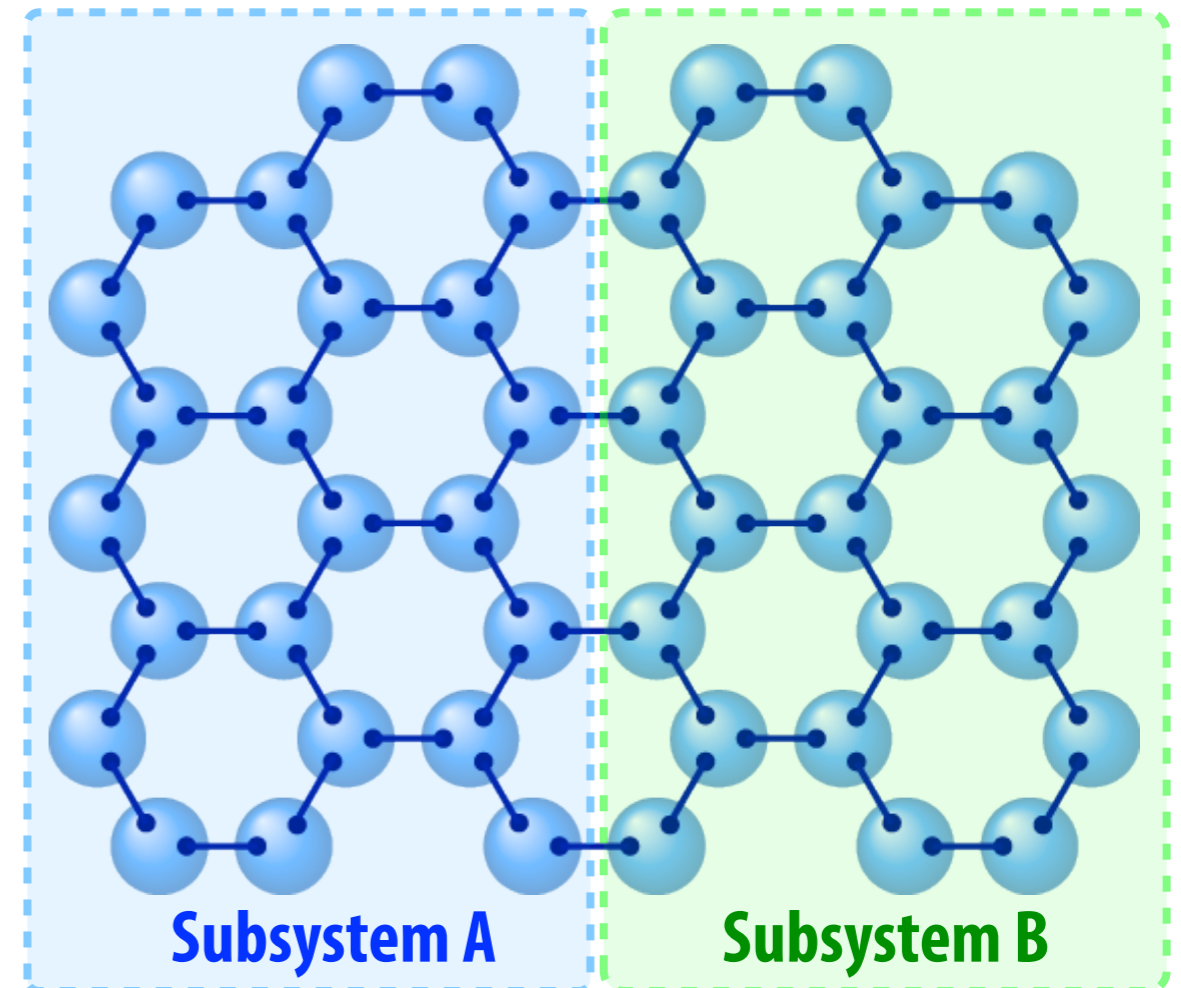
Reflection symmetry



2D square lattice



2D hexagonal lattice



VBS (Valence-Bond-Solid) state

$$|VBS\rangle = \prod_{\langle k,l\rangle} \left(a_k^\dagger b_l^\dagger - b_k^\dagger a_l^\dagger \right) |\text{vac}\rangle$$

$$= \sum_{\{\alpha\}} |\phi_\alpha^{[A]}\rangle \otimes |\phi_\alpha^{[B]}\rangle$$

- Local gauge transformation
- Reflection symmetry

$$\{\alpha\} = \{\alpha_1, \dots, \alpha_{|\Lambda_A|}\}$$

Auxiliary spin: $\alpha_i = \pm 1/2$

#bonds on edge: $|\Lambda_A|$

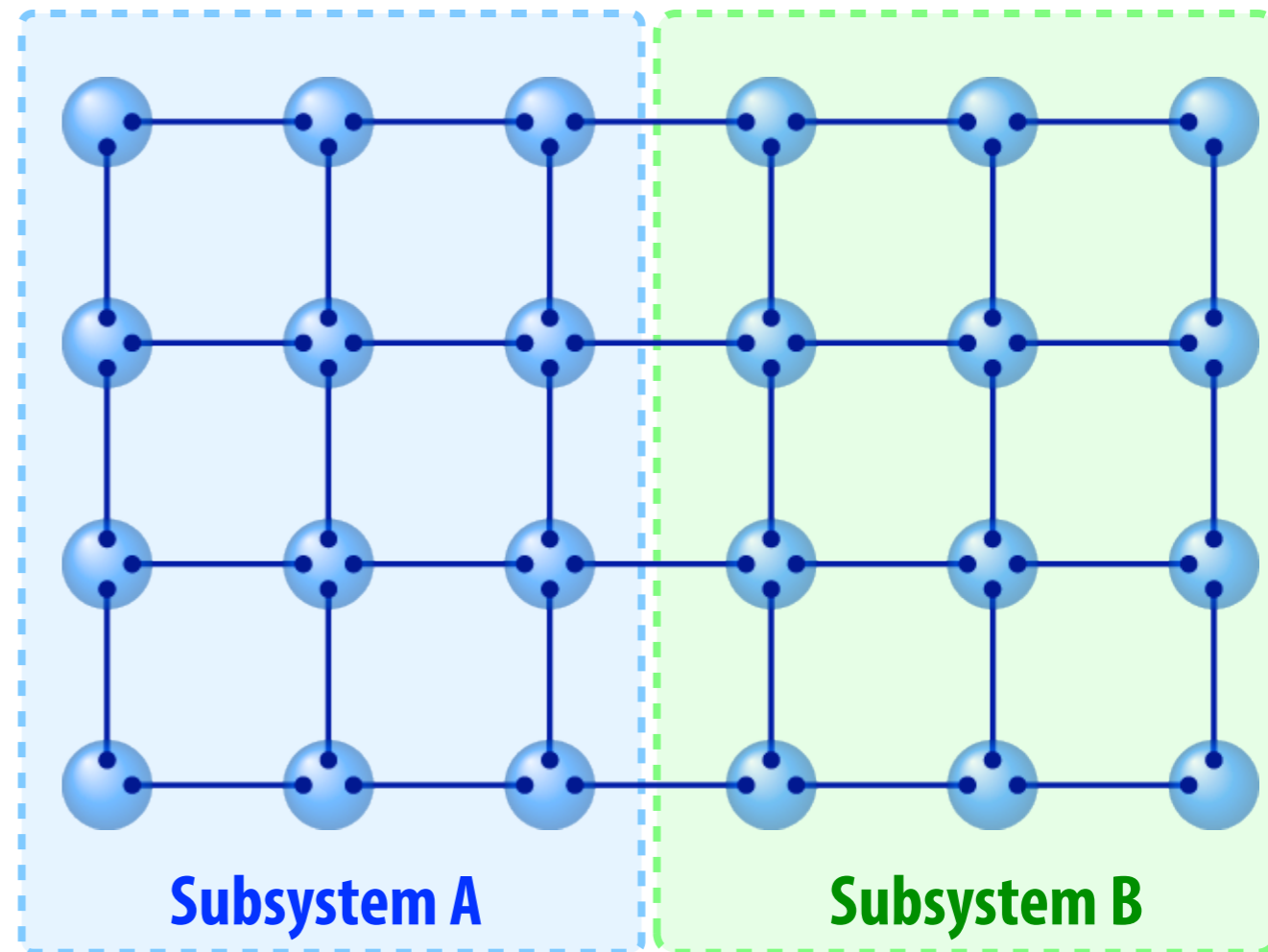


Overlap matrix

$M_{\{\alpha\},\{\beta\}}$: $2^{|\Lambda_A|} \times 2^{|\Lambda_A|}$ matrix

Each element can be obtained by Monte Carlo calculation!!

SU(N) case can be also calculated.



Phys. Rev. B, 84, 245128 (2011)

cf. H. Katsura, arXiv:1407.4262

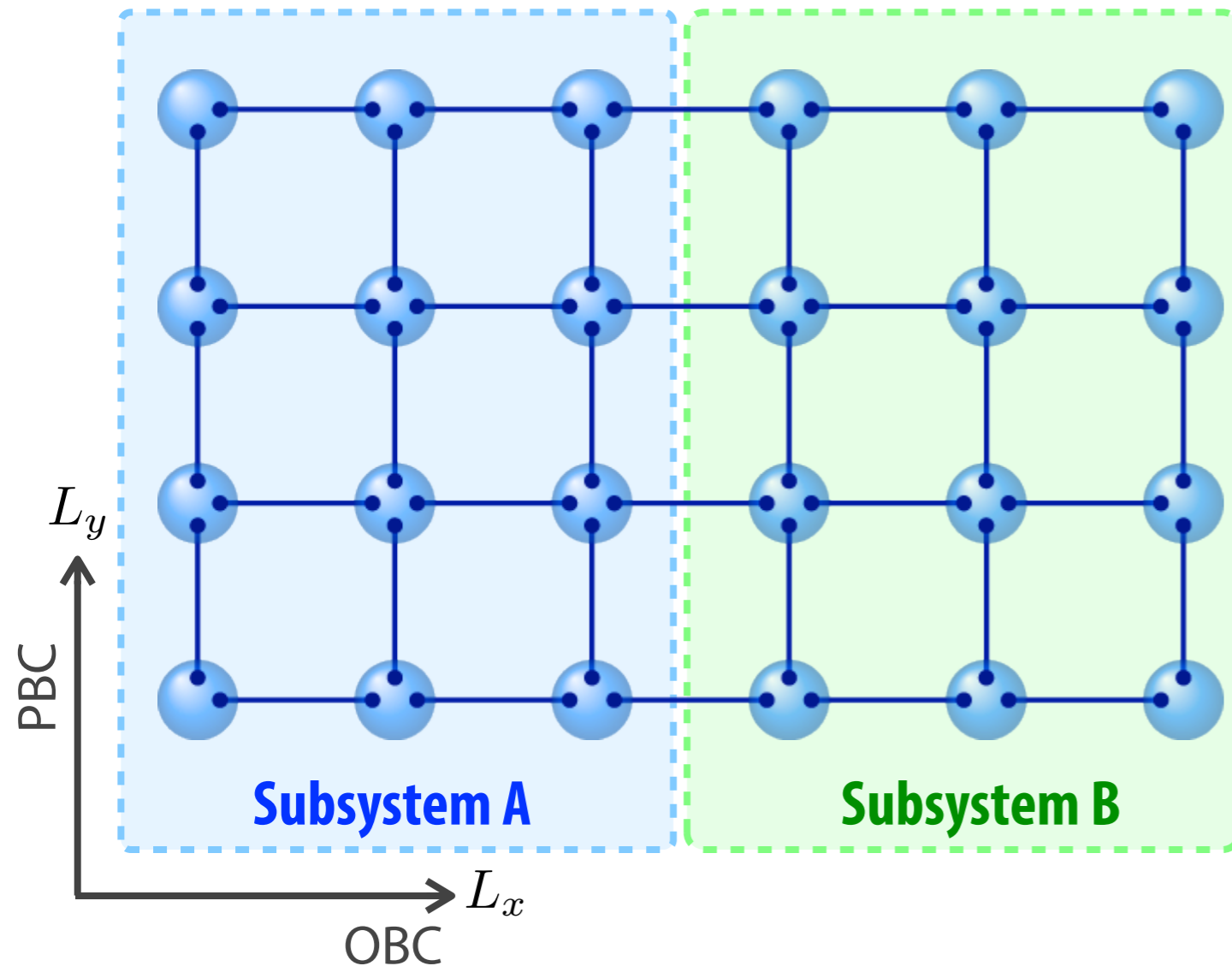
Entanglement properties

- Entanglement entropy***
- Entanglement spectrum***
- Nested entanglement entropy***

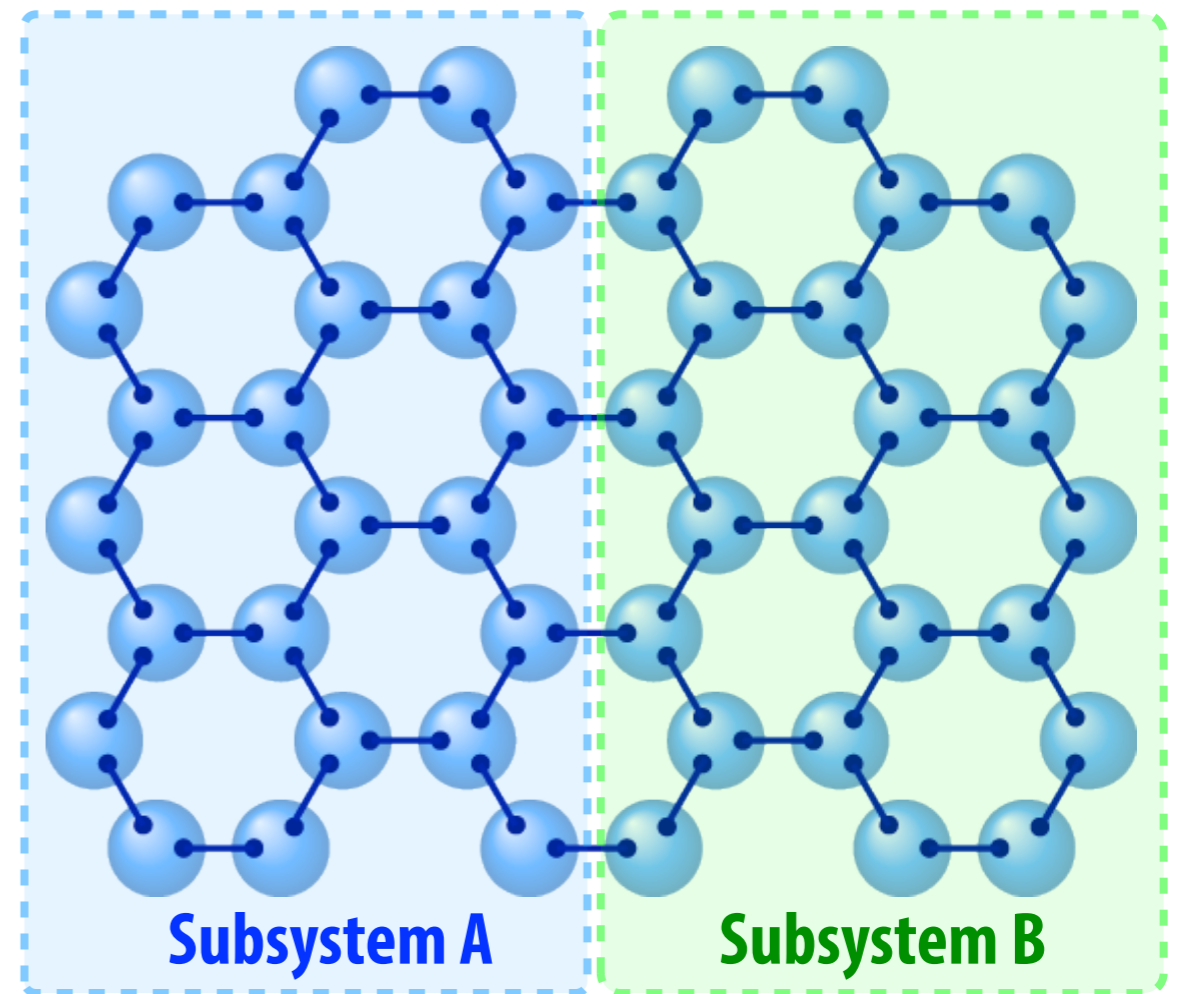
Entanglement properties of 2D VBS states

VBS state = Singlet-covering state

2D square lattice



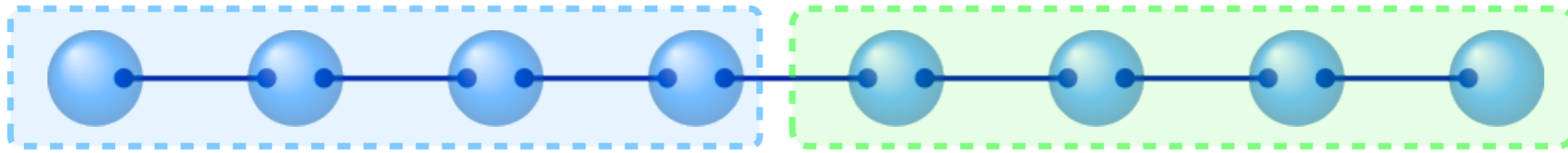
2D hexagonal lattice



Entanglement entropy of 2D VBS states

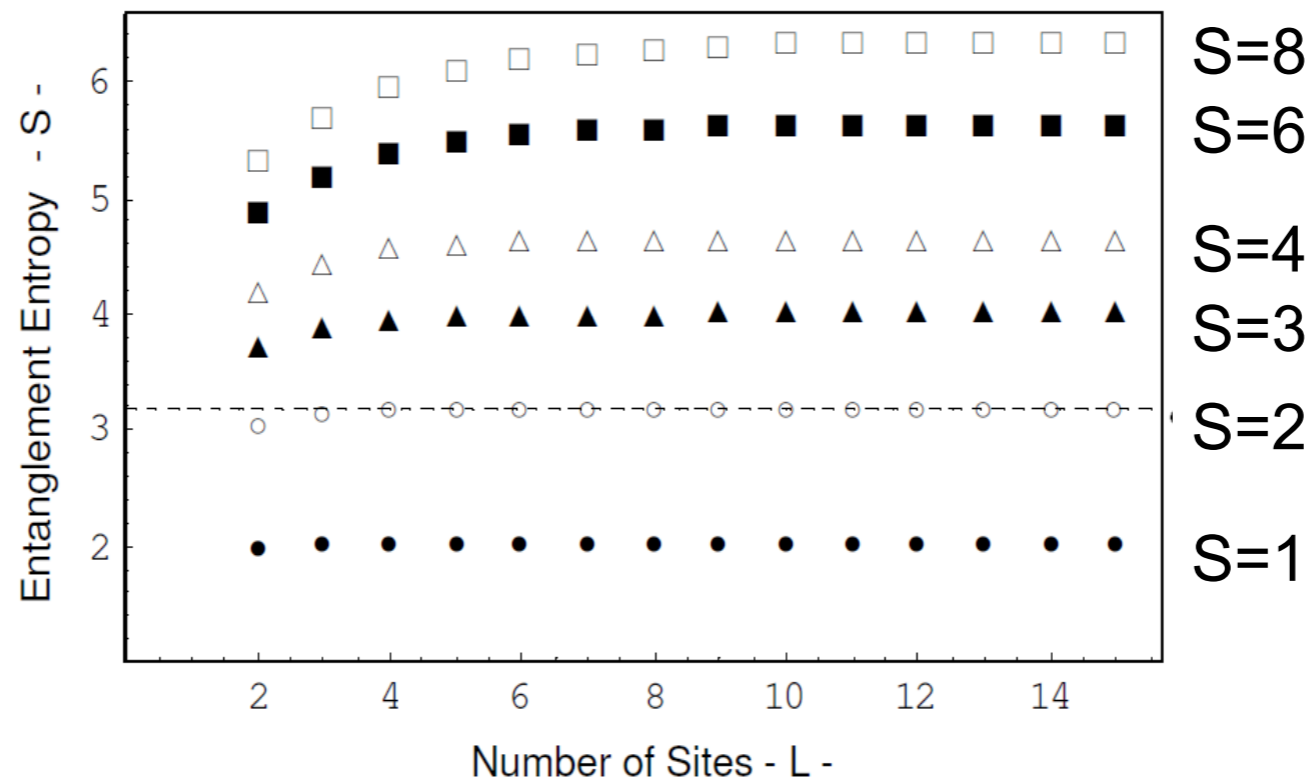
cf. Entanglement entropy of **1D** VBS states

$$|\text{VBS}\rangle = \prod_{i=0}^N \left(a_i^\dagger b_{i+1}^\dagger - b_i^\dagger a_{i+1}^\dagger \right)^S |\text{vac}\rangle$$



Subsystem A

Subsystem B



H. Katsura, T. Hirano, and Y. Hatsugai, PRB 76, 012401 (2007).

$$S = \ln (\# \text{ Edge states})$$

Entanglement entropy of 2D VBS states

$$\frac{\mathcal{S}}{|\Lambda_A|} = \ln 2 - \sigma$$

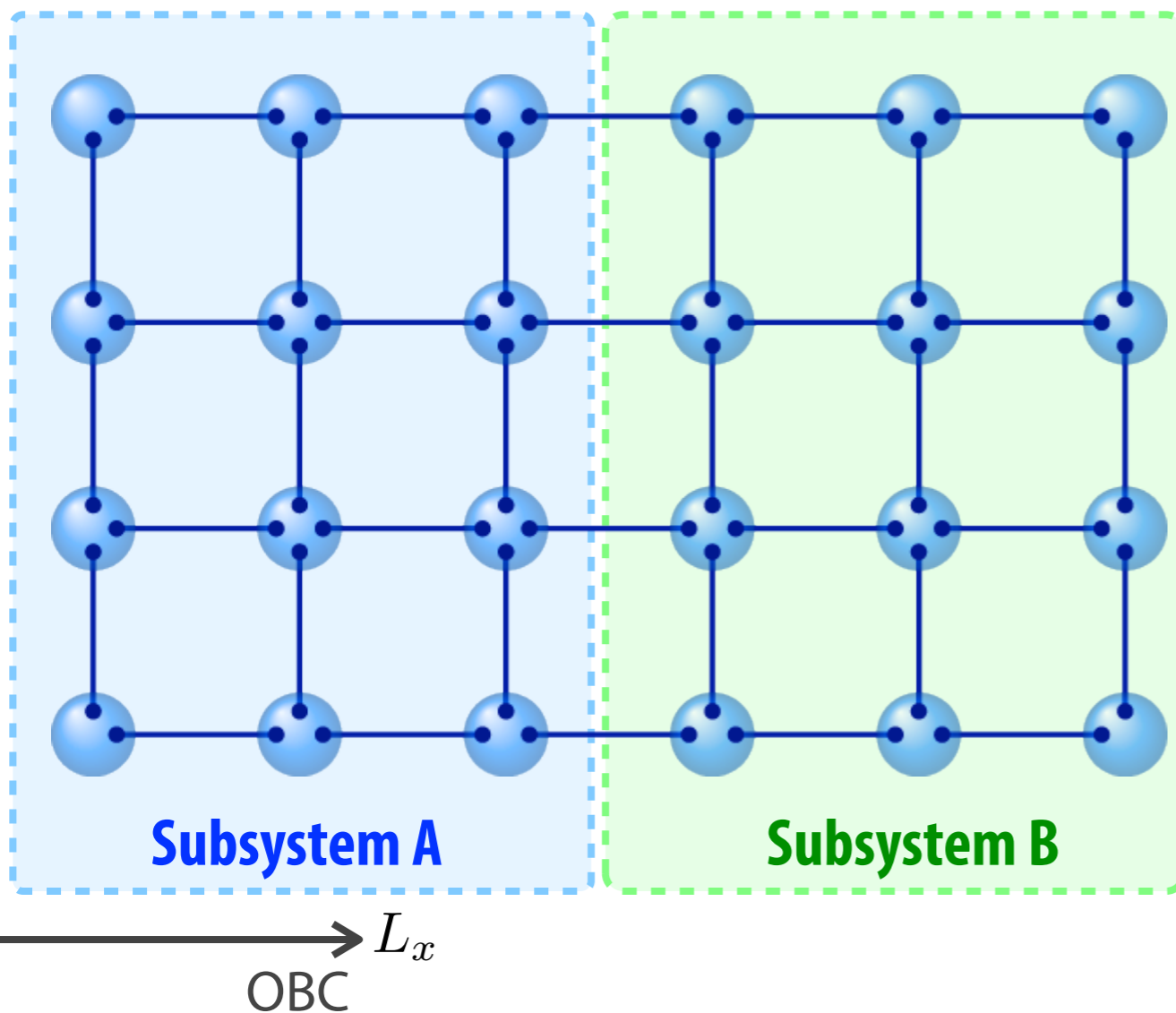
#bonds on edge: $|\Lambda_A|$

$$\sigma \geq 0$$

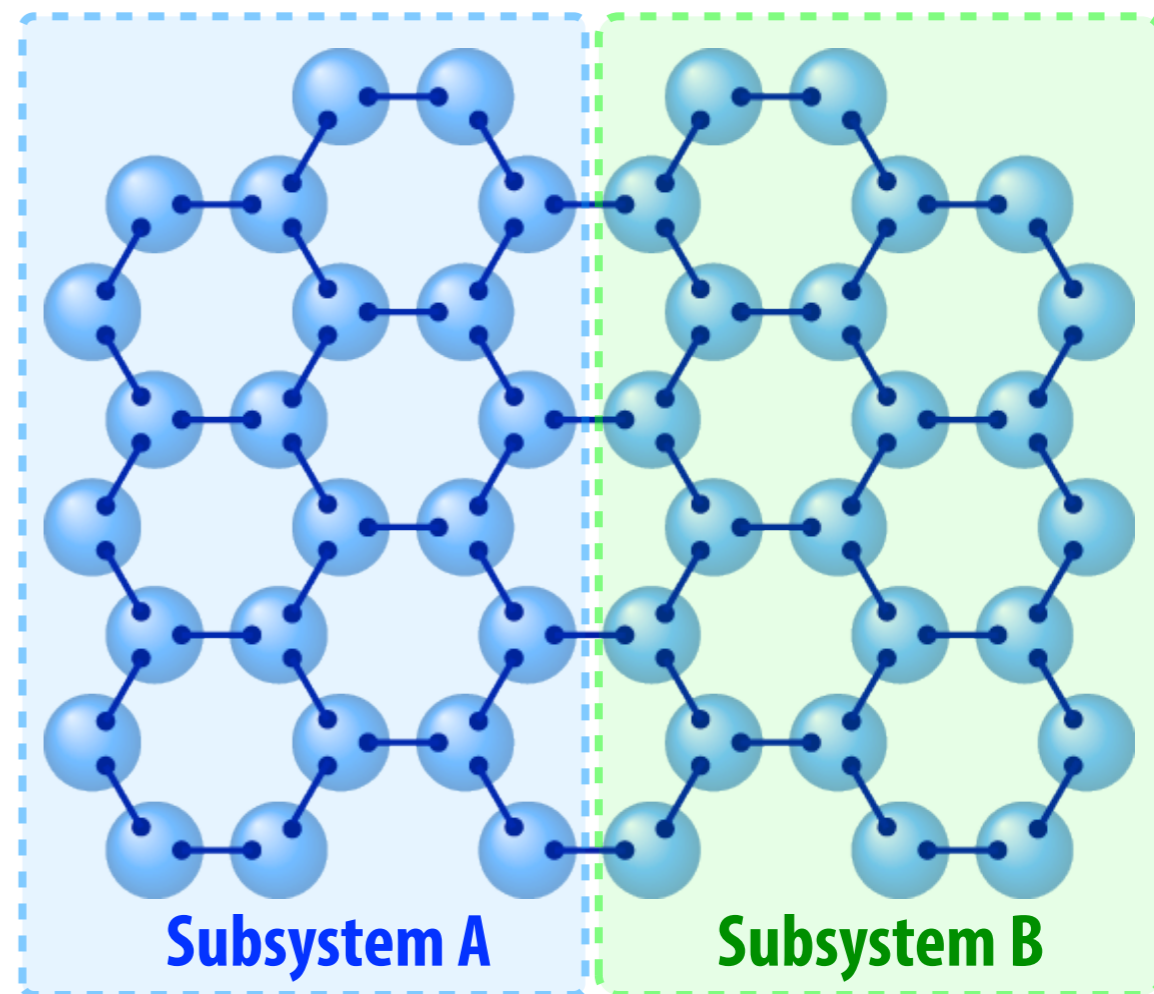
$$\sigma_{1D} = 0$$

$$\sigma_{\text{square}} > \sigma_{\text{hexagonal}} \quad \leftarrow \quad \xi_{\text{square}} > \xi_{\text{hexagonal}}$$

2D square lattice



2D hexagonal lattice



Entanglement spectra of 2D VBS states

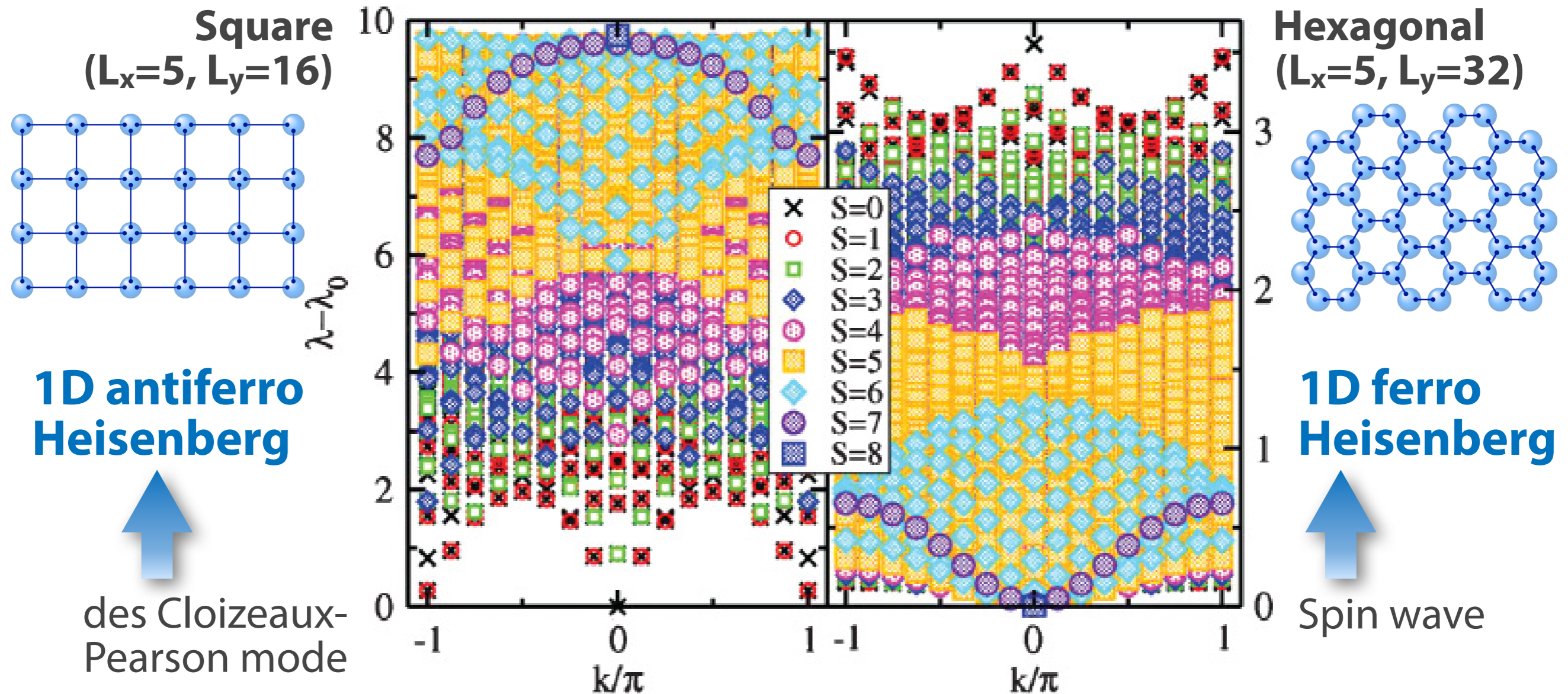
H. Li and F. D. M. Haldane, Phys. Rev. Lett. 101, 010504 (2008).

Reduced density matrix

$$\rho_A = \sum_{\alpha} e^{-\lambda_{\alpha}} |\phi_{\alpha}^{[A]}\rangle \langle \phi_{\alpha}^{[A]}|$$

Entanglement Hamiltonian

$$\rho_A = e^{-\mathcal{H}_E} \quad (\mathcal{H}_E = -\ln \rho_A)$$



cf. J. I. Cirac, D. Poilbranc, N. Schuch, and F. Verstraete, Phys. Rev. B 83, 245134 (2011).

Nested entanglement entropy

“Entanglement” ground state := g.s. of \mathcal{H}_E : $|\Psi_0\rangle$

$$\mathcal{H}_E = -\ln \rho_A$$

$$\mathcal{H}_E |\Psi_0\rangle = E_{\text{gs}} |\Psi_0\rangle \longleftrightarrow \rho_A |\Psi_0\rangle = \overline{\overline{\rho_0}} |\Psi_0\rangle$$

Maximum eigenvalue

Nested reduced density matrix

$$\rho(\ell) := \text{Tr}_{\ell+1, \dots, L} [|\Psi_0\rangle\langle\Psi_0|]$$

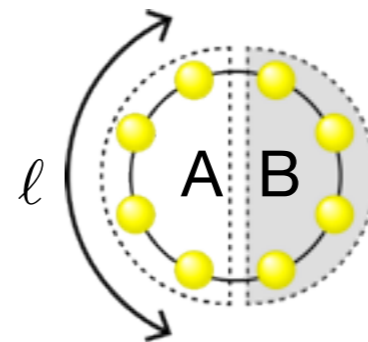
Nested entanglement entropy

$$S(\ell, L) = -\text{Tr}_{1, \dots, \ell} [\rho(\ell) \ln \rho(\ell)]$$

1D quantum critical system (periodic boundary condition)

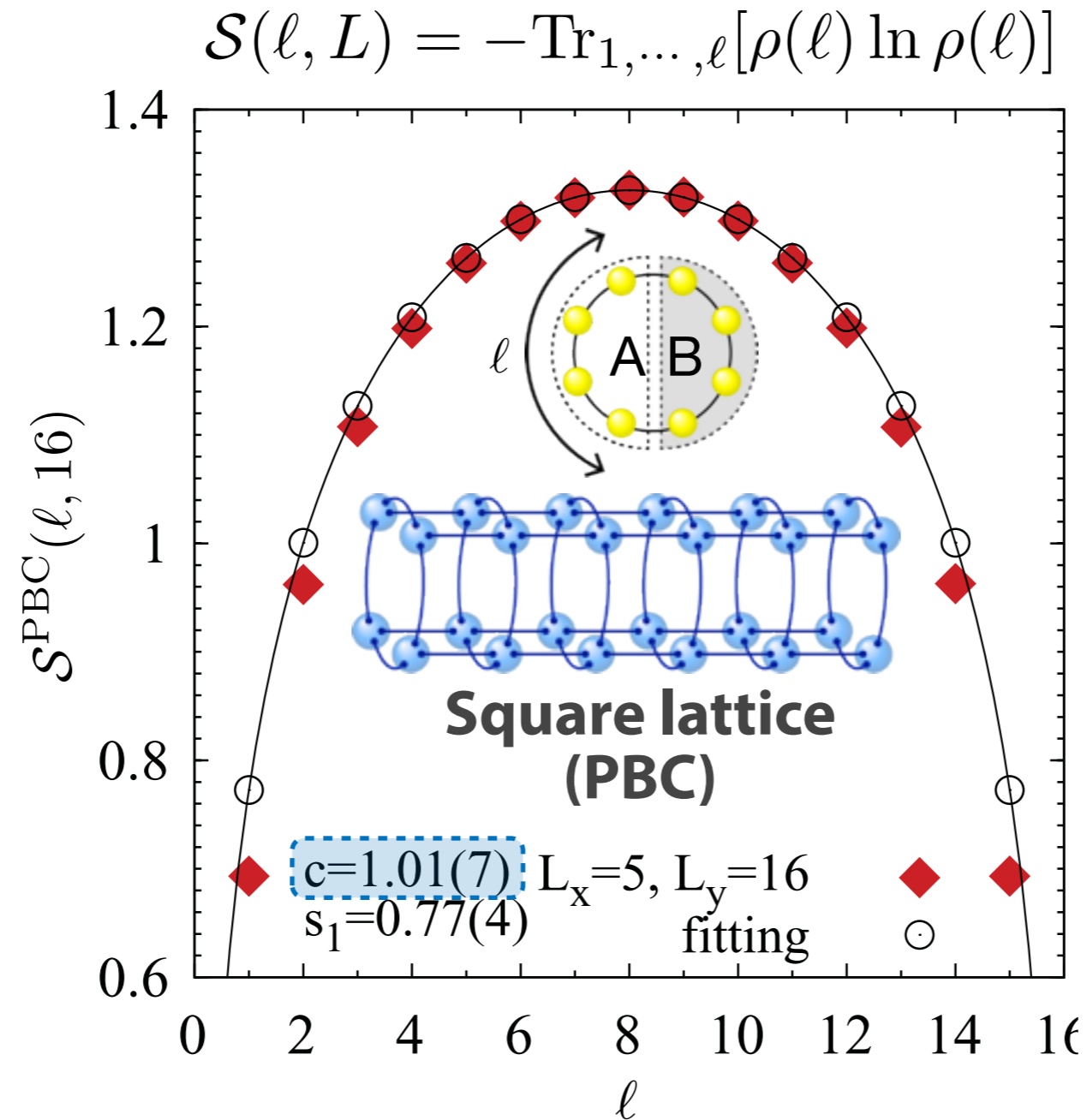
$$S^{\text{PBC}}(\ell, L_y) = \frac{c}{3} \ln[f(\ell)] + s_1$$

$$f(\ell) = \frac{L_y}{\pi} \sin\left(\frac{\pi\ell}{L_y}\right)$$



P. Calabrese and J. Cardy, J. Stat. Mech. (2004) P06002.

Nested entanglement entropy



Central charge: $c = 1$ \rightarrow **1D antiferromagnetic Heisenberg**
des Cloizeaux-Pearson mode in ES supports this result.

VBS/CFT correspondence

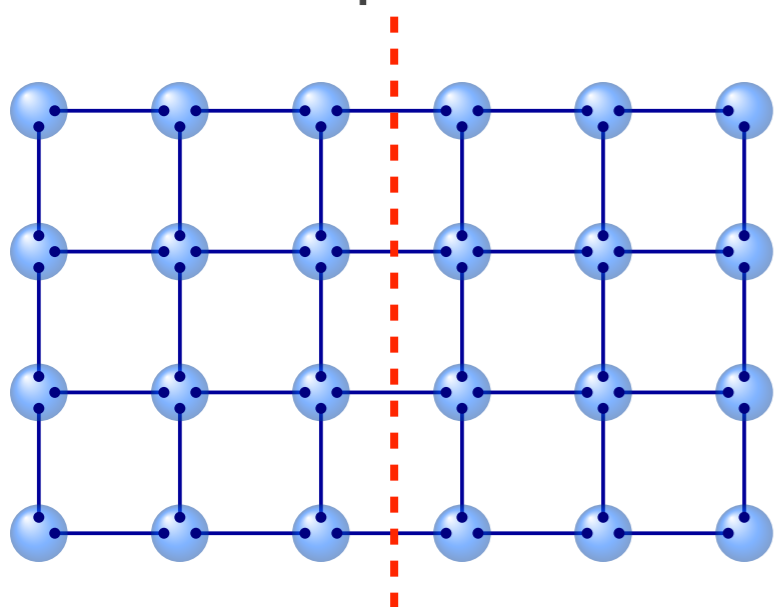
Digest

Entanglement properties of 2D quantum systems

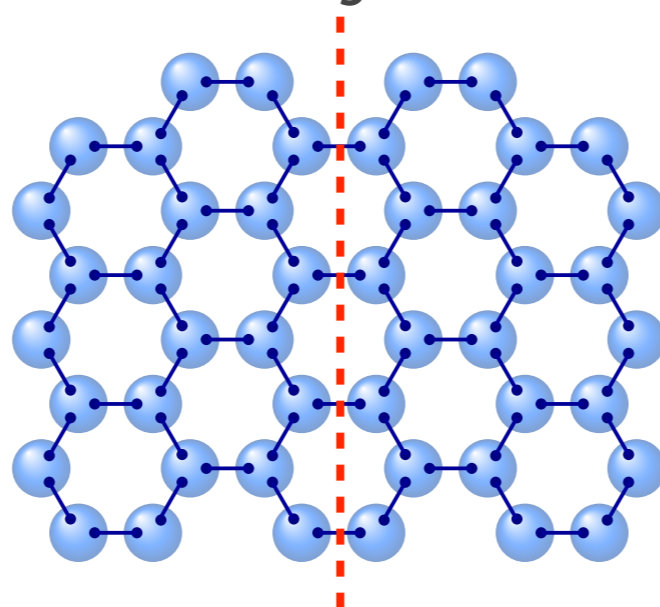


Physical properties of 1D quantum systems

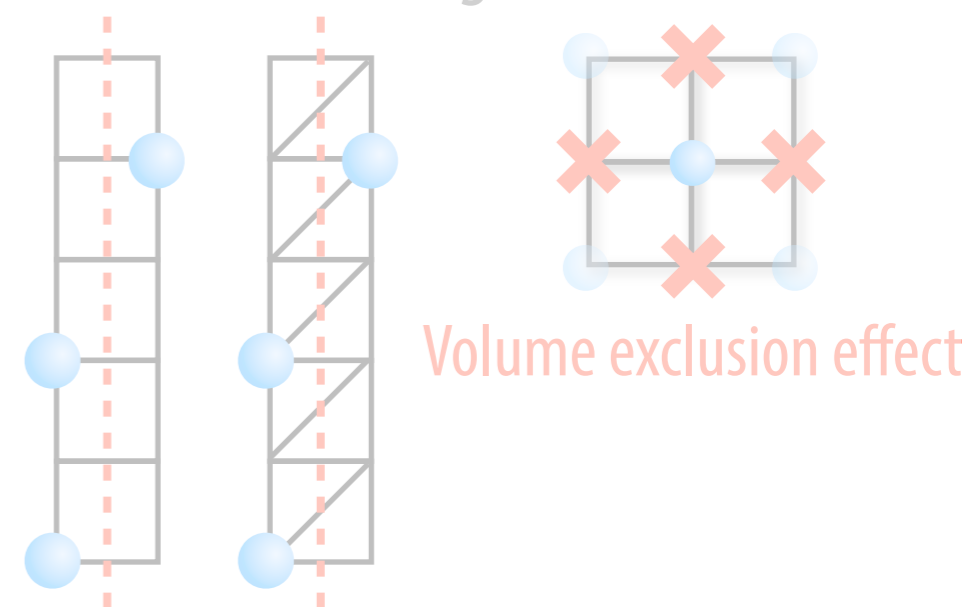
VBS on square lattice



VBS on hexagonal lattice



Quantum lattice gas on ladder



VBS state on 2D lattice

Total system	Entanglement Hamiltonian
Square lattice	1D AF Heisenberg
Hexagonal lattice	1D F Heisenberg

Quantum lattice gas on ladder

Total system	Entanglement Hamiltonian
Square ladder	2D Ising
Triangle ladder	2D 3-state Potts

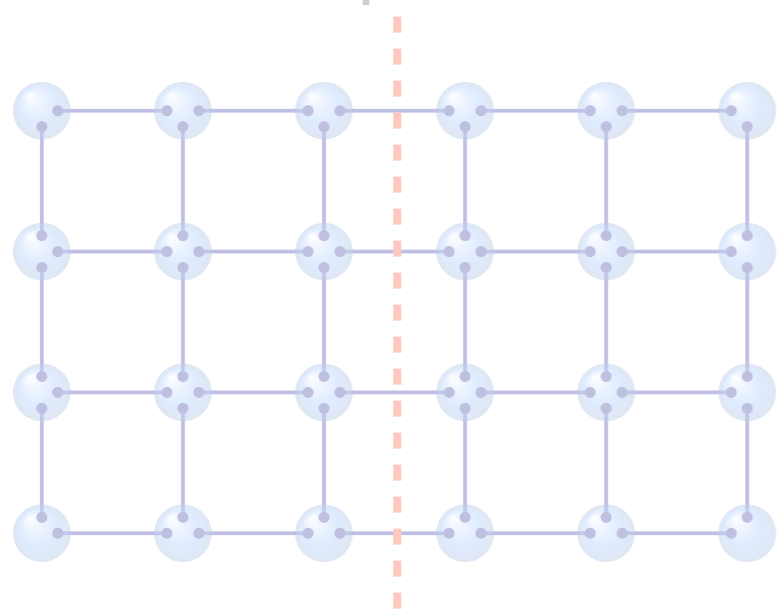
Digest

Entanglement properties of 2D quantum systems

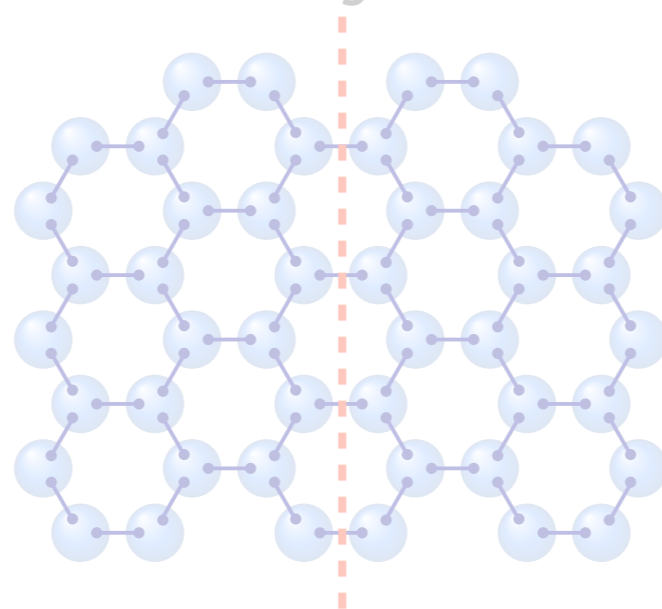


Physical properties of 1D quantum systems

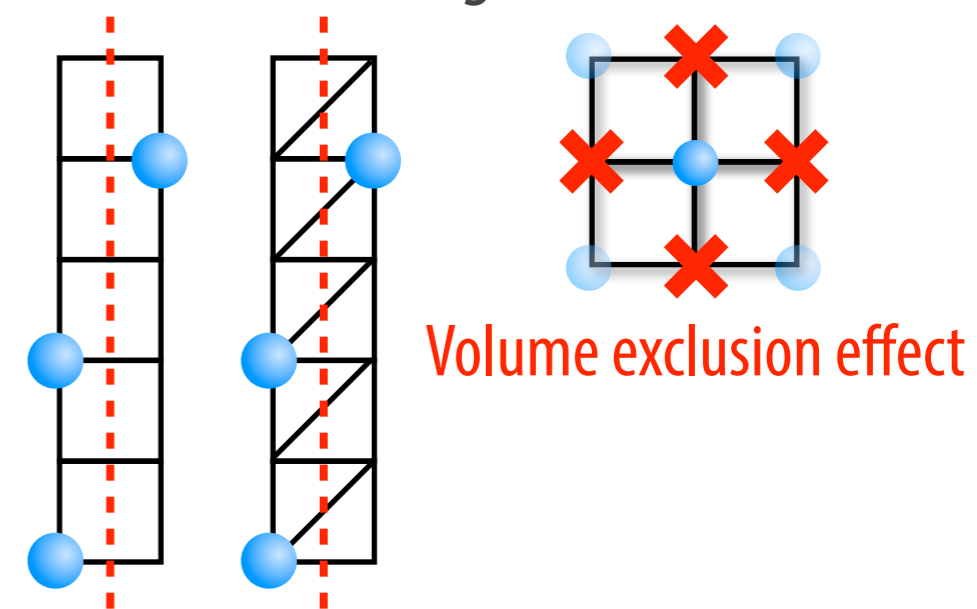
VBS on square lattice



VBS on hexagonal lattice



Quantum lattice gas on ladder



VBS state on 2D lattice

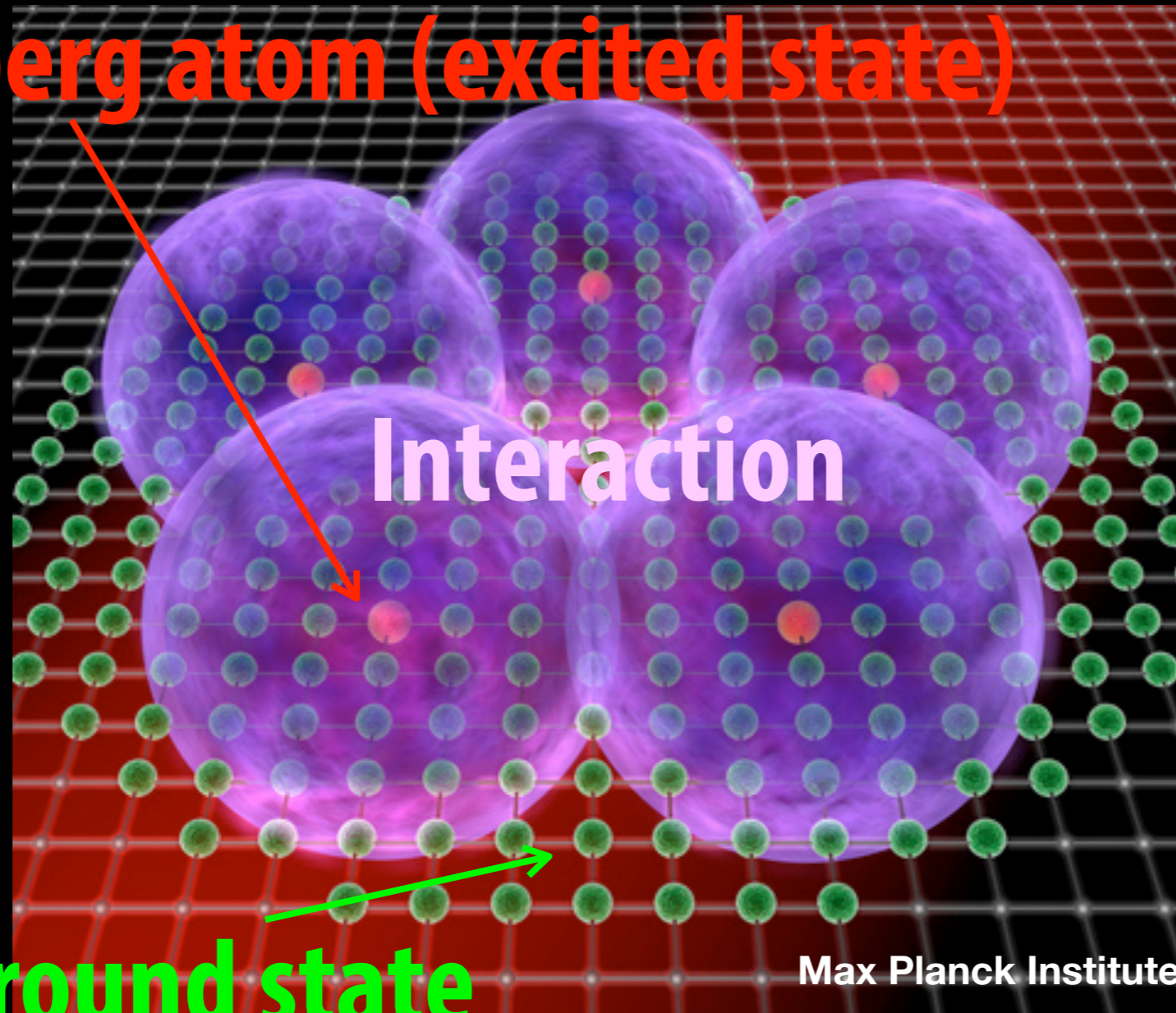
Total system	Entanglement Hamiltonian
Square lattice	1D AF Heisenberg
Hexagonal lattice	1D F Heisenberg

Quantum lattice gas on ladder

Total system	Entanglement Hamiltonian
Square ladder	2D Ising
Triangle ladder	2D 3-state Potts

Rydberg Atom

Rydberg atom (excited state)



$$\mathcal{H} = \Omega \sum_{i \in \Lambda} \sigma_i^x + \Delta \sum_{i \in \Lambda} n_i + V \sum_{i,j} \frac{n_i n_j}{|\mathbf{r}_i - \mathbf{r}_j|^\gamma}$$

Quantum hard-core lattice gas model

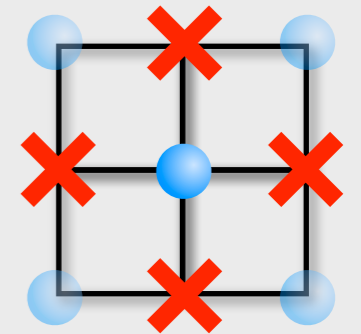
Construct a solvable model

$$\mathcal{H}_{\text{sol}} = \sum_{i \in \Lambda} h_i^\dagger(z) h_i(z), \quad h_i(z) := [\sigma_i^- - \sqrt{z}(1 - n_i)] \mathcal{P}_{\langle i \rangle}$$

$$\mathcal{H}_{\text{sol}} = \underbrace{-\sqrt{z} \sum_{i \in \Lambda} (\sigma_i^+ + \sigma_i^-) \mathcal{P}_{\langle i \rangle}}_{\text{Creation/annihilation}} + \underbrace{\sum_{i \in \Lambda} [(1 - z)n_i + z] \mathcal{P}_{\langle i \rangle}}_{\text{Interaction btw particles \& chemical potential}}$$

$$n_i = \frac{\sigma_i^z + 1}{2}$$

$$\mathcal{P}_{\langle i \rangle} := \prod_{j \in G_i} (1 - n_j)$$



1-dim chain

$$\mathcal{H} = \sum_{i=1}^L \mathcal{P} \left[-\sqrt{z} \sigma_i^x + (1 - 3z)n_i + zn_{i-1}n_{i+1} + z \right] \mathcal{P} \quad \text{Transverse Ising model with constraint}$$

Hamiltonian is positive semi-definite. \longrightarrow Eigenenergies are non-negative.

Zero-energy state (ground state)

$$|z\rangle = \frac{1}{\sqrt{\Xi(z)}} \prod_{i \in \Lambda} \exp(\sqrt{z} \sigma_i^+ \mathcal{P}_{\langle i \rangle}) |\downarrow\downarrow \cdots \downarrow\rangle \quad |\downarrow\downarrow \cdots \downarrow\rangle : \text{Vacuum state}$$

Unique (Perron-Frobenius theorem)

GS of the quantum hard-core lattice gas model

unnormalized ground state: $|\Psi(z)\rangle := \sqrt{\Xi(z)}|z\rangle = \sum_{\mathcal{C} \in \mathcal{S}} z^{n_{\mathcal{C}}/2} |\mathcal{C}\rangle$

\mathcal{C} : classical configuration of particle on Λ

$\langle \mathcal{C} | \mathcal{C}' \rangle = \delta_{\mathcal{C}, \mathcal{C}'}$ ($|\mathcal{C}\rangle$ is orthonormal basis)

\mathcal{S} : set of classical configurations with "constraint"

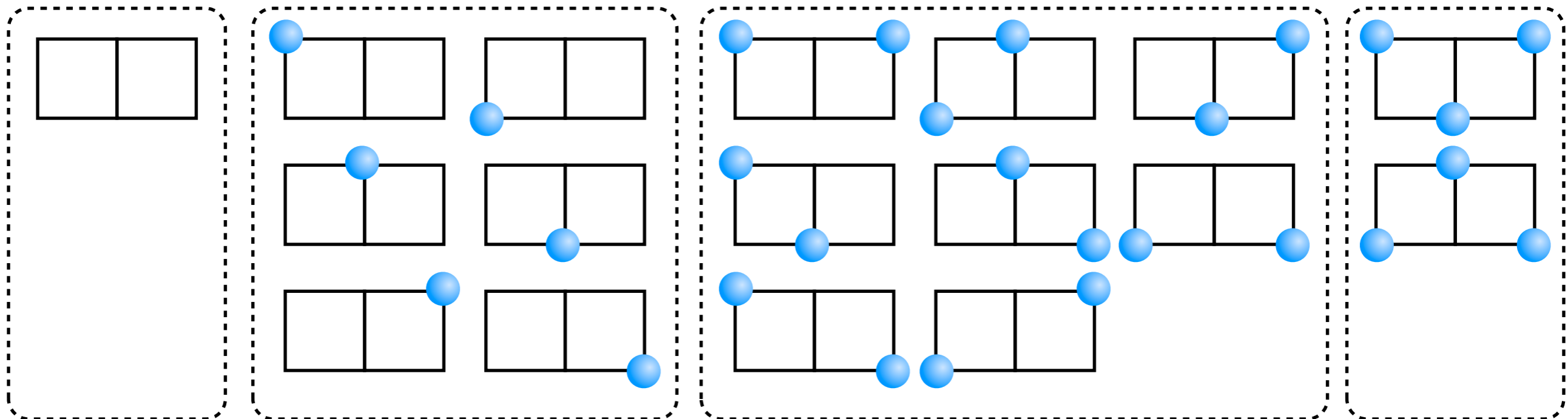
$n_{\mathcal{C}}$: number of particles in the state \mathcal{C}

Normalization factor

= Partition function of classical hard-core lattice gas model

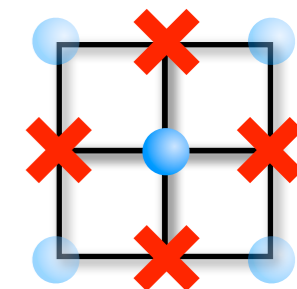
$$\Xi(z) = \langle \Psi(z) | \Psi(z) \rangle = \sum_{\mathcal{C} \in \mathcal{S}} z^{n_{\mathcal{C}}}$$

z : chemical potential

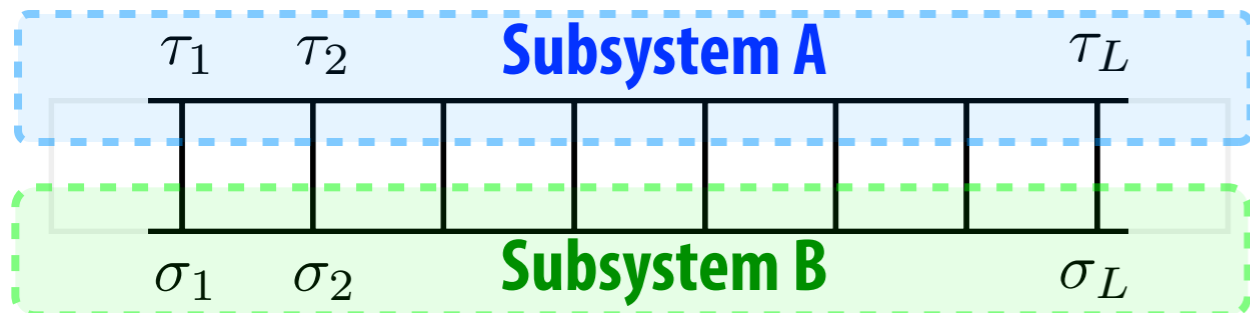


GS of the quantum hard-core lattice gas model

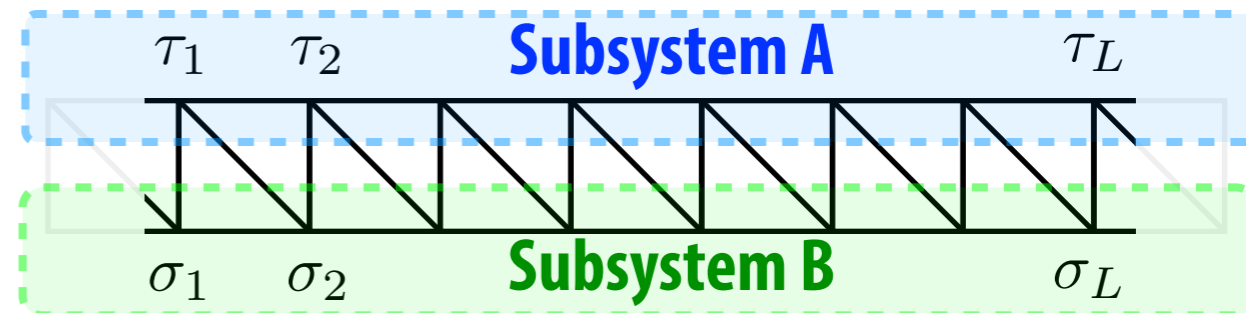
Periodic boundary condition is imposed in the leg direction.



Square ladder

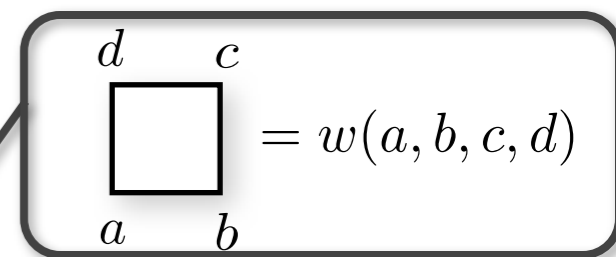


Triangle ladder

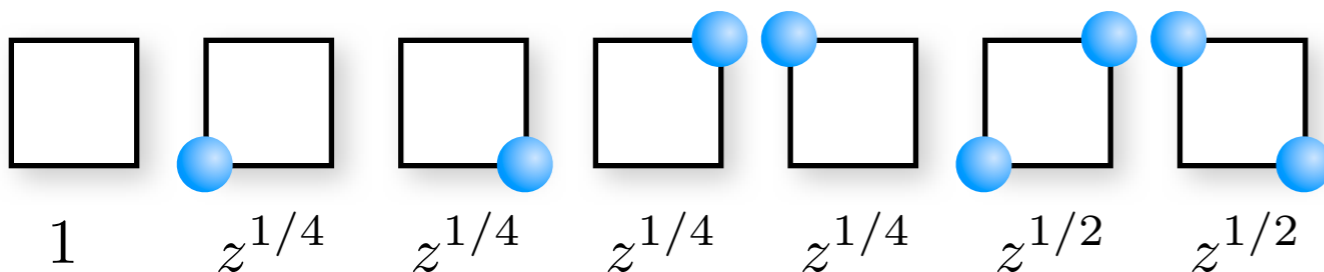


unnormalized ground state:

$$|\Psi(z)\rangle = \sum_{\sigma} \sum_{\tau} [T(z)]_{\tau, \sigma} |\tau\rangle \otimes |\sigma\rangle, \quad [T(z)]_{\tau, \sigma} := \prod_{i=1}^L w(\sigma_i, \sigma_{i+1}, \tau_{i+1}, \tau_i)$$

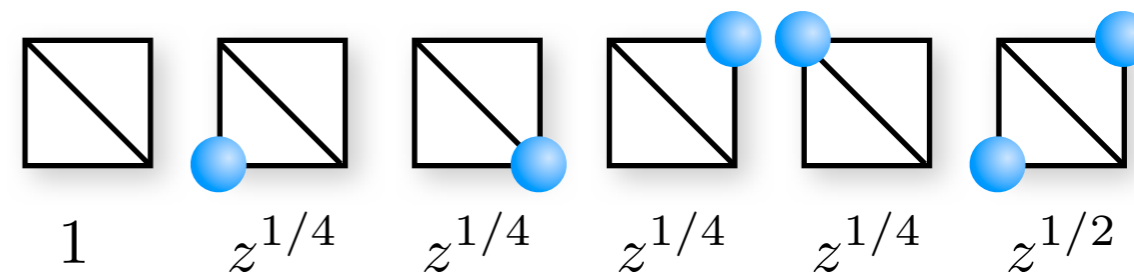


Square ladder



$$[T(z)]_{\tau, \sigma} = \prod_{i=1}^L z^{(\sigma_i + \tau_i)/2} (1 - \sigma_i \tau_i) (1 - \sigma_i \sigma_{i+1}) (1 - \tau_i \tau_{i+1})$$

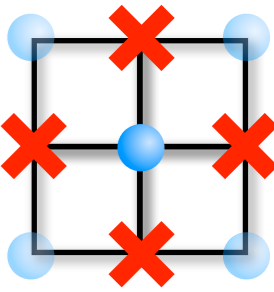
Triangle ladder



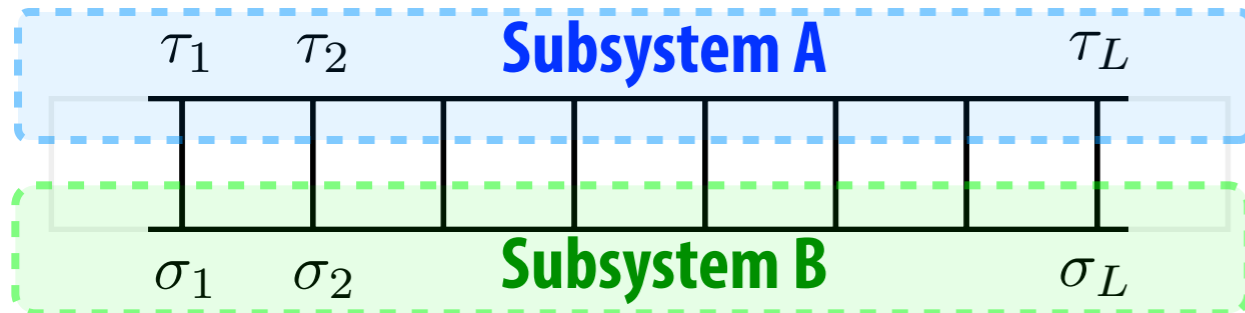
$$[T(z)]_{\tau, \sigma} = \prod_{i=1}^L z^{(\sigma_i + \tau_i)/2} (1 - \sigma_i \tau_i) (1 - \sigma_i \sigma_{i+1}) (1 - \tau_i \tau_{i+1}) (1 - \tau_i \sigma_{i+1})$$

GS of the quantum hard-core lattice gas model

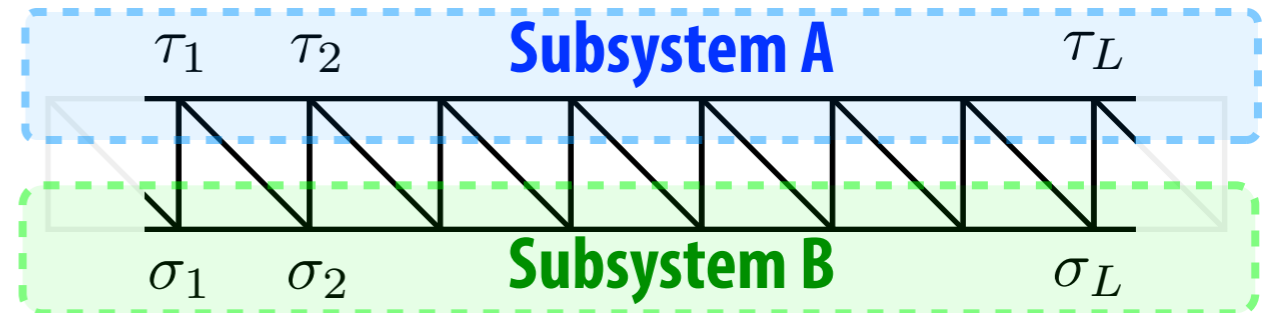
Periodic boundary condition is imposed in the leg direction.



Square ladder

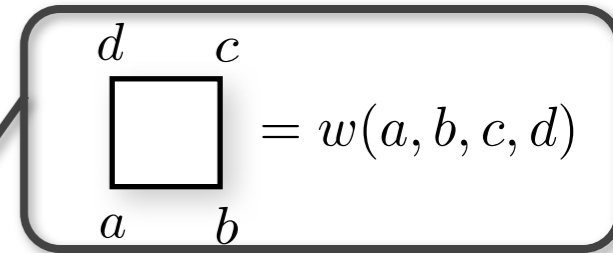


Triangle ladder



unnormalized ground state:

$$|\Psi(z)\rangle = \sum_{\sigma} \sum_{\tau} [T(z)]_{\tau,\sigma} |\tau\rangle \otimes |\sigma\rangle, \quad [T(z)]_{\tau,\sigma} := \prod_{i=1}^L w(\sigma_i, \sigma_{i+1}, \tau_{i+1}, \tau_i)$$



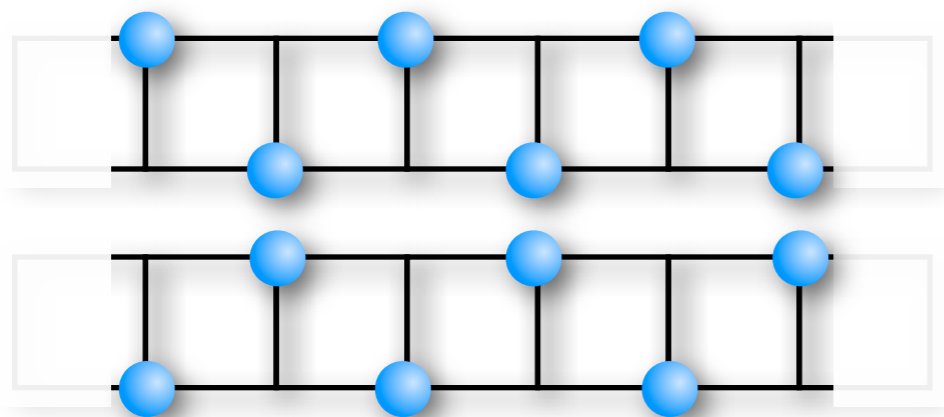
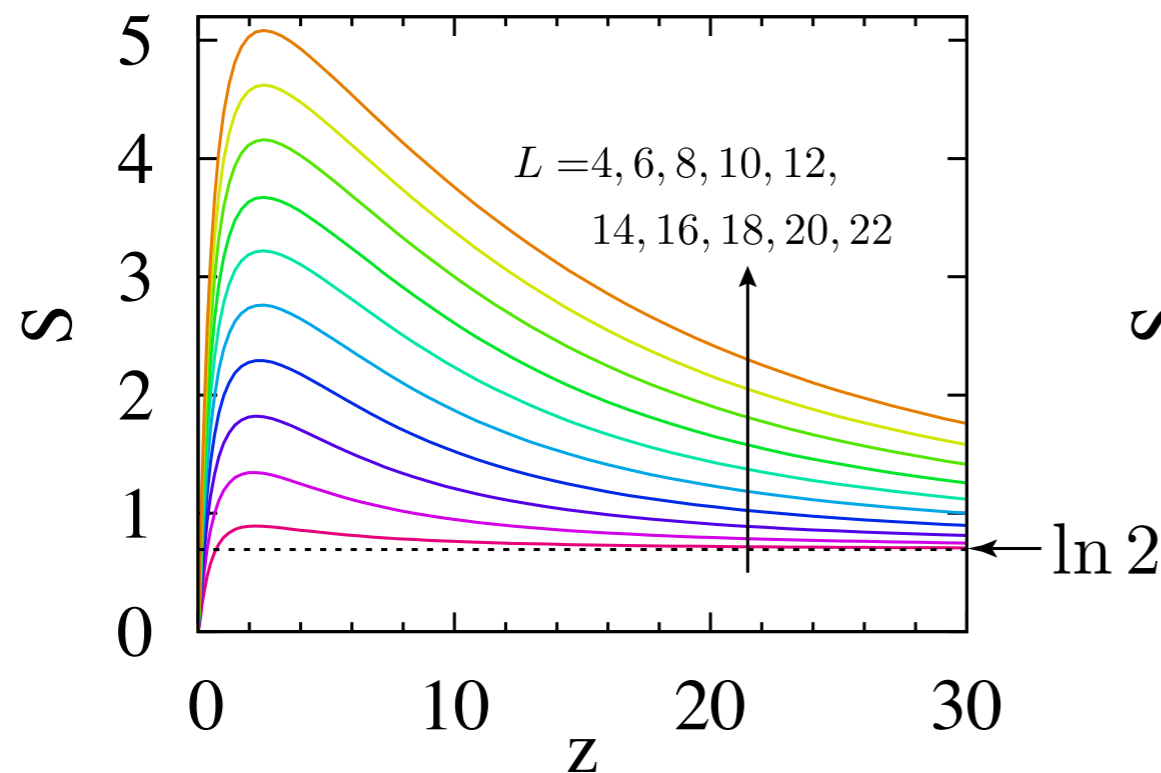
$$|z\rangle = \frac{1}{\sqrt{\Xi(z)}} \sum_{\sigma} \sum_{\tau} [T(z)]_{\tau,\sigma} |\tau\rangle \otimes |\sigma\rangle \quad \longrightarrow \quad \text{Overlap matrix}$$

$$M = \frac{1}{\Xi(z)} [T(z)]^T T(z)$$

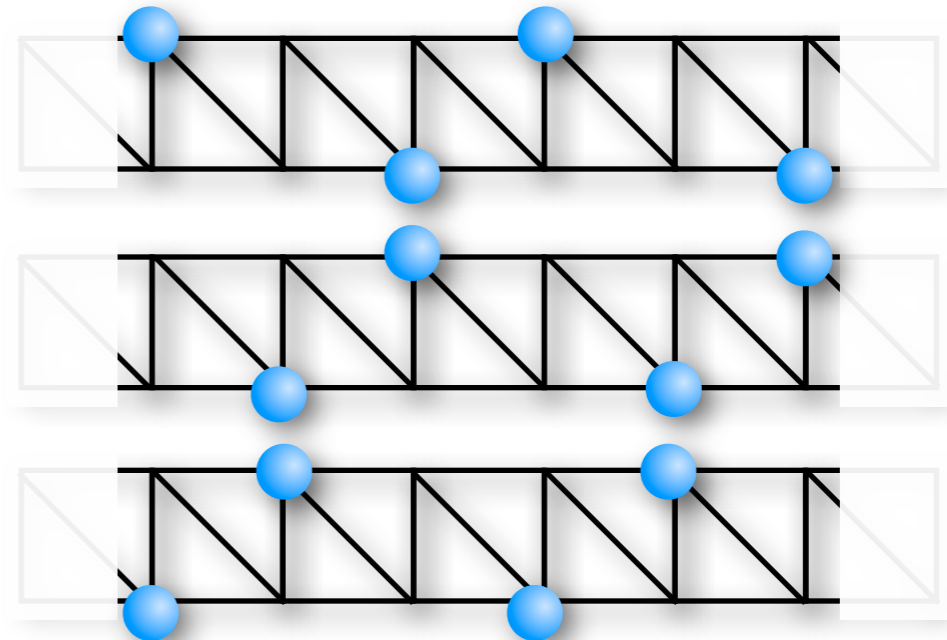
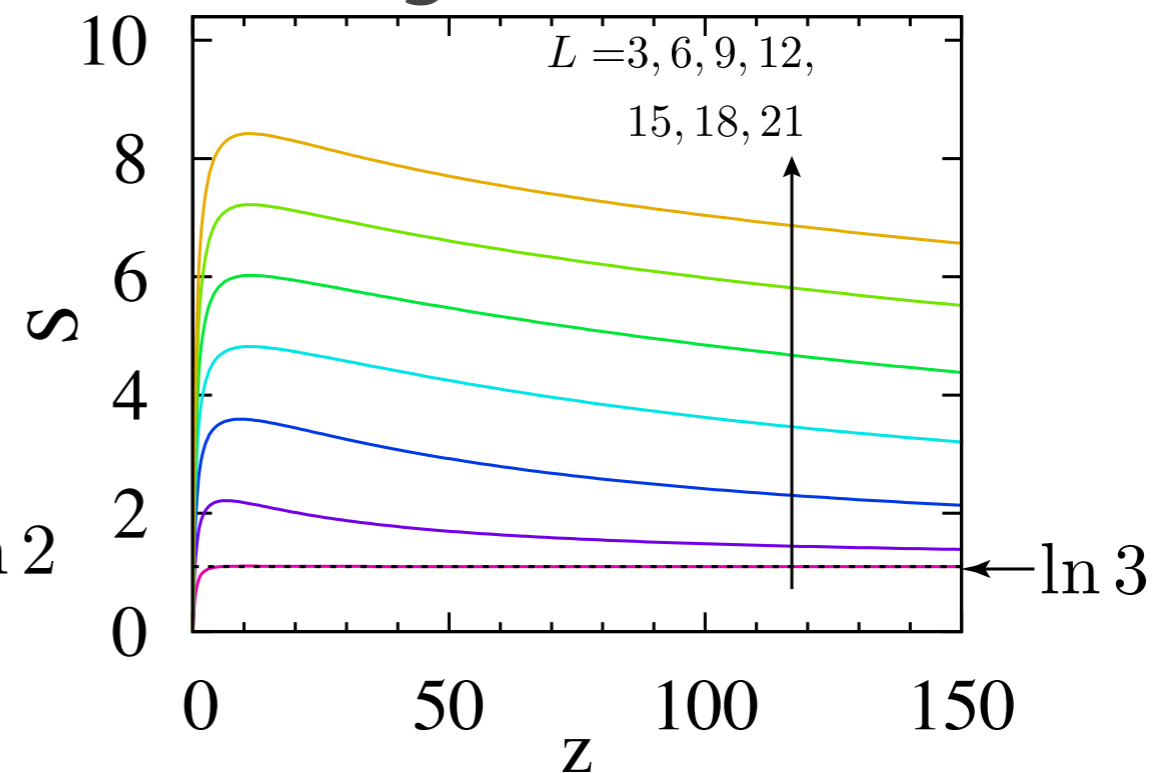
Entanglement entropy

$$\mathcal{S} = -\text{Tr} [M \ln M] = - \sum_{\alpha} p_{\alpha} \ln p_{\alpha} \quad p_{\alpha} (\alpha = 1, 2, \dots, \underline{\underline{N_L}}) \quad \# \text{ of states}$$

Square ladder



Triangle ladder



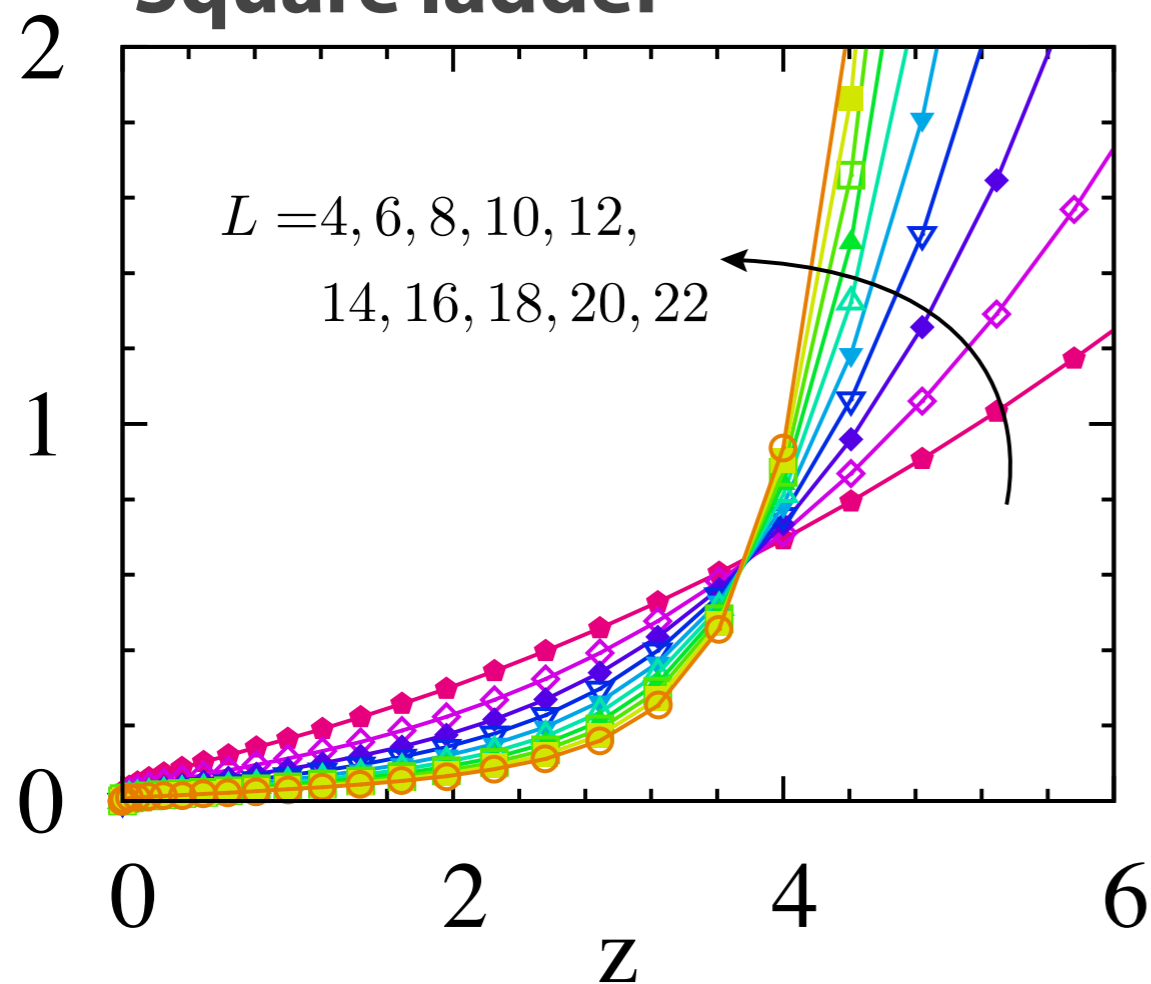
Estimation of z_c

$$\xi(z) := \frac{1}{\ln[p^{(1)}(z)/p^{(2)}(z)]}$$

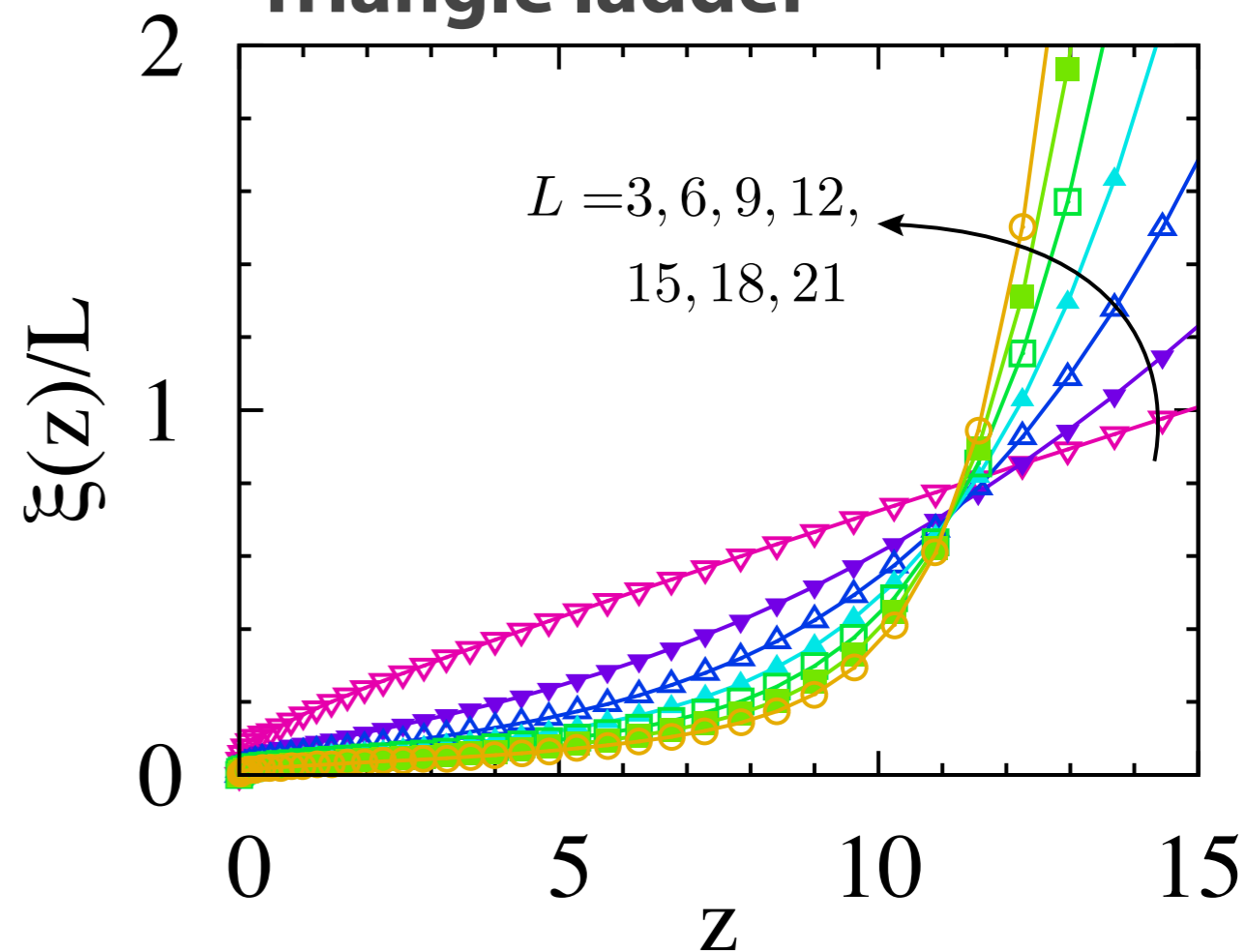
$p^{(1)}(z)$: the largest eigenvalue of M

$p^{(2)}(z)$: the second-largest eigenvalue of M

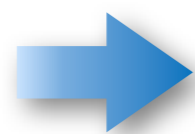
Square ladder



Triangle ladder



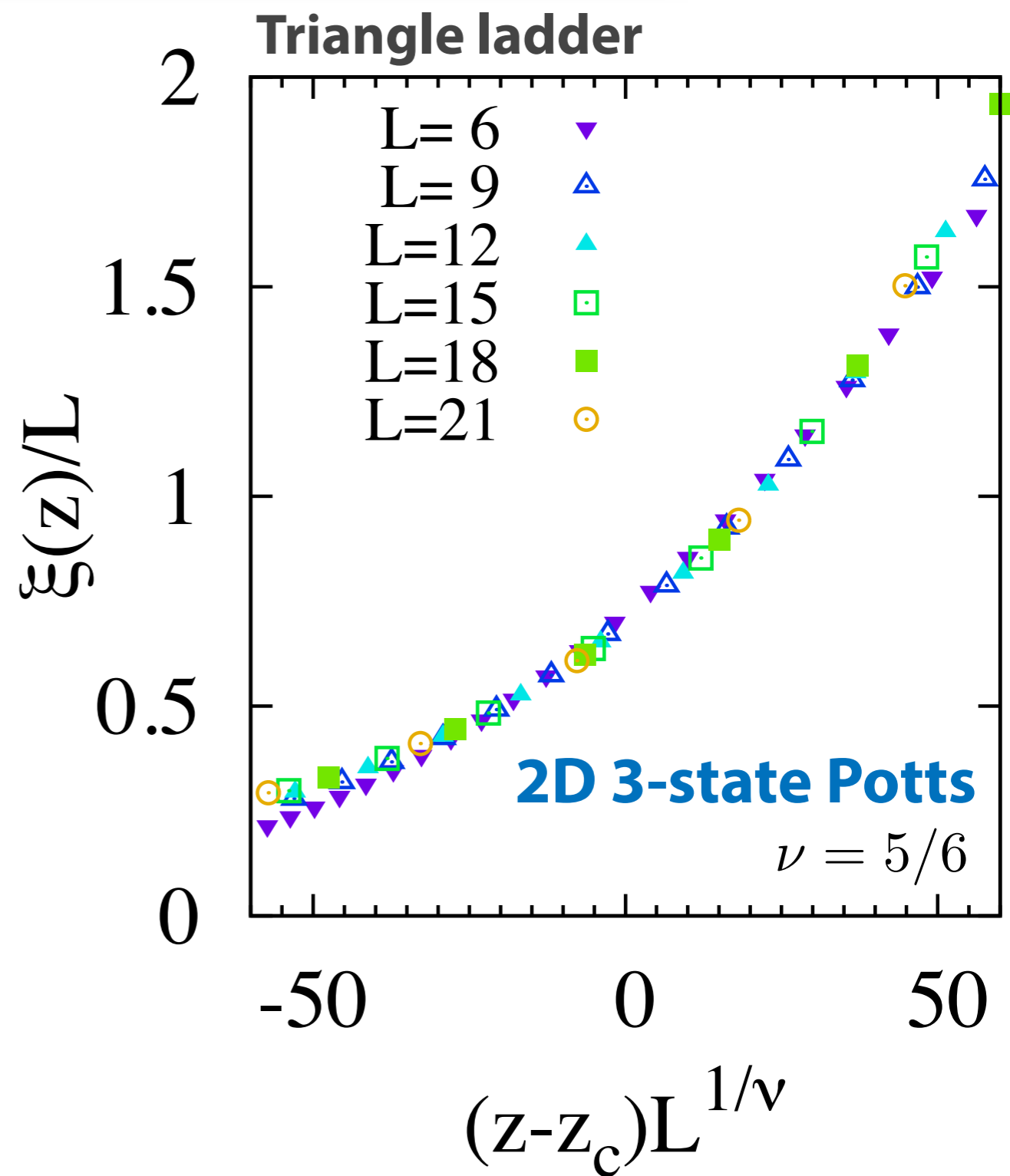
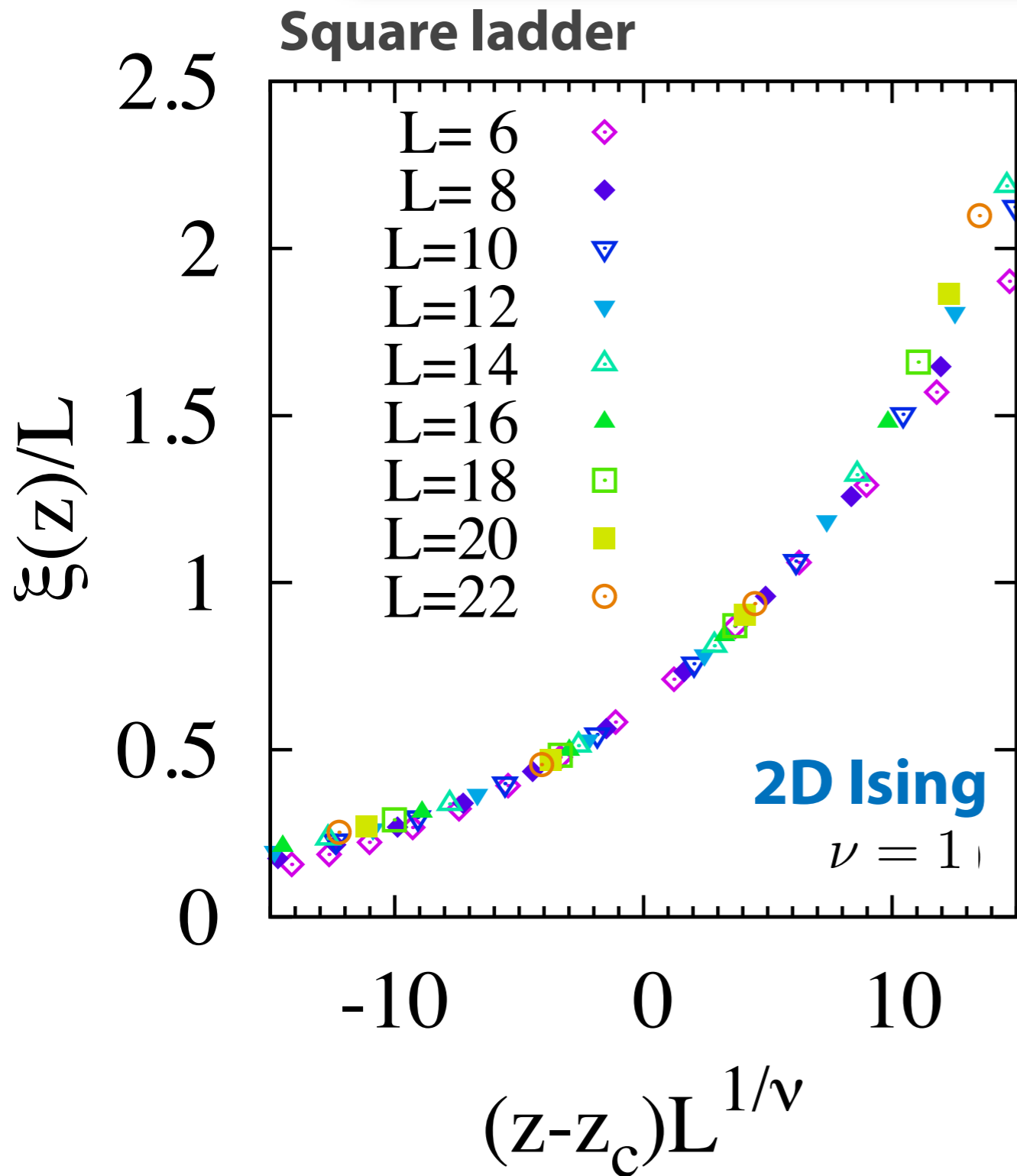
correlation length crosses at z_c



Finite-size scaling for correlation length

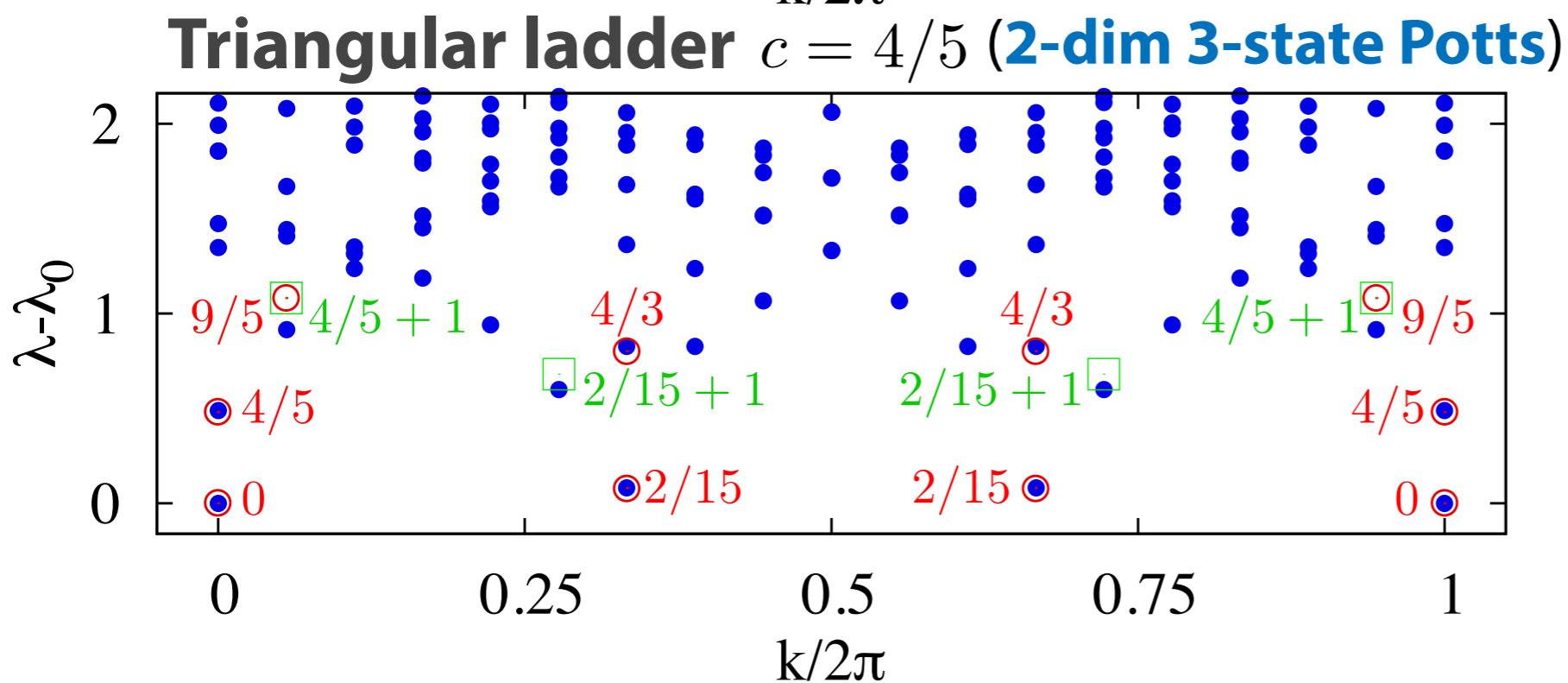
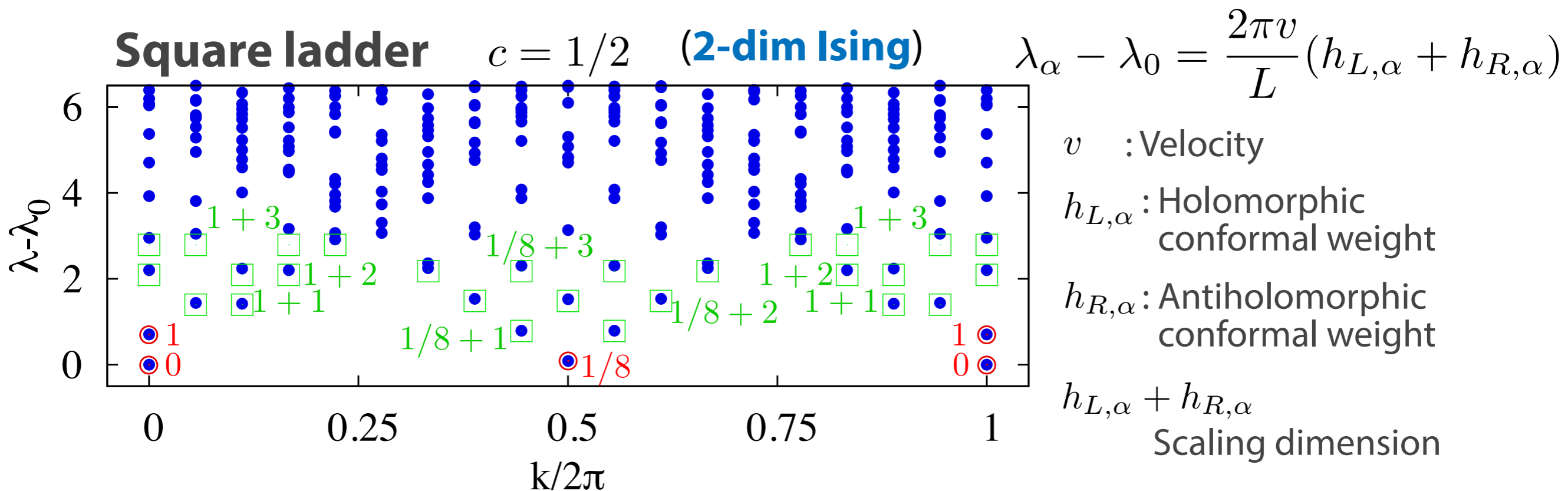
Finite-size scaling

Finite-size scaling relation: $\xi(z)/L = f[(z - z_c)L^{1/\nu}]$



Entanglement spectra at $z=z_c$

Eigenvalues of entanglement Hamiltonian at $z = z_c$



Primary field
Descendant field



M. Henkel "Conformal invariance and critical phenomena" (Springer)

Nested entanglement entropy at $z=z_c$

$|\psi_0\rangle$: Ground state of entanglement Hamiltonian ($z = z_c$)

nested reduced density matrix:

$$\rho(\ell) := \text{Tr}_{\ell+1, \dots, L} [|\psi_0\rangle\langle\psi_0|]$$

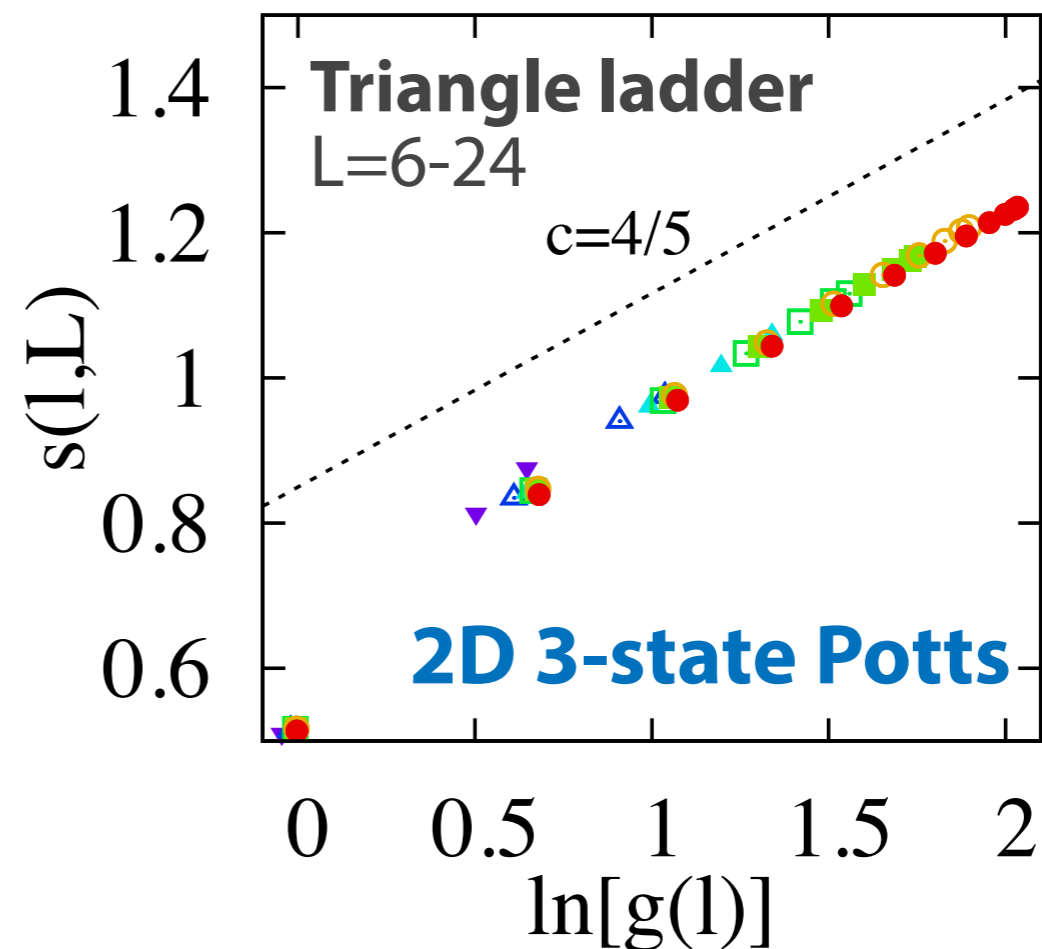
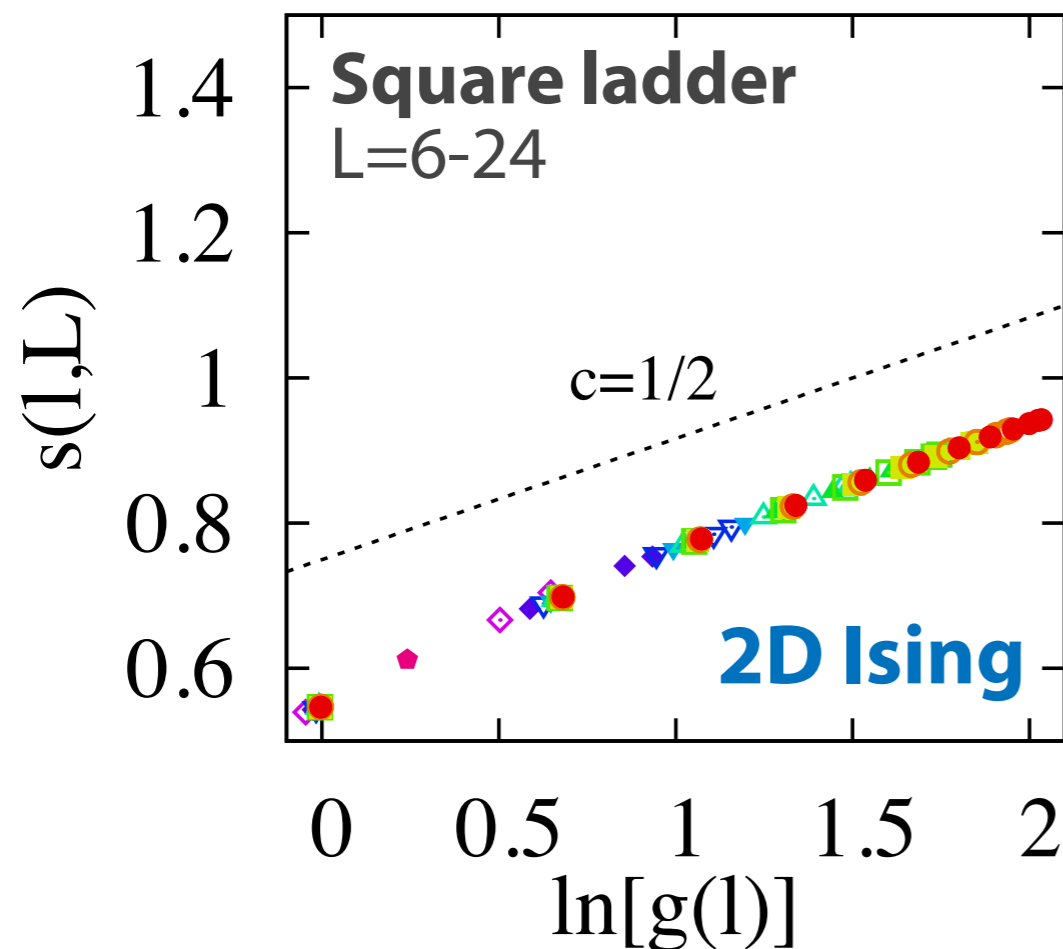
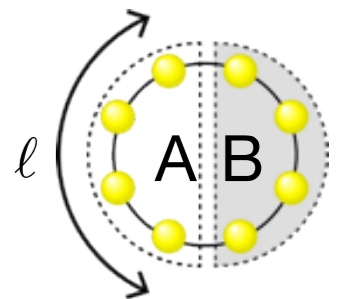
nested entanglement entropy:

$$s(\ell, L) := -\text{Tr}_{1, \dots, \ell} [\rho(\ell) \ln \rho(\ell)]$$

Phys. Rev. B **84**, 245128 (2011).

Interdisciplinary Information Sciences, **19**, 101 (2013)

$$s(\ell, L) = \frac{c}{3} \ln[g(\ell)] + s_1, \quad g(\ell) = \frac{L}{\pi} \sin\left(\frac{\pi\ell}{L}\right)$$



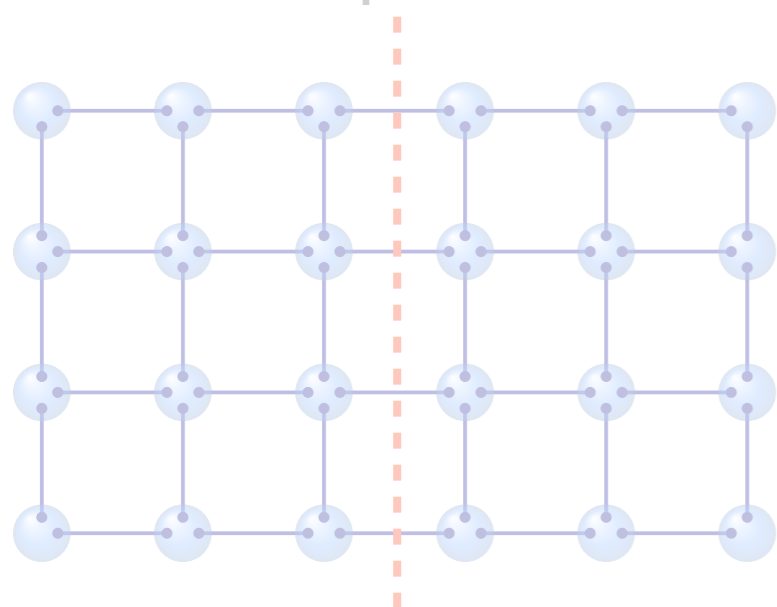
Digest

Entanglement properties of 2D quantum systems

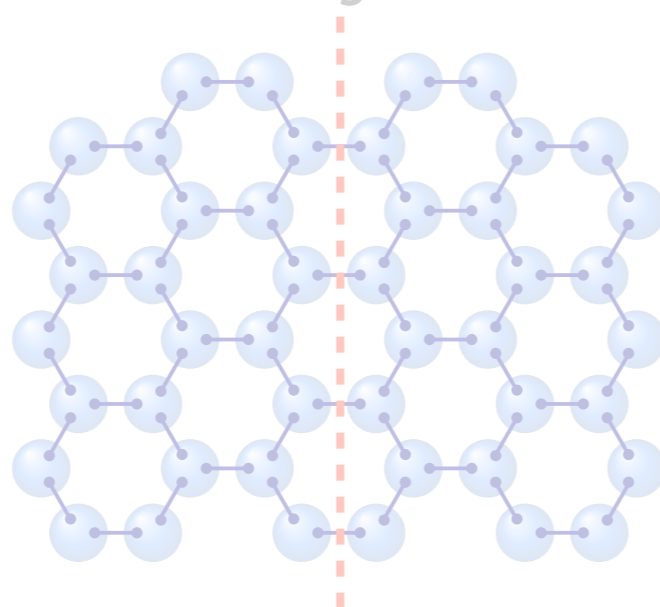
Physical properties of 1D quantum systems



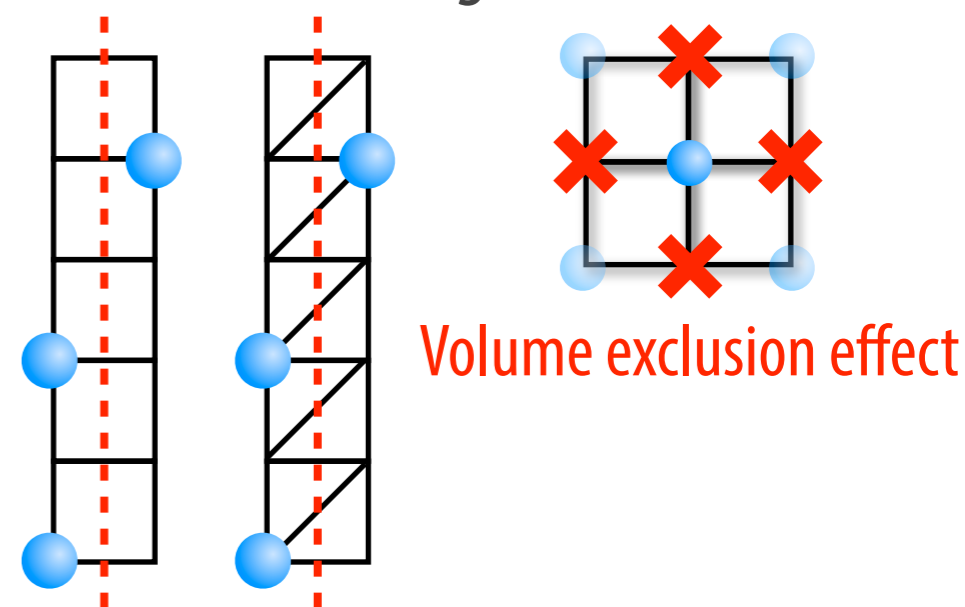
VBS on square lattice



VBS on hexagonal lattice



Quantum lattice gas on ladder



VBS state on 2D lattice

Total system	Entanglement Hamiltonian
Square lattice	1D AF Heisenberg
Hexagonal lattice	1D F Heisenberg

Quantum lattice gas on ladder

Total system	Entanglement Hamiltonian
Square ladder	2D Ising
Triangle ladder	2D 3-state Potts

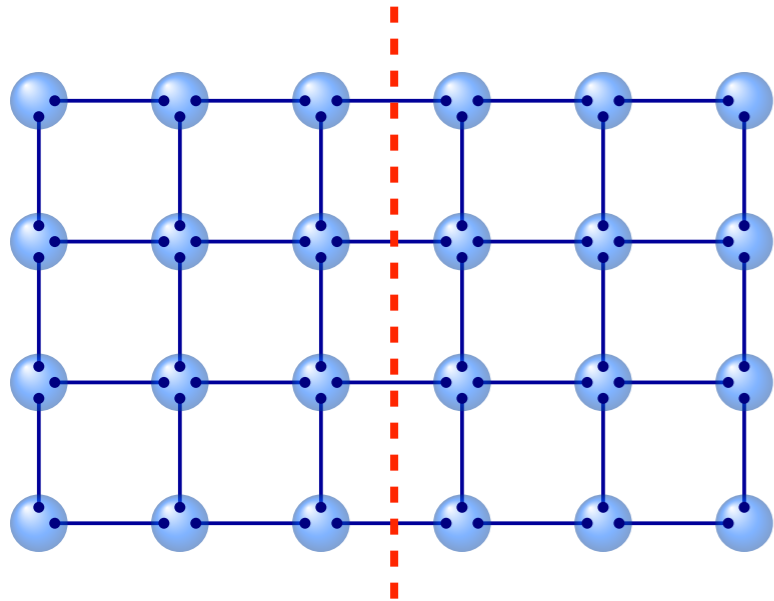
Conclusion

Entanglement properties of 2D quantum systems

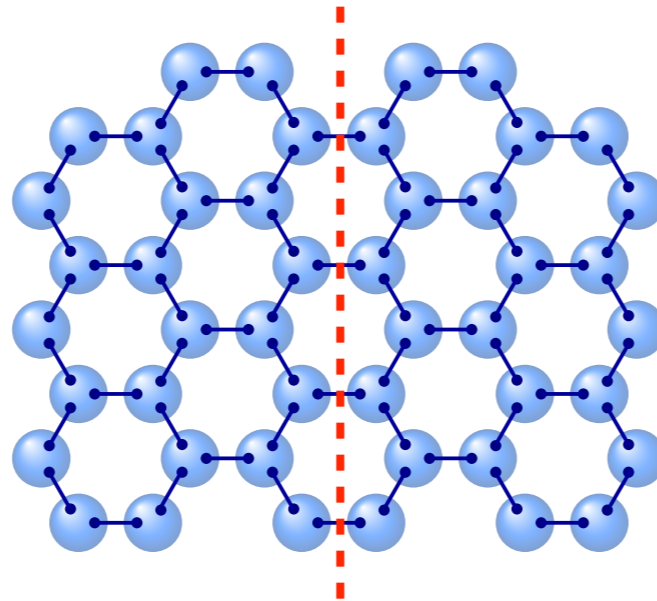


Physical properties of 1D quantum systems

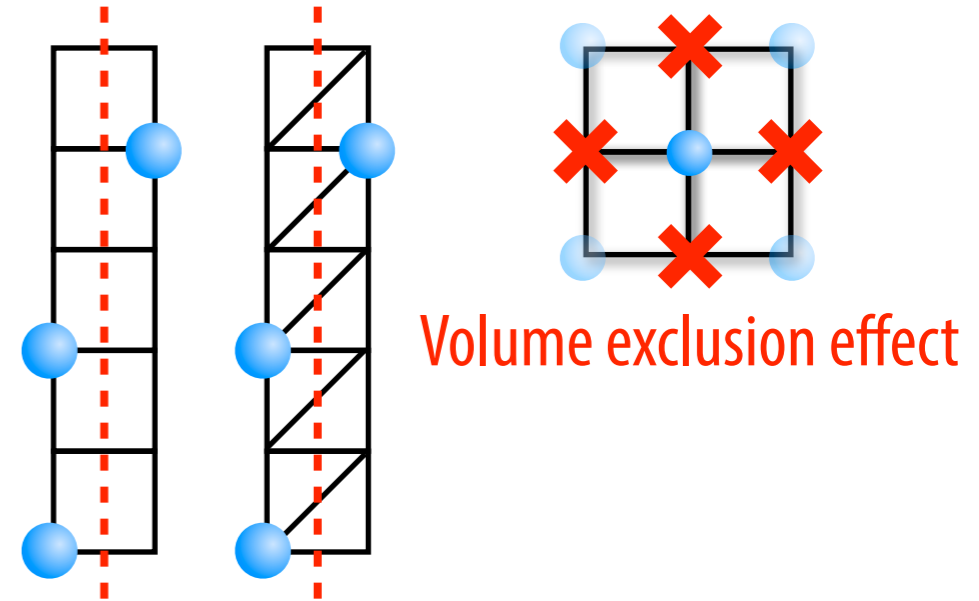
VBS on square lattice



VBS on hexagonal lattice



Quantum lattice gas on ladder



VBS state on 2D lattice

Total system	Entanglement Hamiltonian
Square lattice	1D AF Heisenberg
Hexagonal lattice	1D F Heisenberg

Quantum lattice gas on ladder

Total system	Entanglement Hamiltonian
Square ladder	2D Ising
Triangle ladder	2D 3-state Potts

Thank you for your attention!!

*VBS on symmetric graphs, J. Phys. A, **43**, 255303 (2010)*

*“VBS/CFT correspondence”, Phys. Rev. B, **84**, 245128 (2011)*

*Quantum hard-square model, Phys. Rev. A, **86**, 032326 (2012)*

*Nested entanglement entropy, Interdisciplinary Information Sciences, **19**, 101 (2013)*