

Summer Institute 2010

G-inflation

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What's "G"?

Galileon field

The Galileon field

$\mathcal{L}_1 = \phi$ Field equations are 2nd order

$$\mathcal{L}_2 = (\nabla\phi)^2$$

$$\mathcal{L}_3 = (\nabla\phi)^2 \square\phi$$

$$\mathcal{L}_4 = (\nabla\phi)^2 \left[2(\square\phi)^2 \right.$$

$$\left. -2(\nabla_\mu\nabla_\nu\phi)^2 - \frac{R}{2}(\nabla\phi)^2 \right]$$

$$\mathcal{L}_5 = (\nabla\phi)^2 \left[(\square\phi)^3 + \dots \right]$$

Galilean shift symmetry in flat space

$$\partial_\mu\phi \rightarrow \partial_\mu\phi + b_\mu$$

$$\mathcal{L}_n \sim \partial^{2(n-1)}\phi^n$$

Nicolis et al. '09;
Deffayet et al. '09

Our Lagrangian

$$\mathcal{L} = \frac{R}{2} + K(\phi, X) - F(\phi, X)\square\phi$$

where $X := -\frac{1}{2}(\nabla\phi)^2$

Field equations are 2nd order

Simple motivation

The Galileon field has been used to explain
current cosmic acceleration

Antonio De Felice talk yesterday

Why don't we use the Galileon field to
drive inflation in the early universe?

Talk plan

- I. Introduction
- II. G-inflation
- III. Primordial perturbations
- IV. Pre-G-inflation?
- V. Summary & Outlook

G-inflation

The image features a clear blue sky with scattered white cumulus clouds. A vertical strip of textured, light-colored material, possibly paper or fabric, runs down the left side of the frame. The text 'G-inflation' is centered in a white, serif font with a subtle drop shadow.

Standard picture of inflation

One (or more) canonical scalar field(s) rolling slowly down a nearly flat potential

$$\mathcal{L} = X - V(\phi), \quad X = -\frac{1}{2}(\partial\phi)^2$$

$$3M_{\text{Pl}}^2 H^2 \simeq V(\phi)$$



Kinematically driven inflation

$$\mathcal{L} = K(\phi, X)$$

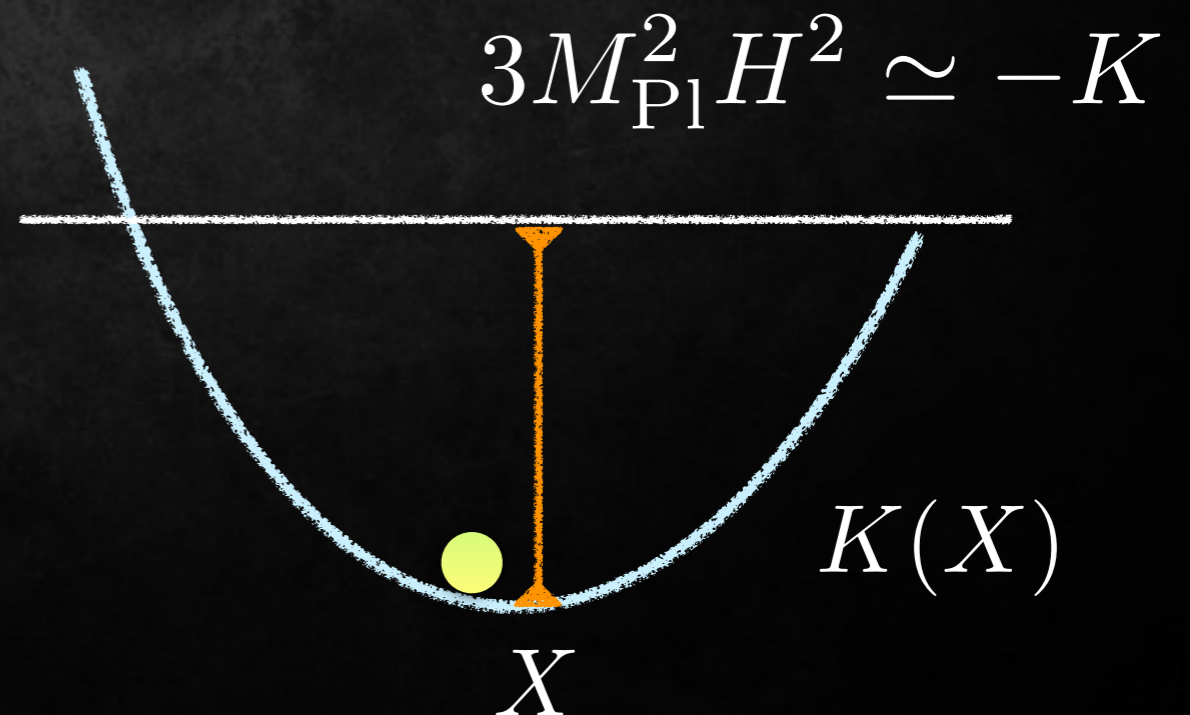
$$\xrightarrow{K = K(X)} \frac{d}{dt} \left(a^3 K_X \dot{\phi} \right) = 0$$

“k-inflation”

Armendariz-Picon et al. '99;

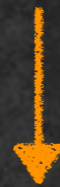
“Ghost condensate”

Arkani-Hamed et al. '04



G-inflation: background

$$\mathcal{L}_\phi = K(\phi, X) - F(\phi, X)\square\phi$$



$$\begin{aligned} 3H^2 &= \rho \\ -3H^2 - 2\dot{H} &= p \end{aligned} + \text{Scalar field EOM is automatically satisfied}$$

$$\rho = 2XK_X - K + 3F_X H \dot{\phi}^3 - 2F_\phi X$$

$$p = K - 2\left(F_\phi + F_X \ddot{\phi}\right) X$$

de Sitter G-inflation

$$K = K(X), \quad F = fX, \quad f = \text{const}$$

Look for **exactly de Sitter** solution:

$$H = \text{const}$$

$$\dot{\phi} = \text{const}$$

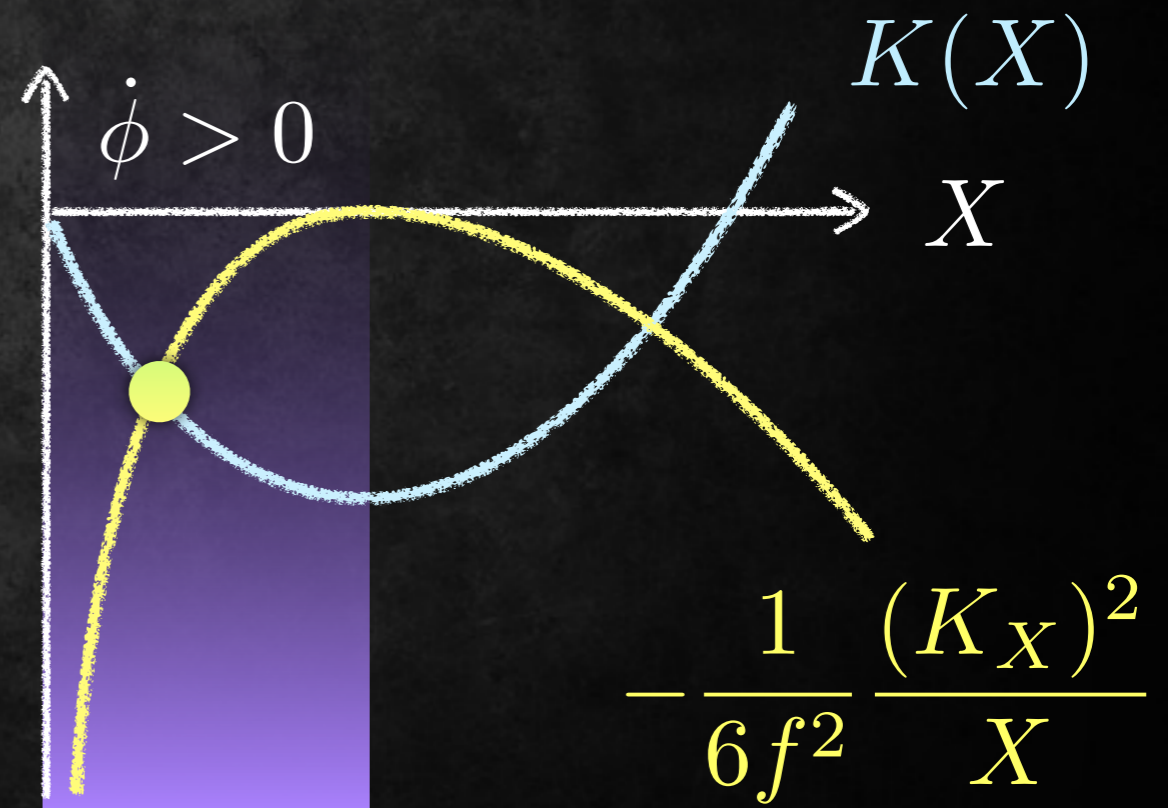
satisfying:

$$3H^2 = -K$$

$$K_X = -3fH\dot{\phi}$$



$$K = -\frac{1}{6f^2} \frac{(K_X)^2}{X}$$



Example

$$K = -X + \frac{X^2}{2M^3\mu}, \quad F = \frac{X}{M^3}$$



$$\mu \ll M_{\text{Pl}}$$

$$\frac{H^2}{M_{\text{Pl}}^2} \simeq \frac{1}{6} \frac{M^3}{M_{\text{Pl}}^3} \frac{\mu}{M_{\text{Pl}}}$$

$$X \simeq \left(1 - \frac{\sqrt{3}\mu}{M_{\text{Pl}}} \right) \mu M^3$$

Quasi-dS G-inflation

$$K = K(X), \quad F = f(\phi)X$$

Quasi-de Sitter solution:

Required to get $n_s - 1 \neq 0$

$$H = H(t), \quad \dot{\phi} = \dot{\phi}(t)$$



Small rate of change

satisfying:

$$\epsilon = -\frac{\dot{H}}{H^2} \ll 1$$

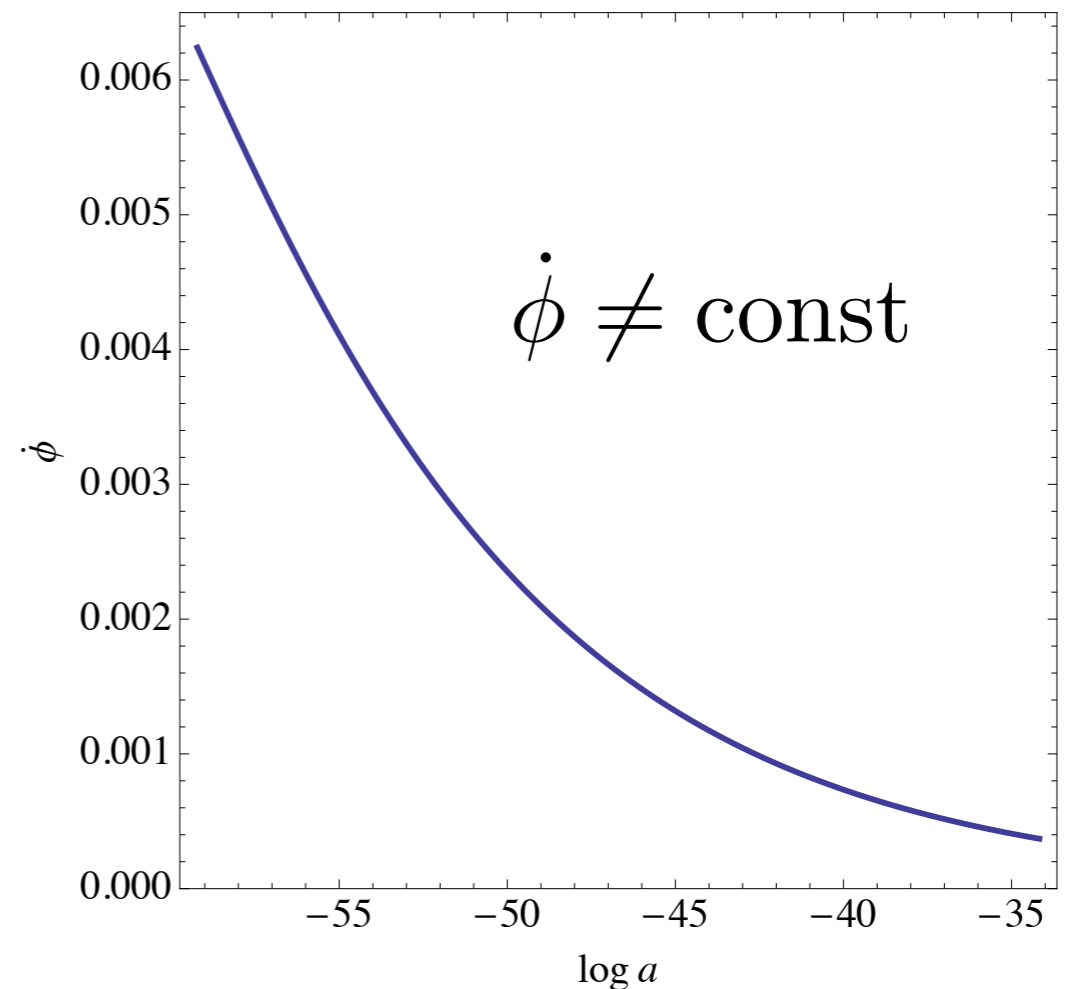
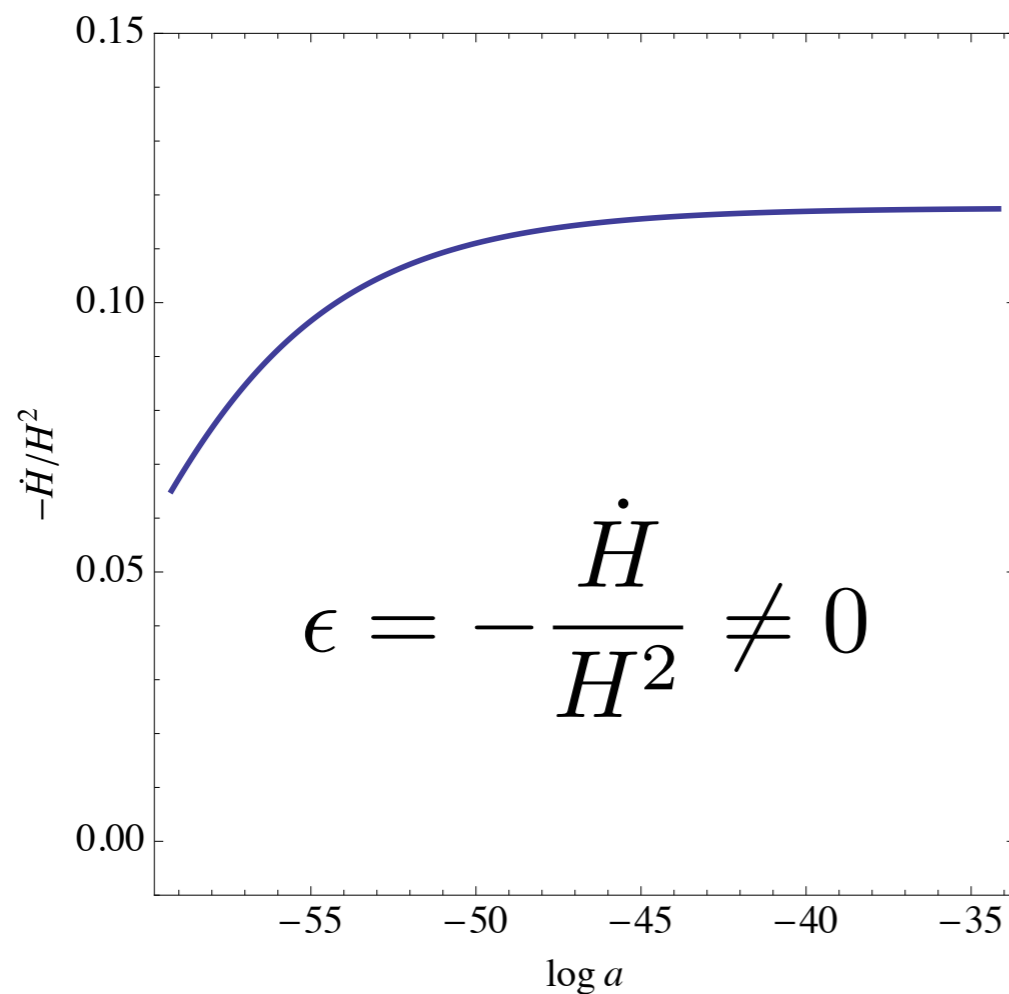
$$H^2 \simeq -K(X)$$

$$K_X \simeq -3f(\phi)H\dot{\phi}$$

$$\eta = -\frac{\ddot{\phi}}{H\dot{\phi}} \ll 1$$

Numerical Example

$$K = -X + \frac{X^2}{2M^3\mu}, \quad F = \frac{e^{\alpha\phi}}{M^3}X$$



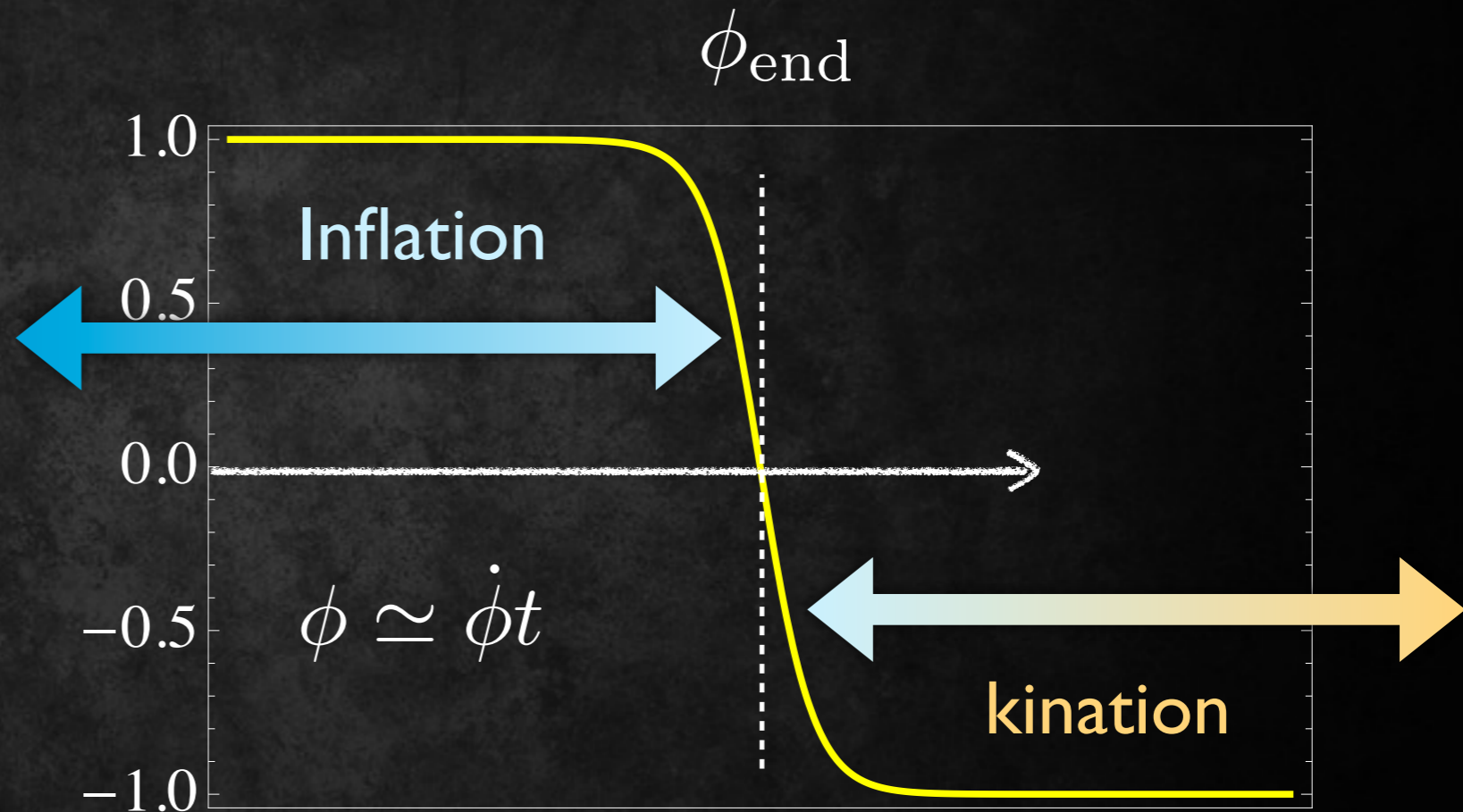
Graceful exit & Reheating

Basic idea

$$K = -A(\phi)X + \dots$$

Example:

$$A = \tanh[\lambda(\phi_{\text{end}} - \phi)]$$



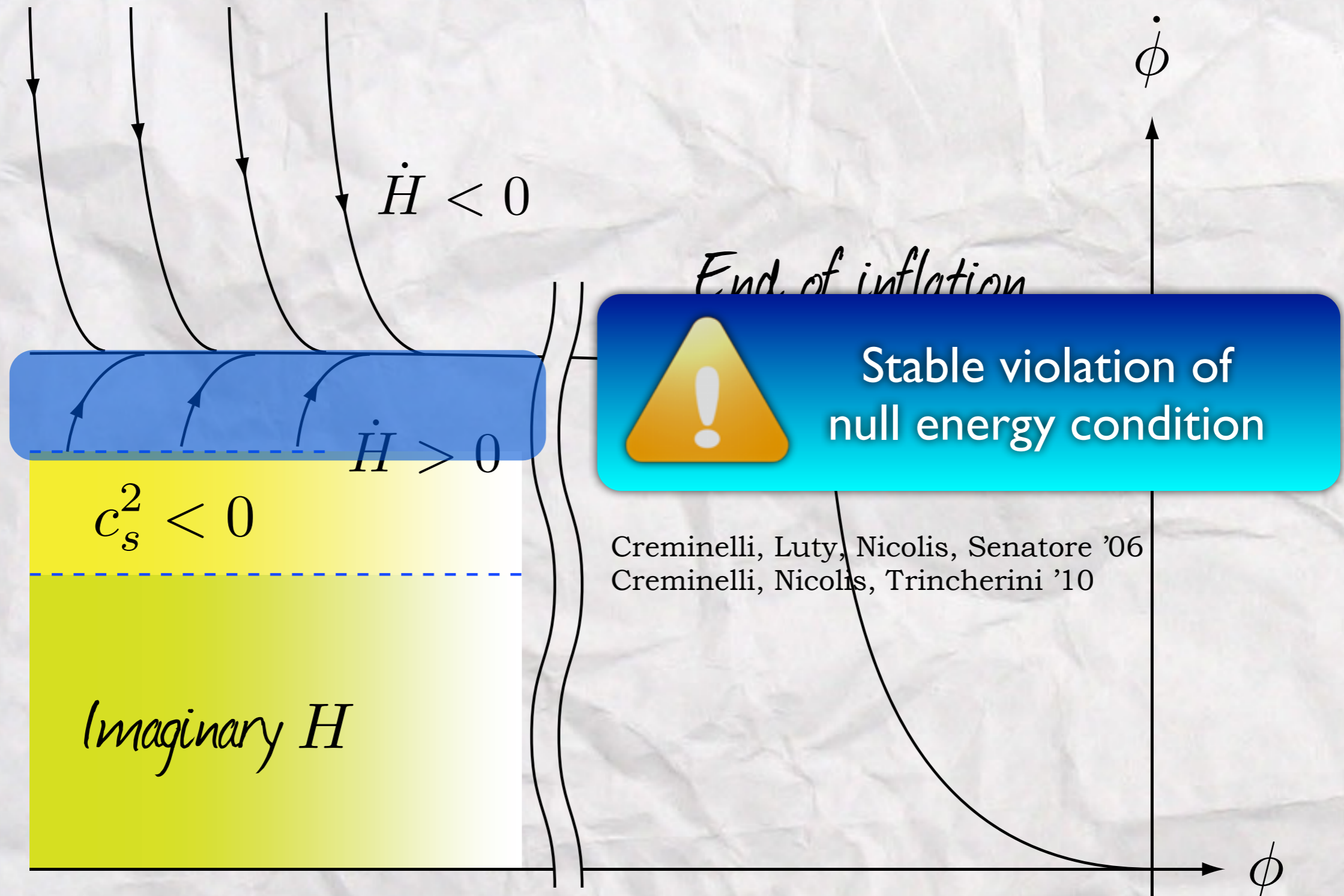
$$\rho \simeq p \simeq X \propto a^{-6}$$

Reheating through gravitational
particle production

*~ massless, canonical field
(normal sign)*

Ford '87

Phase diagram





Primordial perturbations

Cosmological perturbations

$$ds^2 = -(1 + 2\alpha)dt^2 + 2a^2 \beta_{,i} dt dx^i + a^2 (1 + 2\mathcal{R}) \delta_{ij} dx^i dx^j$$

$$\phi = \phi(t)$$

$$\text{Unitary gauge: } \delta\phi = 0$$

1. Expand the action to 2nd order
2. Eliminate α and β using constraint eqs
3. Quadratic action for \mathcal{R}



$$\delta T_i^0 = -F_X \dot{\phi}^3 \alpha_{,i}$$

Uniform ϕ hypersurfaces
 \neq comoving hypersurfaces

Quadratic action

$$S^{(2)} = \frac{1}{2} \int d\tau d^3x z^2 \left[\mathcal{G}(\mathcal{R}')^2 - \mathcal{F}(\vec{\nabla}\mathcal{R})^2 \right]$$

where

No ghost and gradient instabilities if

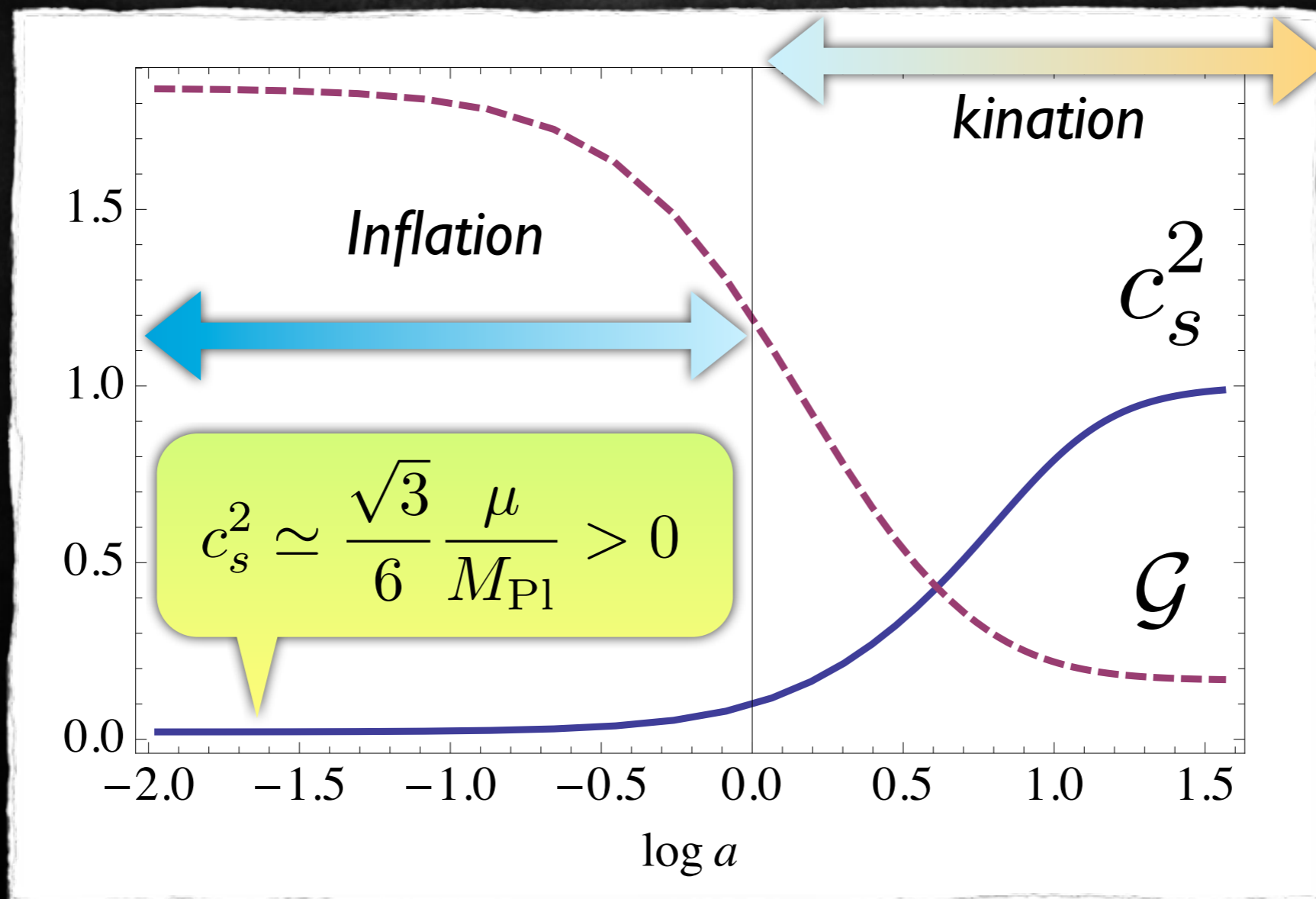
$$z \quad \mathcal{G} > 0, \quad c_s^2 = \mathcal{F}/\mathcal{G} > 0$$

$$\mathcal{F} = K_X + 2F_X \left(\ddot{\phi} + 2H\dot{\phi} \right) - 2F_X^2 X^2 \\ + 2F_{XX} X \ddot{\phi} - 2(F_\phi - X F_{\phi X})$$

$$\mathcal{G} = K_X + 2X K_{XX} + 6F_X H \dot{\phi} + 6F_X^2 X^2 \\ - 2(F_\phi + X F_{\phi X}) + 6F_{XX} H X \dot{\phi}$$

Stable example

$$K = -A(\phi)X + \frac{X^2}{2M^3\mu}, \quad F = \frac{X}{M^3}$$



Primordial spectrum

Consider G-inflation with:

$$K = K(X), \quad F = f(\phi)X$$

New variables:

$$\begin{aligned} dy &= c_s d\tau \\ \tilde{z} &= (\mathcal{F}\mathcal{G})^{1/4} z \\ u &= \tilde{z}\mathcal{R} \end{aligned}$$



$$\frac{\dot{c}_s}{H c_s} = \mathcal{O}(\epsilon)$$

$$\frac{d^2 u}{dy^2} + \left(k^2 - \frac{\tilde{z}_{,yy}}{\tilde{z}} \right) u = 0$$

$$\frac{\tilde{z}_{,yy}}{\tilde{z}} \simeq \frac{1}{(-y)^2} [2 + 3\epsilon\mathcal{C}(X)]$$

$$\mathcal{C}(X) = \frac{K}{K_X} \frac{Q_X}{Q}$$

$$Q(X) = \frac{(K - X K_X)^2}{18 X c_s^2 \sqrt{\mathcal{F}\mathcal{G}}}$$

Primordial spectrum

Normalized mode: $u = \frac{\sqrt{\pi}}{2} \sqrt{-y} H_{3/2+\epsilon\mathcal{C}}^{(1)}(-ky)$



$$\mathcal{P}_{\mathcal{R}} = \frac{Q}{4\pi^2} \Big|_{c_s k=1/(-\tau)}, \quad n_s - 1 = -2\epsilon\mathcal{C}$$

$\propto f, \phi$

\mathcal{R} can be generated even from exact de Sitter

where

$$Q(X) = \frac{(K - XK_X)^2}{18M_{\text{Pl}}^4 X c_s^2 \sqrt{\mathcal{F}\mathcal{G}}}$$

* Tensor mode dynamics: unchanged

Tensor-to-scalar ratio

$$K = -X + \frac{X^2}{2M^3\mu}, \quad F = fX \quad \longrightarrow \quad H^2 \sim \frac{\mu M^3}{M_{\text{Pl}}^2}$$

$(f \simeq M^{-3})$

$$r \simeq \frac{16\sqrt{6}}{3} \left(\frac{\sqrt{3}\mu}{M_{\text{Pl}}} \right)^{3/2}$$

Standard consistency relation is violated

$$r \neq 16\epsilon$$

$$\uparrow \propto f, \phi$$

$$M = 0.00435 \times M_{\text{Pl}}, \quad \mu = 0.032 \times M_{\text{Pl}}$$

$$\longrightarrow \mathcal{P}_{\mathcal{R}} = 2.4 \times 10^{-9}, \quad r = 0.17$$

r can be large!

Definition:

$$r = \frac{\mathcal{P}_T}{\mathcal{P}_{\mathcal{R}}}, \quad \mathcal{P}_T = \frac{8}{M_{\text{Pl}}^2} \left(\frac{H}{2\pi} \right)^2$$



Summary & Outlook

Summary & Outlook

- **G-inflation**: A novel class of inflation models

$$\mathcal{L}_\phi = K(\phi, X) - F(\phi, X)\square\phi$$

- $n_s - 1 \simeq 0$
- Large r
- ~~Consistency relation~~

- **Non-Gaussianity?**

TK, Yamaguchi, Yokoyama *in progress*



A scenic landscape photograph featuring a bright blue sky filled with large, white, fluffy clouds. In the foreground, a snow-covered field is visible, bordered by a wooden fence with wire mesh. The background shows a silhouette of a residential area with houses and trees under a clear sky. The text "Thank you!" is overlaid in a large, bold, orange font across the center of the image.

Thank you!