# "Elliptic genus of K3 surface and Mathieu group M24" 

T.E. and K.Hikami

T.E., H.Ooguri and Y.Tachikawa
\& Elliptic genus of K3

Elliptic genus in string theory is expressed as

$$
Z_{\text {elliptic }}(z ; \tau)=\operatorname{Tr}_{\mathcal{H}_{L} \times \mathcal{H}_{R}(-1)^{F_{L}+F_{R}} e^{4 \pi i z J_{L, 0}^{3}} q^{L_{0}-\frac{c}{24}} \overline{\boldsymbol{q}}_{0}-\frac{c}{24}}
$$

and describes topological invariants of the manifold. Here $L_{0}$ denotes the zero mode of the Visasoro operators and $F_{L}$ and $F_{R}$ are left and right moving fermion numbers. $q=e^{2 \pi i \tau}$. In elliptic genus the right moving sector is frozen to the supersymmetric ground states (BPS states) while in the left moving sector all the states in the Hilbert space $\mathcal{H}_{L}$ contribute.

Elliptic genus of K3 surface is known:

$$
Z_{K 3}(z ; \tau)=8\left[\left(\frac{\theta_{2}(z ; \tau)}{\theta_{2}(0 ; \tau)}\right)^{2}+\left(\frac{\theta_{3}(z ; \tau)}{\theta_{3}(0 ; \tau)}\right)^{2}+\left(\frac{\theta_{4}(z ; \tau)}{\theta_{2}(0 ; \tau)}\right)^{2}\right]
$$

$$
\begin{aligned}
& Z_{K 3}(z=0)=24, \quad Z_{K 3}\left(z=\frac{1}{2}\right)=16+O(q) \\
& Z_{K 3}\left(z=\frac{1+\tau}{2}\right)=2 q^{-\frac{1}{2}}+O\left(q^{\frac{1}{2}}\right)
\end{aligned}
$$

Elliptic genus of a complex D-dimensional manifold is a Jacobi form of weight= 0 and index=D/2. When $D=2$, space of Jacobi form is one-dimensional and given by the above formula.

String theory on K3 has an $\mathrm{N}=4$ superconformal symmetry and its states fall into representations of $\mathrm{N}=4$ superconformal algebra (SCA). $\mathrm{N}=4$ SCA contains an affine $S U(2)_{k}$ symmetry
and has a central charge $c=6 k$. $k=n$ case decsribes complex-2n dimensional hyperKähler manifolds.

We would like to study the decomposition of the elliptic genus in terms of irreducible characters of $\mathrm{N}=4$ SCA. In $\mathrm{N}=4$ SCA, hightest-weight states $|h, I\rangle$ are charactered by

$$
L_{0}|h, \ell\rangle=|h, \ell\rangle, \quad J_{0}^{3}|h, \ell\rangle=\ell|h, \ell\rangle
$$

and the theory possesses two different type of representations, BPS and non-BPS representations. In the case of $k=1$
there are representations (in Ramond sector)


The index of a representation is given by the value of its character at $z=0,\left.\operatorname{Tr}_{\mathcal{R}}(-1)^{F} q^{L_{0}} e^{4 \pi i z J_{0}^{3}}\right|_{z=0}$. BPS representations have a non-vanishing index

$$
\begin{aligned}
& \text { index }(B P S, \ell=0)=1 \\
& \text { index }\left(B P S, \ell=\frac{1}{2}\right)=-2
\end{aligned}
$$

Character function of $\ell=0$ BPS representation has the form

$$
c h_{k=1, h=\frac{1}{4}, \ell=0}^{\tilde{R}}(z ; \tau)=\frac{\theta_{1}(z ; \tau)^{2}}{\eta(\tau)^{3}} \mu(z ; \tau)
$$

where

$$
\mu(z ; \tau)=\frac{-i e^{\pi i z}}{\theta_{1}(z ; \tau)} \sum_{n}(-1)^{n} \frac{q^{\frac{1}{2} n(n+1)} e^{2 \pi i n z}}{1-q^{n} e^{2 \pi i z}}
$$

On the other hand the character of non-BPS representations are given by

$$
c h_{k=1, h>\frac{1}{4}, \ell=\frac{1}{2}}^{\tilde{R}}=q^{h-\frac{3}{8}} \frac{\theta_{1}(z ; \tau)^{2}}{\eta(\tau)^{3}}
$$

These have a vanishing index

$$
\text { index (non-BPS rep) }=0
$$

At the unitarity bound non-BPS representation splits into a sum of BPS representations

$$
q^{-\frac{1}{8}} \frac{\theta_{1}(z ; \tau)^{2}}{\eta(\tau)^{3}}=\operatorname{ch}_{k=1, h=\frac{1}{4}, \ell=\frac{1}{2}}^{\tilde{R}}+2 c h_{k=1, h=\frac{1}{4}, \ell=0}^{\tilde{R}}
$$

Function $\mu(z ; \tau)$ is a typical example of the so-called Mock theta functions (Lerch sum or Appell function). Mock theta functions look like theta functions but they have anomalous modular transformation laws. Recently there has been a development in the understanding the nature of Mock theta
functions by Zwegers and he has developed a way to improve their modular properties. We will adopt his method of handling Mock theta functions.

It is possible to derive the following idenities

$$
\begin{aligned}
& \operatorname{ch}_{k=1, h=\frac{1}{4}, \ell=0}^{\tilde{R}}(z ; \tau)=\left(\frac{\theta_{2}(z ; \tau)}{\theta_{2}(0 ; \tau)}\right)^{2}+\mu_{2}(\tau) \frac{\theta_{1}(z ; \tau)^{2}}{\eta(\tau)^{3}} \\
& c h_{k=1, h=\frac{1}{4}, \ell=0}^{\tilde{R}}(z ; \tau)=\left(\frac{\theta_{3}(z ; \tau)}{\theta_{3}(0 ; \tau)}\right)^{2}+\mu_{3}(\tau) \frac{\theta_{1}(z ; \tau)^{2}}{\eta(\tau)^{3}} \\
& c h_{k=1, h=\frac{1}{4}, \ell=0}^{\tilde{R}}(z ; \tau)=\left(\frac{\theta_{4}(z ; \tau)}{\theta_{4}(0 ; \tau)}\right)^{2}+\mu_{4}(\tau) \frac{\theta_{1}(z ; \tau)^{2}}{\eta(\tau)^{3}}
\end{aligned}
$$

where

$$
\begin{gathered}
\mu_{2}(\tau)=\mu\left(z=\frac{1}{2} ; \tau\right), \mu_{3}(\tau)=\mu\left(z=\frac{1+\tau}{2} ; \tau\right), \mu_{4}(\tau)=\mu\left(z=\frac{\tau}{2} ; \tau\right) \\
\mu(z ; \tau)=\frac{-i e^{\pi i z}}{\theta_{1}(z ; \tau)} \sum_{n}(-1)^{n} \frac{q^{\frac{1}{2} n(n+1)} e^{2 \pi i n z}}{1-q^{n} e^{2 \pi i z}}
\end{gathered}
$$

Then by making use of these formulas we can rewrite the ellitpic genus as

$$
Z_{K 3}=24 c h_{k=1, h=\frac{1}{4}, \ell=0}^{\tilde{R}}(z ; \tau)-8 \sum_{i=2}^{4} \mu_{i}(\tau) \frac{\theta_{1}(z ; \tau)^{2}}{\eta(\tau)^{3}}
$$

q-expansion of functions $\mu_{i}$

$$
8\left(\mu_{2}(\tau)+\mu_{3}(\tau)+\mu_{4}(\tau)\right)=2 q^{-\frac{1}{8}}-2 \sum_{n=1} A_{n} q^{n-\frac{1}{8}}
$$

## polar term

$A_{n}(n=1,2, \cdots)$ are positive integers. Using

$$
q^{-\frac{1}{8}} \frac{\theta_{1}(z ; \tau)^{2}}{\eta(\tau)^{3}}=\operatorname{ch}_{k=1, \ell=\frac{1}{2}}^{\tilde{R}}+2 \operatorname{ch}_{k=1, \ell=0}^{\tilde{R}}
$$

we can rewrite

$$
Z_{K 3}=20 \operatorname{ch}_{k=1, \ell=0}^{\tilde{R}}-2 c h_{k=1, \ell=\frac{1}{2}}^{\tilde{R}}+2 \sum_{n=1} A_{n} q^{n-\frac{1}{8}} \frac{\theta_{1}(z ; \tau)^{2}}{\eta(\tau)^{3}}
$$

This is the character expansion of K3 elliptic genus. Type II string theory compactified on K3 possesses

1 : graviton multiplet
20 : tensor, hyper multiplets

At smaller values of $n$, Fourier coefficients $A_{\boldsymbol{n}}$ may be obtained by direct expansion. We find

| $n$ | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | $\ldots$ |
| :---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | ---: | :--- |
| $A_{n}$ | 45 | 231 | 770 | 2277 | 5796 | 13915 | 30843 | 65550 | $\ldots$ |

A surprize: Dimensions of some irreducible reps. of Mathieu group $M_{24}$ appear
dimension | $45 \quad 231 \quad 770 \quad 2277 \quad 5796 \quad 3520 \quad 10395 \quad 5544 \ldots$

$$
\begin{aligned}
& A_{6}=3520+10395 \\
& A_{7}=10395+5796+5544+5313+3312+2024+1774
\end{aligned}
$$

Mathieu moonshine?
T.E.-Ooguri-Tachikawa

Monsterous moonshine:

$$
J(q)=\frac{1}{q}+196884+221493760 q^{2}+\cdots
$$

q-expansion coeffcients of J-function are decomposed into a sum of irred. reps. of monster simple group.

Mukai: enumeration of eleven K3 surfaces with finite nonAbelian automorphism group. All these groups are sugbgroups of $M_{24}$.
fantasy: Is it possible that these automorphism groups at isoletd points in K3 moduli space are enhanced to $M_{24}$ over the whole moduli space when one cosnider the elliptic genus?

On the other hand, using the method of Rademacher expansion adapted to the case of Mock theta functions (BringmannOno ) we can determine the asymptotic behavior of coefficients $A_{n}$ as

$$
A_{n} \approx \frac{2}{\sqrt{8 n-1}} e^{2 \pi \sqrt{\frac{1}{2}\left(n-\frac{1}{8}\right)}}
$$

T.E.-Hikami

More about Mathieu group

The Mathieu group $M_{24}$ is a subgroup of $A_{24}$ and an automorphism group of binary Golay code. There are 26 conjugacy classes and 26 irreducible representations. Characters
are is given in the table. ( $e_{p}^{ \pm}$stands for $e_{p}^{ \pm}=( \pm \sqrt{-p}-1) / 2$ ).
The dimensions of the irreducible representations are

$$
\begin{aligned}
& 1,23,45,45^{*}, 231,231^{*}, 252,253,483,770,770^{*} \\
& 990,990^{*}, 1035,1035^{*}, 1035^{\prime}, 1265,1771,2024 \\
& 2277,3312,3520,5313,5796,5544,10395 .
\end{aligned}
$$

Here irreducible representations of dimensions

$$
45,231,770,990,1035
$$

come in complex conjugate pairs. There is in addition an extra real 1035-dimensional irreducible representation.

## Character table I

| 1 A | 2A | 3A | 5A | 4B | 7A | 8A | 6A | 11A | 15A | 14A | 23A | 12B | 6B | 4C | 3B | 2B | 10A | 21A | 4A |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | 1 | I | 1 | 1 |
| 23 | 7 | 5 | 3 | 3 | 2 | 1 | 1 | 1 | 0 | 0 | 0 | -1 | -1 | -1 | -1 | -1 | -1 | -1 | -1 |
| 252 | 28 | 9 | 2 | 4 | 0 | 0 | 1 | -1 | -1 | 0 | -1 | 0 | 0 | 0 | 0 | 12 | 2 | 0 | 4 |
| 253 | 13 | 10 | 3 | 1 | 1 | -1 | -2 | 0 | 0 | -1 | 0 | 1 | 1 | 1 | 1 | -11 | -1 | 1 | -3 |
| 1771 | -21 | 16 | 1 | -5 | 0 | -1 | 0 | 0 | 1 | 0 | 0 | -1 | -1 | -1 | 7 | 11 | 1 | 0 | 3 |
| 3520 | 64 | 10 | 0 | 0 | -1 | 0 | -2 | 0 | 0 | 1 | 1 | 0 | 0 | 0 | -8 | 0 | 0 | -1 | 0 |
| 45 | -3 | 0 | 0 | 1 | $e_{7}^{+}$ | -1 | 0 | 1 | 0 | $-e_{7}^{+}$ | -1 | 1 | -1 | 1 | 3 | 5 | 0 | $e_{7}^{-}$ | -3 |
| 990 | -18 | 0 | 0 | 2 | $e_{7}^{+}$ | 0 | 0 | 0 | 0 | $e_{7}^{+}$ | 1 | 1 | -1 | -2 | 3 | -10 | 0 | $e_{7}^{-}$ | 6 |
| 1035 | -21 | 0 | 0 | 3 | $2 e_{7}^{+}$ | -1 | 0 | 1 | 0 | 0 | 0 | -1 | 1 | -1 | -3 | -5 | 0 | $-e_{7}^{-}$ | 3 |
| 1035 | 27 | 0 | 0 | -1 | -1 | 1 | 0 | 1 | 0 | -1 | 0 | 0 | 2 | 3 | 6 | 35 | 0 | -1 | 3 |
| 231 | 7 | -3 | 1 | -1 | 0 | -1 | 1 | 0 | $e_{15}^{+}$ | 0 | 1 | 0 | 0 | 3 | 0 | -9 | 1 | 0 | -1 |
| 770 | -14 | 5 | 0 | -2 | 0 | 0 | 1 | 0 | 0 | 0 | $e_{23}^{+}$ | 1 | 1 | -2 | -7 | 10 | 0 | 0 | 2 |
| 483 | 35 | 6 | -2 | 3 | 0 | -1 | 2 | -1 | 1 | 0 | 0 | 0 | 0 | 3 | 0 | 3 | -2 | 0 | 3 |
| 1265 | 49 | 5 | 0 | 1 | -2 | 1 | 1 | 0 | 0 | 0 | 0 | 0 | 0 | -3 | 8 | -15 | 0 | 1 | -7 |
| 2024 | 8 | -1 | -1 | 0 | 1 | 0 | -1 | 0 | -1 | 1 | 0 | 0 | 0 | 0 | 8 | 24 | -1 | 1 | 8 |
| 2277 | 21 | 0 | -3 | 1 | 2 | -1 | 0 | 0 | 0 | 0 | 0 | 0 | 2 | -3 | 6 | -19 | 1 | -1 | -3 |
| 3312 | 48 | 0 | -3 | 0 | 1 | 0 | 0 | 1 | 0 | -1 | 0 | 0 | -2 | 0 | -6 | 16 | 1 | 1 | 0 |
| 5313 | 49 | -15 | 3 | -3 | 0 | -1 | 1 | 0 | 0 | 0 | 0 | 0 | 0 | -3 | 0 | 9 | -1 | 0 | 1 |
| 5796 | -28 | -9 | 1 | 4 | 0 | 0 | -1 | -1 | 1 | 0 | 0 | 0 | 0 | 0 | 0 | 36 | 1 | 0 | -4 |
| 5544 | -56 | 9 | -1 | 0 | 0 | 0 | 1 | 0 | -1 | 0 | 1 | 0 | 0 | 0 | 0 | 24 | -1 | 0 | -8 |
| 10395 | -21 | 0 | 0 | -1 | 0 | 1 | 0 | 0 | 0 | 0 | -1 | 0 | 0 | 3 | 0 | -45 | 0 | 0 | 3 |

## Character table II

| conj. class $\backslash n$ | 0 | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | 9 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $M_{23}$ |  |  |  |  |  |  |  |  |  |  |
| 1A, $1^{24}$ | 24 | 2x45 | 2x231 | 2x770 | 2x2277 | 2x5796 | 2x13915 | 2x30843 | 2x65550 | 2x132825 |
| 2A, $1^{8} 2^{8}$ | 16 | -6 | 14 | -28 | 42 | -56 | 86 | -138 | 188 | -238 |
| 3A, $1^{6} 3^{6}$ | 6 | 0 | -6 | 10 | 0 | -18 | 20 | 0 | -30 | 42 |
| 5A, $1^{4} 5^{4}$ | 4 | 0 | 2 | 0 | -6 | 2 | 0 | 6 | 0 | -10 |
| 7A, $1^{3} 7^{3}$ | 3 | -1 | 0 | 0 | 4 | 0 | -2 | 2 | -3 | 0 |
| 4B, $1^{4} 2^{2} 4^{4}$ |  | 2 | -2 | -4 | 2 | 8 | -2 | -10 | 4 | 10 |
| $6 \mathrm{~A}, 1^{2} 2^{2} 3^{2} 6^{2}$ |  | 0 | 2 | 2 | 0 | -2 | -4 | 0 | 2 | 2 |
| 8A, $1^{2} \cdot 2 \cdot 4 \cdot 8^{2}$ |  | -2 | -2 | 0 | -2 | 0 | 2 | -2 | 0 | -2 |
| 11A, $1^{2} 11^{2}$ |  | 2 | 0 | 0 | 0 | -2 | 0 | -2 | 2 | 0 |
| 14A, $1 \cdot 2 \cdot 7 \cdot 14$ |  | 1 | 0 | 0 | 0 | 0 | 2 | 2 | -1 | 0 |
| 15A, $1 \cdot 3 \cdot 5 \cdot 15$ |  | 0 | -1 | 0 | 0 | 2 | 0 | 0 | 0 | 2 |
| 23A, $1 \cdot 23$ |  | -2 | 2 | -1 | 0 | 0 | 0 | 0 | 0 | 0 |
| $M_{24}$ |  |  |  |  |  |  |  |  |  |  |
| 2B, $2^{12}$ |  | 10 | -18 | 20 | -38 | 72 | -90 | 118 | -180 | 258 |
| 4A, $2^{4} 4^{4}$ |  | -6 | -2 | 4 | -6 | -8 | 6 | 6 | -4 | -14 |
| 4C, $4^{6}$ |  | 1 | 3 | -2 | -3 | 0 | 3 | -1 | -6 | 5 |
| 3B, $3^{8}$ |  | 6 | 0 | -14 | 12 | 0 | -16 | 30 | 0 | -42 |
| 6B, $6^{4}$ |  | -2 | 0 | 2 | 4 | 0 | 0 | -2 | 0 | 6 |
| 12B, $12{ }^{2}$ |  | 2 | 0 | 2 | 0 | 0 | 0 | -2 | 0 | -2 |
| 10A, $2^{2} 10^{2}$ |  | 0 | 2 | 0 | 2 | 2 | 0 | -2 | 0 | -2 |
| 12A, $2 \cdot 4 \cdot 6 \cdot 12$ |  | 0 | -2 | -2 | 0 | -2 | 0 | 0 | 2 | -2 |
| 21A, 3.21 |  | -1 | 0 | 0 | -2 | 0 | -2 | 2 | 0 | 0 |

## Twisted elliptic genera

$$
\begin{aligned}
& Z_{p A}(\tau, z)=\frac{2}{p+1} \phi_{0,1}(\tau, z)+\frac{2 p}{p+1} \phi_{-2,1}(\tau, z) \phi_{2}^{(p)}(\tau), \quad p=2,3,5,7 \\
& Z_{4 B}(\tau, z)=\frac{1}{3} \phi_{0,1}(\tau, z)+\phi_{-2,1}(\tau, z)\left(-\frac{1}{3} \phi_{2}^{(2)}(\tau)+2 \phi_{2}^{(4)}(\tau)\right), \\
& Z_{6 A}(\tau, z)= \frac{1}{6} \phi_{0,1}(\tau, z)+\phi_{-2,1}(\tau, z)\left(-\frac{1}{6} \phi_{2}^{(2)}(\tau)-\frac{1}{2} \phi_{2}^{(4)}(\tau)+\frac{5}{2} \phi_{2}^{(6)}(\tau)\right), \\
& Z_{8 A}(\tau, z)=\frac{1}{6} \phi_{0,1}(\tau, z)+\phi_{-2,1}(\tau, z)\left(-\frac{1}{2} \phi_{2}^{(4)}(\tau)+\frac{7}{3} \phi_{2}^{(8)}(\tau)\right), \\
& Z_{11 A}(\tau, z)=\frac{1}{6} \phi_{0,1}(\tau, z)+\phi_{-2,1}(\tau, z)\left(\frac{11}{6} \phi_{2}^{(11)}(\tau)-\frac{22}{5} f_{11}(\tau)\right), \\
& Z_{14 A}(\tau, z)=\frac{1}{12} \phi_{0,1}(\tau, z)+\phi_{-2,1}(\tau, z)\left(-\frac{1}{36} \phi_{2}^{(2)}(\tau)-\frac{7}{12} \phi_{2}^{(7)}(\tau)+\frac{91}{36} \phi_{2}^{(14)}(\tau)-\frac{14}{3} f_{14}(\tau)\right), \\
& \quad f_{14}(\tau)=\eta(\tau) \eta(2 \tau) \eta(7 \tau) \eta(14 \tau), \\
& Z_{15 A}(\tau, z)=\frac{1}{12} \phi_{0,1}(\tau, z)+\phi_{-2,1}(\tau, z)\left(-\frac{1}{16} \phi_{2}^{(3)}(\tau)-\frac{5}{24} \phi_{2}^{(5)}(\tau)+\frac{35}{16} \phi_{2}^{(15)}(\tau)-\frac{15}{4} f_{15}(\tau)\right), \\
& f_{15}(\tau)=\eta(\tau) \eta(3 \tau) \eta(5 \tau) \eta(15 \tau),
\end{aligned}
$$

$$
\phi_{0,1}=1 / 2 \times(\text { elliptic genus of } \mathrm{K} 3), \phi_{-2,1}=\frac{\theta_{1}^{2}}{\eta^{6}}, \quad \phi_{2}^{(\mathrm{N})}=\frac{24}{\mathrm{~N}-1} \mathrm{q} \partial_{\mathrm{q}} \log \frac{\eta(\mathrm{~N} \tau)}{\eta(\tau)}
$$

There is a strong evidence in favor of the Mathieu moonshine conjecture. M.Cheng,Gaberdiel-Hohennegger-Volpato
© Entropy of HyperKähler manifold

Since in the elliptic genus the right-moving sector is frozen to the ground (BPS) state, left-moving non-BPS states are halfBPS states when we consider both left-right moving sectors simultaneously. Thus $A_{\boldsymbol{n}}$ describes the multiplicity of half-BPS states at higher levels. When the multiplicity of states protected by supersymmetry increases like an exponential, we may interpret it as the entropy of the system.

We want to compute the entropy for higher dimensional hyperKähler manifolds, in particular symmetric products of K3 surfaces, $K 3^{[k]}$ and check if the entropy considered here agrees with the entropy of black holes of string theory compactified on K3 surface.

Character expansion of elliptic genera at level $k$ is given by

$$
\begin{aligned}
& Z\left(H_{2 k}\right)=\sum_{\ell=0}^{\frac{k}{2}} n_{\ell} c h_{k, h=\frac{k}{4}, \ell}^{\tilde{R}}(z ; \tau) \\
& +\sum_{\ell=0}^{\frac{k-1}{2}} \sum_{n=1}^{\infty} A_{n}^{\ell} q^{n-\frac{\left(\ell+\frac{1}{2}\right)^{2}}{k+1}} \chi_{k-1, \ell}(z ; \tau) \frac{\theta_{1}(z ; \tau)^{2}}{\eta(\tau)^{3}}
\end{aligned}
$$

Here $\chi_{k-1, \ell}$ is the character of affine $S U(2)$ at level- $(k-1)$.

Asymptotic growth of $A_{n}^{\ell}$ is determined by the polar parts of the series. Results are given as (for large $k$ )

$$
A_{n}^{\ell} \approx \exp \left(\pi \sqrt{4 k n-\ell^{2}}\right)
$$

## T.E.-Hikami

Dominant multiplicity $A_{n}$ behaves as

$$
A_{n} \approx \exp (\pi \sqrt{4 k n})
$$

and agrees with the entropy $S=2 \pi \sqrt{Q_{1} Q_{5} n}$ of the standard $D_{1} D_{5}$ black hole with $k=Q_{1} Q_{5}$.

K3 surface corresponds to the mini black hole $Q_{1}=Q_{5}=$

1. It will be very intersting if one can interprete the Mathieu group as acting on the micro-states of mini black holes.
