

Bell's spin-entangled electron pair generated by local-Fermi-liquid exchange interaction

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Collaboration:

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Discussion:

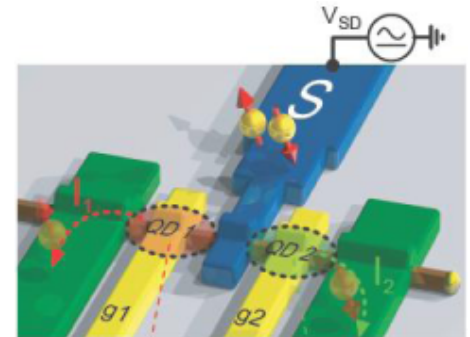
Shiro Kawabata (AIST)

Bell's pair generator by many-body effect

Cooper pair splitter

Theory: Chtchelkatchev *et al.*, PRB **66**, 161320(R) (2002)

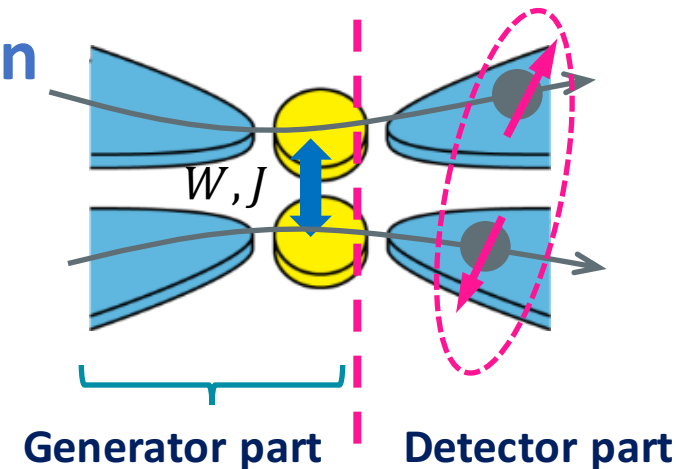
Experiment: Hofstetter *et al.*, Nature **461**, 960-963 (2009).



New idea!

Pair by Local-Fermi-liquid interaction

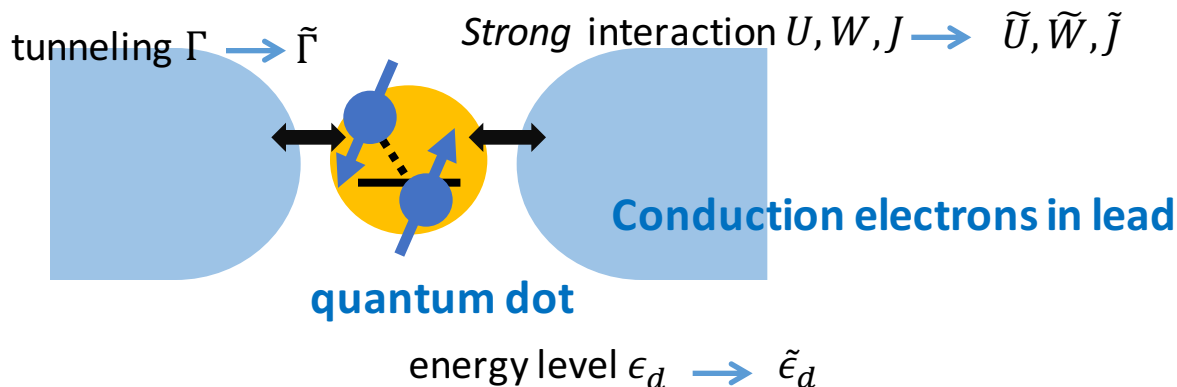
Bell's pair in current through a double quantum dot



Local Fermi-liquid

Landau's Fermi liquid theory \longrightarrow Quantum dot (impurity) systems

Ground state of “the Kondo effect”



Outline

Introduction

Electric transport and Kondo effect in quantum dot

Charge pair creation mechanism

shot noise measurement for Kondo effect in quantum dot

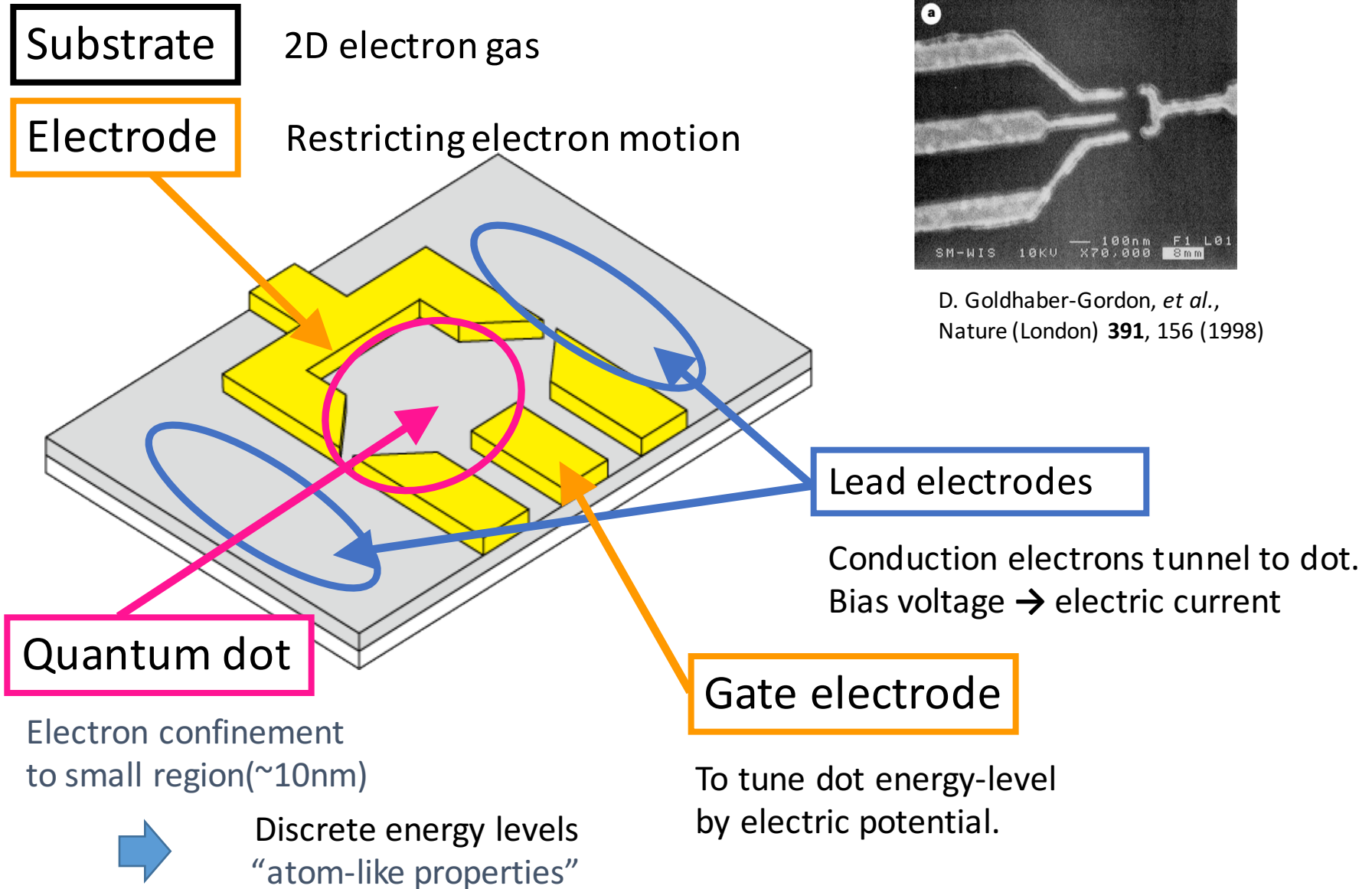
Theory: Sela *et al.*, PRL (2006), Gogolin and Komnik PRL (2006), RS *et al.*, PRL (2012)

Experiment: Ferrier, RS *et al.*, Nat. Phys. (2015)

Bell's pair creation in double quantum dot

Electric transport and Kondo effect in quantum dot

Quantum dot

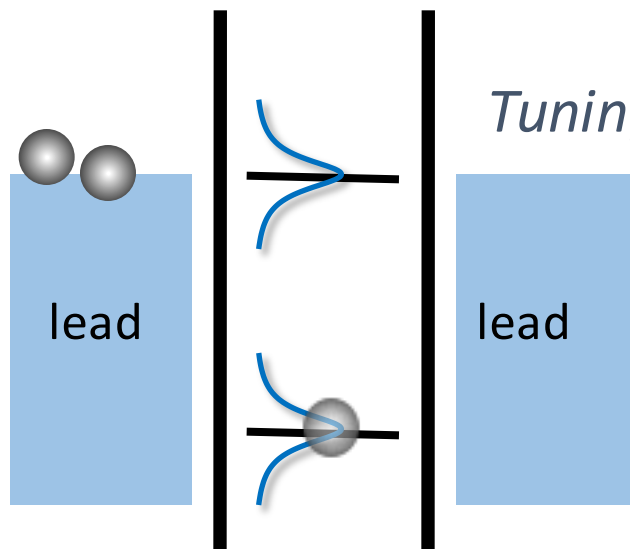


Transport through quantum dot

Tunnel effect → Electron transport

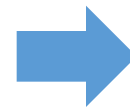
〈Energy schematic of quantum dot〉

Coulomb blockade



Tuning gate electrode

One-body effect



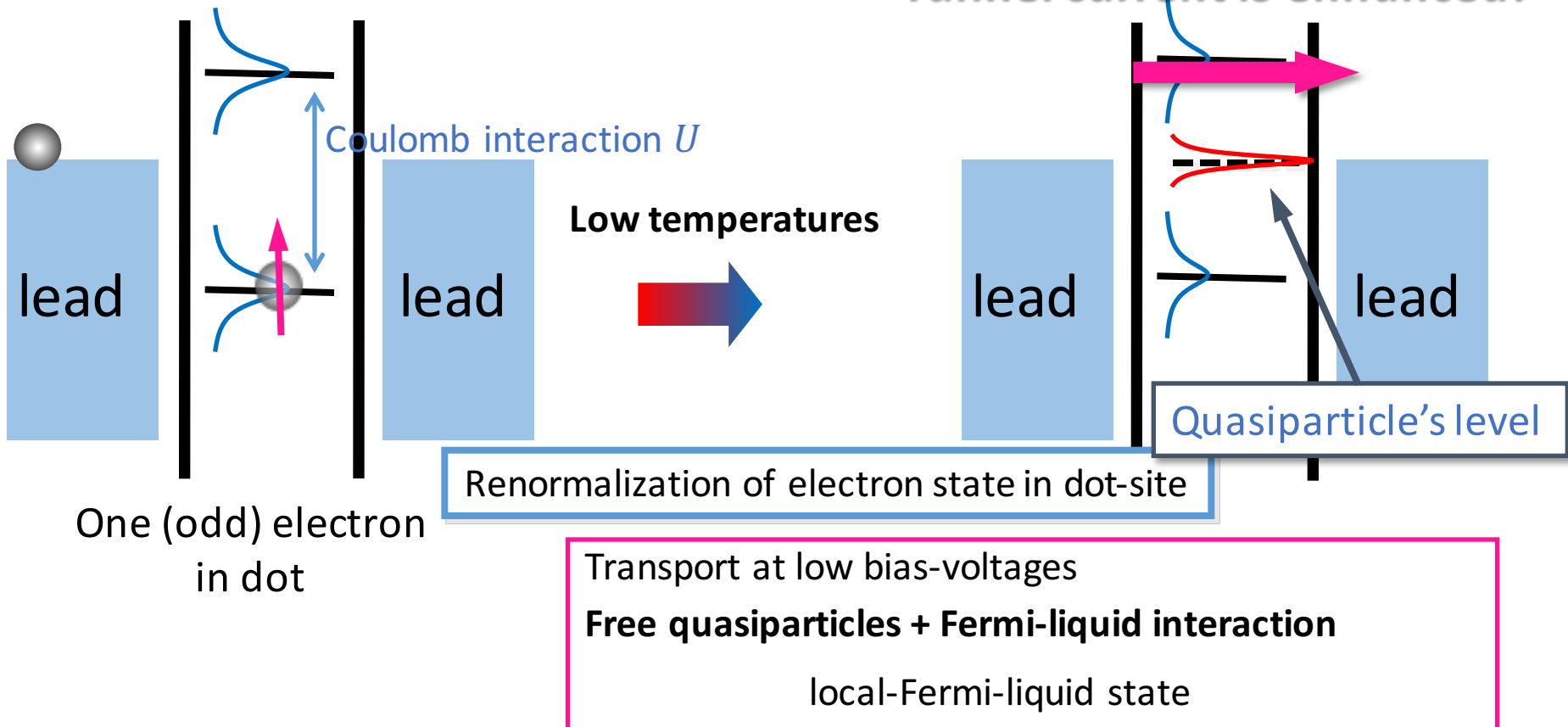
Many-body effect

Quantum dot

Kondo effect in quantum dot

Coulomb blockade

No transport occurs.



Detection of pair charge creation in local Fermi-liquid

Shot noise measurement

Current generated by rare event (Poissonian process)

At zero temperature where no thermal noise

Discretized charge state  Current noise

$$\frac{S}{\text{Average current}} = 2e^*$$

The equation shows the spectral density S divided by the average current $\langle I \rangle$ is equal to $2e^*$. A blue arrow points from the text 'Current-current correlation (current noise)' below to the S in the numerator of the fraction.

Current-current correlation (current noise)

$$S := \int dt \langle \delta I(t) \delta I(0) + \delta I(0) \delta I(t) \rangle$$

$$\text{Current fluctuation: } \delta I(t) := I(t) - \langle I \rangle$$

The proportional constant e^*

Effective charge of current carrying states

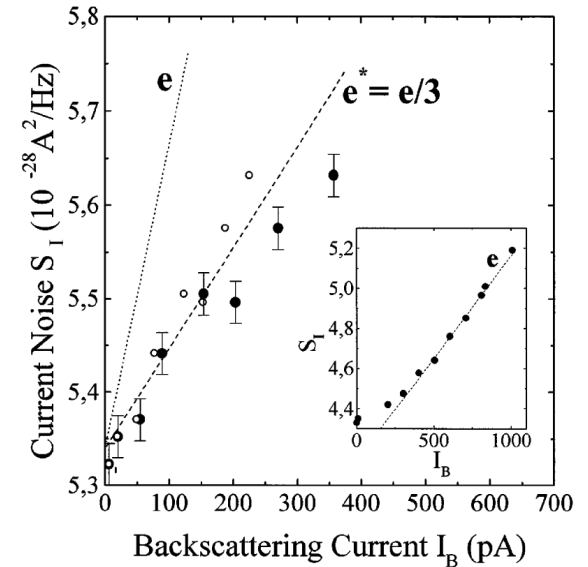
Effective charges of many-body systems

Fractional quantum Hall system

Fractional charges of Laughlin's quasi-particle:
 $e^* = e/3$, etc.

Picciotto, *et al.*, Nature **389**, 162 (1997).

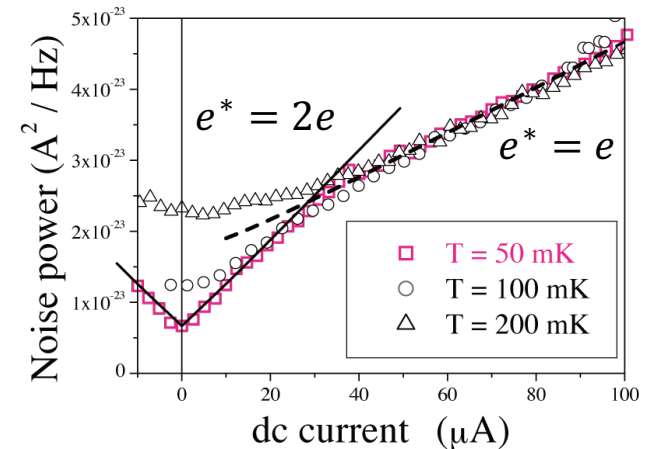
Saminadayar, *et al.*, PRL **79**, 2526 (1997).



Superconductor/normal metal junction

Cooper pair charge: $e^* = 2e$

Lefloch, *et al.*, PRL **90**, 067002 (2003).



Question

Effective charge of Kondo-correlated quantum dots?

Our model

Impurity Anderson model

$$\mathcal{H}_A = \mathcal{H}_0 + \mathcal{H}_I$$

- lead part ($\alpha = L, R$) and quantum dot

$$\mathcal{H}_0 = \sum_{k\alpha\sigma} \epsilon_k c_{k\alpha\sigma}^\dagger c_{k\alpha\sigma} + \sum_{\sigma} \epsilon_d d_{\sigma}^\dagger d_{\sigma}$$

σ : spin

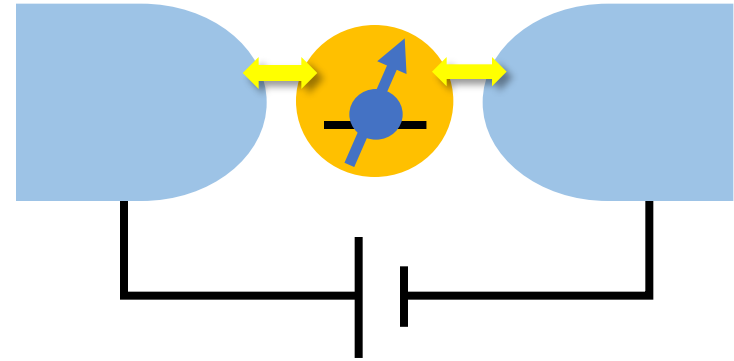
$$+ \sum_{k\alpha} \left(\frac{v_{\alpha}}{\sqrt{\mathcal{N}}} d_{\sigma}^\dagger c_{k\alpha\sigma} + \text{H.c.} \right)$$

lead-dot tunneling

- Coulomb interaction

$$\mathcal{H}_I = U d_{\uparrow}^\dagger d_{\uparrow} d_{\downarrow}^\dagger d_{\downarrow}$$

lead electrode L quantum dot lead electrode R
 $\mu_L = eV/2$ $\mu_R = -eV/2$



Applied bias-voltages eV

Dot level linewidth: $\Gamma = \frac{1}{2}(\Gamma_L + \Gamma_R)$

$$\Gamma_{\alpha} := 2\pi\rho_c v_{\alpha}^2$$

ρ_c : Density of state for conduction electrons in lead electrodes

Reorganized perturbation

Original model

Perturbation expansion up to *infinite* order in interaction U

➡ Exact transport quantities up to V^3 (Local Fermi liquid)



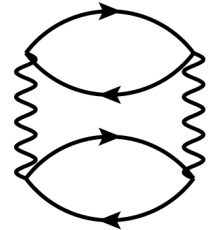
We know how the parameters are renormalized.

Quasiparticle model (local Fermi liquid)

Perturbation expansion up to *2nd* order in Fermi-liquid interaction \tilde{U}

+ counter term to avoid overcounting

➡ Exact transport quantities up to V^3 (Local Fermi liquid)



good point

Contribution decomposed to perturbation processes in Fermi-liquid interaction \tilde{U} .

Quasiparticle Hamiltonian

Replacement **electron in dot** $d_\sigma \rightarrow$ **quasiparticle** \tilde{d}_σ

interaction $U \rightarrow \tilde{U}$

energy level $\epsilon_d \rightarrow \tilde{\epsilon}_d$

tunneling $v_\alpha \rightarrow \tilde{v}_\alpha$

$$\tilde{\mathcal{H}}_{qp} = \tilde{\mathcal{H}}_0 + \tilde{\mathcal{H}}_I$$

- **lead part ($\alpha = L, R$) and quasiparticle**

$$\tilde{\mathcal{H}}_0 = \sum_{k\alpha\sigma} \epsilon_k c_{k\alpha\sigma}^\dagger c_{k\alpha\sigma} + \sum_{\sigma} \tilde{\epsilon}_d \tilde{d}_\sigma^\dagger \tilde{d}_\sigma + \sum_{k\alpha} \left(\frac{\tilde{v}_\alpha}{\sqrt{\mathcal{N}}} \tilde{d}_\sigma^\dagger c_{k\alpha\sigma} + \text{H.c.} \right)$$

σ : spin

lead-quasiparticle tunneling

- **Fermi liquid interaction**

$$\mathcal{H}_I = \tilde{U} \tilde{d}_\uparrow^\dagger \tilde{d}_\uparrow \tilde{d}_\downarrow^\dagger \tilde{d}_\downarrow$$

quasiparticle's linewidth: $\tilde{\Gamma} = \frac{1}{2} (\tilde{\Gamma}_L + \tilde{\Gamma}_R)$

$$\tilde{\Gamma}_\alpha := 2\pi\rho_c \tilde{v}_\alpha^2$$

Electric current

Oguri, PRB **64**, 153305 (2001)

Low-bias current through spin Kondo dot

$$I = \frac{2e^2}{h} V \left[\mathbf{1} - \left[\frac{\mathbf{1}}{\mathbf{12}} + \frac{\mathbf{5}}{\mathbf{12}} \left(\frac{\tilde{U}}{\pi\tilde{\Gamma}} \right)^2 \right] \left(\frac{eV}{\tilde{\Gamma}} \right)^2 \right] + O(V^4)$$

$=: I_b$

V -linear term:

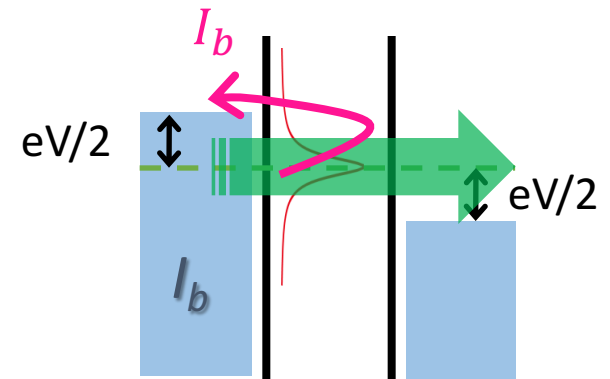
Free quasiparticle contribution

Perfect transmission

V^3 term:

Fermi-liquid interaction enhances the backscattering current

with small probability “**Poissonian process**”.



Noise-current ratio $\frac{S}{2eI_b}$ for the backscattering current

Noise-current ratio

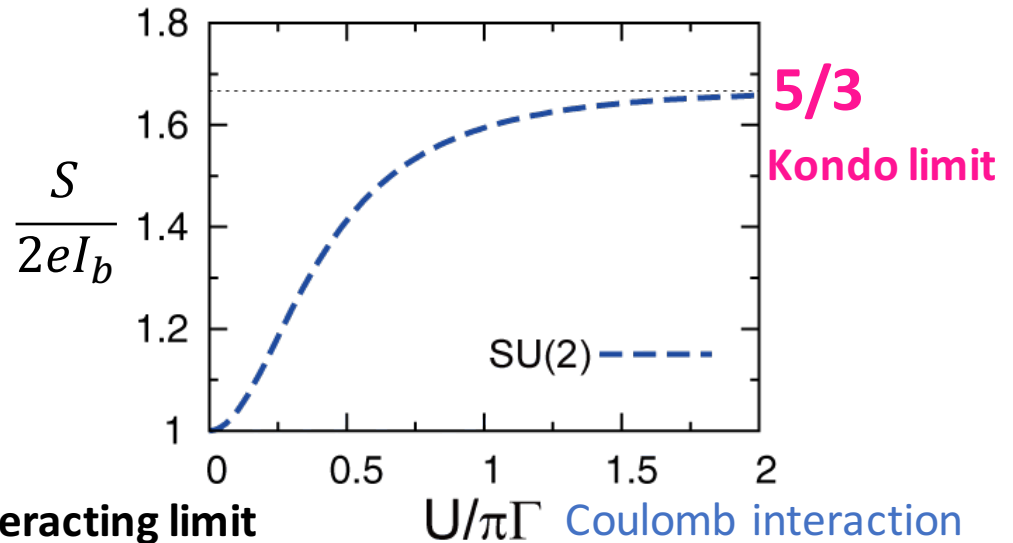
Local-Fermi-liquid calculation

$$\rightarrow \frac{S}{2eI_b} = \frac{1+9\left(\frac{\tilde{U}}{\pi\tilde{\Gamma}}\right)^2}{1+5\left(\frac{\tilde{U}}{\pi\tilde{\Gamma}}\right)^2}$$

Fermi-liquid interaction $\frac{\tilde{U}}{\pi\tilde{\Gamma}}$

Exact solution

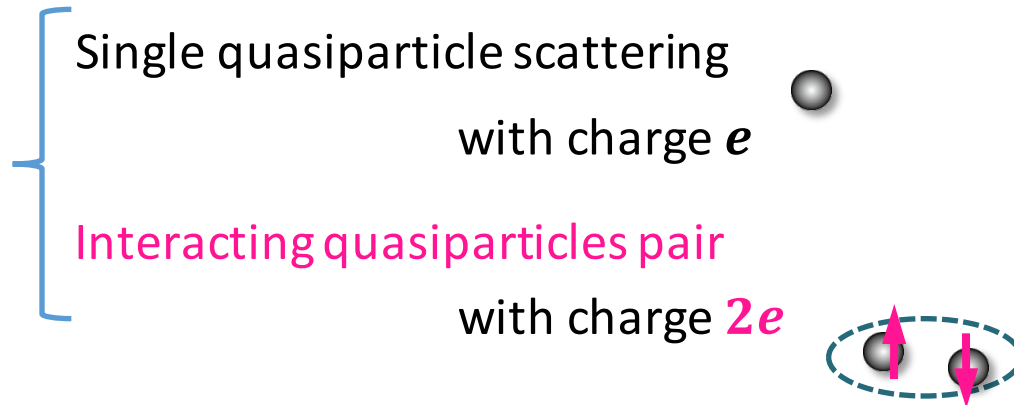
Non-interacting limit



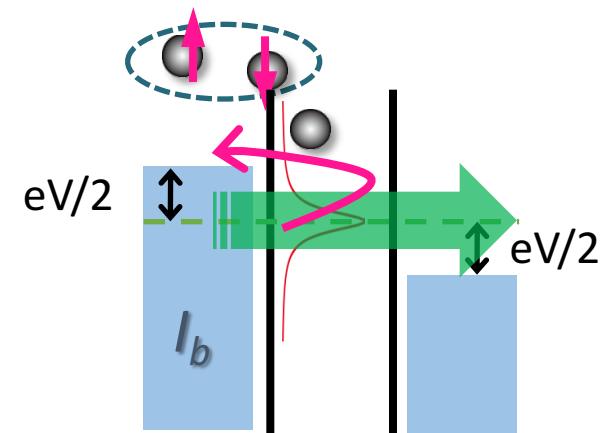
The meaning of the values of $\frac{S}{2eI_b}$

Analysis of scattering process in quasiparticle picture

Backscattering current



➔ $1 \leq \frac{S}{2eI_b} < 2$

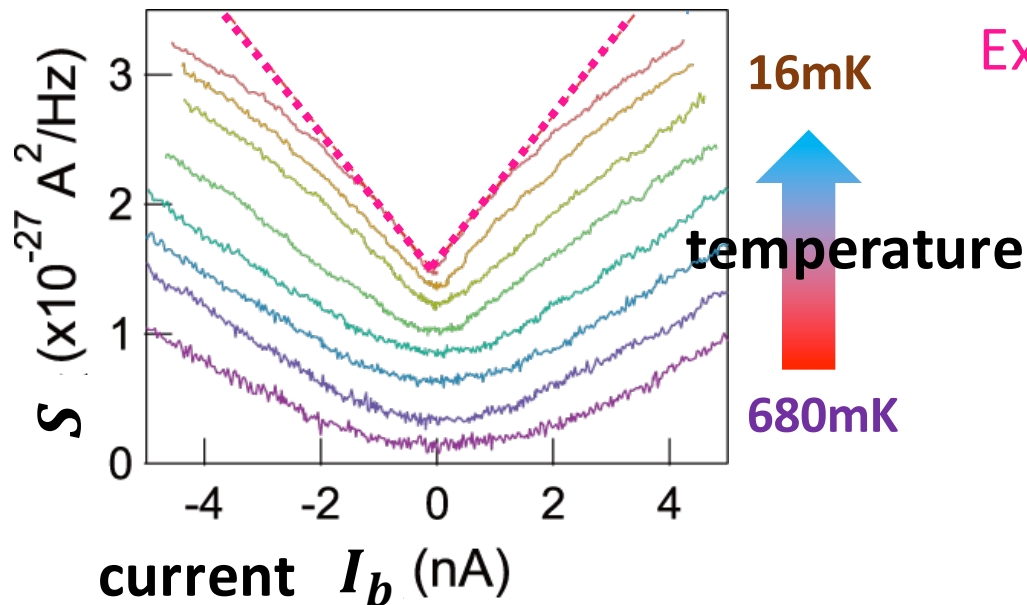


**“Not elemental charge of excited states,
but evidence of the charge pair creation by the Fermi-liquid interaction.**

Experiment in SU(2) Kondo dot

M. Ferrier, RS *et al.*, Nat. Phys. (2015)

Shot noise measurement



Experiment: $\frac{S}{2eI_b} = 1.7 \pm 0.1$

Good agreement with

theory: $\frac{S}{2eI_b} = 1.66 \dots$

Direct measurement on “charge pair” creation.

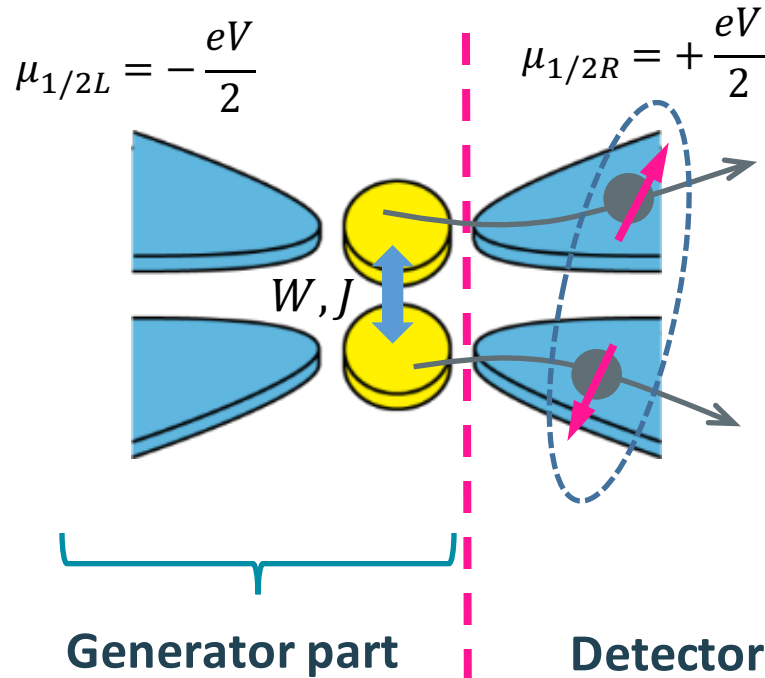
This is not our goal, yet.

**Bell's pair creation
by local-Fermi-liquid interaction**

Double quantum dot device

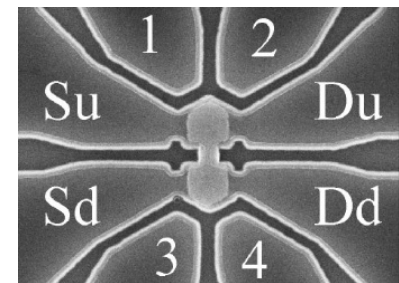
Motivation

Testing “quantum” correlation of quasiparticle pair



- Coulomb interaction
- Exchange coupling

Correlation between channels



Bell's correlation in current

Chtchelkatchev *et al.*, PRB **66**, 161320(R) (2002)

Experimental observable quantity

Electric current



Counting number of electrons

coming from lead $\alpha = \text{Left, Right}$ to the dot
with spin orientating θ
through channel $m = 1, 2$
in time interval $[t, t + \tau]$

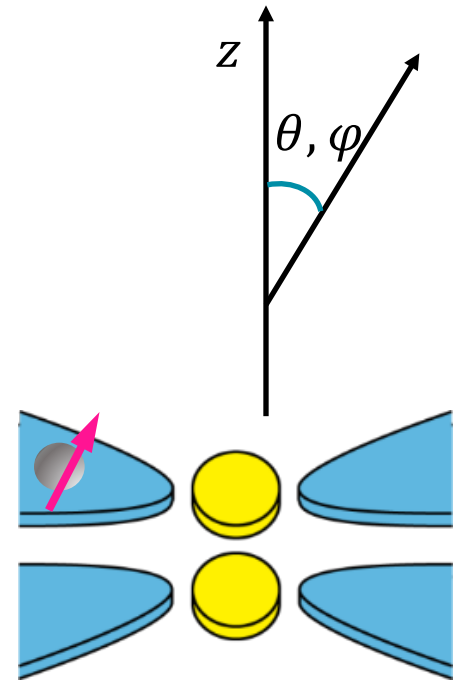
$$N_{Rm\theta}(t, t + \tau) = \frac{1}{e} \int_t^{t+\tau} dt' I_{Rm\theta}(t')$$

Generated spin per a particle

$$\frac{N_{\alpha m \uparrow \theta} - N_{\alpha m \downarrow \theta}}{N_{\alpha m \uparrow \theta} + N_{\alpha m \downarrow \theta}}$$



Bell's correlation

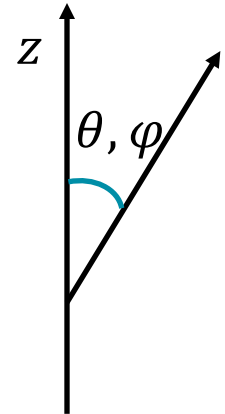


CHSH type Bell's correlation

Current-current correlation with twisted spin orientation

$$S(\theta, \varphi) := \int dt \langle \delta I_{R1\theta}(t) \delta I_{R2\varphi}(0) \rangle$$

Current fluctuation: $\delta I_{\alpha m \theta}(t) := I_{\alpha m \theta}(t) - \langle I_{\alpha m \theta} \rangle$

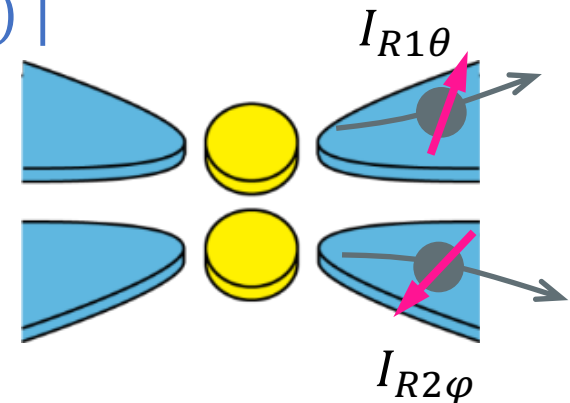


CHSH type Bell's correlator

Four spin angle: $\theta, \theta', \varphi, \varphi'$

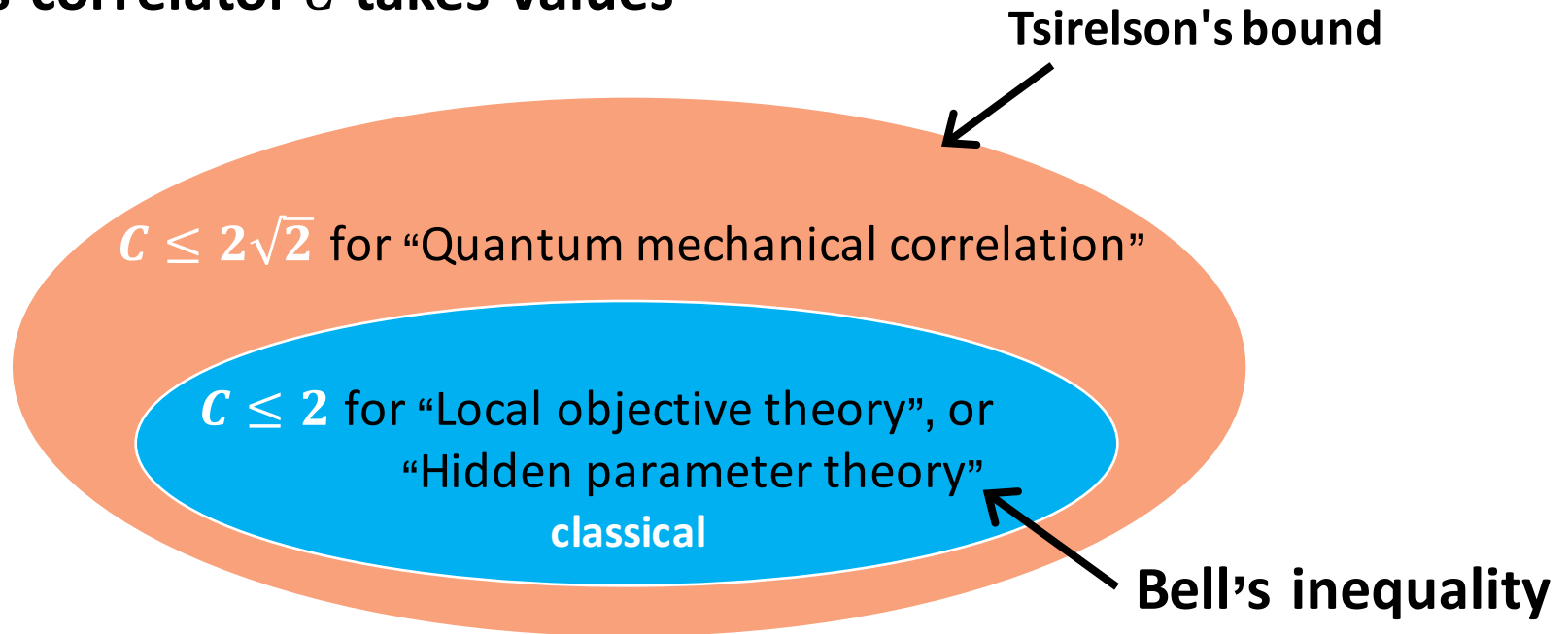
$$C = |F(\theta, \varphi) - F(\theta', \varphi) + F(\theta, \varphi') + F(\theta', \varphi')|$$

$$F(\theta, \varphi) := \frac{S(\theta, \varphi) - S(\theta + \pi, \varphi) - S(\theta, \varphi + \pi) + S(\theta + \pi, \varphi + \pi)}{S(\theta, \varphi) + S(\theta + \pi, \varphi) + S(\theta, \varphi + \pi) + S(\theta + \pi, \varphi + \pi)}$$



Upper bound of the CHSH Bell's correlator

Bell's correlator C takes values



Sufficient condition for quantum correlation

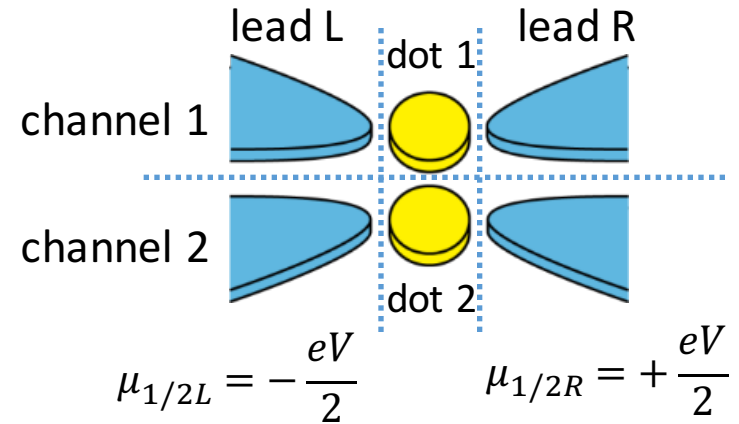
$$2 < C \leq 2\sqrt{2}$$

Orbital degenerate impurity Anderson model

$$\mathcal{H}_A = \mathcal{H}_0 + \mathcal{H}_T + \mathcal{H}_I$$

$$\mathcal{H}_0 = \sum_{k\alpha m\sigma} \epsilon_k c_{k\alpha m\sigma}^\dagger c_{k\alpha m\sigma} + \sum_m \epsilon_d d_{m\sigma}^\dagger d_{m\sigma}$$

Electric lead $\alpha = L, R$ **dot**
 $m=1,2$: channel
 $\sigma = \uparrow, \downarrow$: spin



$$\mathcal{H}_T = \sum_{k\alpha m} \left(\frac{v_\alpha}{\sqrt{N}} d_{m\sigma}^\dagger c_{k\alpha m\sigma} + \text{H.c.} \right) \quad \text{Tunneling between lead and dot}$$

“Conservation of spin and channel” \rightarrow **Local Fermi liquid state**

$$\mathcal{H}_I = \sum_{m=1,2} \mathbf{U} d_{m\uparrow}^\dagger d_{m\uparrow} d_{m\downarrow}^\dagger d_{m\downarrow} + \sum_{\sigma\sigma'} \mathbf{W} d_{1\sigma}^\dagger d_{1\sigma} d_{2\sigma'}^\dagger d_{2\sigma'}$$

Intra- Coulomb repulsion *Inter-dot Coulomb repulsion*

$$+ 2\mathbf{J} \mathbf{S}_{d1} \cdot \mathbf{S}_{d2'} \quad \text{“Exchange coupling } J$$

- **particle-hole symmetric:** $\epsilon_d = -U/2 - (M - 1)W$
symmetric coupling: $v_L = v_R$

Many-body effect

Quasiparticle Hamiltonian

Replacement

Electron in dot $d_{m\sigma}$ \rightarrow Quasiparticle $\tilde{d}_{m\sigma}$

interaction U, W, J \rightarrow $\tilde{U}, \tilde{W}, \tilde{J}$

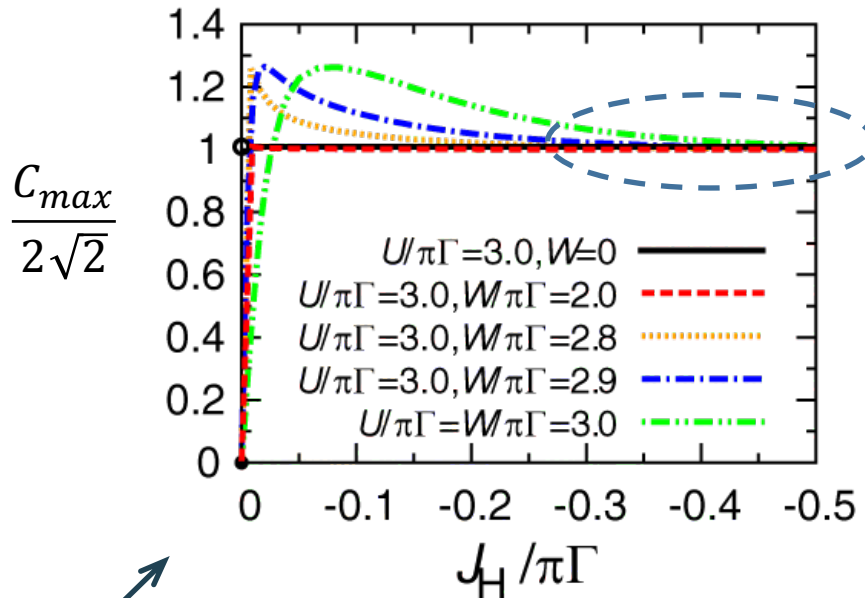
energy level ϵ_d \rightarrow $\tilde{\epsilon}_d$

tunneling v_α \rightarrow \tilde{v}_α

Second order perturbation in local-Fermi-liquid interactions $\tilde{U}, \tilde{W}, \tilde{J}$

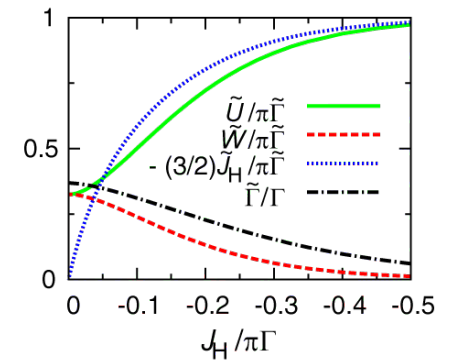
+ counter term to avoid overcounting

Maximum value of Bell's correlator

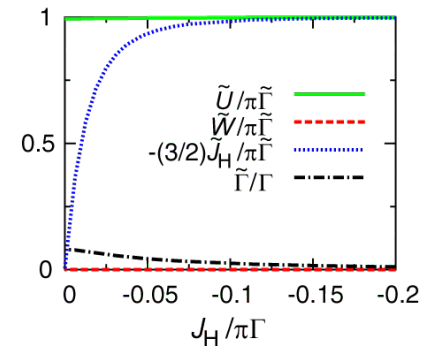


Numerical renormalization Group for Fermi liquid interactions

$$U/\pi\Gamma = W/\pi\Gamma = 3$$



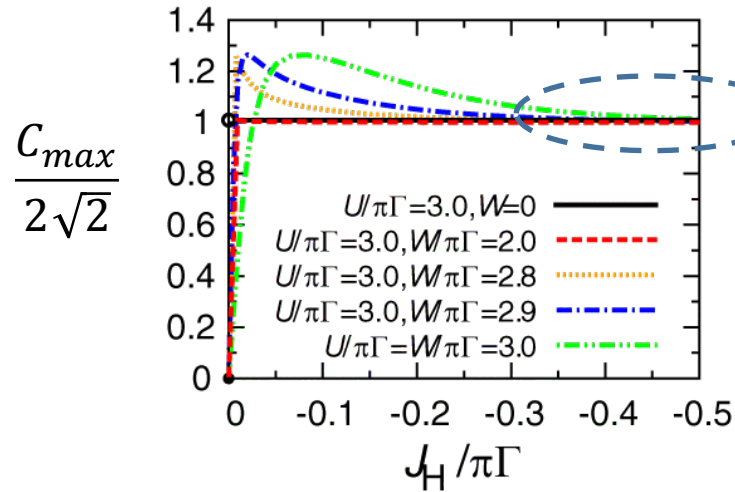
$$U/\pi\Gamma = 3, W = 0$$



$$J = 0$$

Coulomb interactions U, W are independent of spin orientation.

Large ferromagnetic coupling



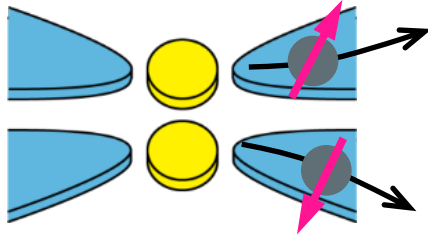
Suppression of charge fluctuation between channels $\tilde{W} \rightarrow 0$

Maximum value of the Bell correlator

$$C_{max} \rightarrow 2\sqrt{2}$$

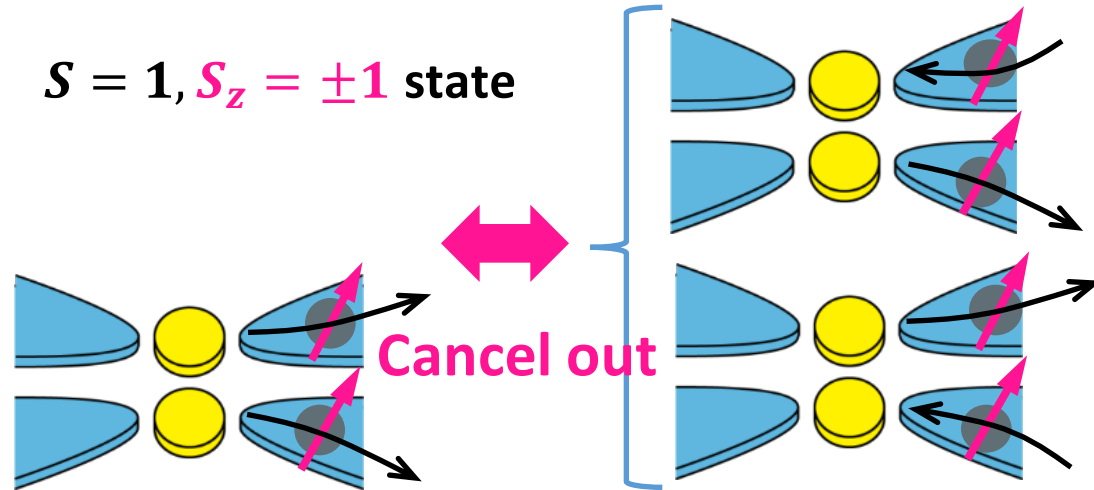
Spin triplet Bell's pair

$S = 1, S_z = 0$ state



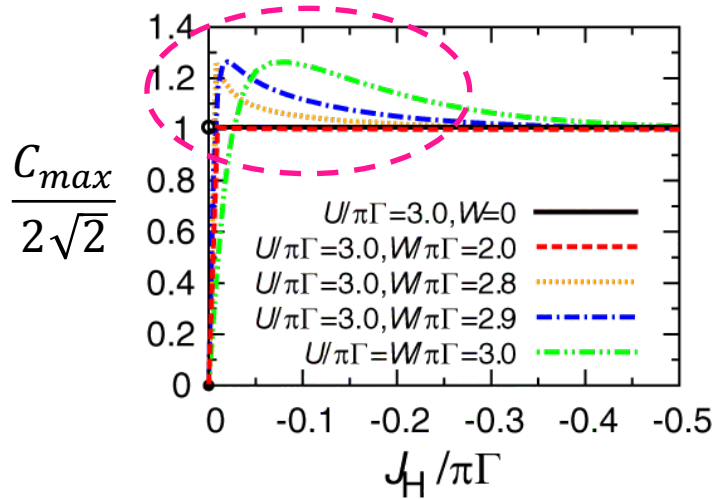
Only this state is observed!

$S = 1, S_z = \pm 1$ state



sign of current-current correlation opposite

Larger than quantum upper bound?

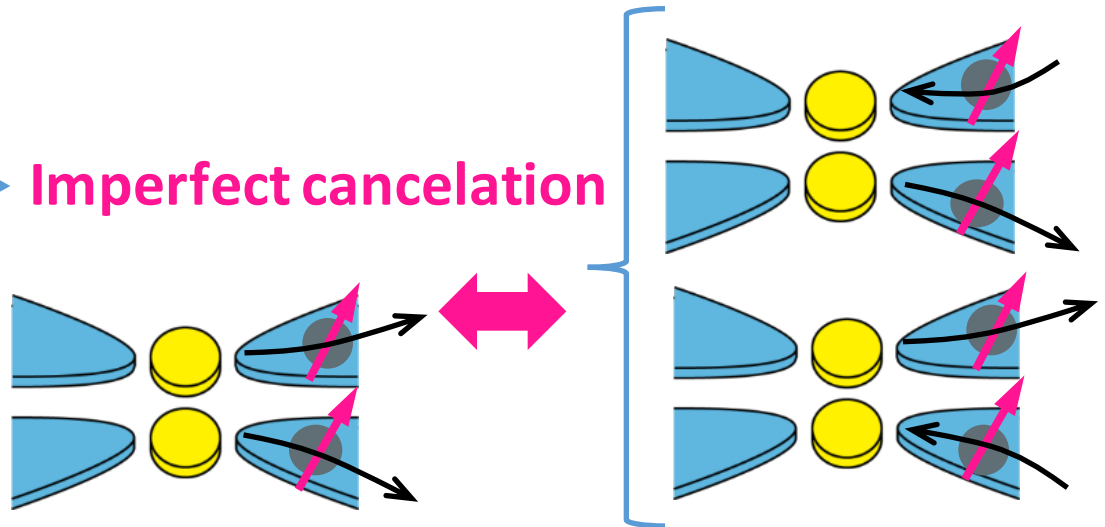


Maximum value of the Bell correlator

$$C_{max} > 2\sqrt{2} \quad \text{why?}$$

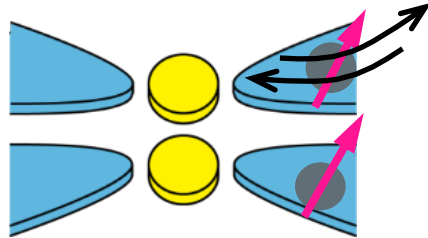
$S = 1, S_z = \pm 1$ state

intermediate \tilde{W} \longrightarrow Imperfect cancelation



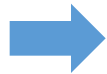
Triplet pair $S = 1, S_z = \pm 1$ state

Particles are generated in two directions in each channel.



Simple integration

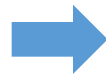
$$N_{Rm\theta}(t, t + \tau) = \frac{1}{e} \int_t^{t+\tau} dt' I_{Rm\theta}(t')$$



Underestimated number of generated particles.

Conjecture

Normalization of Bell's correlator is failed.



$$C_{max} > 2\sqrt{2} \text{ (Tsirelson's bound)}$$

Conclusion

Double quantum dot in local Fermi liquid state (Kondo state)

Exchange-type interaction

➔ Quasiparticle triplet pair

Observable Bell pair

$$S = 1, S_z = 0$$

