

# Galileon and, Modifications of gravity

Antonio De Felice,

Tokyo University of Science, 東京理科大学

Summer Institute 2010 – 10 August 2010

with Shinji Mukohyama, Takahiro Tanaka, Shinji Tsujikawa



# Mt. Fuji, 富士山

# Mt. Fuji, 富士山

- Mazinger-Z was built with 超合金 (chougokin)

# Mt. Fuji, 富士山

- Mazinger-Z was built with 超合金 (chougokin)
- Only found in Mount Fuji

# Mt. Fuji, 富士山

- Mazinger-Z was built with 超合金 (chougokin)
- Only found in Mount Fuji



# Introduction

- Mystery of Dark Energy

# Introduction

- Mystery of Dark Energy
- Complex problem

# Introduction

- Mystery of Dark Energy
- Complex problem
- Motivated different approaches



# Three approaches

- Introduction of new matter, e.g. quintessence

# Three approaches

- Introduction of new matter, e.g. quintessence
- Introduction of new gravity

# Three approaches

- Introduction of new matter, e.g. quintessence
- Introduction of new gravity
  1. DGP

# Three approaches

- Introduction of new matter, e.g. quintessence
- Introduction of new gravity
  1. DGP
  2.  $f(R)$ ,  $f(R, G)$  . . . .

# Three approaches

- Introduction of new matter, e.g. quintessence
- Introduction of new gravity
  1. DGP
  2.  $f(R)$ ,  $f(R, G)$  . . .
- New gravity and new matter

# Three approaches

- Introduction of new matter, e.g. quintessence
- Introduction of new gravity
  1. DGP
  2.  $f(R)$ ,  $f(R, G)$  . . .
- New gravity and new matter
  1. Extended Brans-Dicke

# Three approaches

- Introduction of new matter, e.g. quintessence
- Introduction of new gravity
  1. DGP
  2.  $f(R)$ ,  $f(R, G)$  . . . .
- New gravity and new matter
  1. Extended Brans-Dicke
  2. Galileon . . . .

# Part I – Modified gravity

- Consider  $\mathcal{L} = f(R)$  theories



# Part I – Modified gravity

- Consider  $\mathcal{L} = f(R)$  theories
- Equivalent to ST theory  $\mathcal{L} = f'(\phi) (R - \phi) + f(\phi)$

# Part I – Modified gravity

- Consider  $\mathcal{L} = f(R)$  theories
- Equivalent to ST theory  $\mathcal{L} = f'(\phi) (R - \phi) + f(\phi)$
- In a consistent way? **No extra tensors**,  $f' = f_R > 0$

# Part I – Modified gravity

- Consider  $\mathcal{L} = f(R)$  theories
- Equivalent to ST theory  $\mathcal{L} = f'(\phi) (R - \phi) + f(\phi)$
- In a consistent way? **No extra tensors**,  $f' = f_R > 0$
- What new degrees of freedom? **1 SF**

# Part I – Modified gravity

- Consider  $\mathcal{L} = f(R)$  theories
- Equivalent to ST theory  $\mathcal{L} = f'(\phi) (R - \phi) + f(\phi)$
- In a consistent way? **No extra tensors**,  $f' = f_R > 0$
- What new degrees of freedom? **1 SF**
- Interesting phenomenology [Review by ADF, Tsujikawa '10]

# General MGM

- Introduce the gravity action  $\mathcal{L} = f(R, G)$   
[ADF, Carroll, Duvvuri, Easson, Trodden, Turner '05]

# General MGM

- Introduce the gravity action  $\mathcal{L} = f(R, G)$   
[ADF, Carroll, Duvvuri, Easson, Trodden, Turner '05]
- No ghost tensor modes,  $\mathcal{L} = f_\lambda(R - \lambda) + f_\sigma(G - \sigma) + f$

# General MGM

- Introduce the gravity action  $\mathcal{L} = f(R, G)$   
[ADF, Carroll, Duvvuri, Easson, Trodden, Turner '05]
- No ghost tensor modes,  $\mathcal{L} = f_\lambda(R - \lambda) + f_\sigma(G - \sigma) + f$
- Two extra new scalar fields,  $\lambda, \sigma$

# General MGM

- Introduce the gravity action  $\mathcal{L} = f(R, G)$   
[ADF, Carroll, Duvvuri, Easson, Trodden, Turner '05]
- No ghost tensor modes,  $\mathcal{L} = f_\lambda(R - \lambda) + f_\sigma(G - \sigma) + f$
- Two extra new scalar fields,  $\lambda, \sigma$
- Only one on FLRW and  $\omega^2 \propto (k/a)^4$   
[ADF, Suyama '09; ADF, Tanaka '10]



## Degrees of freedom [ADF, Tanaka '10]

- On Kasner background, expand  $S$  at 2nd order, 3 flds

## Degrees of freedom [ADF, Tanaka '10]

- On Kasner background, expand  $S$  at 2nd order, 3 flds
- $\mathcal{L} \propto A_{ij} \dot{\Phi}_i \dot{\Phi}_j - [C_{ij} k^2 + M_{ij}] \Phi_i \Phi_j + \dots$

## Degrees of freedom [ADF, Tanaka '10]

- On Kasner background, expand  $S$  at 2nd order, 3 flds
- $\mathcal{L} \propto A_{ij} \dot{\Phi}_i \dot{\Phi}_j - [C_{ij} k^2 + M_{ij}] \Phi_i \Phi_j + \dots$
- On Kasner  $\det A < 0$ , 1 ghost mode

## Degrees of freedom [ADF, Tanaka '10]

- On Kasner background, expand  $S$  at 2nd order, 3 flds
- $\mathcal{L} \propto A_{ij} \dot{\Phi}_i \dot{\Phi}_j - [C_{ij} k^2 + M_{ij}] \Phi_i \Phi_j + \dots$
- On Kasner  $\det A < 0$ , 1 ghost mode
- On FLRW,  $\det A = 0$ ,  $A_{33} = 0$ ,  $C_{33} = 0$ ,  $M_{33} \neq 0$

## Degrees of freedom [ADF, Tanaka '10]

- On Kasner background, expand  $S$  at 2nd order, 3 flds
- $\mathcal{L} \propto A_{ij} \dot{\Phi}_i \dot{\Phi}_j - [C_{ij} k^2 + M_{ij}] \Phi_i \Phi_j + \dots$
- On Kasner  $\det A < 0$ , **1 ghost mode**
- On FLRW,  $\det A = 0$ ,  $A_{33} = 0$ ,  $C_{33} = 0$ ,  $M_{33} \neq 0$
- $\Phi_3$ , the ghost, integrated out.

## Galileon [Nicolis et al '09]

- Keeping EOM 2nd order necessary but **not** sufficient

## Galileon [Nicolis et al '09]

- Keeping EOM 2nd order necessary but **not** sufficient
- DGP inspired introducing kinetic non-linearities to SF to implement Veinshtein mech. (mass, Cham)

## Galileon [Nicolis et al '09]

- Keeping EOM 2nd order necessary but **not** sufficient
- DGP inspired introducing kinetic non-linearities to SF to implement Veinshtein mech. (mass, Cham)
- Able to freeze the field avoiding SSC



## Galileon [Nicolis et al '09]

- Keeping EOM 2nd order necessary but **not** sufficient
- DGP inspired introducing kinetic non-linearities to SF to implement Veinshtein mech. (mass, Cham)
- Able to freeze the field avoiding SSC
- But need to keep EOM 2nd order

## Galileon [Nicolis et al '09]

- Keeping EOM 2nd order necessary but **not** sufficient
- DGP inspired introducing kinetic non-linearities to SF to implement Veinshtein mech. (mass, Cham)
- Able to freeze the field avoiding SSC
- But need to keep EOM 2nd order
- Possible example for 4D action coming from ED, ST [de Rham, Tolley '10]

## The action [Nicolis et al; Deffayet et al '09]

- Galileon symm in EOM for SF:  $\partial_\mu\phi \rightarrow \partial_\mu\phi + b_\mu$

## The action [Nicolis et al; Deffayet et al '09]

- Galileon symm in EOM for SF:  $\partial_\mu\phi \rightarrow \partial_\mu\phi + b_\mu$
- Building blocks

$$\mathcal{L}_1 = M^3\phi,$$

## The action [Nicolis et al; Deffayet et al '09]

- Galileon symm in EOM for SF:  $\partial_\mu\phi \rightarrow \partial_\mu\phi + b_\mu$
- Building blocks

$$\mathcal{L}_1 = M^3\phi,$$
$$\mathcal{L}_2 = \partial_\mu\phi\partial^\mu\phi,$$

## The action [Nicolis et al; Deffayet et al '09]

- Galileon symm in EOM for SF:  $\partial_\mu\phi \rightarrow \partial_\mu\phi + b_\mu$
- Building blocks

$$\mathcal{L}_1 = M^3\phi,$$

$$\mathcal{L}_2 = \partial_\mu\phi\partial^\mu\phi,$$

$$\mathcal{L}_3 = (\square\phi)(\nabla\phi)^2/M^3,$$

## The action [Nicolis et al; Deffayet et al '09]

- Galileon symm in EOM for SF:  $\partial_\mu\phi \rightarrow \partial_\mu\phi + b_\mu$
- Building blocks

$$\mathcal{L}_1 = M^3\phi,$$

$$\mathcal{L}_2 = \partial_\mu\phi\partial^\mu\phi,$$

$$\mathcal{L}_3 = (\square\phi)(\nabla\phi)^2/M^3,$$

$$\mathcal{L}_4 = (\nabla\phi)^2 [2(\square\phi)^2 - 2\phi_{;\mu\nu}\phi^{;\mu\nu} - R(\nabla\phi)^2/2] /M^6,$$

## The action [Nicolis et al; Deffayet et al '09]

- Galileon symm in EOM for SF:  $\partial_\mu\phi \rightarrow \partial_\mu\phi + b_\mu$
- Building blocks

$$\mathcal{L}_1 = M^3\phi,$$

$$\mathcal{L}_2 = \partial_\mu\phi\partial^\mu\phi,$$

$$\mathcal{L}_3 = (\square\phi)(\nabla\phi)^2/M^3,$$

$$\mathcal{L}_4 = (\nabla\phi)^2 [2(\square\phi)^2 - 2\phi_{;\mu\nu}\phi^{;\mu\nu} - R(\nabla\phi)^2/2] /M^6,$$

$$\begin{aligned} \mathcal{L}_5 = & (\nabla\phi)^2 [(\square\phi)^3 - 3(\square\phi)\phi_{;\mu\nu}\phi^{;\mu\nu} \\ & + 2\phi_{;\mu}{}^\nu\phi_{;\nu}{}^\rho\phi_{;\rho}{}^\mu - 6\phi_{;\mu}\phi^{;\mu\nu}\phi^{;\rho}G_{\nu\rho}] /M^9. \end{aligned}$$



- Similar to Lovelock scalars [Deffayet et al '09]

- Similar to Lovelock scalars [Deffayet et al '09]
- The Galileon symmetry is broken in general

- Similar to Lovelock scalars [Deffayet et al '09]
- The Galileon symmetry is broken in general
- Is it needed for Veinshtein mech?

- Similar to Lovelock scalars [Deffayet et al '09]
- The Galileon symmetry is broken in general
- Is it needed for Veinshtein mech?
- Study started with  $\mathcal{L}_3$ 
  - [Silva and Koyama '09]
  - [Chow and Khoury '09]
  - [Kobayashi, Tashiro, Suzuki '10]
  - [Kobayashi, Yamaguchi, Yokoyama '10]
  - [ADF and Tsujikawa '10]

# Generalized Brans-Dicke theory

[ADF, Tsujikawa '10; ADF, Mukohyama, Tsujikawa '10]

- Lagrangian

$$\mathcal{L} = \frac{1}{2}F(\phi)R - \frac{1}{2}B(\phi)\partial\phi^2 + \xi(\phi)\mathcal{L}_3.$$

# Generalized Brans-Dicke theory

[ADF, Tsujikawa '10; ADF, Mukohyama, Tsujikawa '10]

- Lagrangian

$$\mathcal{L} = \frac{1}{2}F(\phi)R - \frac{1}{2}B(\phi)\partial\phi^2 + \xi(\phi)\mathcal{L}_3.$$

- Non-minimal coupling and GL symm broken

# Generalized Brans-Dicke theory

[ADF, Tsujikawa '10; ADF, Mukohyama, Tsujikawa '10]

- Lagrangian

$$\mathcal{L} = \frac{1}{2}F(\phi)R - \frac{1}{2}B(\phi)\partial\phi^2 + \xi(\phi)\mathcal{L}_3.$$

- Non-minimal coupling and GL symm broken

- $F = M_p^2(\phi/M_p)^{3-n}$ ,  $B = \omega(\phi/M_p)^{1-n}$ ,  
 $\xi = (\lambda/\mu^3)(\phi/M_p)^{-n}$

# Generalized Brans-Dicke theory

[ADF, Tsujikawa '10; ADF, Mukohyama, Tsujikawa '10]

- Lagrangian

$$\mathcal{L} = \frac{1}{2}F(\phi)R - \frac{1}{2}B(\phi)\partial\phi^2 + \xi(\phi)\mathcal{L}_3.$$

- Non-minimal coupling and GL symm broken
- $F = M_p^2(\phi/M_p)^{3-n}$ ,  $B = \omega(\phi/M_p)^{1-n}$ ,  
 $\xi = (\lambda/\mu^3)(\phi/M_p)^{-n}$
- Existence of late dS era



# Results

- $\mu = (M_p H_{\text{dS}}^2)^{1/3}$

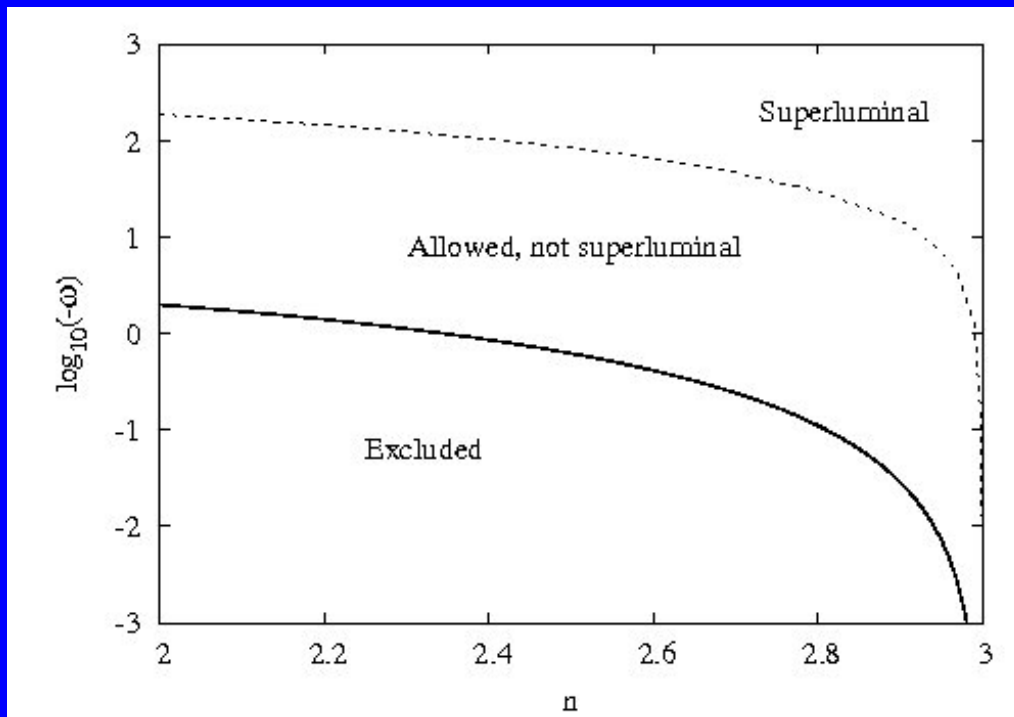
# Results

- $\mu = (M_p H_{\text{dS}}^2)^{1/3}$
- Constraints from BG ( $n < 3$ ) + Lapl inst. ( $n > 2$ ) + no-ghost  
[ADF, Mukohyama, Tsujikawa '10]

# Results

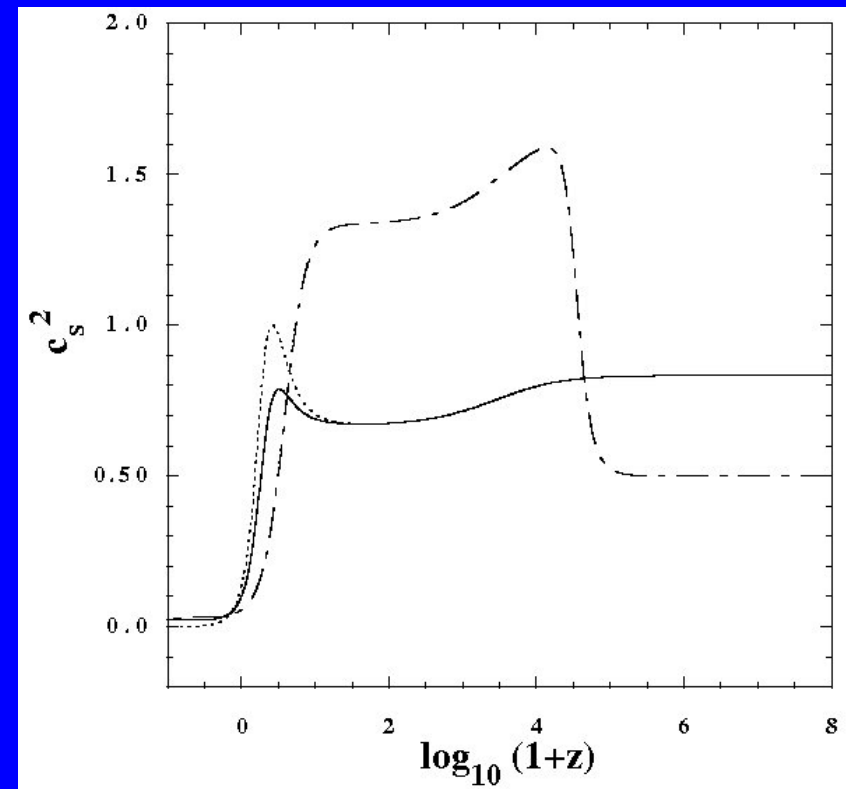
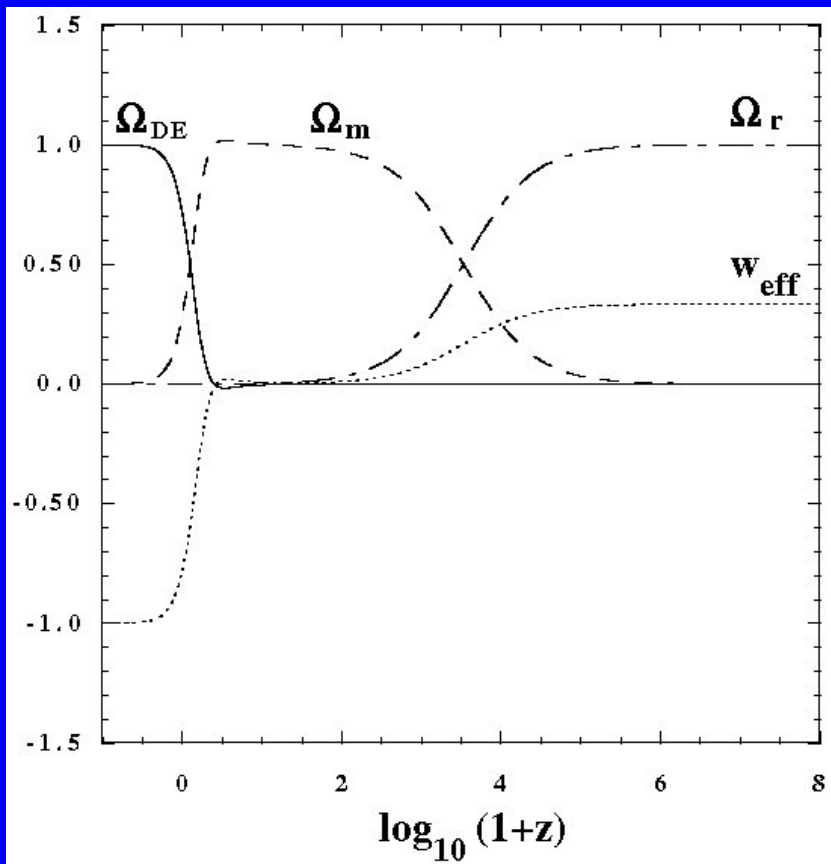
- $\mu = (M_p H_{\text{dS}}^2)^{1/3}$
- Constraints from BG ( $n < 3$ ) + Lapl inst. ( $n > 2$ ) + no-ghost  
[ADF, Mukohyama, Tsujikawa '10]
- $2 \leq n \leq 3, \dot{\phi}/(H\phi) > 0 \rightarrow \omega < -n(n-3)^2, \lambda < 0$

# Allowed parameter space



Allowed parameter space in terms of  $\omega$  vs  $n$ .  $n < 3$  for DM era,  $n > 2$  for  $c_{s,\text{dS}}^2 > 0$ .  $\dot{\phi}/\phi > 0$  for no scalar ghost.  $F > 0$  for no tensor ghost.

# Typical background evolution



## Pure Galileon [ADF, Tsujikawa '10]

- Let us come back to the Galileon Lagrangian

## Pure Galileon [ADF, Tsujikawa '10]

- Let us come back to the Galileon Lagrangian
- Building blocks

$$\mathcal{L}_1 = M^3 \phi ,$$

$$\mathcal{L}_2 = \partial_\mu \phi \partial^\mu \phi ,$$

$$\mathcal{L}_3 = (\square \phi)(\nabla \phi)^2 / M^3 ,$$

$$\mathcal{L}_4 = (\nabla \phi)^2 [2(\square \phi)^2 - 2\phi_{;\mu\nu}\phi^{;\mu\nu} - R(\nabla \phi)^2/2] / M^6 ,$$

$$\begin{aligned} \mathcal{L}_5 = & (\nabla \phi)^2 [(\square \phi)^3 - 3(\square \phi) \phi_{;\mu\nu}\phi^{;\mu\nu} \\ & + 2\phi_{;\mu}{}^\nu \phi_{;\nu}{}^\rho \phi_{;\rho}{}^\mu - 6\phi_{;\mu}\phi^{;\mu\nu}\phi^{;\rho}G_{\nu\rho}] / M^9 . \end{aligned}$$

# Total Lagrangian

$$\mathcal{L} = \frac{1}{2}M_p^2 R + \frac{1}{2} \sum_i c_i \mathcal{L}_i + \mathcal{L}_{\text{mat}},$$



# Total Lagrangian

$$\mathcal{L} = \frac{1}{2}M_p^2 R + \frac{1}{2} \sum_i c_i \mathcal{L}_i + \mathcal{L}_{\text{mat}},$$

- 2nd order EOM

# Total Lagrangian

$$\mathcal{L} = \frac{1}{2}M_p^2 R + \frac{1}{2} \sum_i c_i \mathcal{L}_i + \mathcal{L}_{\text{mat}},$$

- 2nd order EOM
- Cosmological background

# Total Lagrangian

$$\mathcal{L} = \frac{1}{2}M_p^2 R + \frac{1}{2} \sum_i c_i \mathcal{L}_i + \mathcal{L}_{\text{mat}},$$

- 2nd order EOM
- Cosmological background
- Linear perturbation analysis on FLRW

# Total Lagrangian

$$\mathcal{L} = \frac{1}{2}M_p^2 R + \frac{1}{2} \sum_i c_i \mathcal{L}_i + \mathcal{L}_{\text{mat}},$$

- 2nd order EOM
- Cosmological background
- Linear perturbation analysis on FLRW
- Studied up to  $\mathcal{L}_4$  [Gannouji, Sami '10]

# Field equations

- On FLRW:

$$3M_p^2 H^2 = \rho_{\text{DE}} + \rho_m + \rho_r ,$$

$$3M_p^2 H^2 + 2M_p^2 \dot{H} = -p_{\text{DE}} - \frac{1}{3}\rho_r .$$

# Field equations

- On FLRW:

$$3M_p^2 H^2 = \rho_{\text{DE}} + \rho_m + \rho_r ,$$

$$3M_p^2 H^2 + 2M_p^2 \dot{H} = -p_{\text{DE}} - \frac{1}{3}\rho_r .$$

- Imply  $\dot{\rho}_{\text{DE}} + 3H(\rho_{\text{DE}} + p_{\text{DE}}) = 0$

# Field equations

- On FLRW:

$$3M_p^2 H^2 = \rho_{\text{DE}} + \rho_m + \rho_r ,$$

$$3M_p^2 H^2 + 2M_p^2 \dot{H} = -p_{\text{DE}} - \frac{1}{3}\rho_r .$$

- Imply  $\dot{\rho}_{\text{DE}} + 3H(\rho_{\text{DE}} + p_{\text{DE}}) = 0$
- Admits de Sitter solutions

- Define for convenience

$$r_1 \equiv \frac{\dot{\phi}_{\text{dS}} H_{\text{dS}}}{\dot{\phi} H}, \quad r_2 \equiv \frac{\dot{\phi}^4}{r_1 \dot{\phi}_{\text{dS}}^4}, \quad x_{\text{dS}} \equiv \frac{\dot{\phi}_{\text{dS}}}{H_{\text{dS}} M_p}$$



- Define for convenience

$$r_1 \equiv \frac{\dot{\phi}_{\text{dS}} H_{\text{dS}}}{\dot{\phi} H}, \quad r_2 \equiv \frac{\dot{\phi}^4}{r_1 \dot{\phi}_{\text{dS}}^4}, \quad x_{\text{dS}} \equiv \frac{\dot{\phi}_{\text{dS}}}{H_{\text{dS}} M_p}$$

- From FRD eq

$$\begin{aligned} \Omega_{\text{DE}} \equiv \frac{\rho_{\text{DE}}}{3M_p^2 H^2} &= -\frac{1}{6} c_2 x_{\text{dS}}^2 r_1^3 r_2 + c_3 x_{\text{dS}}^3 r_1^2 r_2 \\ &\quad - \frac{15}{2} c_4 x_{\text{dS}}^4 r_1 r_2 + 7 c_5 x_{\text{dS}}^5 r_2 \end{aligned}$$

- Define for convenience

$$r_1 \equiv \frac{\dot{\phi}_{\text{dS}} H_{\text{dS}}}{\dot{\phi} H}, \quad r_2 \equiv \frac{\dot{\phi}^4}{r_1 \dot{\phi}_{\text{dS}}^4}, \quad x_{\text{dS}} \equiv \frac{\dot{\phi}_{\text{dS}}}{H_{\text{dS}} M_p}$$

- From FRD eq

$$\begin{aligned} \Omega_{\text{DE}} \equiv \frac{\rho_{\text{DE}}}{3M_p^2 H^2} &= -\frac{1}{6} c_2 x_{\text{dS}}^2 r_1^3 r_2 + c_3 x_{\text{dS}}^3 r_1^2 r_2 \\ &\quad - \frac{15}{2} c_4 x_{\text{dS}}^4 r_1 r_2 + 7 c_5 x_{\text{dS}}^5 r_2 \end{aligned}$$

- Cannot forget  $\mathcal{L}_5$

- Define

$$\alpha \equiv c_4 x_{\text{dS}}^4, \quad \beta \equiv c_5 x_{\text{dS}}^5$$

- Define

$$\alpha \equiv c_4 x_{\text{dS}}^4, \quad \beta \equiv c_5 x_{\text{dS}}^5$$

- EOM for de Sitter impose

$$c_2 x_{\text{dS}}^2 = 6 + 9\alpha - 12\beta, \quad c_3 x_{\text{dS}}^3 = 2 + 9\alpha - 9\beta$$

- Define

$$\alpha \equiv c_4 x_{\text{dS}}^4, \quad \beta \equiv c_5 x_{\text{dS}}^5$$

- EOM for de Sitter impose

$$c_2 x_{\text{dS}}^2 = 6 + 9\alpha - 12\beta, \quad c_3 x_{\text{dS}}^3 = 2 + 9\alpha - 9\beta$$

- EOM for  $r'_1$ ,  $r'_2$ , and  $\Omega'_r$ ,  $\Omega_m$  from FRD

## きれいなGalileon

- EOM for  $r_1$ :

$$r_1' = (r_1 - 1) f(r_1, r_2, \Omega_r), \quad r_2' = \dots, \quad \Omega_r' = \dots$$

## きれいなGalileon

- EOM for  $r_1$ :

$$r_1' = (r_1 - 1) f(r_1, r_2, \Omega_r), \quad r_2' = \dots, \quad \Omega_r' = \dots$$

- Existence of equilibrium point

$$r_1 = 1$$

## きれいなGalileon

- EOM for  $r_1$ :

$$r_1' = (r_1 - 1) f(r_1, r_2, \Omega_r), \quad r_2' = \dots, \quad \Omega_r' = \dots$$

- Existence of equilibrium point

$$r_1 = 1, \quad \Omega_{\text{DE}} = r_2, \quad \dot{\phi} \propto H^{-1}$$



## きれいなGalileon

- EOM for  $r_1$ :

$$r'_1 = (r_1 - 1) f(r_1, r_2, \Omega_r), \quad r'_2 = \dots, \quad \Omega'_r = \dots$$

- Existence of equilibrium point

$$r_1 = 1, \quad \Omega_{\text{DE}} = r_2, \quad \dot{\phi} \propto H^{-1}$$

- Small BG perturbations

$$\delta r'_1 = -\frac{9 + \Omega_r + 3r_2}{2(1 + r_2)} \delta r_1$$

# Other variables

- EOM

$$r_2' = \frac{2r_2(3 - 3r_2 + \Omega_r)}{1 + r_2}, \quad \Omega_r' = \frac{\Omega_r(\Omega_r - 1 - 7r_2)}{1 + r_2},$$

## Other variables

- EOM

$$r'_2 = \frac{2r_2(3 - 3r_2 + \Omega_r)}{1 + r_2}, \quad \Omega'_r = \frac{\Omega_r(\Omega_r - 1 - 7r_2)}{1 + r_2},$$

- **Universality:** no  $\alpha$  and  $\beta$  dependence

## Other variables

- EOM

$$r_2' = \frac{2r_2(3 - 3r_2 + \Omega_r)}{1 + r_2}, \quad \Omega_r' = \frac{\Omega_r(\Omega_r - 1 - 7r_2)}{1 + r_2},$$

- **Universality**: no  $\alpha$  and  $\beta$  dependence

- Three fixed points: RAD =  $(r_1 = 1, r_2 = 0, \Omega_r = 1)$ ,  
DM =  $(1, 0, 0)$ , dS =  $(1, 1, 0)$

## Other variables

- EOM

$$r_2' = \frac{2r_2(3 - 3r_2 + \Omega_r)}{1 + r_2}, \quad \Omega_r' = \frac{\Omega_r(\Omega_r - 1 - 7r_2)}{1 + r_2},$$

- **Universality:** no  $\alpha$  and  $\beta$  dependence
- Three fixed points: RAD =  $(r_1 = 1, r_2 = 0, \Omega_r = 1)$ ,  
DM =  $(1, 0, 0)$ , dS =  $(1, 1, 0)$
- Last is stable, others are saddle as  $r_2$  unstable

# Features

- BG Signature

$$w_{\text{DE}} \equiv \frac{p_{\text{DE}}}{\rho_{\text{DE}}} = -\frac{\Omega_r + 6}{3(r_2 + 1)}, \quad w_{\text{eff}} \equiv -1 - \frac{2H'}{3H} = \frac{\Omega_r - 6r_2}{3(r_2 + 1)}.$$

# Features

- BG Signature

$$w_{\text{DE}} \equiv \frac{p_{\text{DE}}}{\rho_{\text{DE}}} = -\frac{\Omega_r + 6}{3(r_2 + 1)}, \quad w_{\text{eff}} \equiv -1 - \frac{2H'}{3H} = \frac{\Omega_r - 6r_2}{3(r_2 + 1)}.$$

- Evolution

$$w_{\text{DE}} = -\frac{7}{3} \rightarrow -2 \rightarrow -1, \quad w_{\text{eff}} = \frac{1}{3} \rightarrow 0 \rightarrow -1$$

# Features

- BG Signature

$$w_{\text{DE}} \equiv \frac{p_{\text{DE}}}{\rho_{\text{DE}}} = -\frac{\Omega_r + 6}{3(r_2 + 1)}, \quad w_{\text{eff}} \equiv -1 - \frac{2H'}{3H} = \frac{\Omega_r - 6r_2}{3(r_2 + 1)}.$$

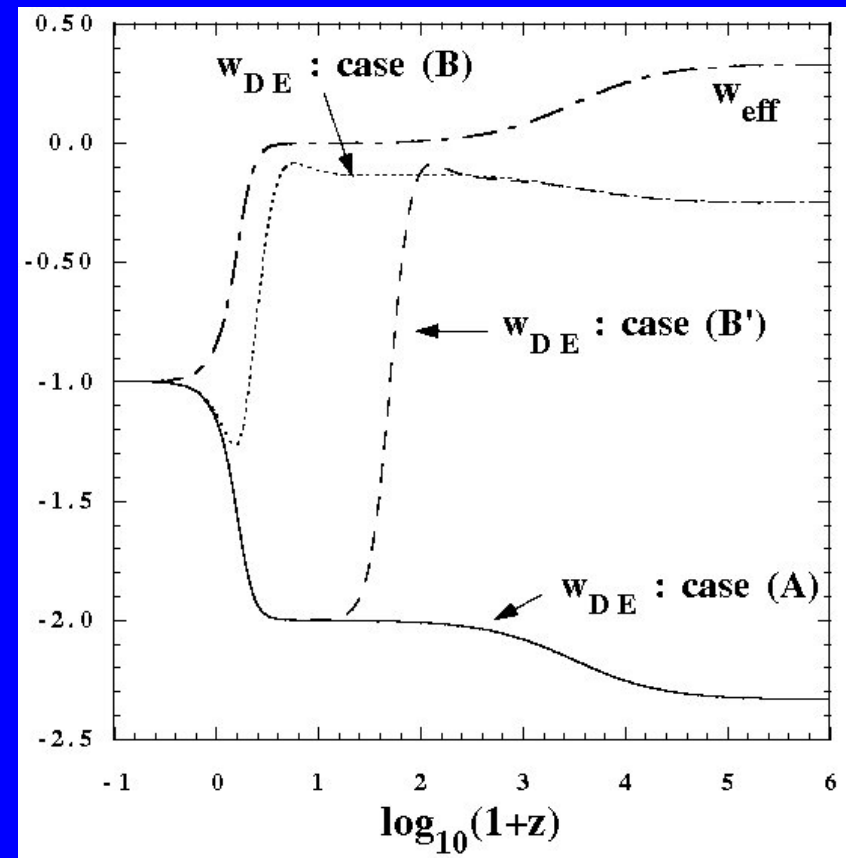
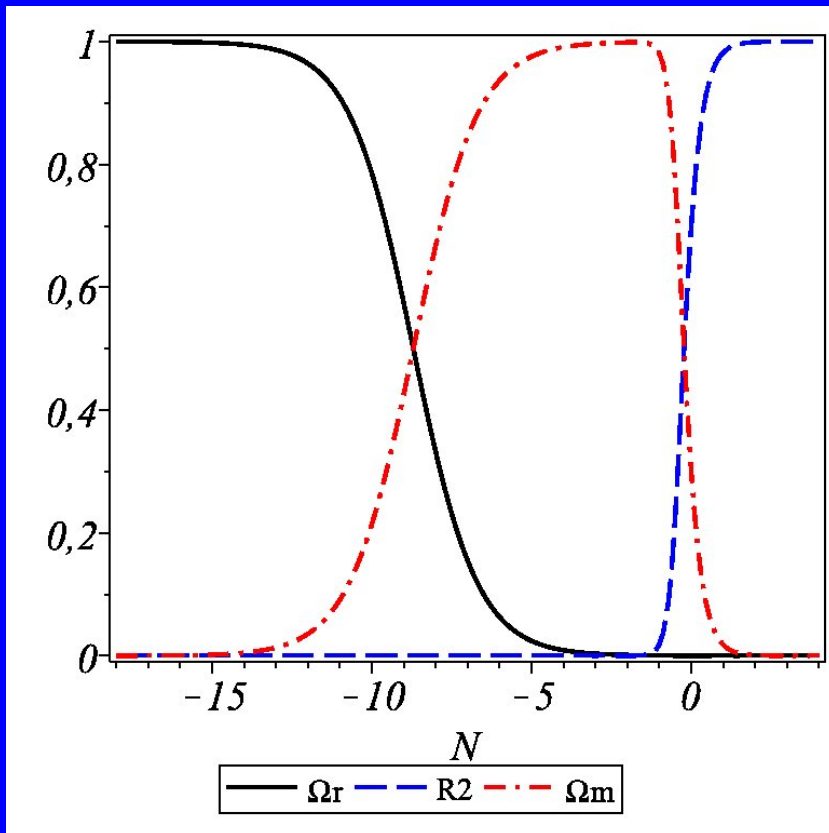
- Evolution

$$w_{\text{DE}} = -\frac{7}{3} \rightarrow -2 \rightarrow -1, \quad w_{\text{eff}} = \frac{1}{3} \rightarrow 0 \rightarrow -1$$

- To be tested against observations



# Background



# Removing degeneracy

- Universality leads to BG degeneracy

# Removing degeneracy

- Universality leads to BG degeneracy
- LPT removes degeneracy

# Removing degeneracy

- Universality leads to BG degeneracy
- LPT removes degeneracy
- Use no-ghost and  $c_S^2 > 0, c_T^2 > 0$  conditions

# LPT results

