

Single-ion anisotropy in Haldane chains
and
the form factor
of the $O(3)$ nonlinear sigma model

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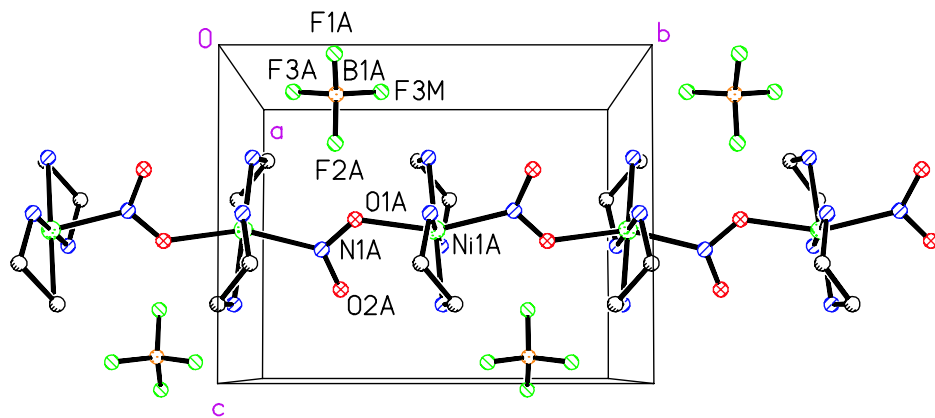
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single-ion anisotropy in S=1 Haldane chains

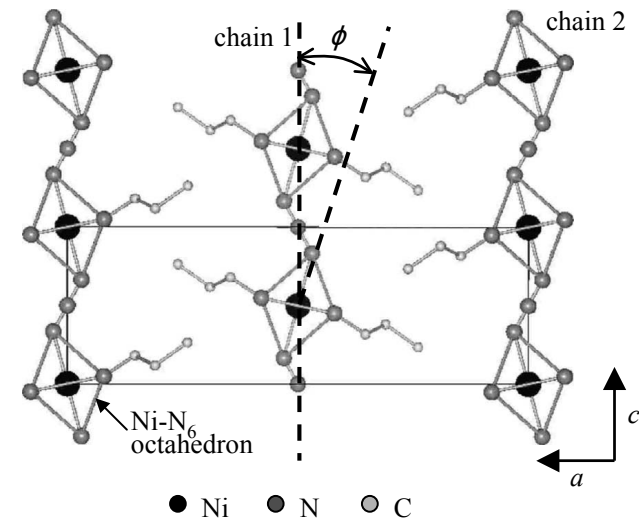
$$\mathcal{H} = J \sum_j \vec{S}_j \cdot \vec{S}_{j+1} + D \sum_j (S_j^z)^2 + E \sum_j \{(S_j^x)^2 - (S_j^y)^2\}$$

✓ NENB ($D/J = 0.17, |E|/J = 0.016$)



E. Čížmár et al., New J. Phys. (2008)

✓ NDMAP ($D/J = 0.25$)



T. Kashiwagi et al., PRB (2009)

S=1 isotropic Haldane chains

→ O(3) nonlinear sigma model $\mathcal{A} = \int dt dx \frac{1}{2g} (\partial_\mu \vec{n})^2 + i\theta Q$

$\theta = 2\pi S$

integrable

exact form factors are available

→ static & dynamical correlations

$$\langle 0 | n^a | Z_b^\dagger(\theta) \rangle = \delta_{ab} \sqrt{Z}$$

E. S. Sørensen and I. Affleck, PRB (1994)

$$\langle 0 | l^a | Z_{a_1}^\dagger(\theta_1) Z_{a_2}^\dagger(\theta_2) \rangle = i \frac{\Delta_0 \pi^2}{4} \epsilon^{aa_1 a_2} \psi(\theta_1 - \theta_2), \quad \psi(\theta) = \frac{\theta - i\pi}{\theta(\theta - 2\pi i)} \tanh^2 \frac{\theta}{2}$$

I. Affleck and R. A. Weston PRB (1992);
R. M. Konik PRB (2003)

“form factor perturbation theory”

D. Controzzi and G. Mussardo, PRL (2004)

cf.:

$$|\Psi_n\rangle = |\Phi_n\rangle + \sum_{m \neq n} \frac{\langle \Phi_m | \mathcal{H}' | \Phi_n \rangle}{E_n - E_m} |\Phi_m\rangle + \dots$$

The unperturbed system is strongly interacting,
but integrable.

Two-particle form factor of the single-ion anisotropy

$$\langle 0 | \mathcal{H}' | Z_{a_1}^\dagger(\theta_1) Z_{a_2}^\dagger(\theta_2) \rangle = ?$$

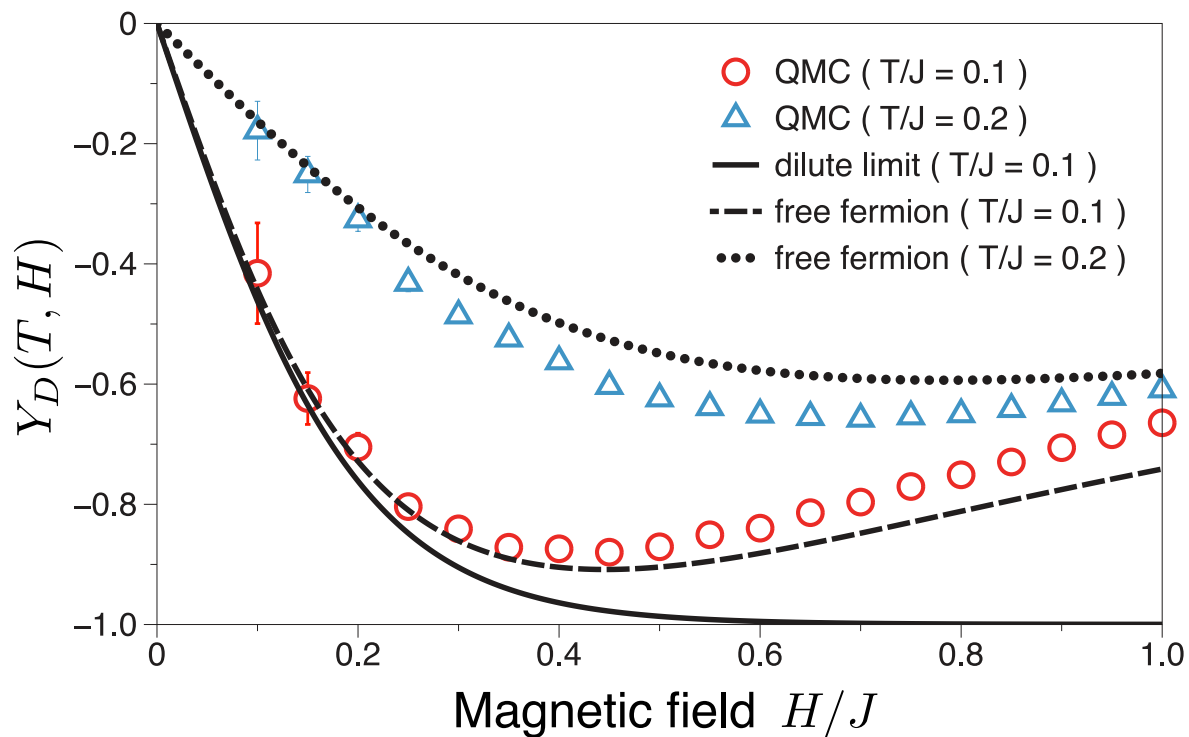
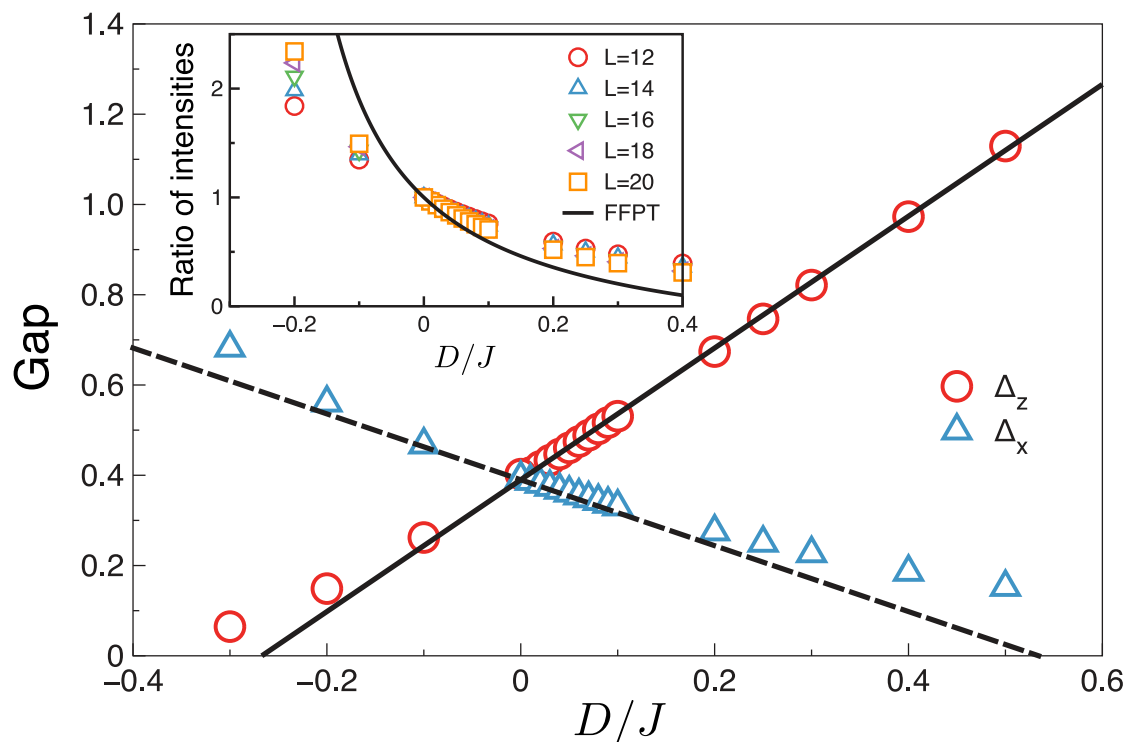
exact solution + iTEBD method

✓ excitation gaps

$$\Delta_a = \Delta_0 + \Delta_a^{(1)} + \dots$$

$$\Delta_a^{(1)} \sim \frac{\langle Z_a(\theta) | \mathcal{H}' | Z_a^\dagger(0) \rangle}{\langle Z_a(\theta) | Z_a(0) \rangle}$$

FFPT vs diagonalization



✓ ESR shifts

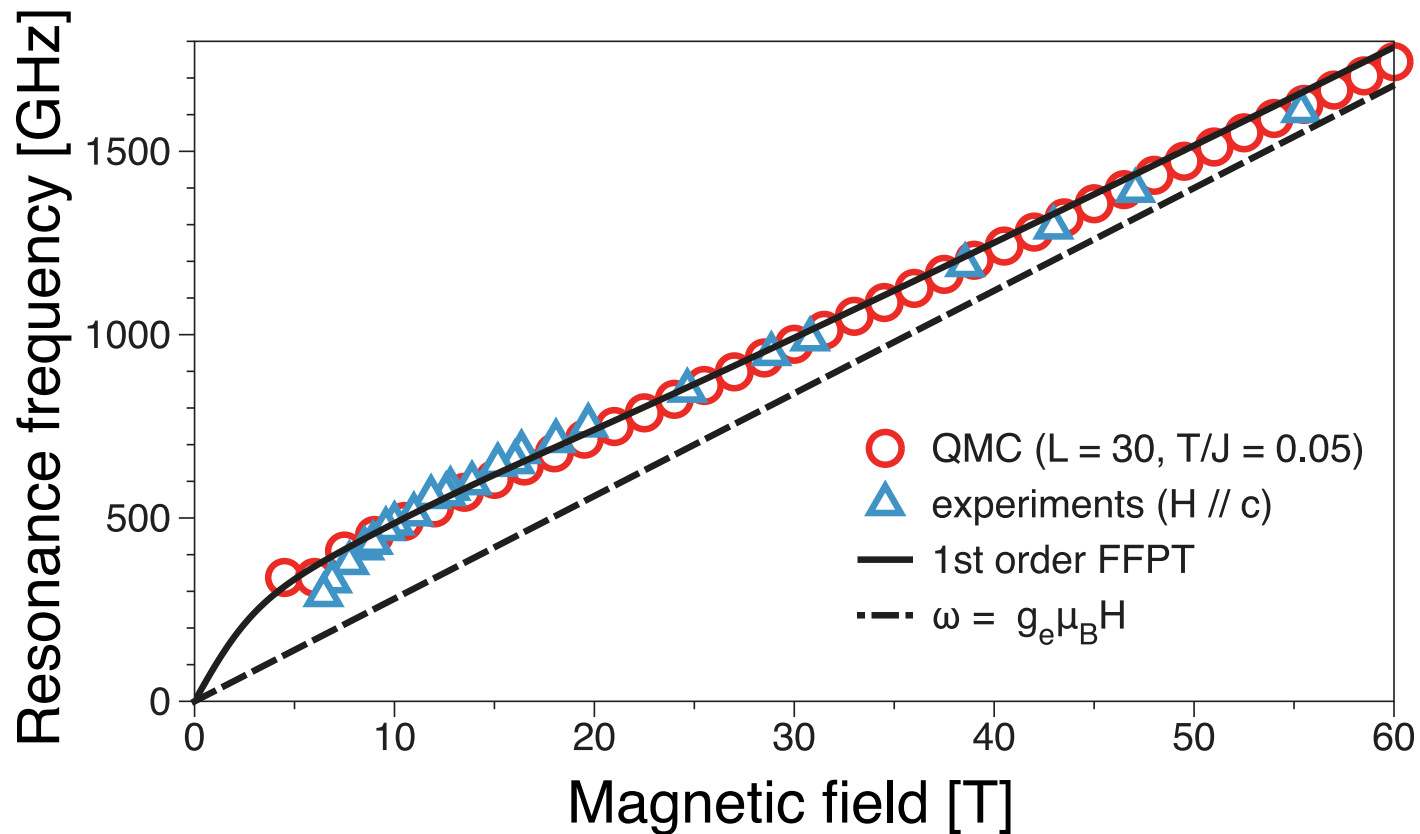
FFPT vs QMC

$$\delta\omega = \omega_r - g_e\mu_B H$$

dynamics under a magnetic field

✓ ESR frequency of NDMAP

FFPT, QMC vs experiments



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