

# Dynamical vertex approximation — a step beyond dynamical mean field theory

K. Held

MPI-FKF Stuttgart → TU Vienna, as of March 2008

*YKIS, Nov 26, 2007*

## 1) Dynamical vertex approximation

- Motivation
- Method
- Results for 3D, 2D, and 1D Hubbard model

## 2) Kinks in the dispersion relation of correlated electrons



# Dynamical vertex approximation — a step beyond dynamical mean field theory

K. Held

MPI-FKF Stuttgart → TU Vienna, as of March 2008

*YKIS, Nov 26, 2007*

## 1) Dynamical vertex approximation (D $\Gamma$ A)

- Motivation
- Method
- Results for 3D, 2D, and 1D Hubbard model

## 2) Kinks in the dispersion relation of correlated electrons

# Thanks to...

## 1) Dynamical vertex approximation ( $D\Gamma A$ )

[A. Toschi](#), [A. Katanin](#) – MPI-FKF Stuttgart

PRB 75, 045118 (2007)

## 2) Kinks

[Y.-F. Yang](#) – MPI-FKF Stuttgart

[K. Byczuk](#), [M. Kollar](#), [D. Vollhardt](#) – Augsburg

[I. A. Nekrasov](#) – Ekaterinburg

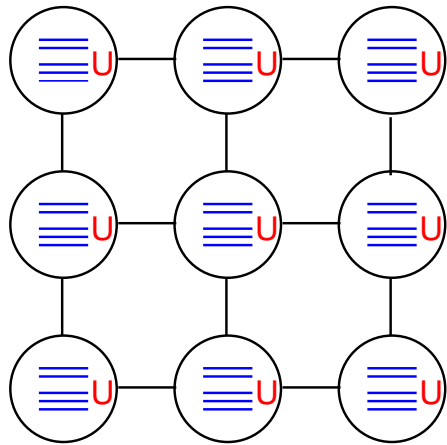
[Th. Pruschke](#) – Göttingen

PRB 73, 155112 (2006)

Nature Physics 3 168 (2007)

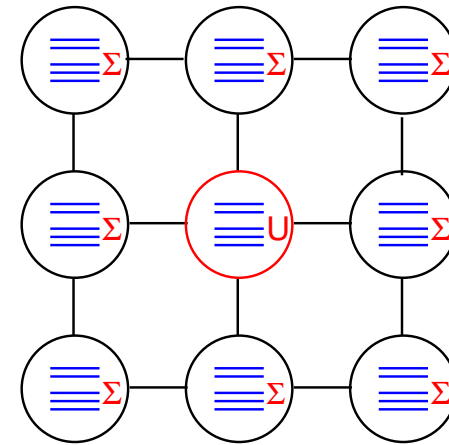
# Motivation

## Dynamical mean field theory



DMFT

(Metzner, Vollhardt'89; Georges, Kotliar'92)



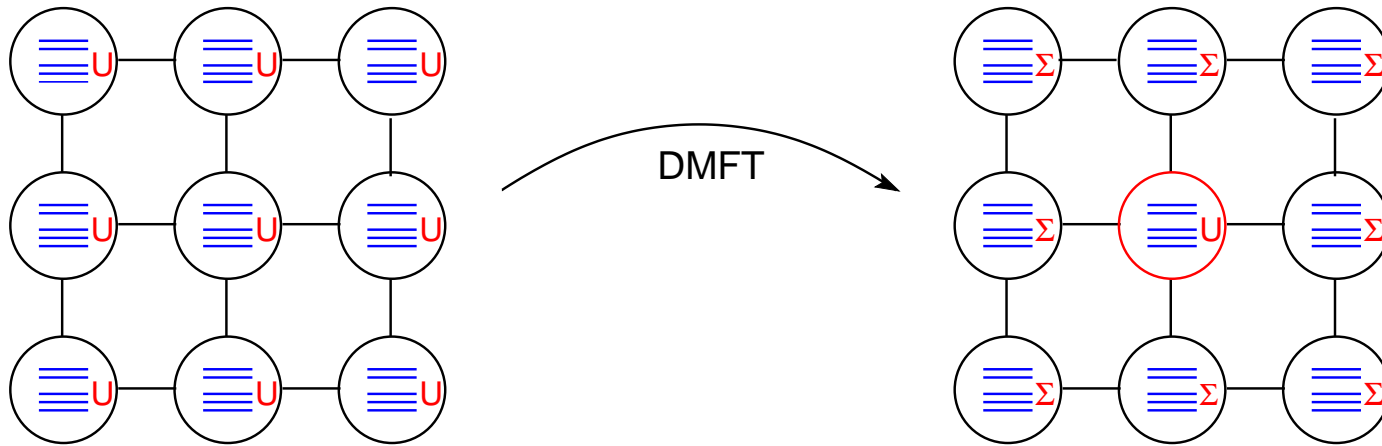
$\Sigma$  all topologically distinct, but **local** diagrams

**Success story:** quasiparticle renormalizations, magnetism, kinks ...

# Motivation

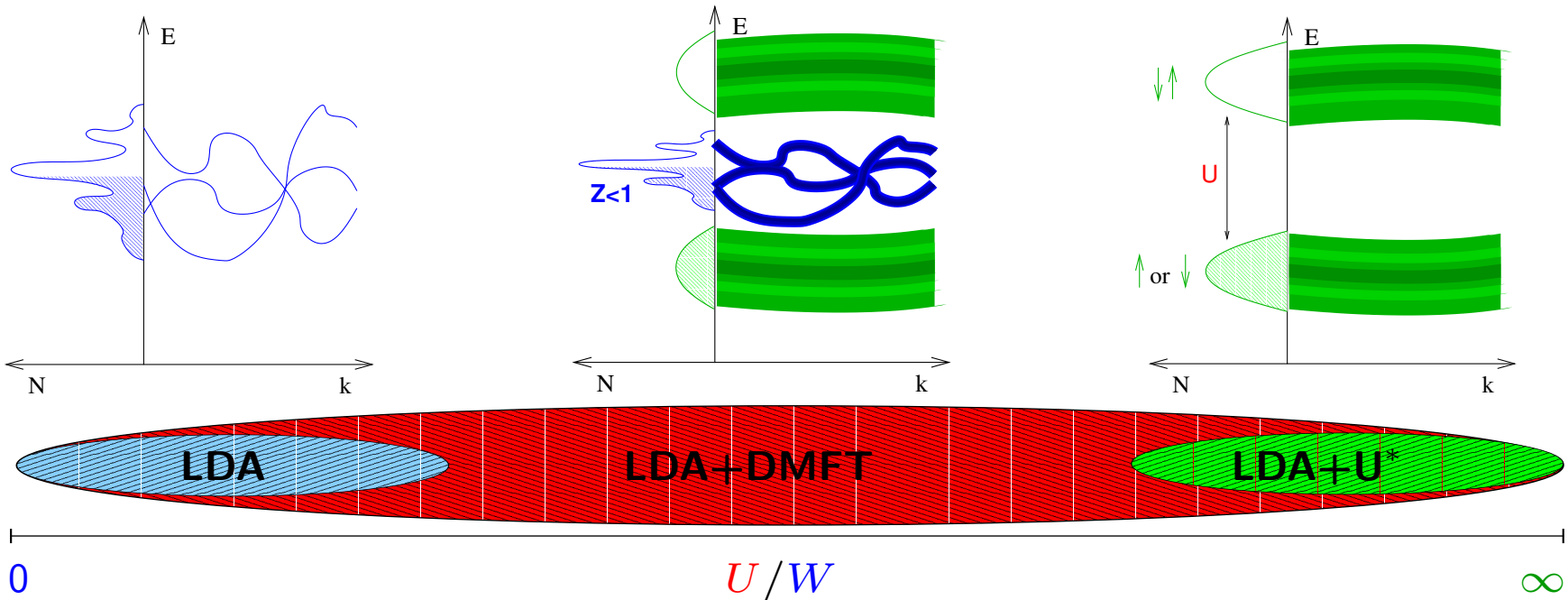
## Dynamical mean field theory

(Metzner, Vollhardt '89; Georges, Kotliar '92)



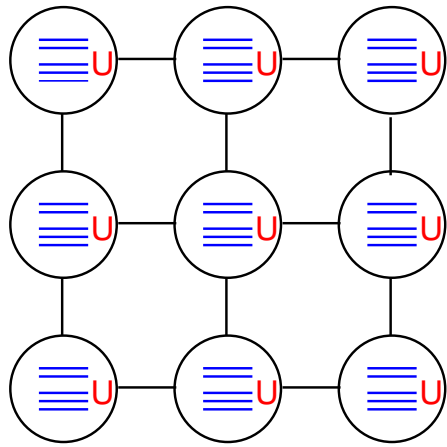
$\Sigma$  all topologically distinct, but **local** diagrams

**Success story:** quasiparticle renormalizations, magnetism, kinks ...



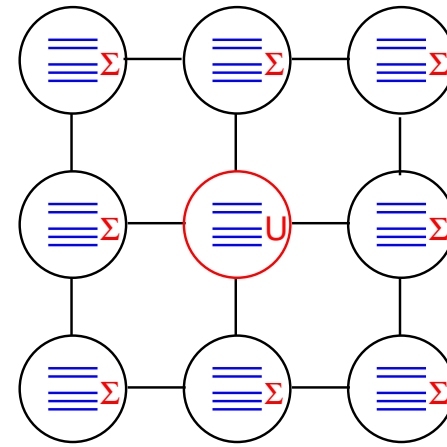
# Motivation

## Dynamical mean field theory



DMFT

(Metzner, Vollhardt'89; Georges, Kotliar'92)



$\Sigma$  all topologically distinct, but **local** diagrams

**Success story:** quasiparticle renormalizations, magnetism, kinks ...

**Not included:**

**non-local correlations**

*p*-, *d*-wave superconductivity, spin Peierls

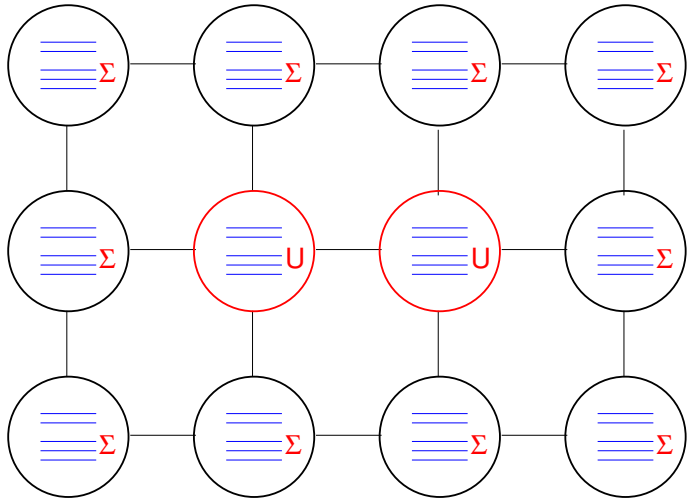
magnons, (quantum) critical behavior ...

**k**-dependence of  $\Sigma$

# beyond DMFT



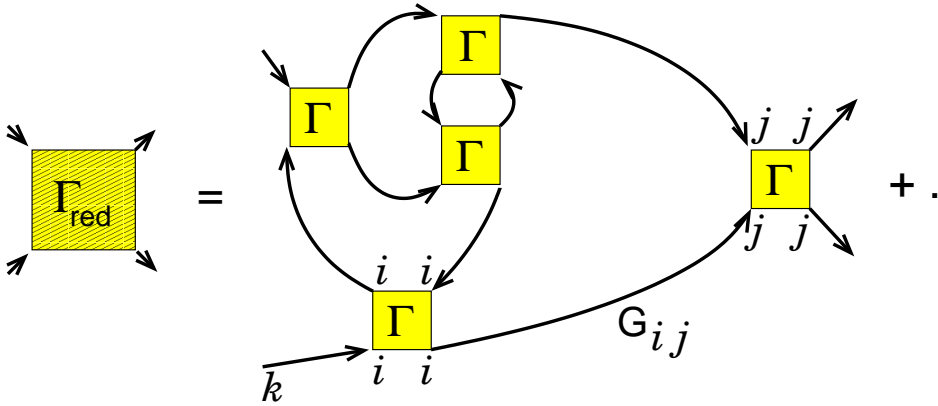
## cluster extensions of DMFT



- non-local **short-range** correlations
- *d/p*-wave superconductivity

Hettler *et al.*'98, Lichtenstein Katsnelson'00,  
Kotliar *et al.*'01, Potthoff'03

## diagrammatic extensions of DMFT



## dynamical vertex approximation

- non-local **long-range** correlations
- (para-)magnons, phase transitions ...

Toschi, Katanin, KH cond-mat/0603100  
cf. Kusunose cond-mat/0602451  
Slezak *et al.* cond-mat/0603421

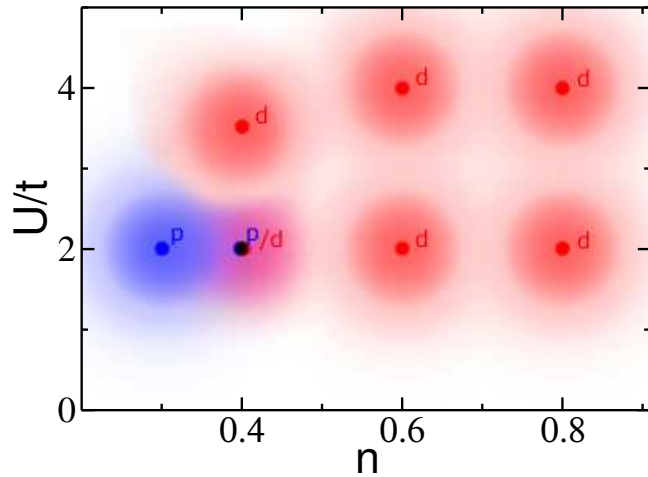


# beyond DMFT

dominant superconducting susceptibility

$t-t'$  2D Hubbard model

Arita, KH PRB'06

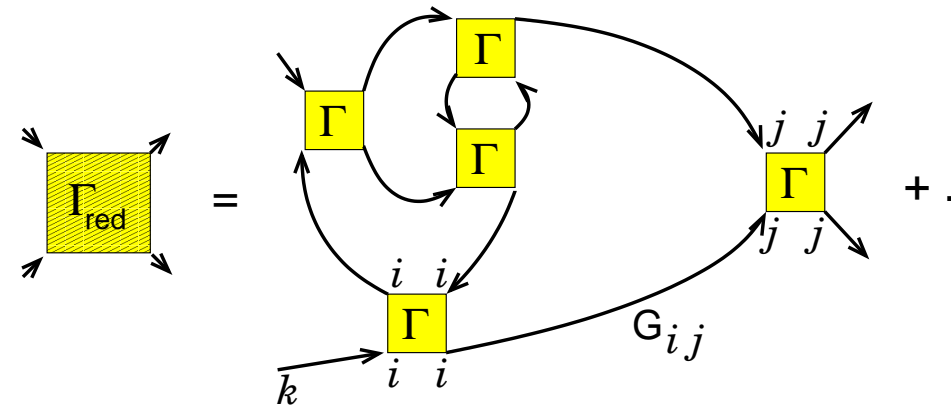


$t' = 0.4t$ ,  $N_c = 4 \times 4 = 16$

- non-local **short-range** correlations
- $d/p$ -wave superconductivity

Hettler *et al.*'98, Lichtenstein Katsnelson'00,  
Kotliar *et al.*'01, Potthoff'03

## diagrammatic extensions of DMFT



## dynamical vertex approximation

- non-local **long-range** correlations
- (para-)magnons, phase transitions ...

Toschi, Katanin, KH cond-mat/0603100  
cf. Kusunose cond-mat/0602451  
Slezak *et al.* cond-mat/0603421

# Dynamical vertex approximation (D $\Gamma$ A)

**DMFT:** all (topological distinct) **local** diagram for  $\Sigma$

**Generalization:** all **local** diagrams for  $n$ -particle fully irreducible vertex  $\Gamma$

# Dynamical vertex approximation (D $\Gamma$ A)

**DMFT:** all (topological distinct) **local** diagram for  $\Sigma$

**Generalization:** all **local** diagrams for  $n$ -particle fully irreducible vertex  $\Gamma$

$n = 1 \rightarrow$  DMFT

# Dynamical vertex approximation (D $\Gamma$ A)

**DMFT:** all (topological distinct) **local** diagrams for  $\Sigma$

**Generalization:** all **local** diagrams for  $n$ -particle fully irreducible vertex  $\Gamma$

$n = 1 \rightarrow$  DMFT

$n = 2 \rightarrow$  D $\Gamma$ A: from **2-particle** irreducible vertex  $\Gamma$   
construct  $\Sigma$  (**local** and **non-local** diagrams)

# Dynamical vertex approximation (D $\Gamma$ A)

DMFT: all (topological distinct) **local** diagram for  $\Sigma$

**Generalization:** all **local** diagrams for  $n$ -particle fully irreducible vertex  $\Gamma$

$n = 1 \rightarrow$  DMFT

$n = 2 \rightarrow$  D $\Gamma$ A: from **2-particle** irreducible vertex  $\Gamma$   
construct  $\Sigma$  (**local** and **non-local** diagrams)

...

$n = \infty \rightarrow$  exact solution

# Dynamical vertex approximation (D $\Gamma$ A)

DMFT: all (topological distinct) **local** diagram for  $\Sigma$

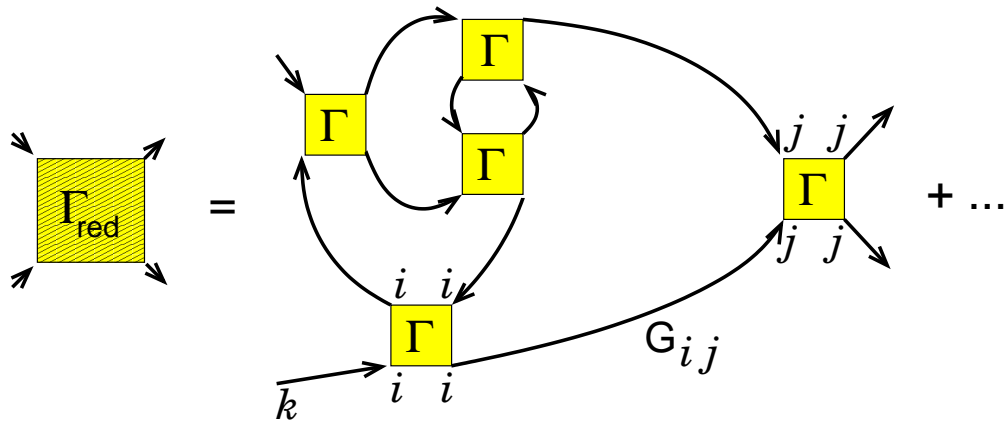
**Generalization:** all **local** diagrams for  $n$ -particle fully irreducible vertex  $\Gamma$

$n = 1 \rightarrow$  DMFT

$n = 2 \rightarrow$  D $\Gamma$ A: from **2-particle** irreducible vertex  $\Gamma$   
construct  $\Sigma$  (**local** and **non-local** diagrams)

...

$n = \infty \rightarrow$  exact solution



**local**  $\Gamma$ , **non-local**  $G$

$\rightarrow$

**non-local** reducible vertex  $\Gamma_{\text{red}}$

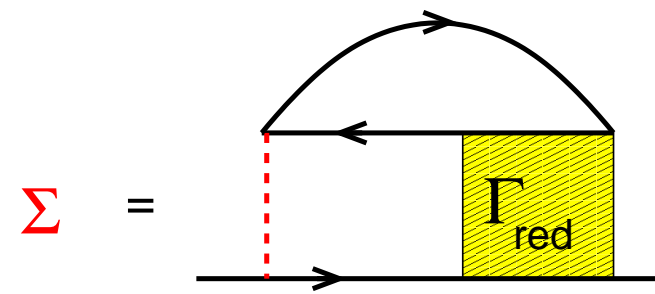
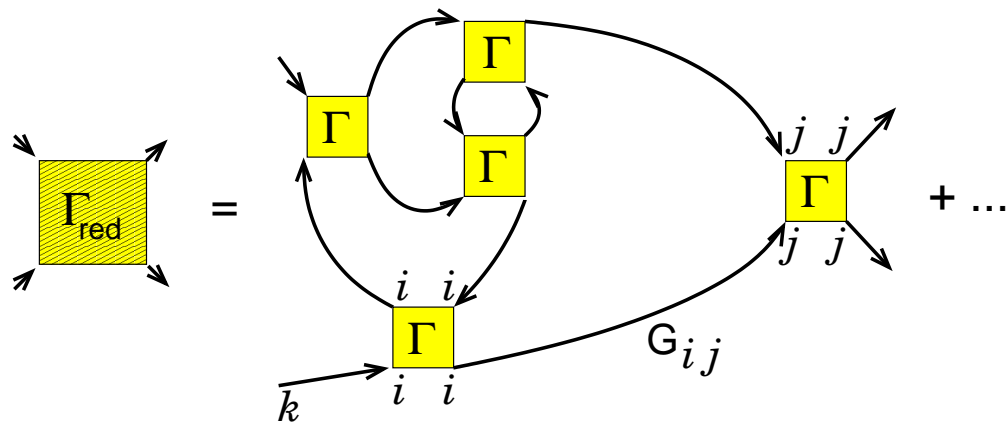
via parquet equations

# Dynamical vertex approximation (D $\Gamma$ A)

DMFT: all (topological distinct) **local** diagram for  $\Sigma$

**Generalization:** all **local** diagrams for  $n$ -particle fully irreducible vertex  $\Gamma$

$n = 1 \rightarrow$  DMFT  
 $n = 2 \rightarrow$  D $\Gamma$ A: from **2-particle** irreducible vertex  $\Gamma$   
 construct  $\Sigma$  (**local** and **non-local** diagrams)  
 ...  
 $n = \infty \rightarrow$  exact solution



**local**  $\Gamma$ , **non-local**  $G$   
 $\rightarrow$   
**non-local** reducible vertex  $\Gamma_{\text{red}}$   
 via parquet equations

$\Gamma_{\text{red}}$   
 $\rightarrow$   
**non-local**  $\Sigma$   
 exact relation (eq. of motion)

# Dynamical vertex approximation (D $\Gamma$ A)

DMFT: all (topological distinct) **local** diagram for  $\Sigma$

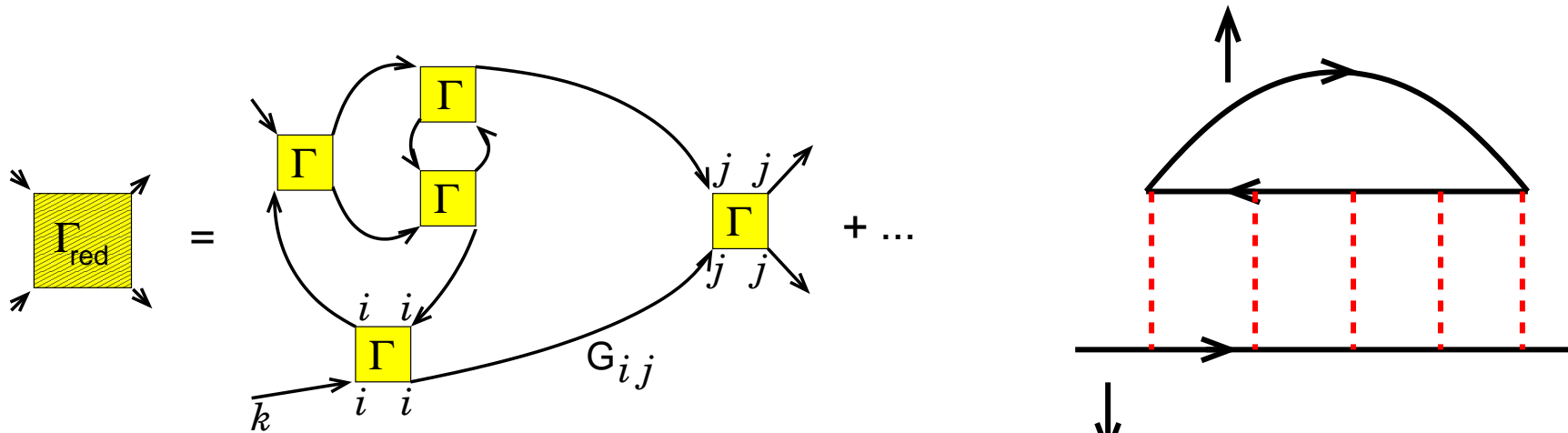
**Generalization:** all **local** diagrams for  $n$ -particle fully irreducible vertex  $\Gamma$

$n = 1 \rightarrow$  DMFT

$n = 2 \rightarrow$  D $\Gamma$ A: from **2-particle** irreducible vertex  $\Gamma$   
construct  $\Sigma$  (**local** and **non-local** diagrams)

...

$n = \infty \rightarrow$  exact solution



**local**  $\Gamma$ , **non-local**  $G$

$\rightarrow$

**non-local** reducible vertex  $\Gamma_{\text{red}}$

via parquet equations

Moriya Edwards-Hertz



# Dynamical vertex approximation (D $\Gamma$ A)

DMFT: all (topological distinct) **local** diagram for  $\Sigma$

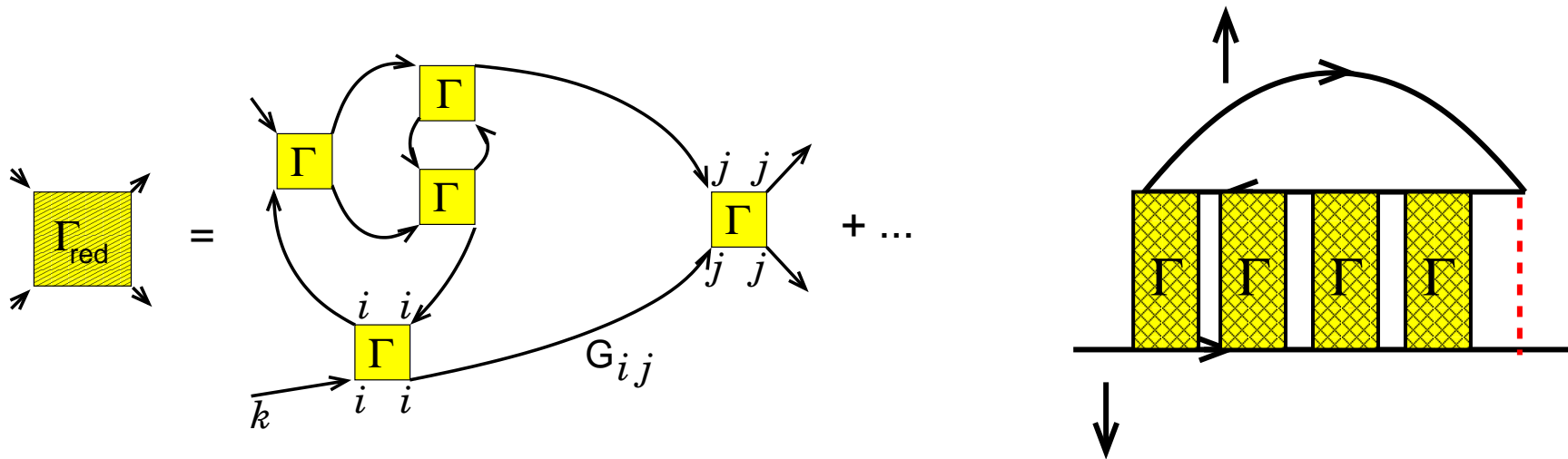
**Generalization:** all **local** diagrams for  $n$ -particle fully irreducible vertex  $\Gamma$

$n = 1 \rightarrow$  DMFT

$n = 2 \rightarrow$  D $\Gamma$ A: from **2-particle** irreducible vertex  $\Gamma$   
construct  $\Sigma$  (**local** and **non-local** diagrams)

...

$n = \infty \rightarrow$  exact solution



**local**  $\Gamma$ , **non-local**  $G$

$\rightarrow$

**non-local** reducible vertex  $\Gamma_{\text{red}}$

via parquet equations

# Dynamical vertex approximation (D $\Gamma$ A)

DMFT: all (topological distinct) **local** diagram for  $\Sigma$

**Generalization:** all **local** diagrams for  $n$ -particle fully irreducible vertex  $\Gamma$

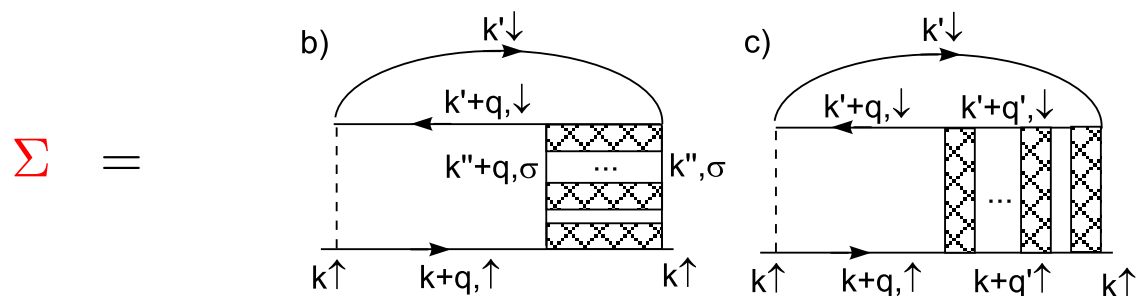
$n = 1 \rightarrow$  DMFT

$n = 2 \rightarrow$  D $\Gamma$ A: from **2-particle** irreducible vertex  $\Gamma$   
construct  $\Sigma$  (**local** and **non-local** diagrams)

...

$n = \infty \rightarrow$  exact solution

**First step: restriction to ladder diagrams**



lines: **non-local**  $G$

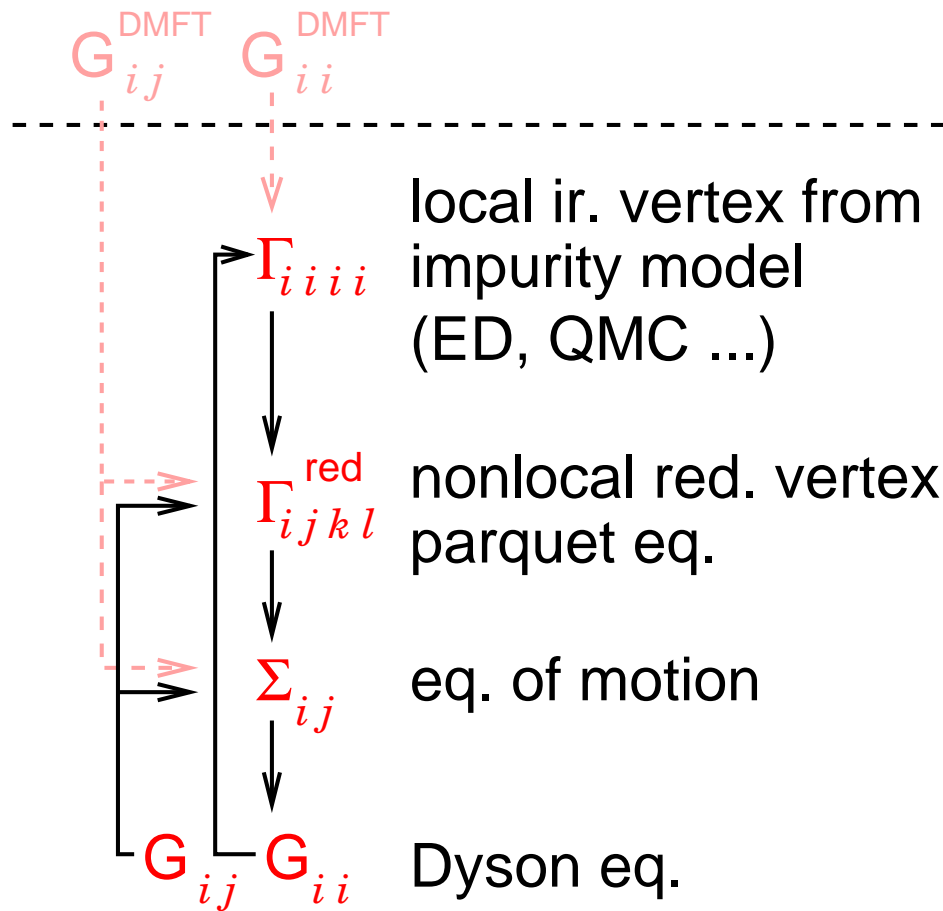
crosshatched: local irreducible vertex in spin/charge channels

$$\Gamma_{S,C}(\nu, \nu', \omega) = \chi_{0,loc}^{-1} - \chi_{S,C}^{-1}$$

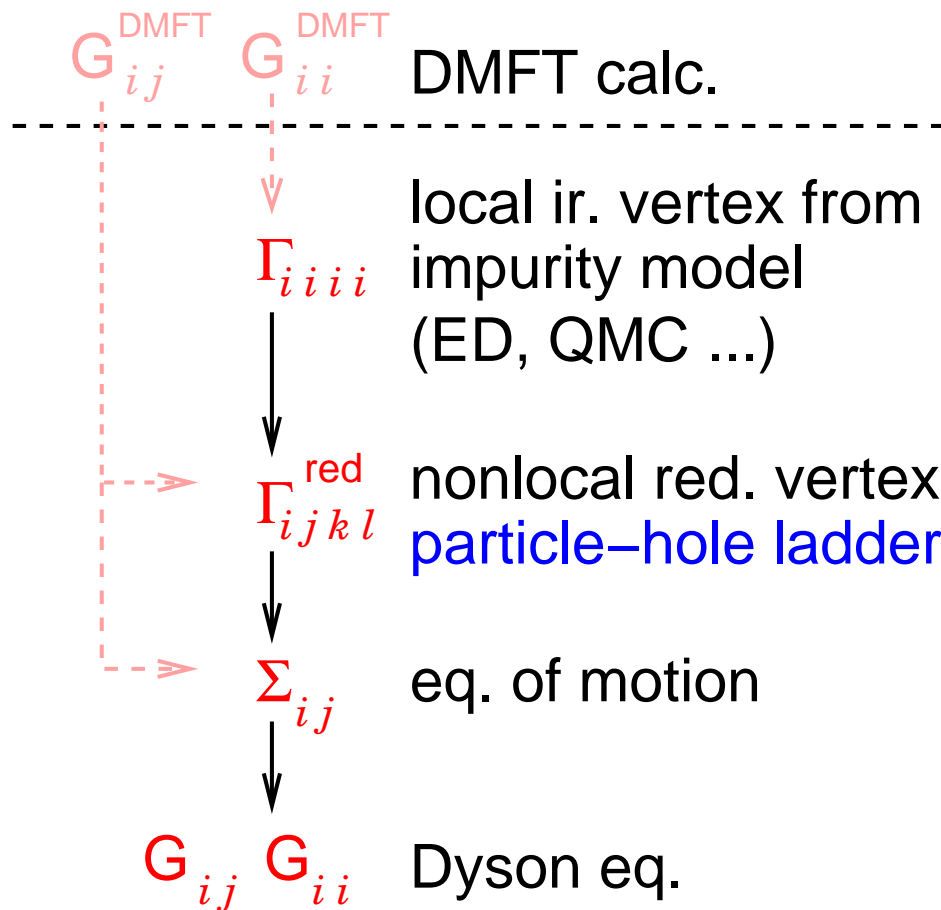
magnons, spin-fluctuations at (A)FM phase transition

$G_{ij}$  from DMFT

# D $\Gamma$ A algorithm



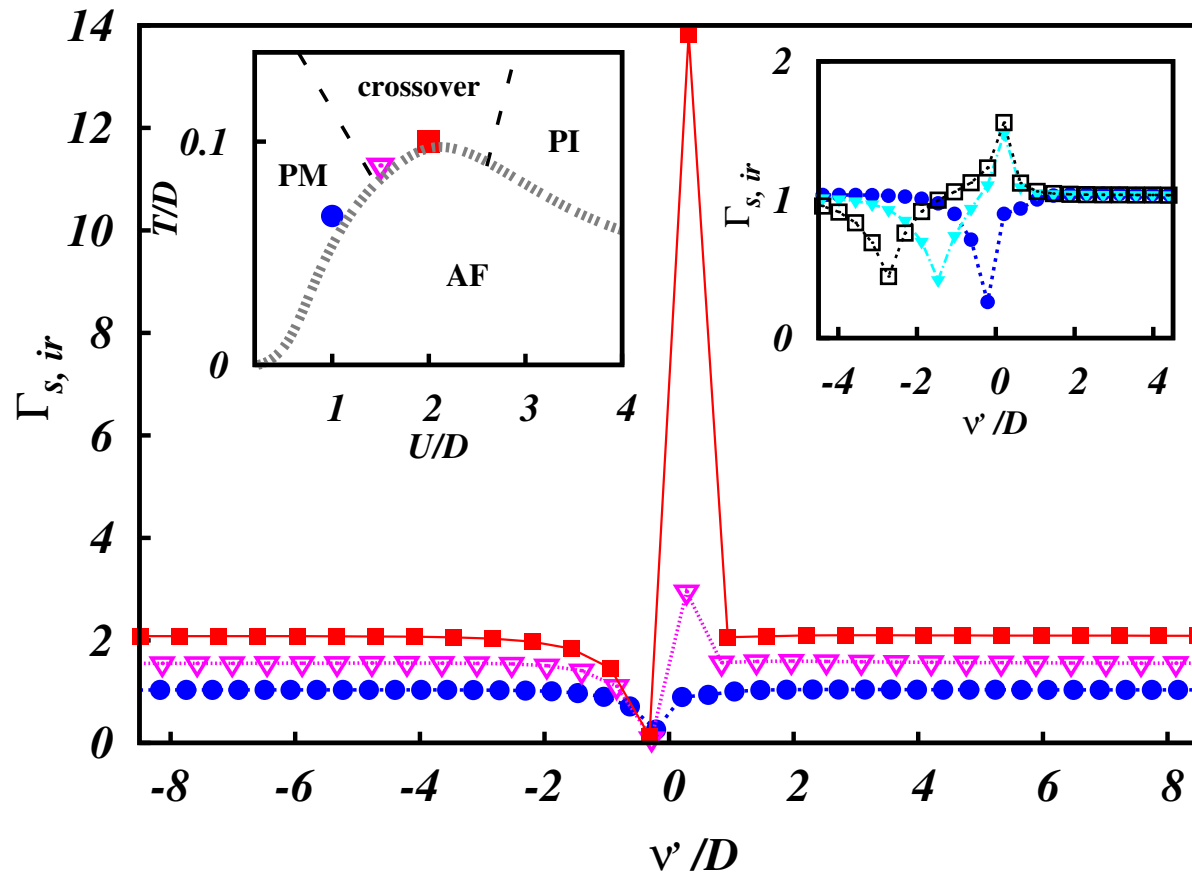
# D $\Gamma$ A algorithm (restriction to ph ladders)



# Results: 3D Hubbard model

$$H = -t \sum_{\langle i,j \rangle \sigma} c_{i\sigma}^\dagger c_{j\sigma} + U \sum_i n_{i\uparrow} n_{i\downarrow}$$

cubic lattice, exact diagonalization as impurity solver

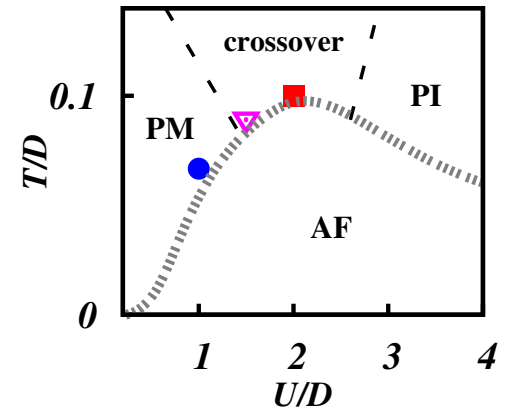


eff. bandwidth  $\equiv 2D$   
 $\omega = 0$   
 $\nu = \pi T$

$\Gamma_{s,ir}(\nu, \nu', \omega)$  strongly frequency dependent

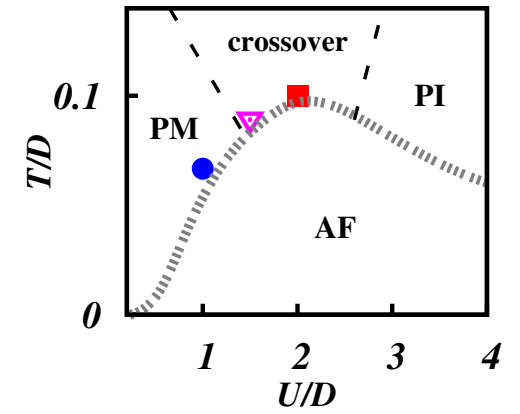
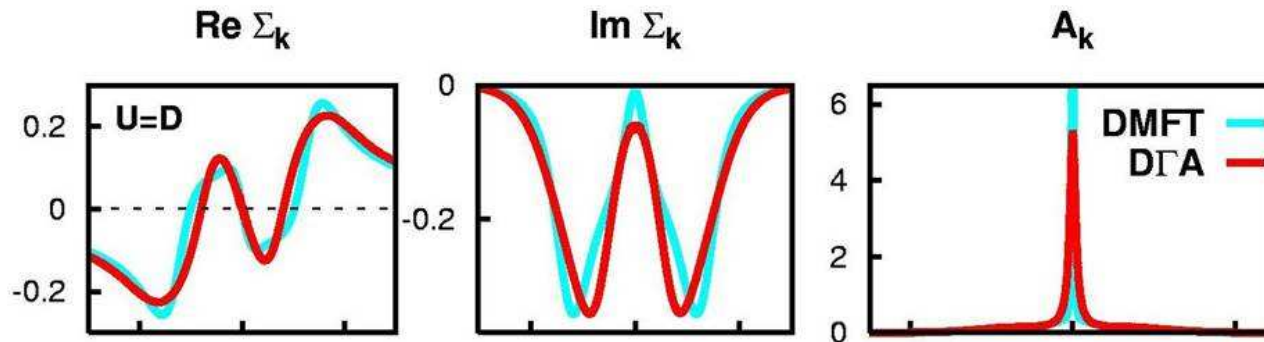
## Results: 3D Hubbard model

$\Sigma$  and  $A$  for  $\mathbf{k} = (\pi/2, \pi/2, \pi/2)$  (on Fermi surface)



# Results: 3D Hubbard model

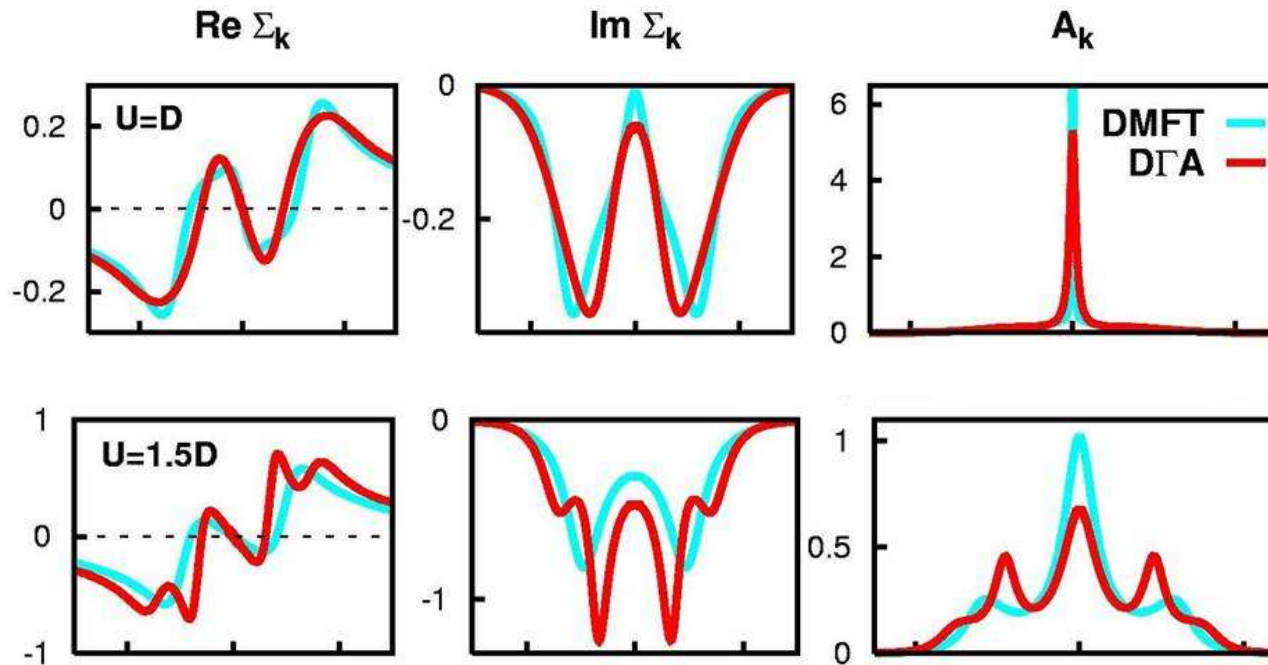
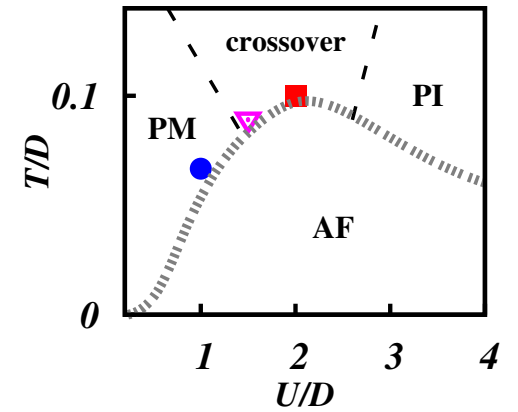
$\Sigma$  and  $A$  for  $\mathbf{k} = (\pi/2, \pi/2, \pi/2)$  (on Fermi surface)



← weak damping of QP peak

# Results: 3D Hubbard model

$\Sigma$  and  $A$  for  $\mathbf{k} = (\pi/2, \pi/2, \pi/2)$  (on Fermi surface)



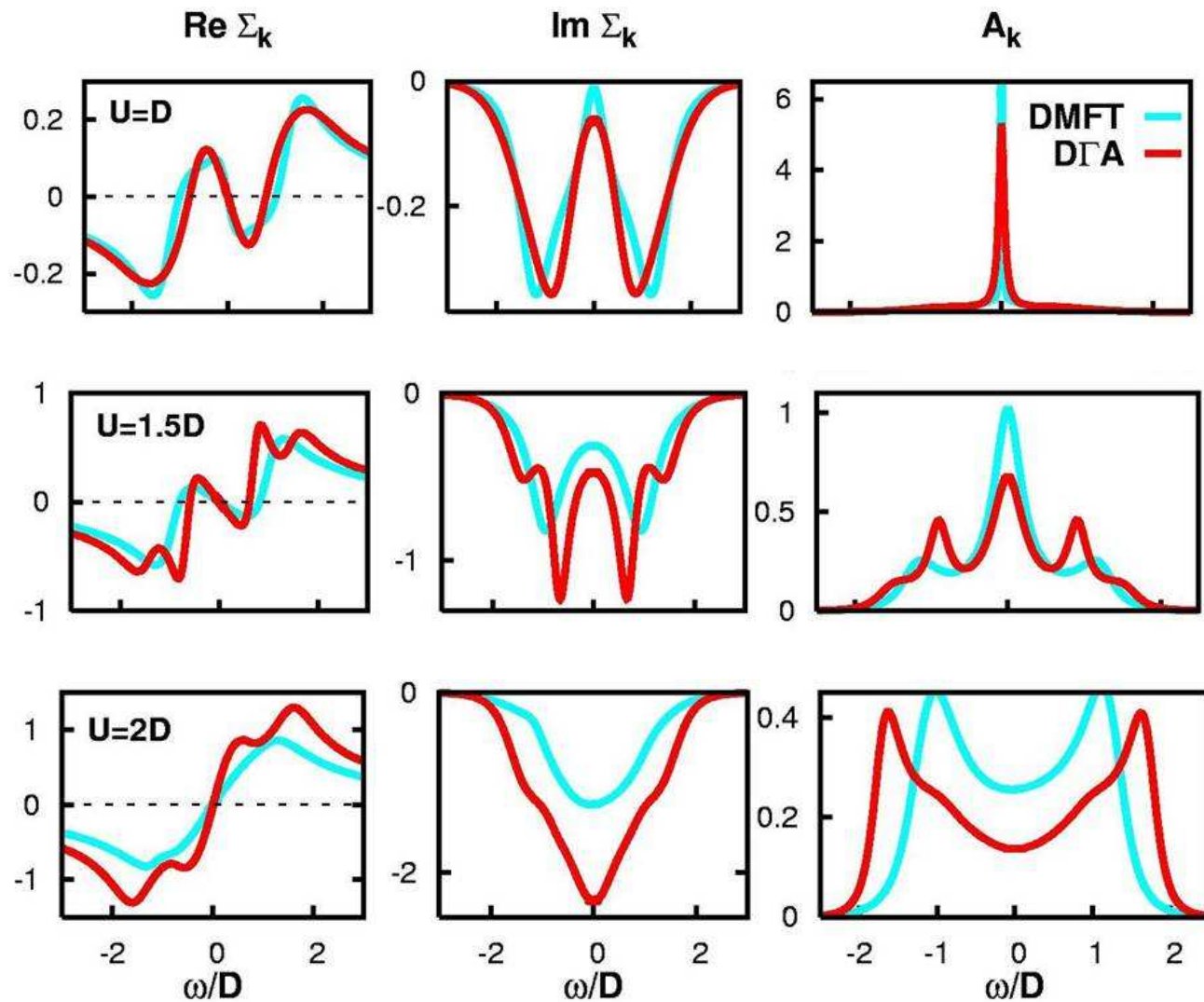
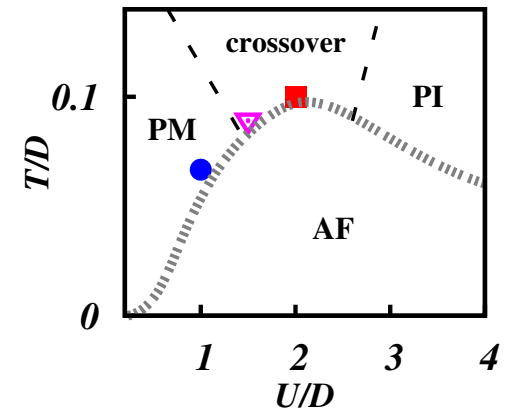
← weak damping of QP peak

← QP-damping strongly enhanced



# Results: 3D Hubbard model

$\Sigma$  and  $A$  for  $\mathbf{k} = (\pi/2, \pi/2, \pi/2)$  (on Fermi surface)

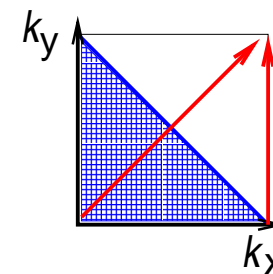


← weak damping of QP peak

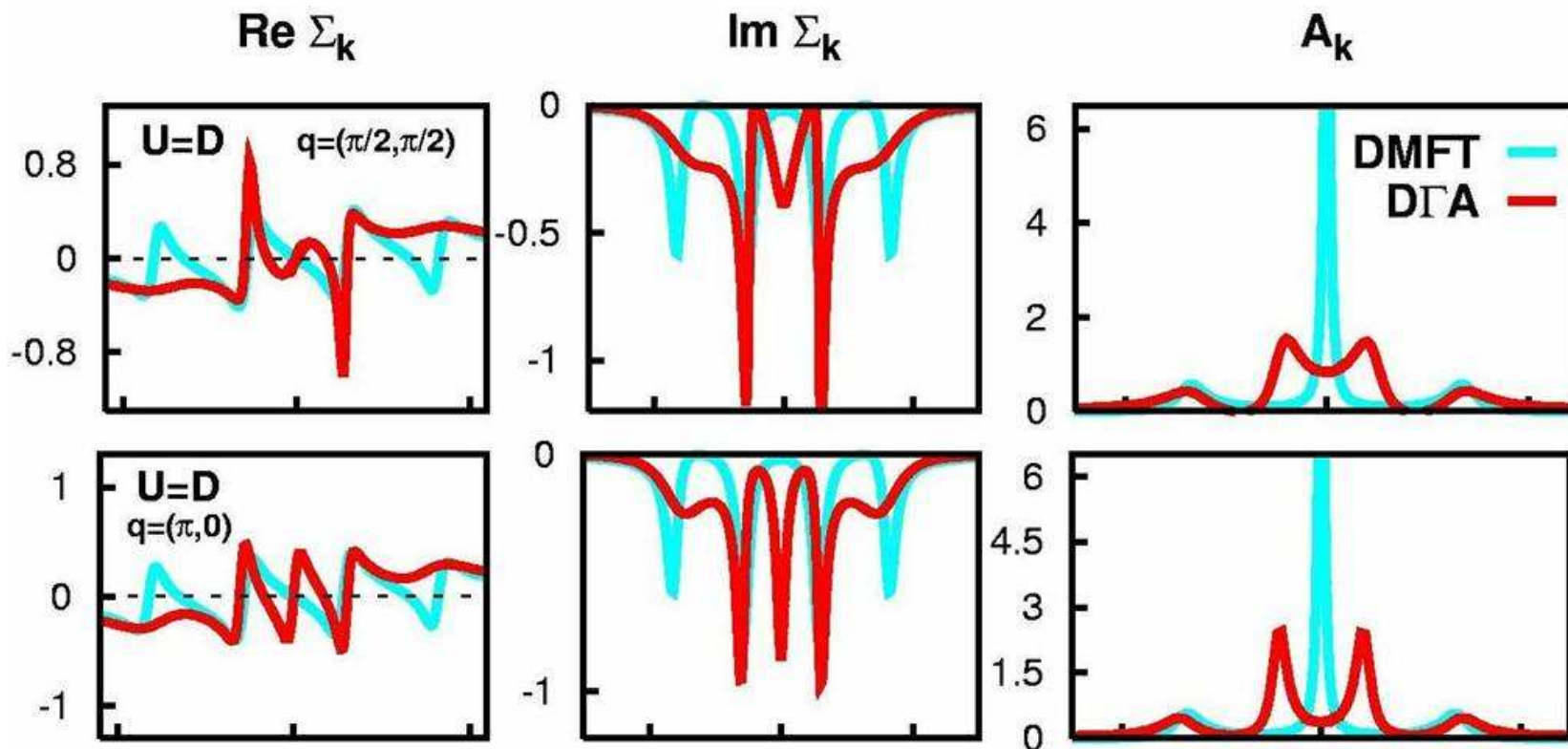
← QP-damping strongly enhanced

← more insulating

# Results: 2D Hubbard model (half-filling)



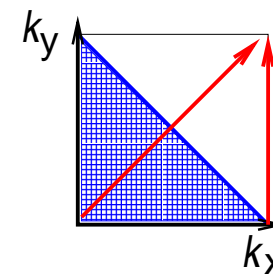
nodal  
 $\mathbf{k} = (\frac{\pi}{2}, \frac{\pi}{2})$



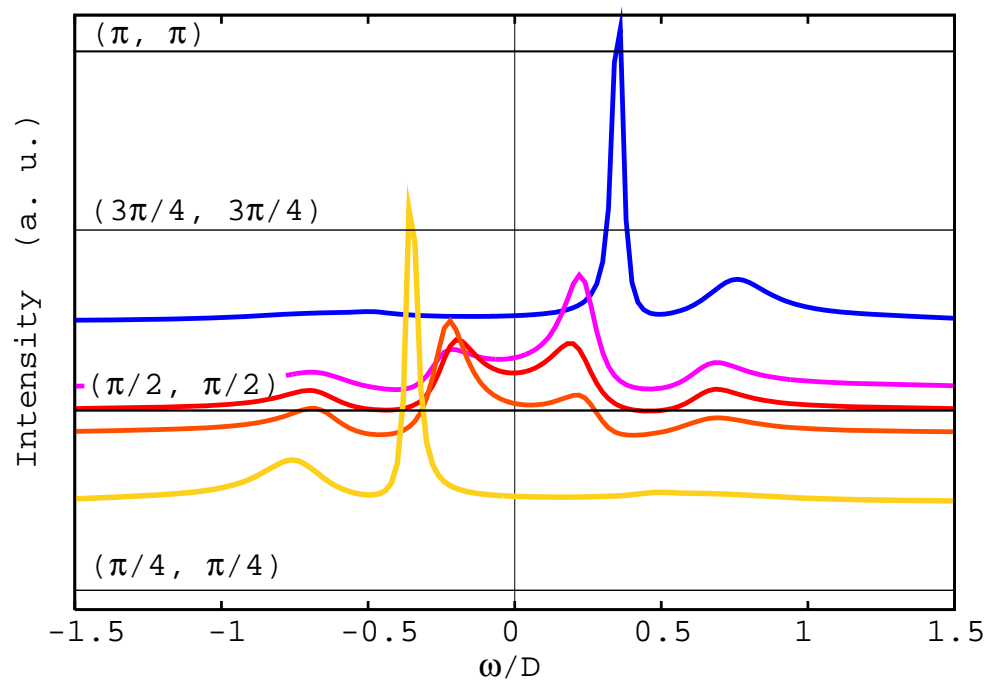
anti-nodal  
 $\mathbf{k} = (\pi, 0)$

anisotropic pseudogap

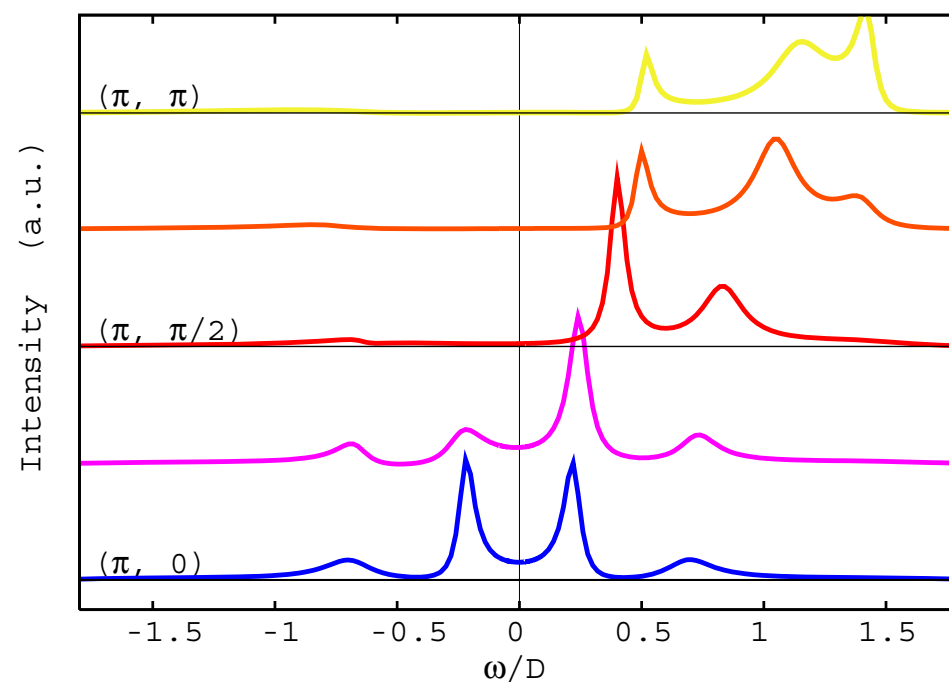
# Results: 2D Hubbard model (half-filling)



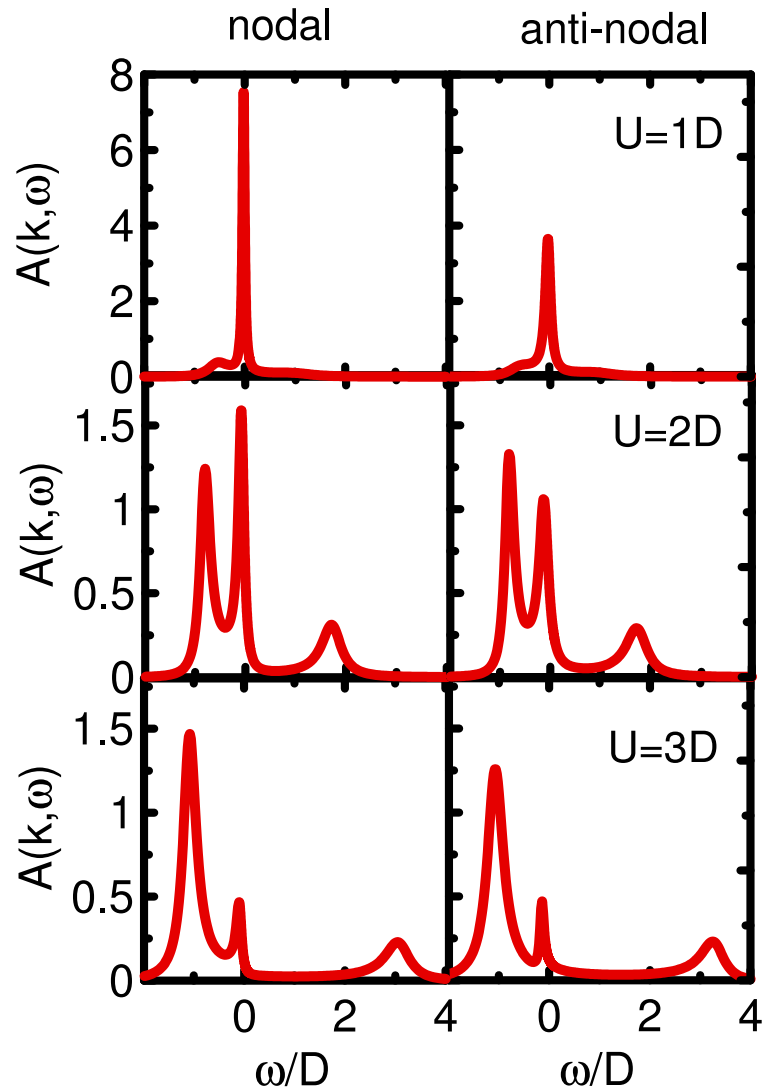
nodal



antinodal



# Results: 2D Hubbard model (off half-filling)



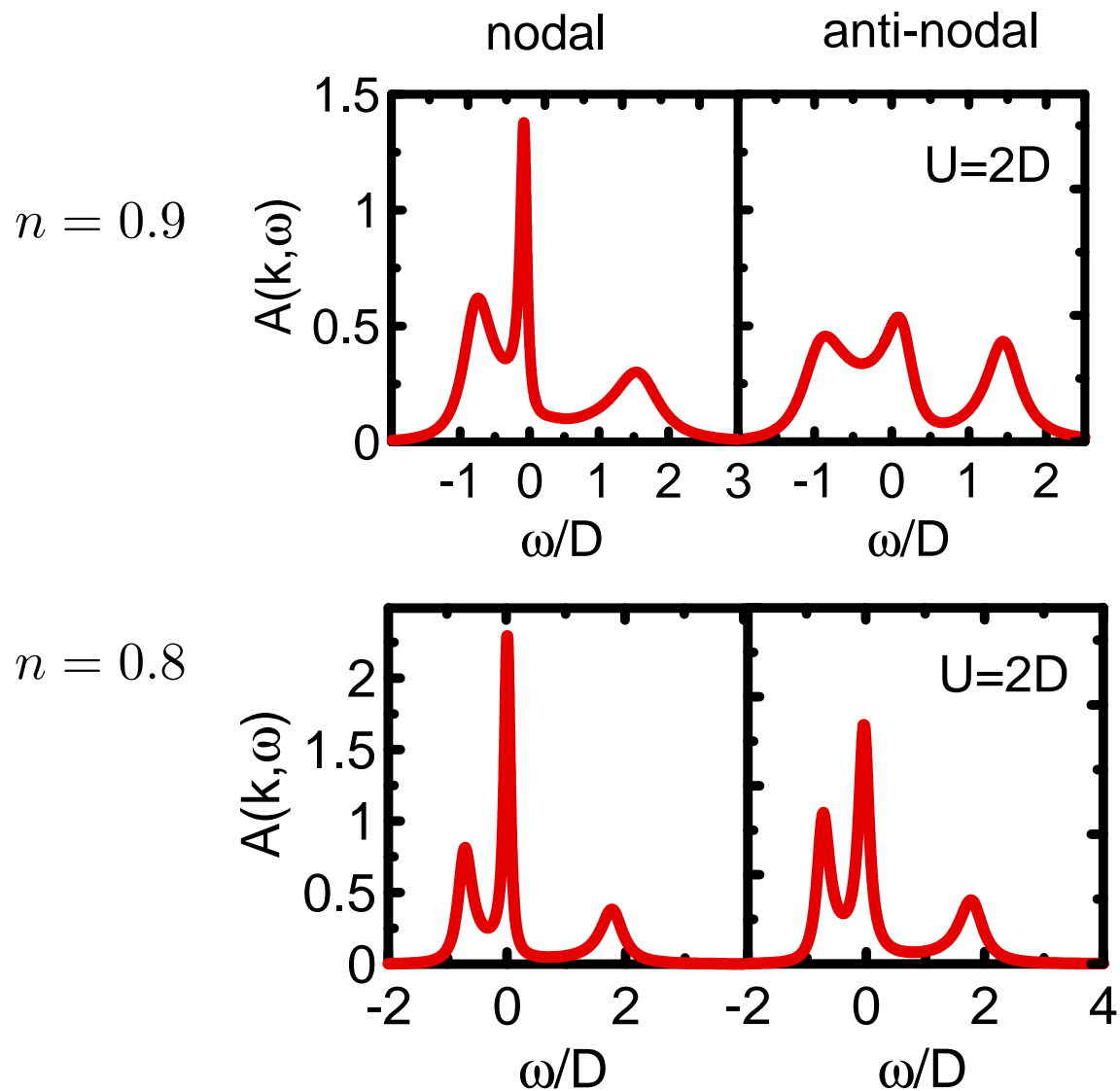
$$t'/t = 0.3$$

$$n = 0.8$$

$$\beta = 100/D$$

less anisotropic  
at strong coupling

# Results: 2D Hubbard model (off half-filling)



less anisotropic  
at larger doping

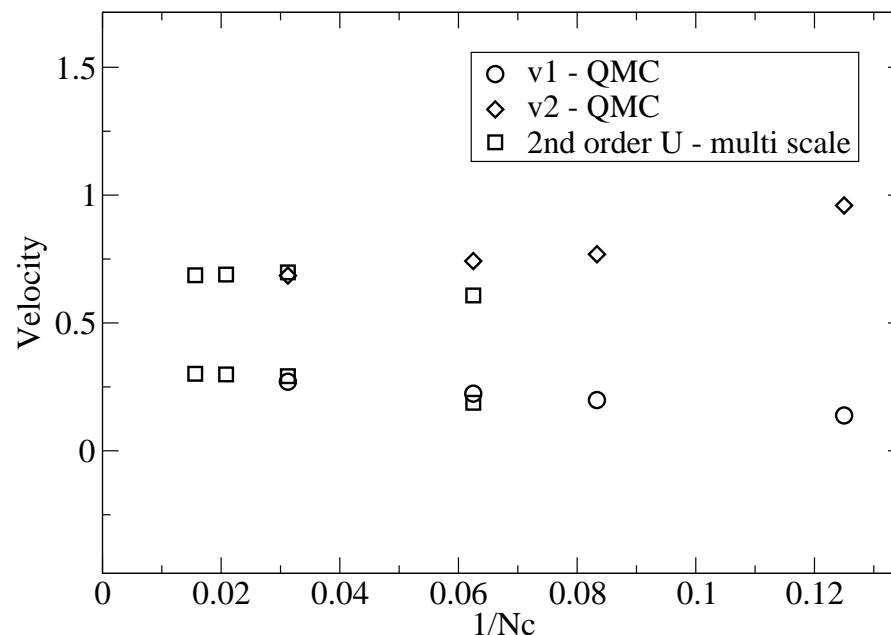
## Spin-charge separation

$$U = W = 1$$

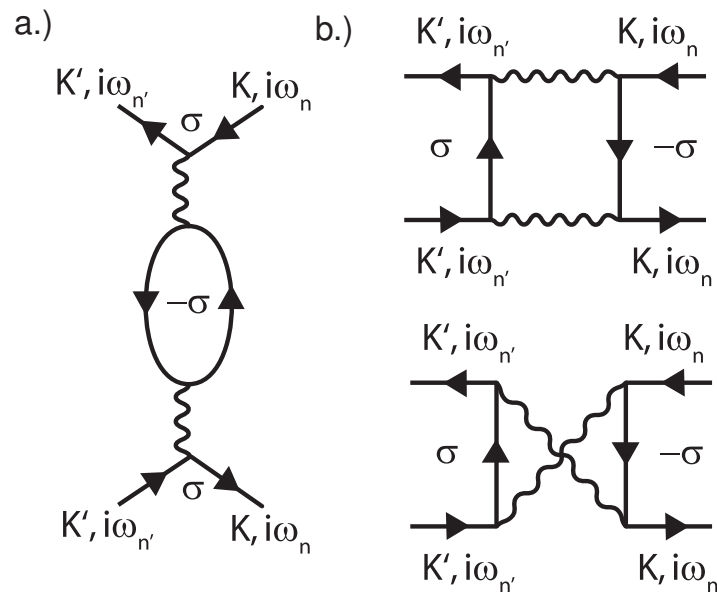
$$k = \pi/2$$

$$\beta = 31$$

$$n = 0.7$$



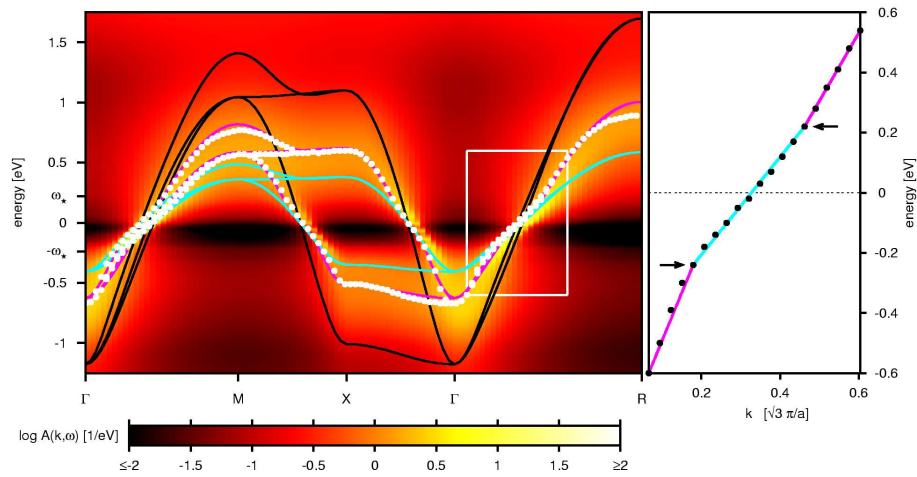
Here, only **2nd order diagrams** for vertex  
 ( $q = 0, \omega = 0$ )  
 but **8-site DCA** for short-range  $\Sigma$



## 2) Kinks — direct consequence of strong correlations

Kinks in  $\text{SrVO}_3$

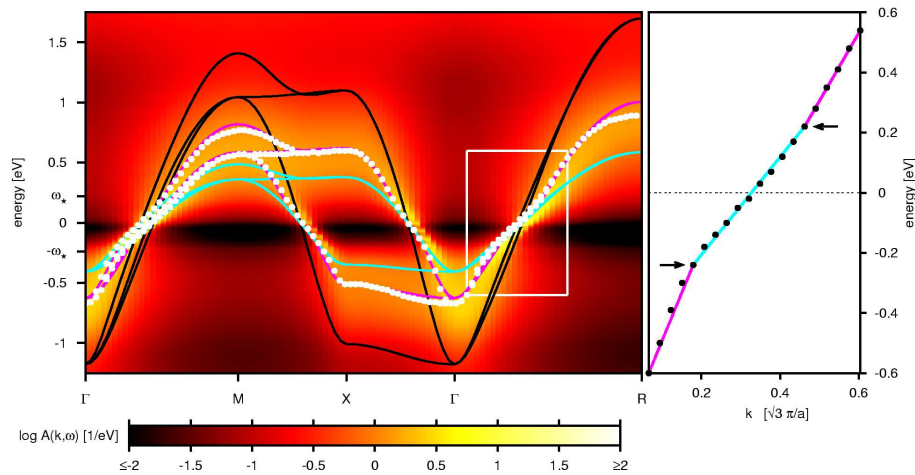
Nekrasov et al. PRB'06



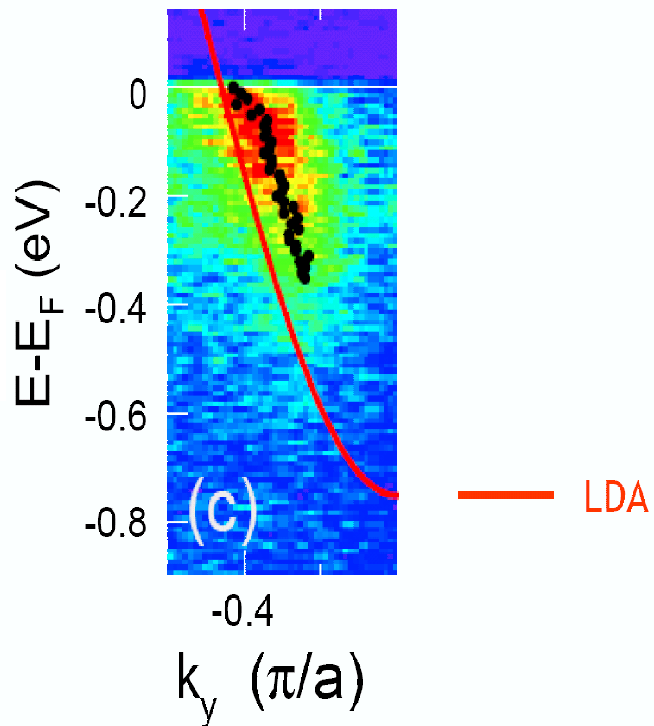
## 2) Kinks — direct consequence of strong correlations

Kinks in  $\text{SrVO}_3$

Nekrasov et al. PRB'06



experimentally observed Fujimori et al.'06

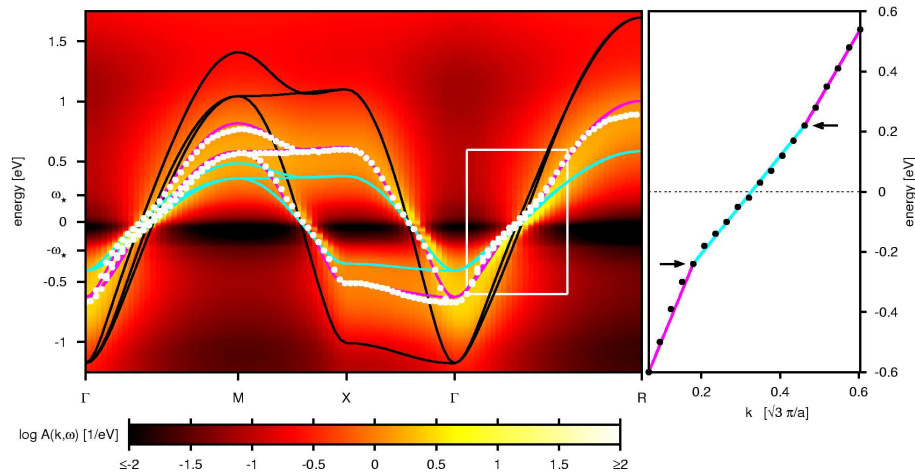




## 2) Kinks — direct consequence of strong correlations

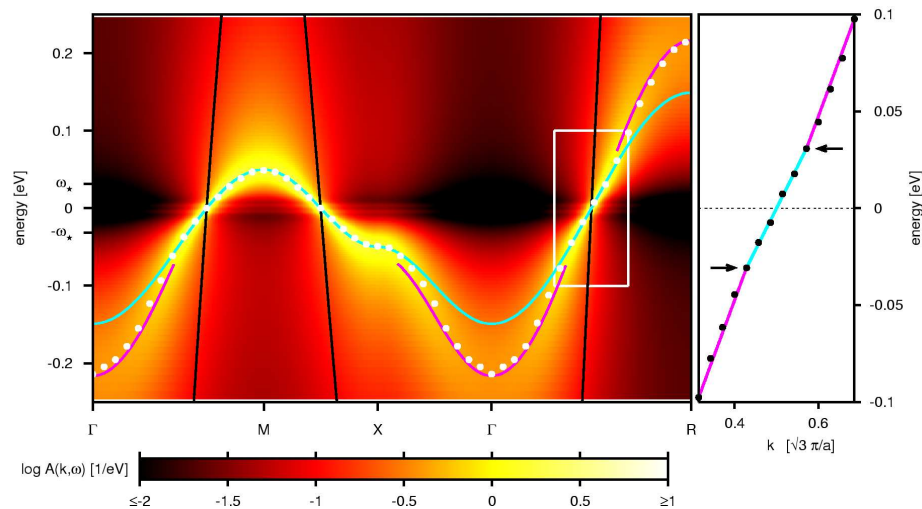
Kinks in  $\text{SrVO}_3$

Nekrasov et al. PRB'06



Kinks in the 3D Hubbard model

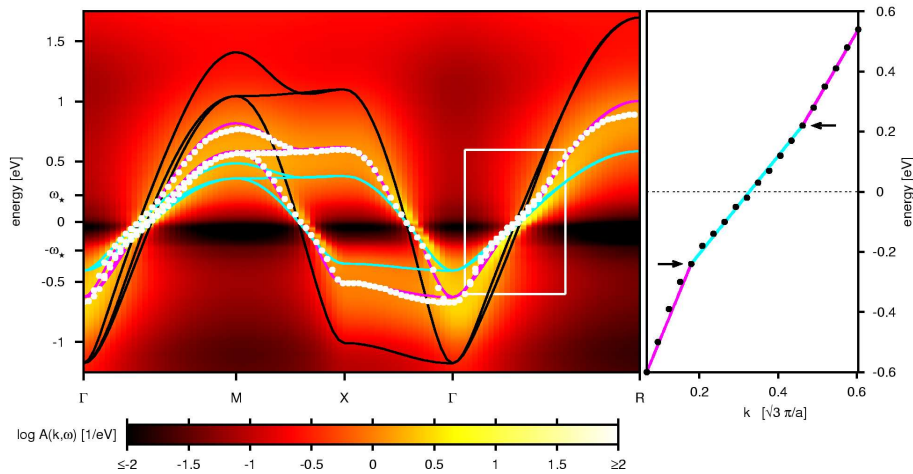
Byczuk, Kollar, KH, Yang, Nekrasov,  
Pruschke, Vollhardt Nature Phys.'07



## 2) Kinks — direct consequence of strong correlations

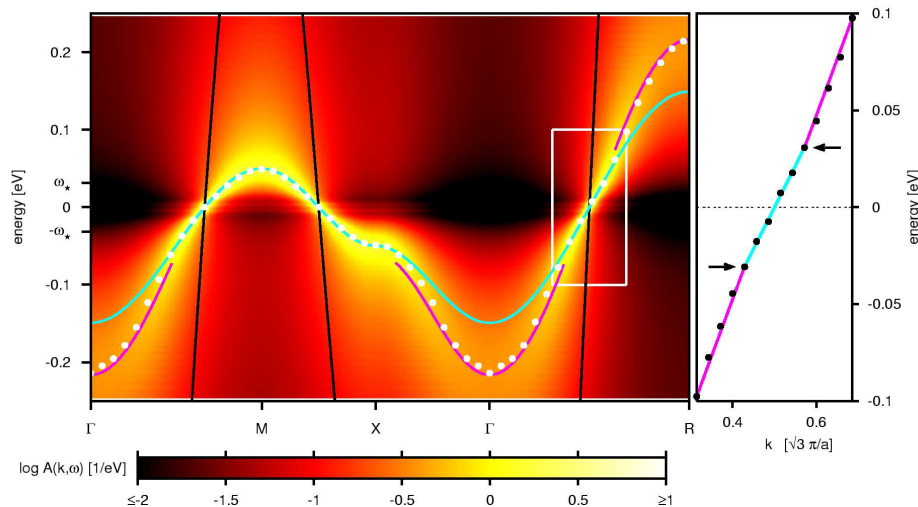
Kinks in  $\text{SrVO}_3$

Nekrasov et al. PRB'06

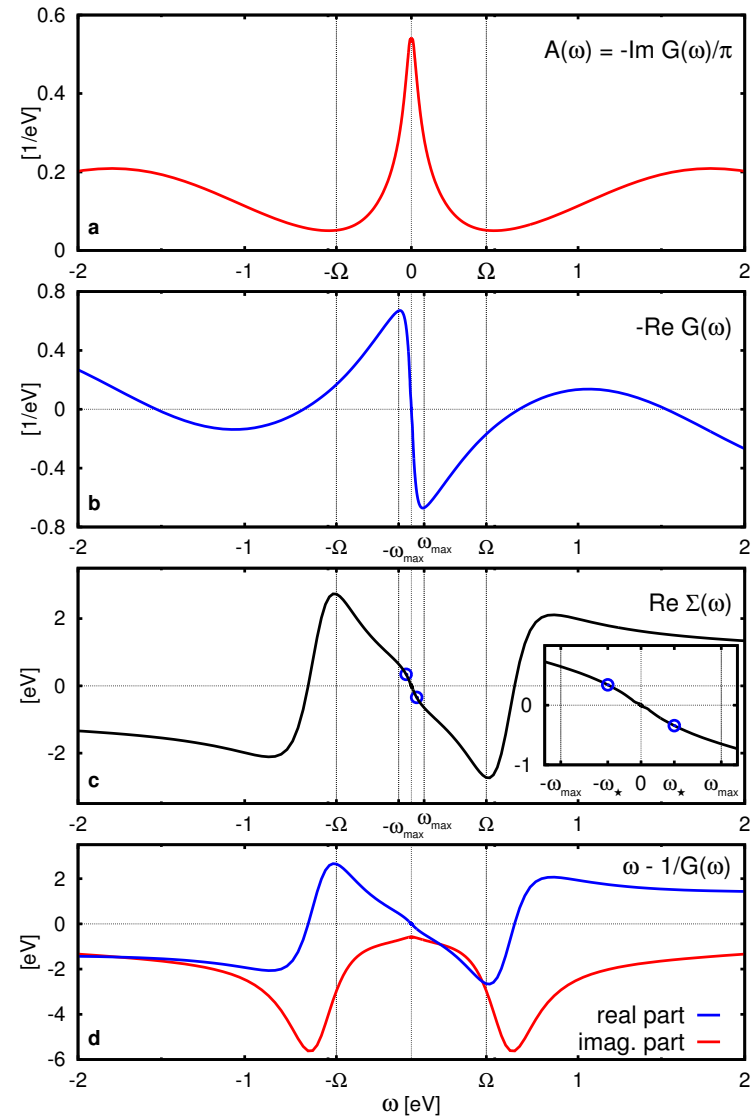


Kinks in the 3D Hubbard model

Byczuk, Kollar, KH, Yang, Nekrasov, Pruschke, Vollhardt Nature Phys.'07



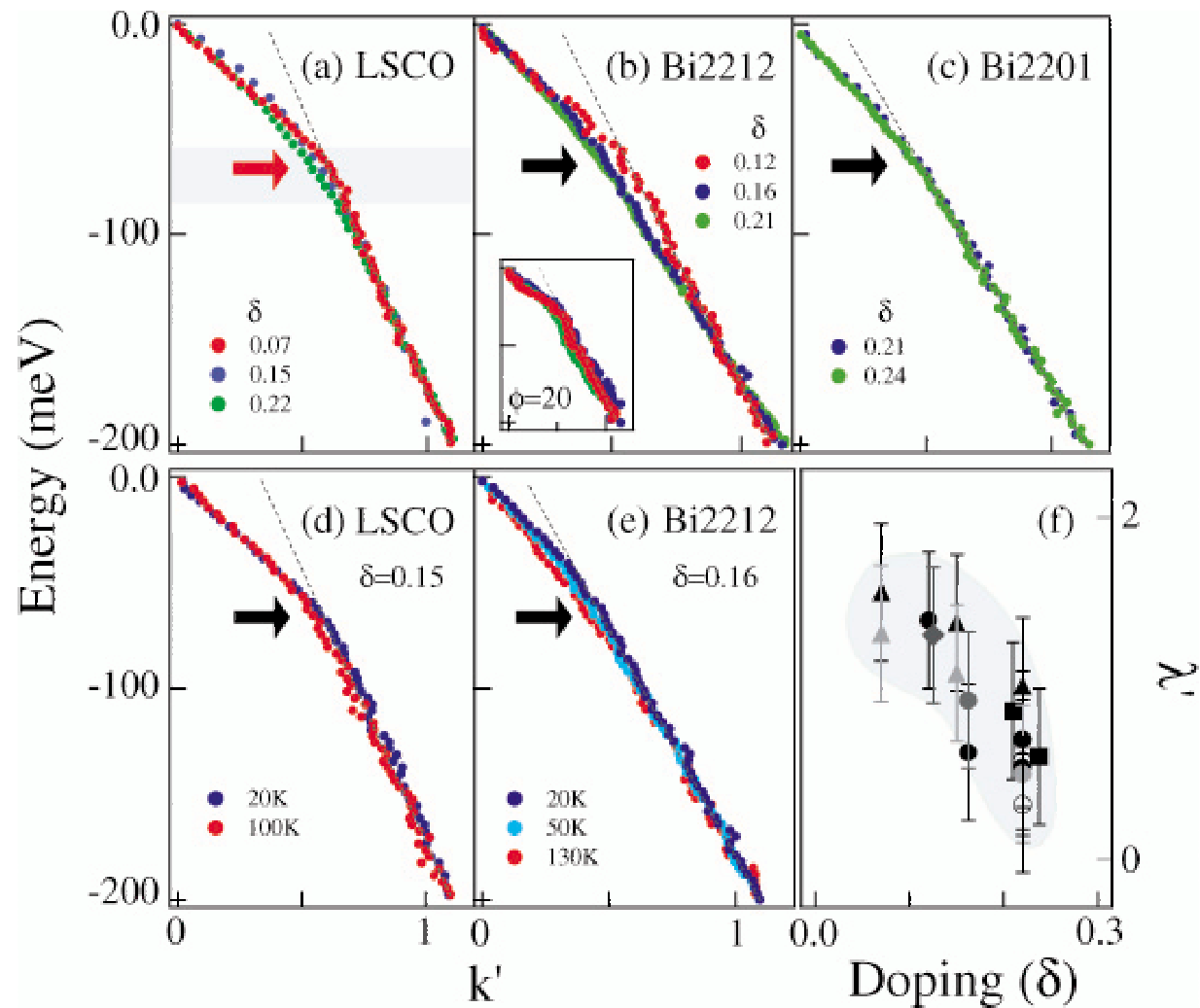
Kinks follow from 3-peak-structure



$$\Sigma(\omega) = \omega + \mu - 1/G(\omega) - \Delta(G(\omega))$$

Fermi-liquid regime:  $E_{\mathbf{k}} = Z_{\text{FL}} \epsilon_{\mathbf{k}}$  for  $|E_{\mathbf{k}}| < \omega_*$   
 Beyond FL regime:  $E_{\mathbf{k}} = Z_{\text{CP}} \epsilon_{\mathbf{k}} \pm c$  for  $|E_{\mathbf{k}}| > \omega_*$

# ARPES: low-energy kinks in cuprates

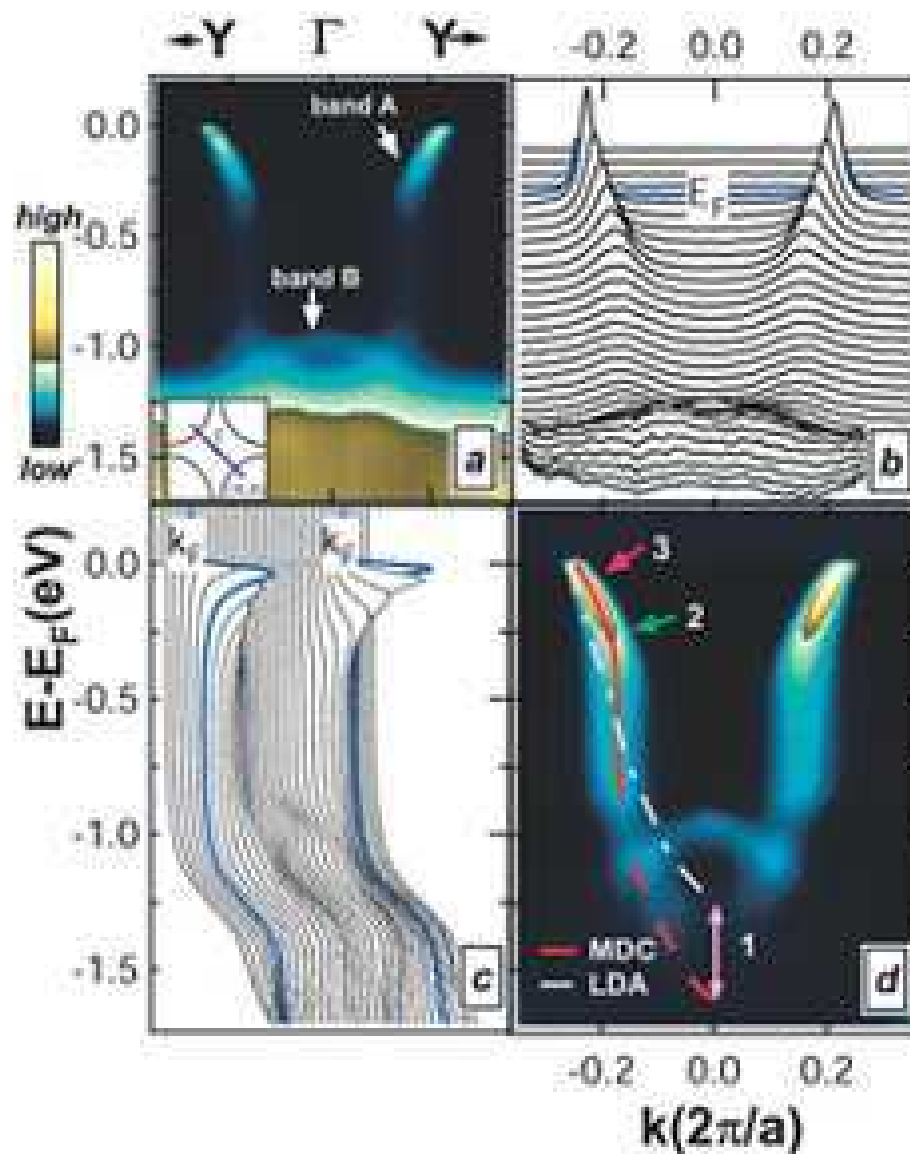


Lanzara *et al.*'01

energy range  $\sim 70$  meV

# ARPES: high-energy kinks in cuprates

Bi2201 at  $T = 30\text{K}$  ( $> T_c$ )



Meevasana *et al.* cond-mat/0612541

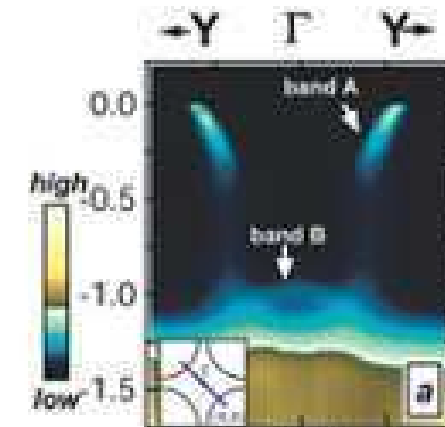
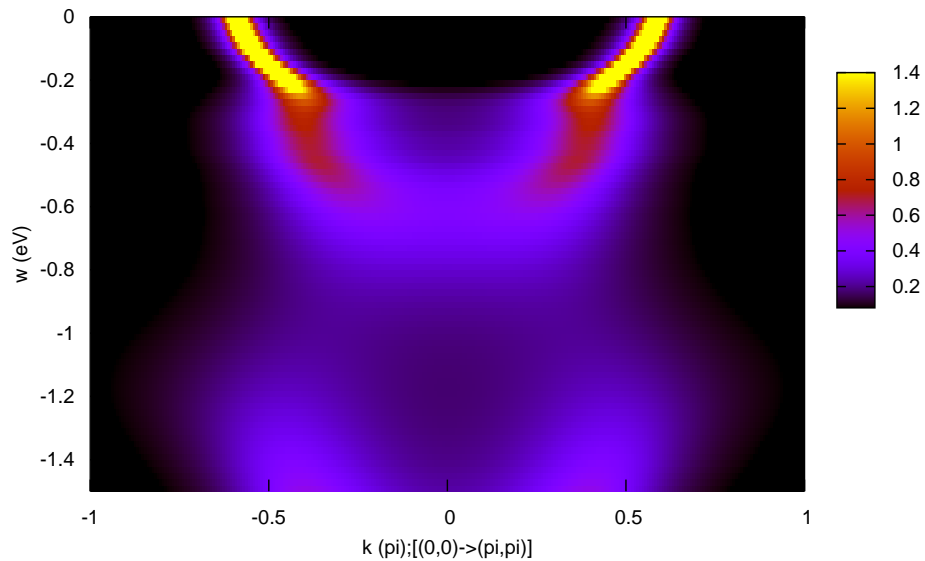
energy range  $\sim 0.3$  eV

# Connection to high-energy kinks

Yang, Held'07

## 2D Hubbard model; DMFT(QMC)

$n = 0.85$ ,  $U = 3$ ,  $t = 0.435$ ,  $t' = -0.1$ ,  $t'' = 0.038$ ,  $T = 1/40$  (eV)

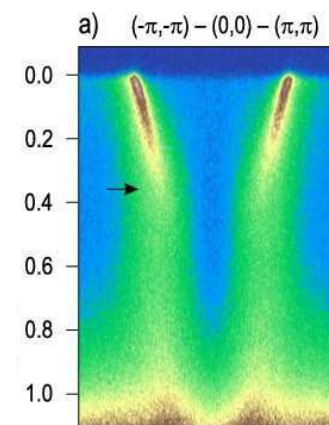


Meevasana *et al.*'06

cf. Macridin *et al.*'07

cf. Byczuk, Kollar, Vollhardt'07

Kink position correct  
but two features kink+waterfall



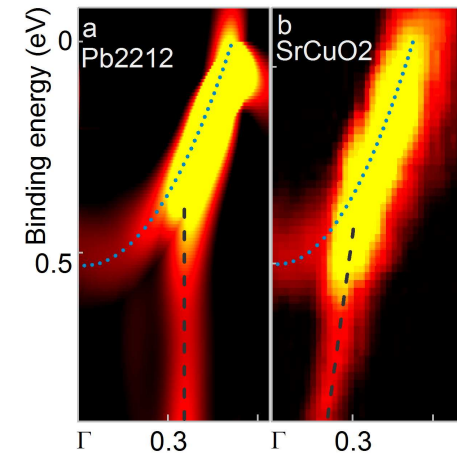
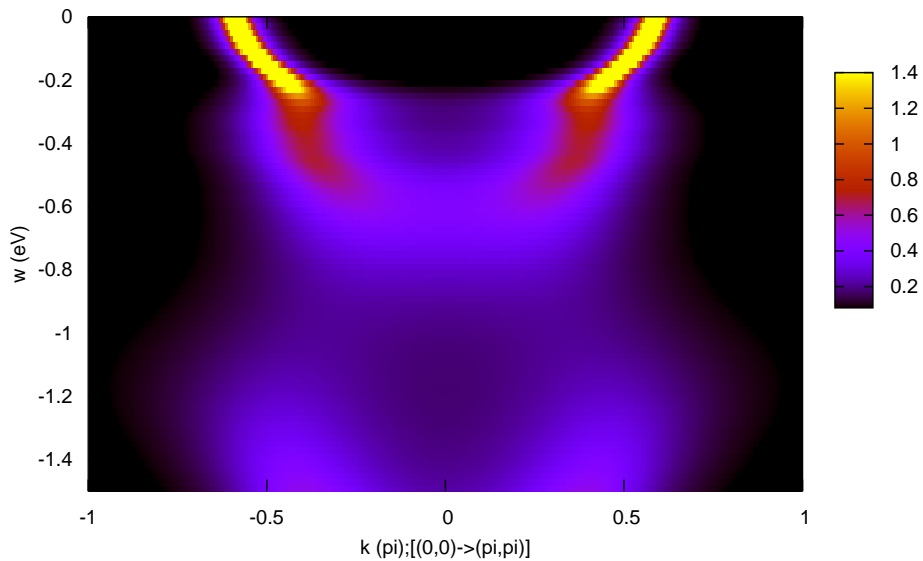
Inosov *et al.*'07

# Connection to high-energy kinks

Yang, Held'07

## 2D Hubbard model; DMFT(QMC)

$n = 0.85$ ,  $U = 3$ ,  $t = 0.435$ ,  $t' = -0.1$ ,  $t'' = 0.038$ ,  $T = 1/40$  (eV)



Graf *et al.*'06

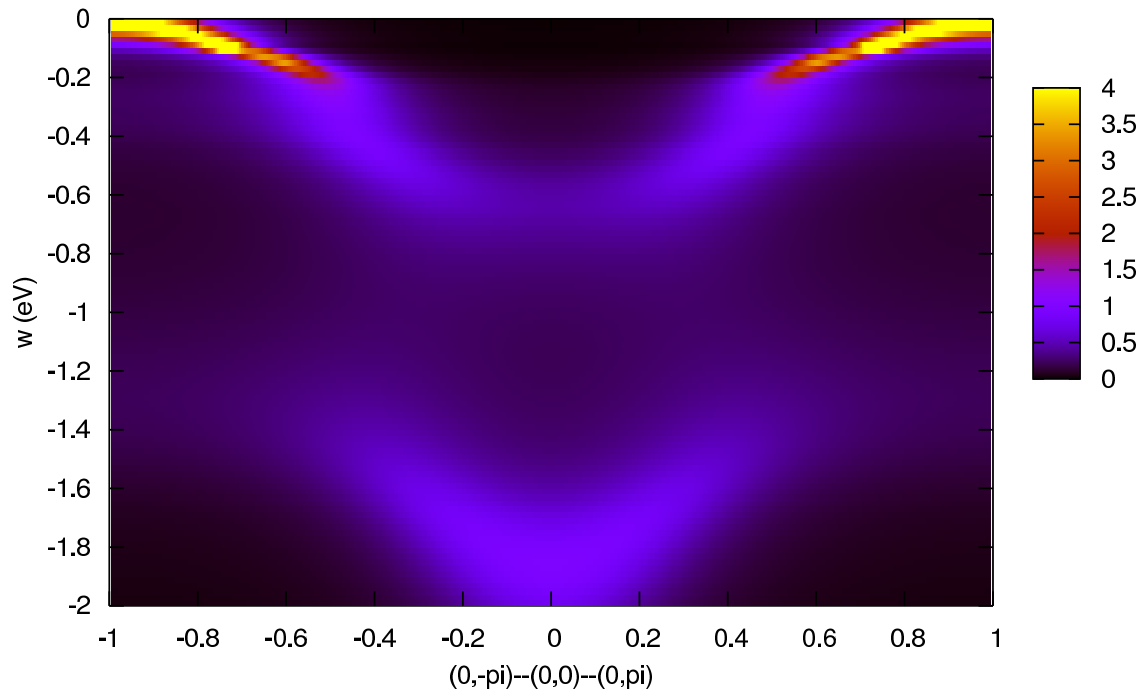
cf. Macridin *et al.*'07

cf. Byczuk, Kollar, Vollhardt'07

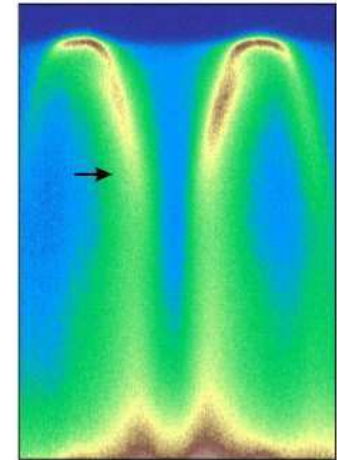
Kink position correct  
but two features kink+waterfall

# High-energy kinks in anti-nodal direction

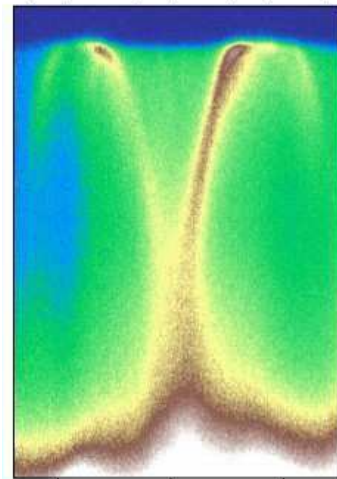
Yang, Held'07



b)  $(0,-\pi) - (0,0) - (0,\pi)$



d)  $(2\pi,-\pi) - (2\pi,0) - (2\pi,\pi)$

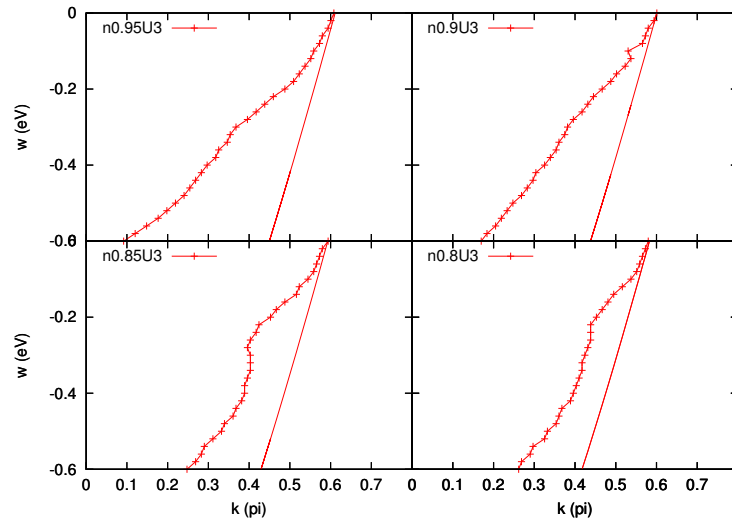


Inosov *et al.*'07

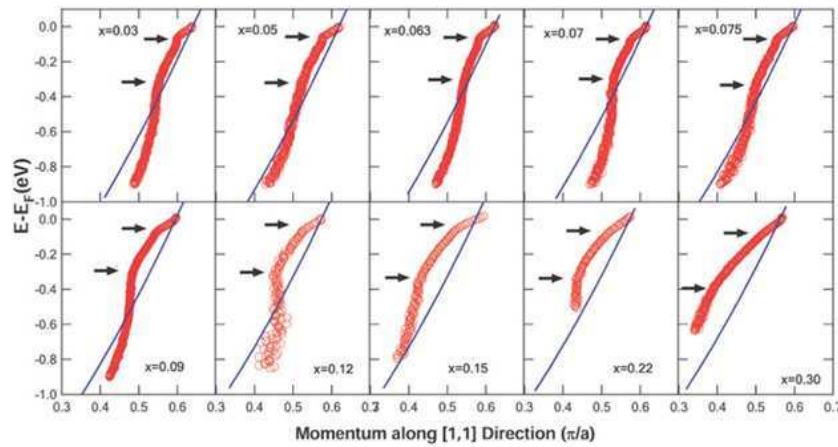
# Kinks more pronounced at higher doping

Yang, Held'07

Theory:



Experiment:

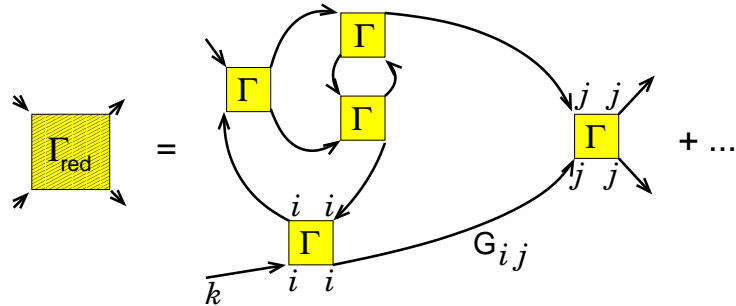


Meevasana *et al.*'06



# Conclusion — D $\Gamma$ A

- D $\Gamma$ A assumption: **local** 2-particle irreducible  $\Gamma$

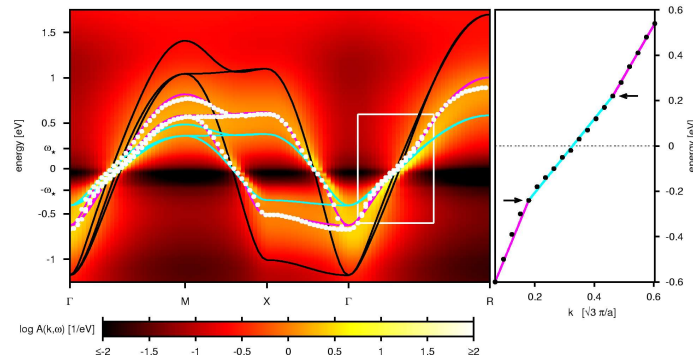


- D $\Gamma$ A can access **short-** and **long-range** correlations
- **Pseudogap** in 2D; **Mott transition** in 3D

## Outlook

- Physics: **magnons**, interplay between **AFM** and **superconductivity**, **QCP**
- Realistic **multi-orbital** calculations possible

# Conclusion – kinks



- Kinks direct consequence of strong correlations  
→ kinks are everywhere (three peak structure)

- Fermi-liquid regime:  $E_{\mathbf{k}} = Z_{\text{FL}} \epsilon_{\mathbf{k}}$  for  $|E_{\mathbf{k}}| < \omega_*$

Beyond Fermi-liquid regime:  $E_{\mathbf{k}} = Z_{\text{CP}} \epsilon_{\mathbf{k}} \pm c$  for  $|E_{\mathbf{k}}| > \omega_*$

- Connection to [high-energy kinks/waterfalls in cuprates](#)