

# 複雑流体の非平衡輸送現象: 分子から細胞集団へ

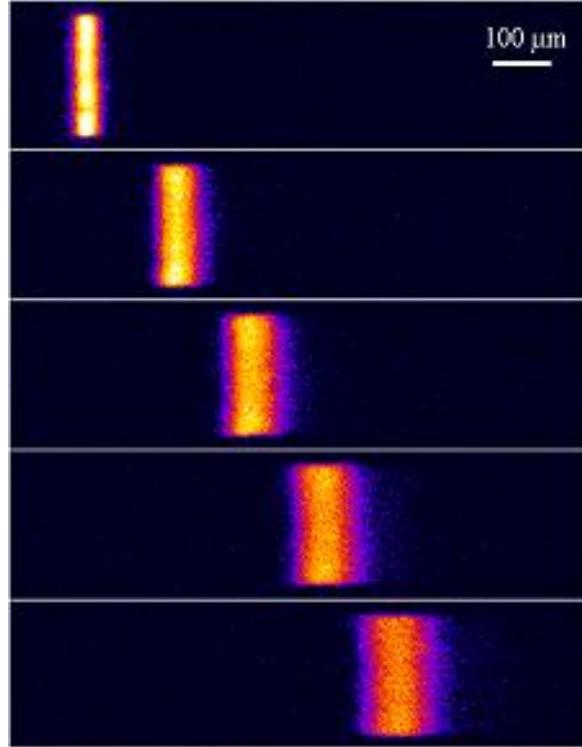
前多裕介<sup>1,2,3</sup> and Albert Libchaber<sup>3</sup>

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<sup>3</sup> The Rockefeller University

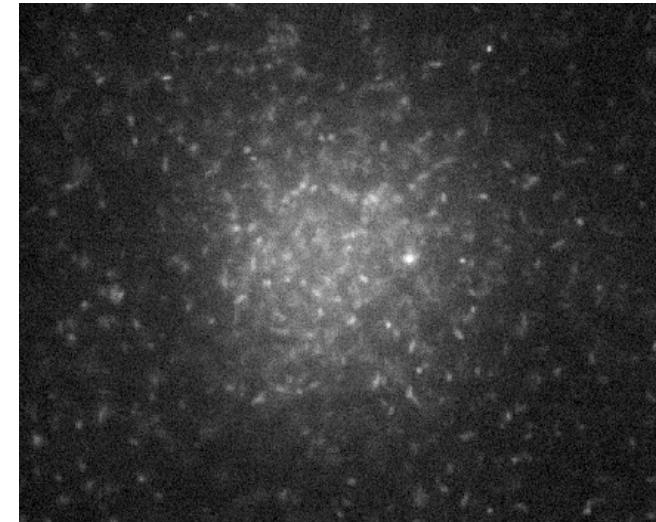
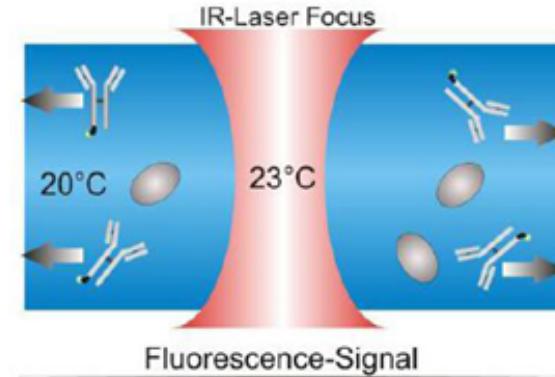
# Transport and Directed motions in microenvironments



Electro-osmotic flow

Directed motions in a gradient of thermodynamic variables

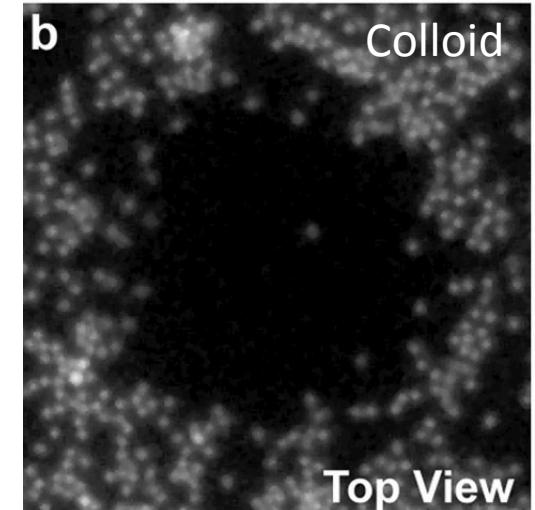
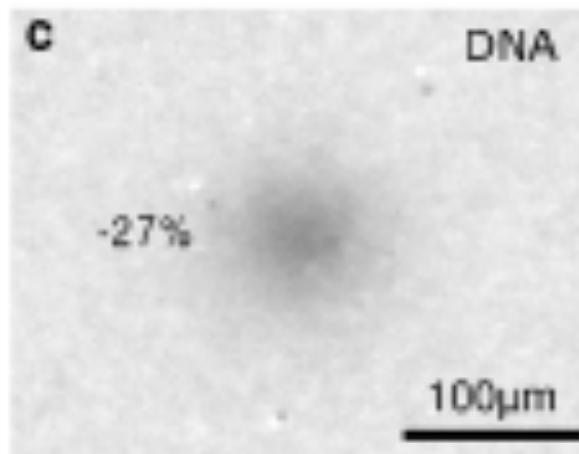
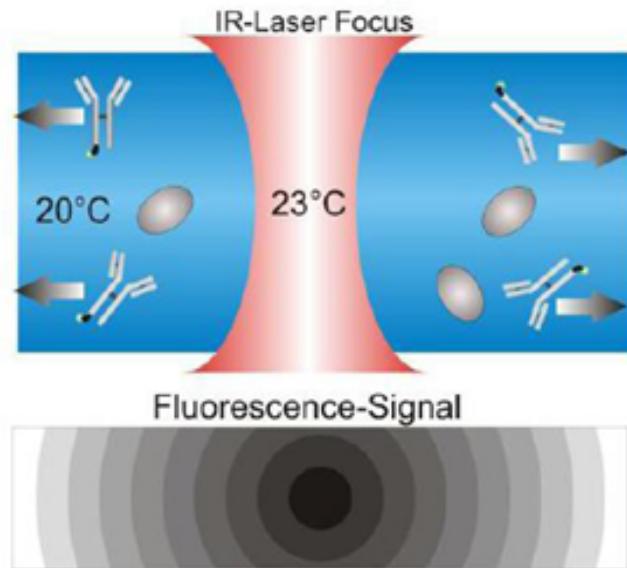
Thermophoresis



# The Soret effect 1/3

**The Ludwig-Soret effect:** a solute move along a temperature gradient

C. Ludwig, S-B Akad. Wiss. Wien. *Nature* **20**, 539 (1856);  
C. Soret. *Arch. Geneve.* **3**, 48 (1879)

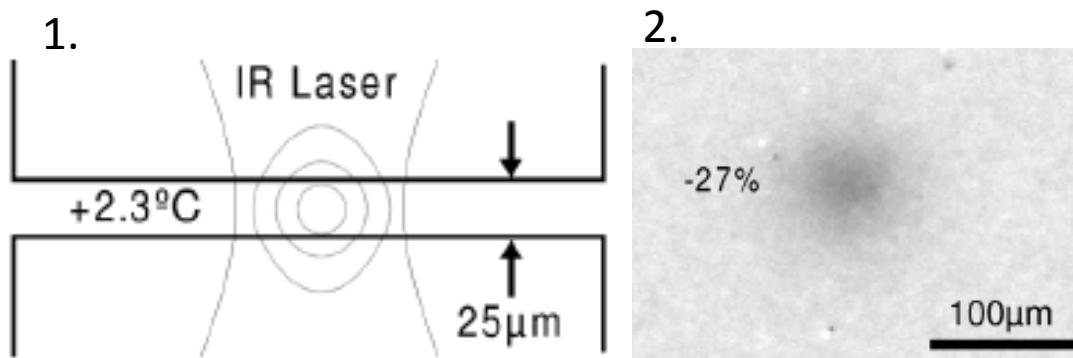


D Braun and A Libchaber. *Phys. Rev. Lett.* **89**, 188103 (2002)  
S Duhr and D Braun. *Phys. Rev. Lett.* **101** 168301 (2008)

Charged molecules move along a temperature gradient

From **HOT** to **COLD**

# The Soret effect 2/3



1. Local heating by infrared light focusing
2. Depletion of DNA from the hot region

$$J = -D \nabla c - c D_T \nabla T$$

Diffusion    Thermal diffusion

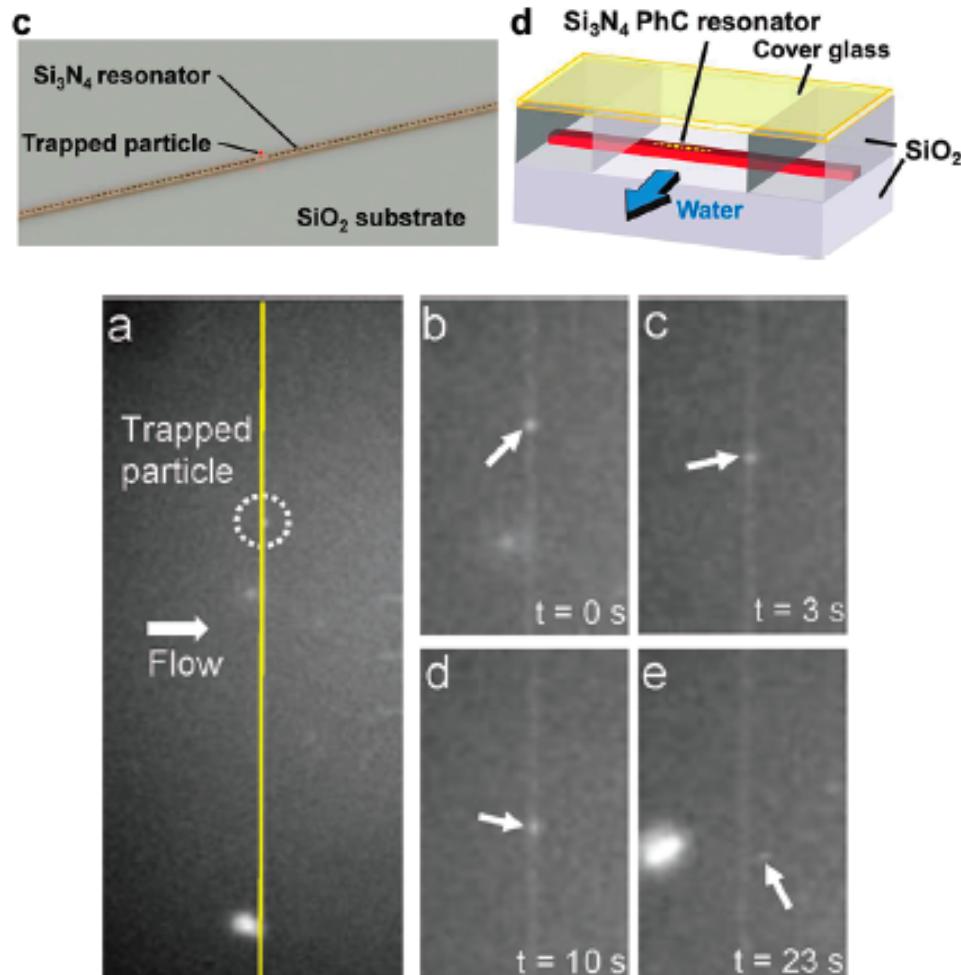
Distribution at the steady state

$$c(z) = c_0 \exp\left[-S_T(T(z) - T_0)\right]$$

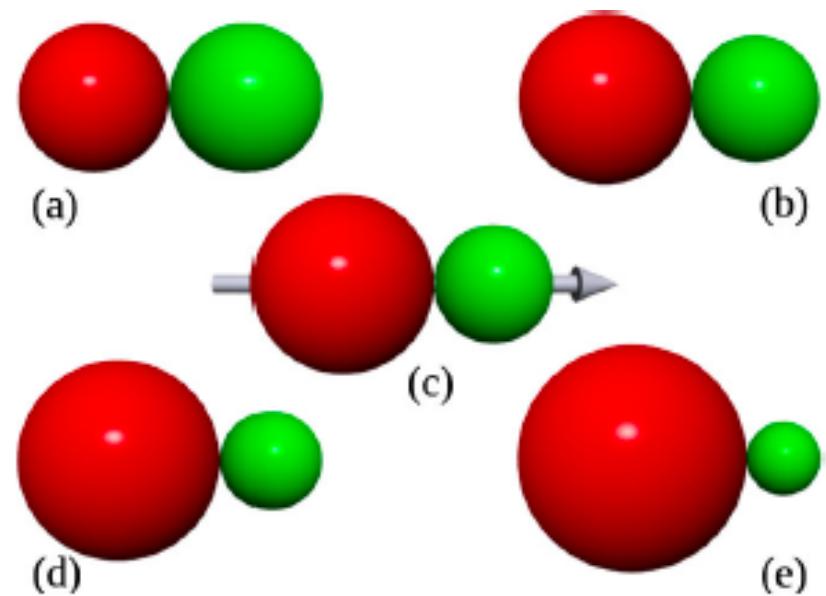
Dieter Braun, Albert Libchaber

*Physical Review Letters* **89**: 188103 (2002)

The Soret coefficient     $S_T = D_T / D$



## Alignment of anisotropic molecules

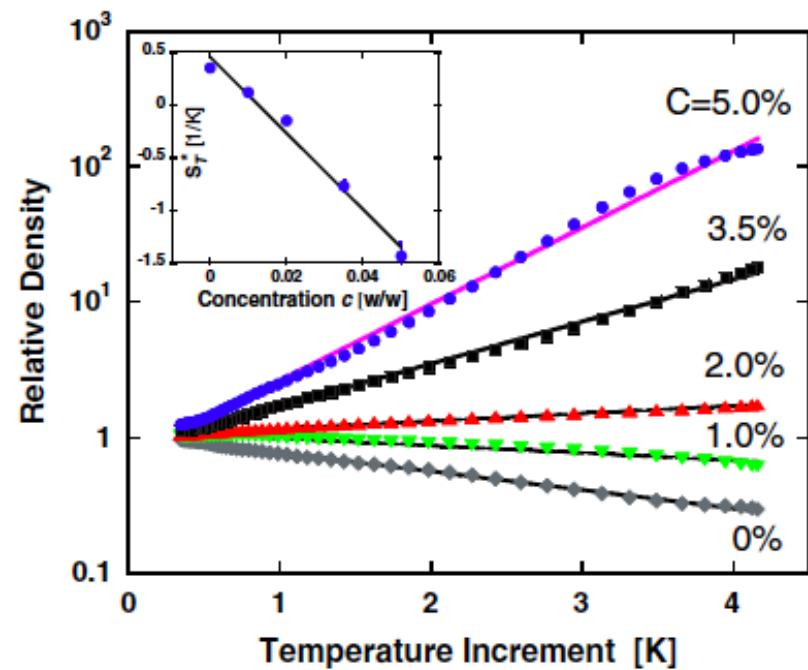
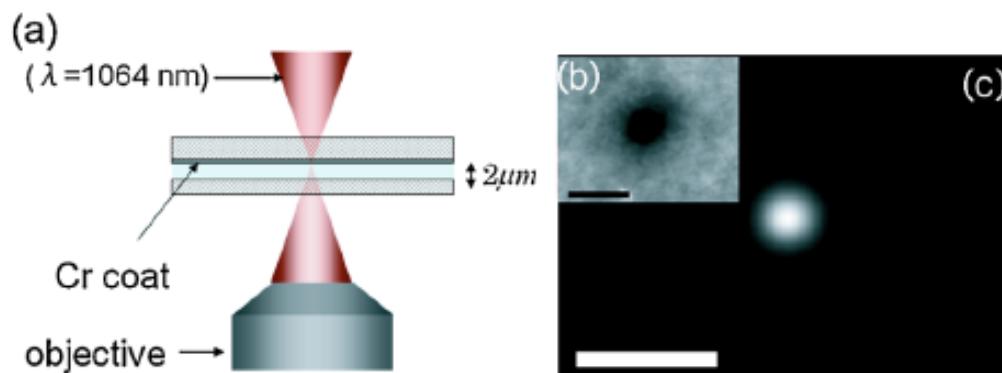


YF Chen et al.  
*Nano Letters* **12**: 1633 (2012)

F Romer et al  
*Physical Review Letters* **108**: 105901 (2012)

# The Soret effect 3/3

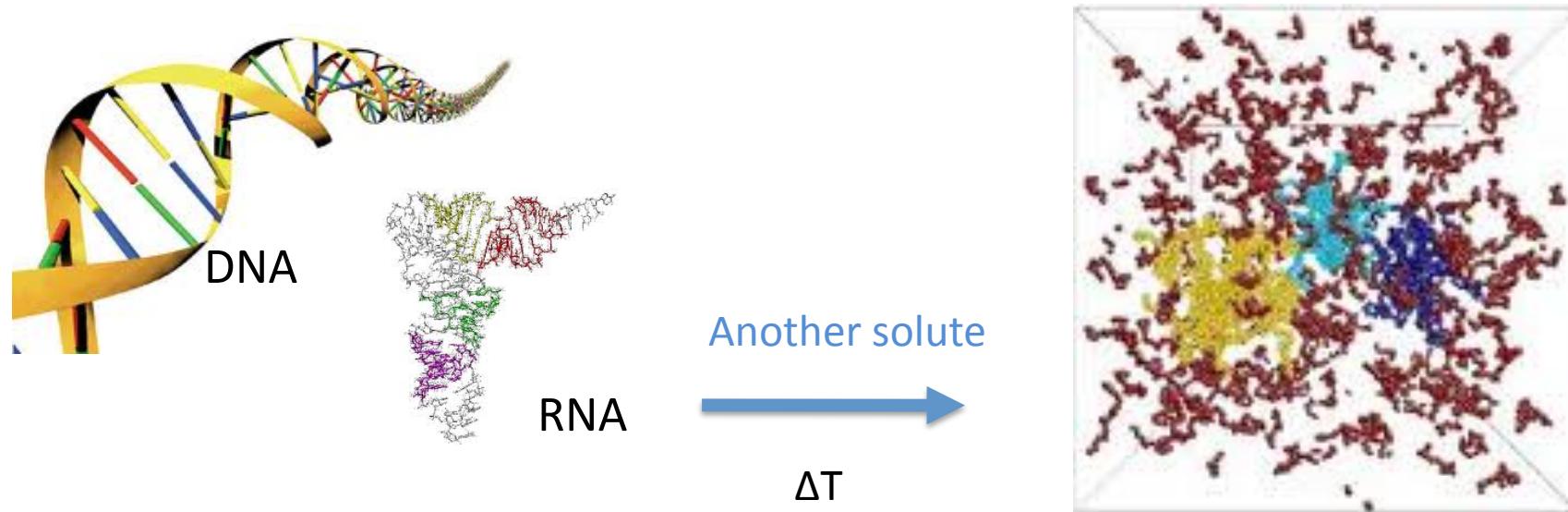
- Thermophoresis of colloidal particles in a polymer solution
- Direction of colloidal transport is flipped.  
→ Little material dependence



H.R.Jiang, H.Wada, N.Yoshinaga, M.Sano  
*Physical Review Letters* **102**: 208301 (2009)

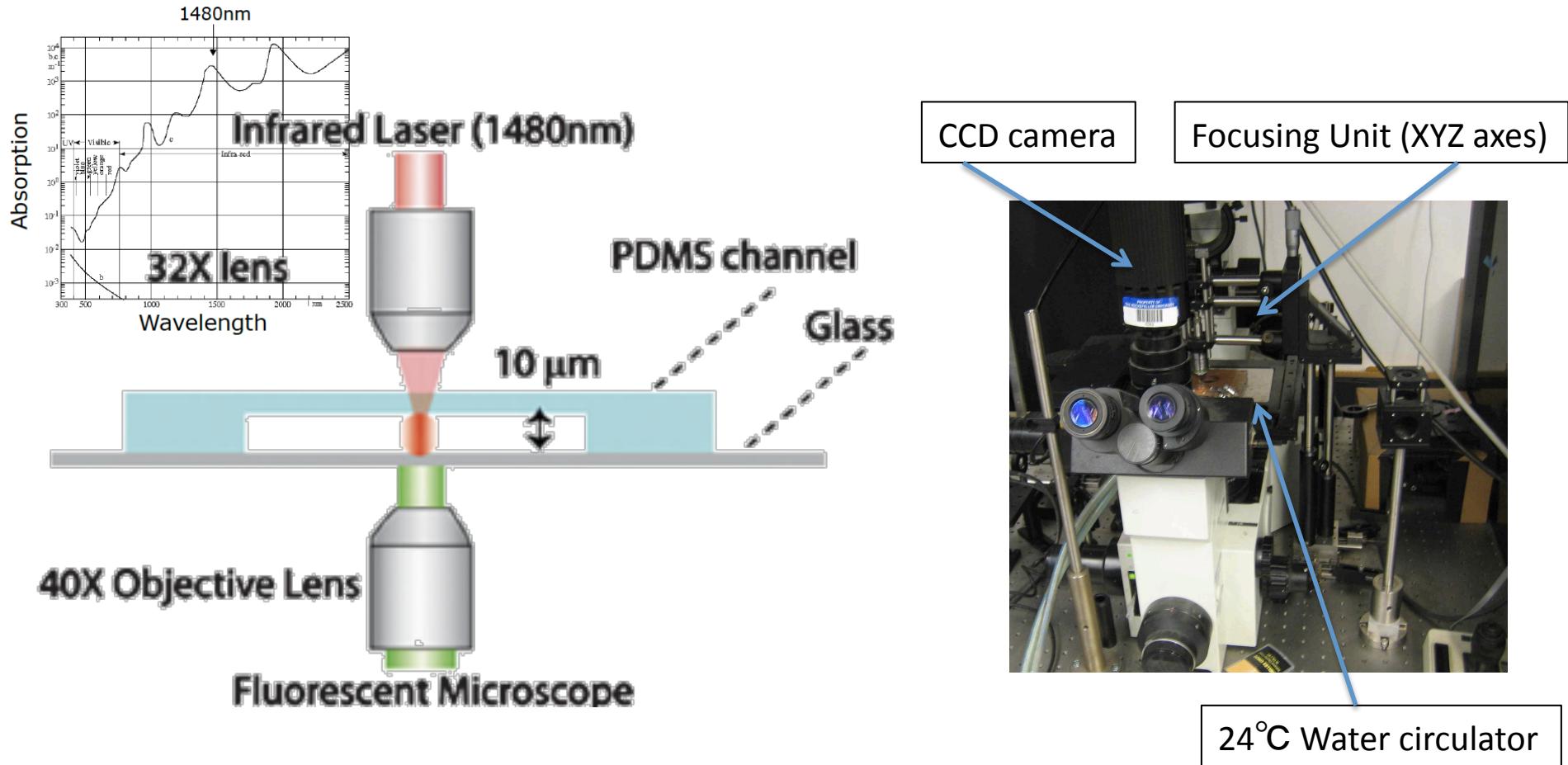
# Focus in this study

## Complex fluids in a temperature gradient



1. Thermophoresis of complex fluids and molecular separation
2. The effects of molecular folding on thermophoresis
3. The origin of the entropic force

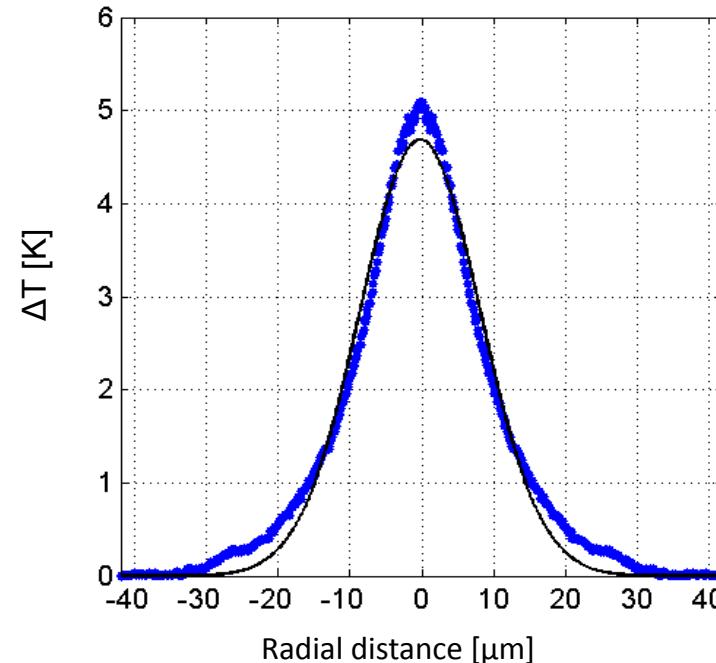
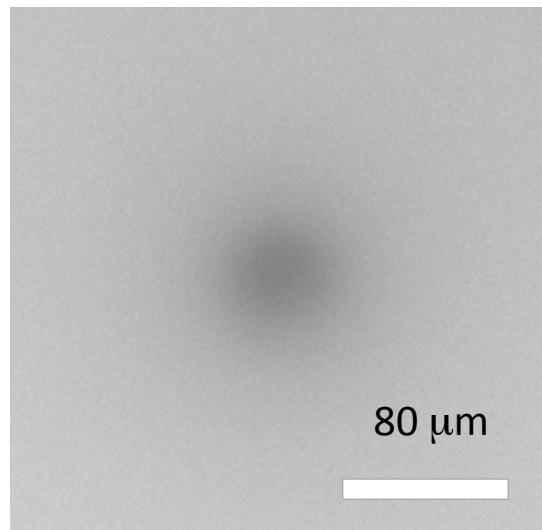
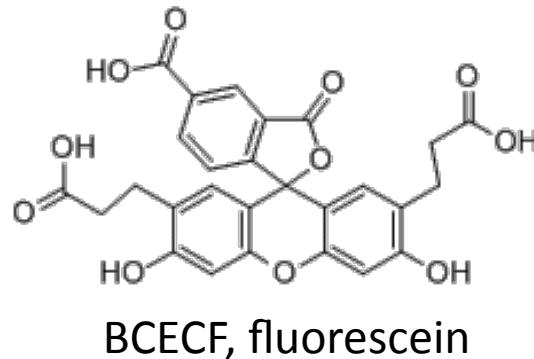
# 1. Thermophoresis of complex fluids and molecular separation

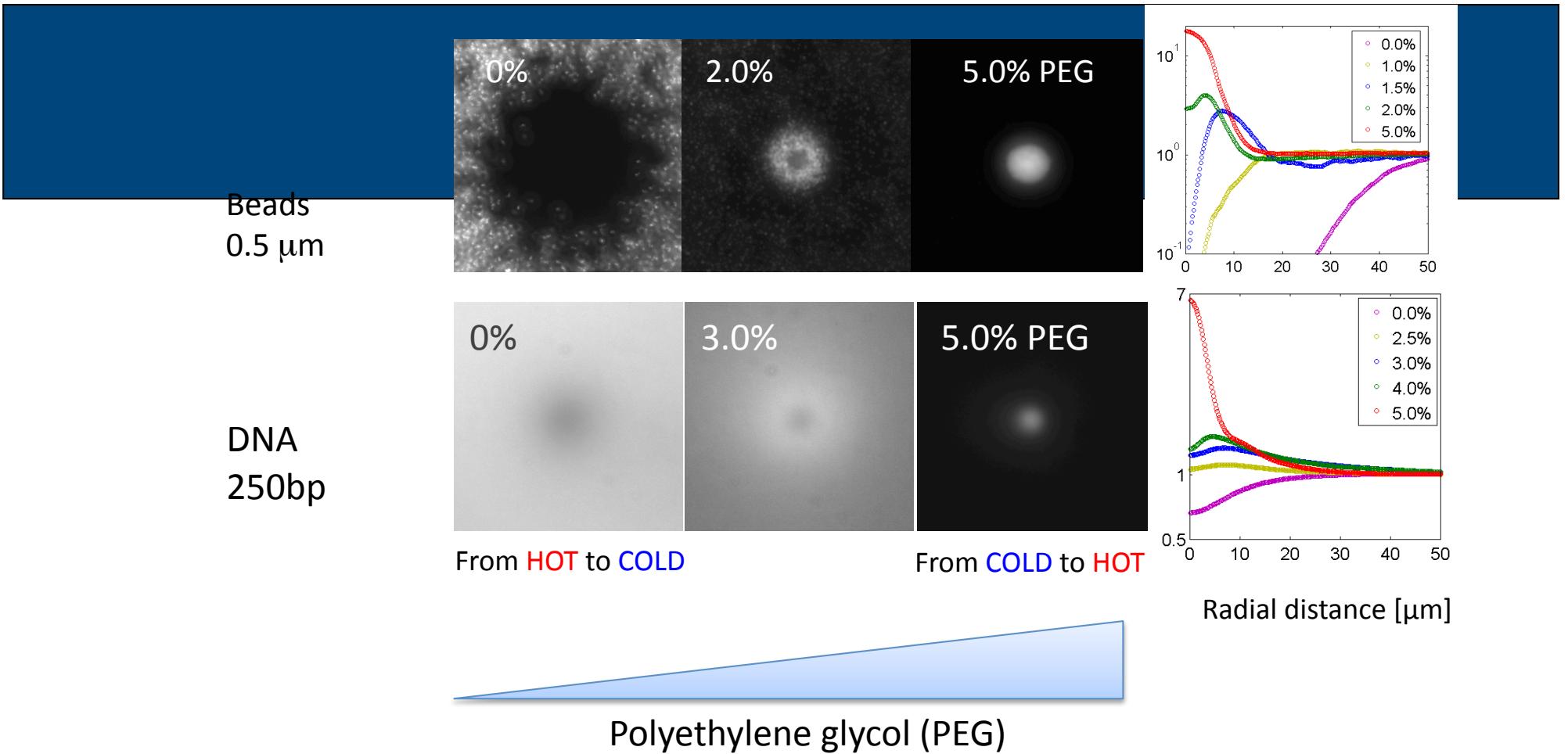


Polymer PEG10,000 (polyethylene glycol 10,000) solution of 0-5% volume fraction is enclosed in a thin chamber.

# Temperature gradient in experiments

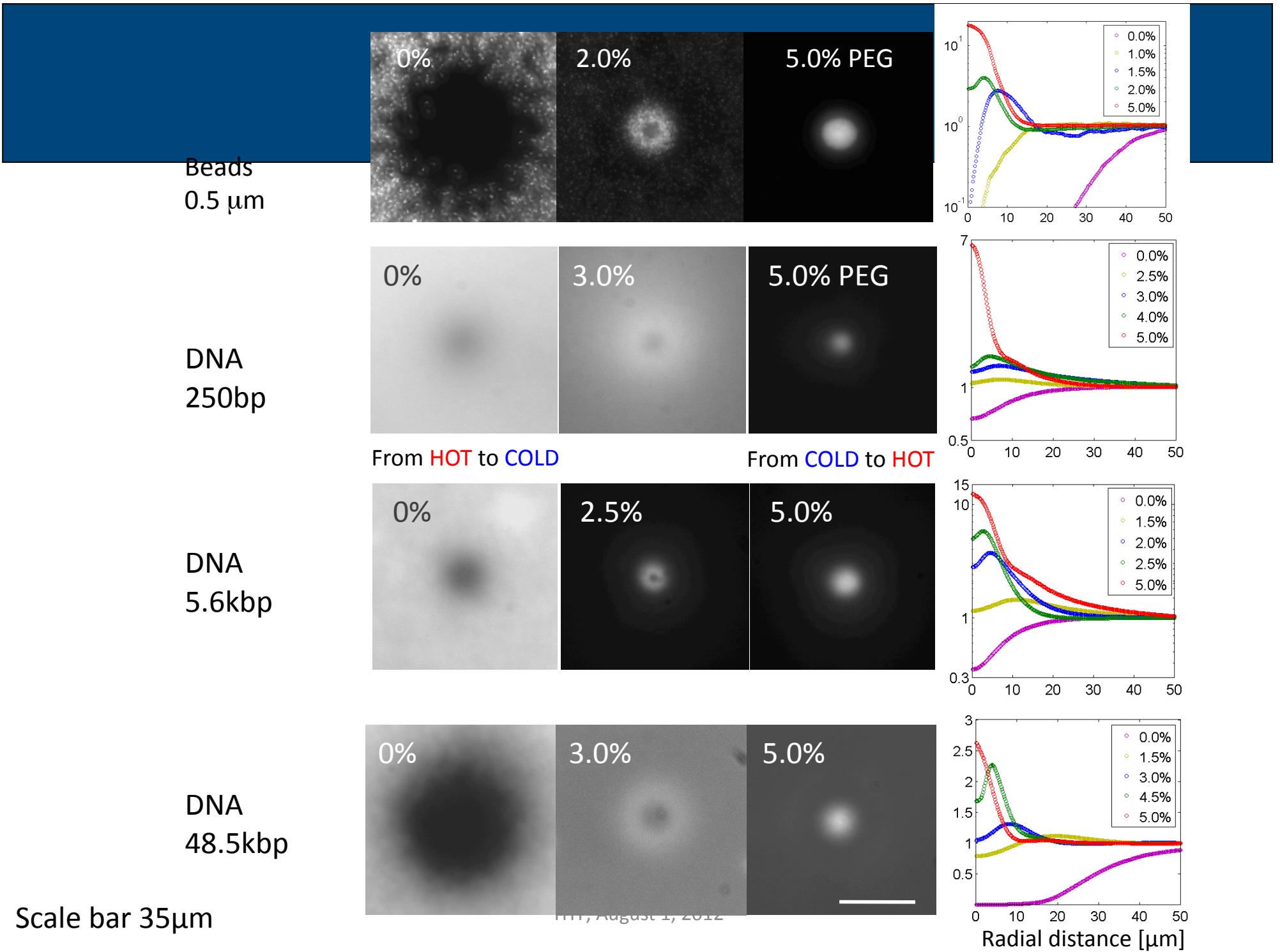
Temperature visualization using the intensity drop of fluorescent dye





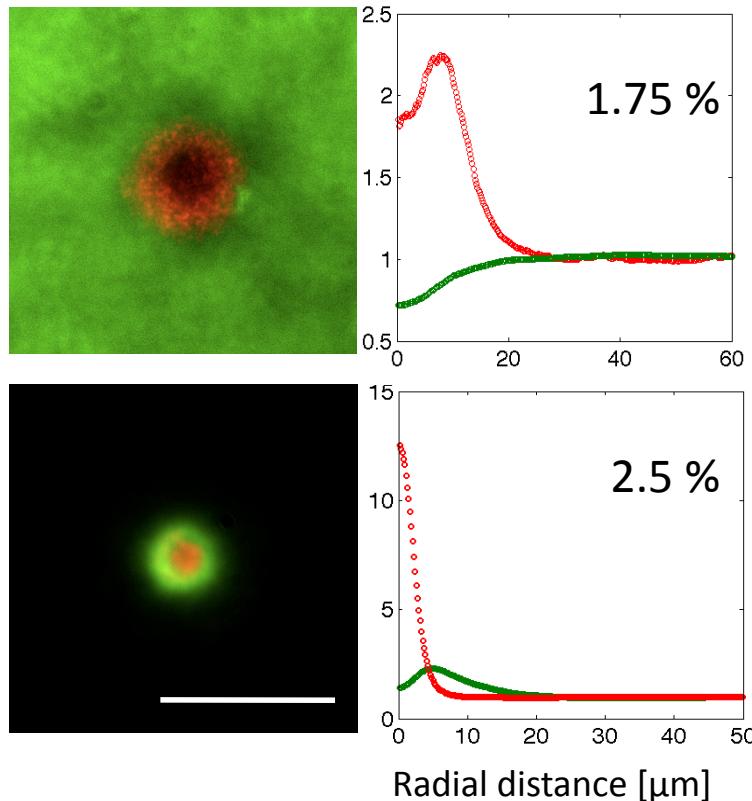
Scale bar 35μm

YITP, August 1, 2012

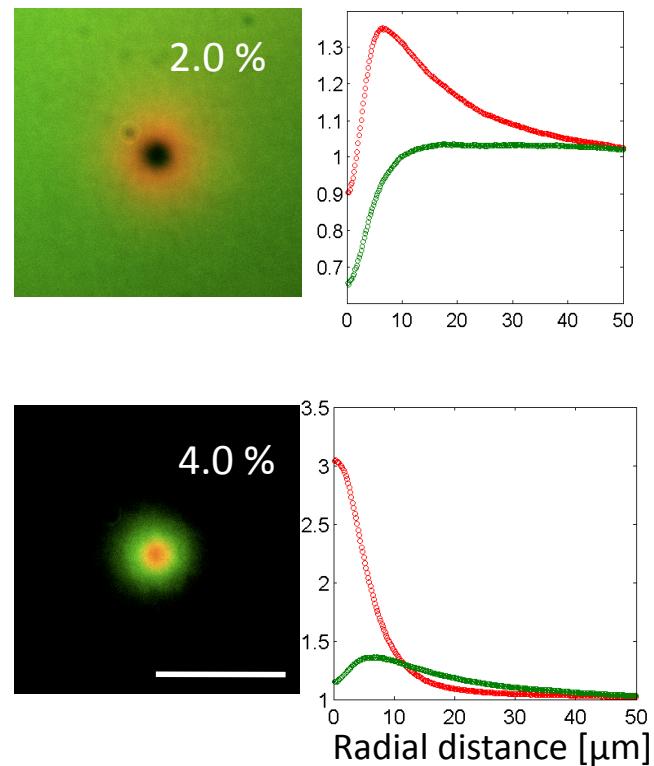


# Size-dependent separation

Beads  $0.5\mu\text{m}$   $0.1\mu\text{m}$



RNA(1.5kb) DNA(250bp)



Yusuke T. Maeda, Axel Buguin, Albert Libchaber

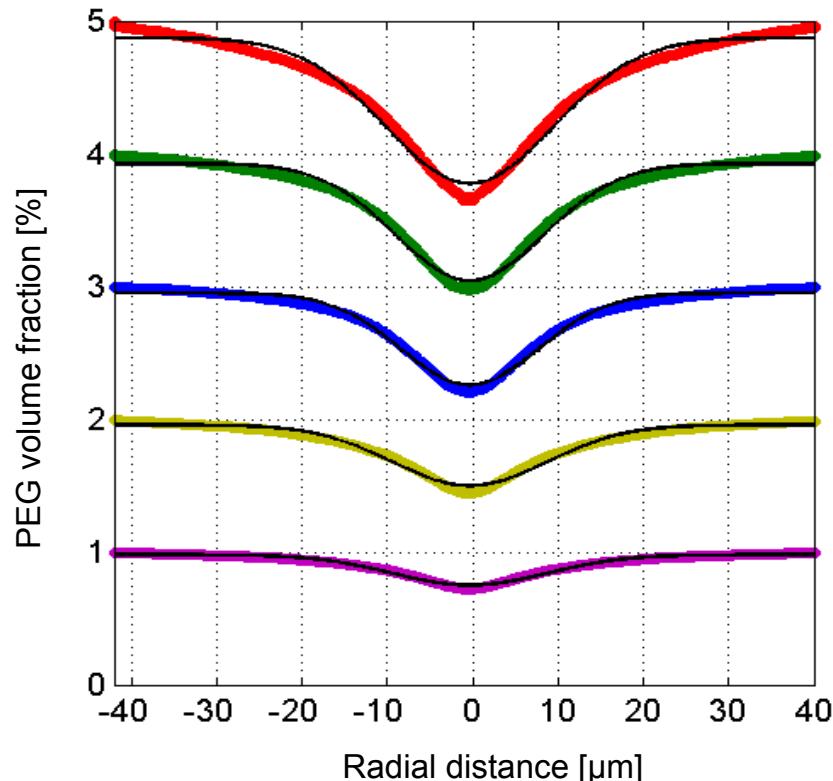
*Physical Review Letters* **107**: 038301 (2011)

P, August 1, 2012

Scale bar  $35\mu\text{m}$

# The mechanism of localization (1/4)

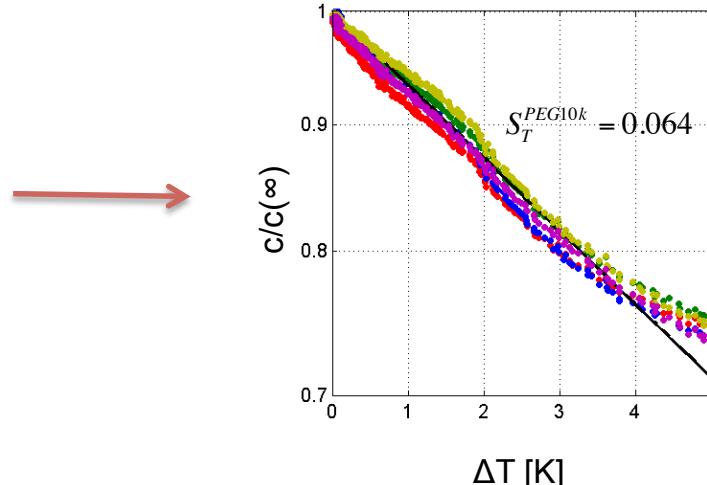
PEG molecules are also subject to thermophoresis



$$J = -D \nabla c - c D_T \nabla T$$

Diffusion      Thermal diffusion

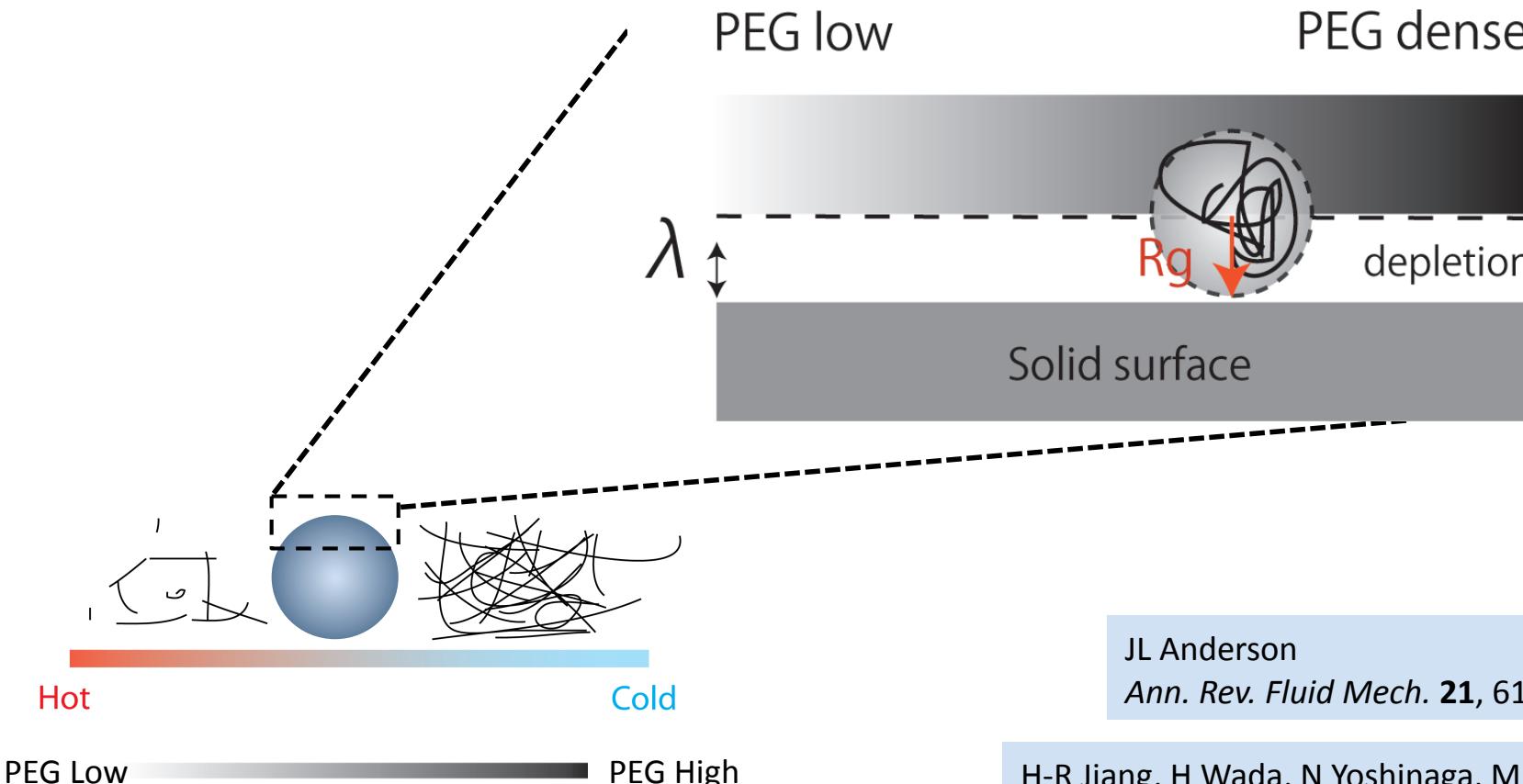
$$J = 0 \Rightarrow c(r) = c(\infty) \exp\left(-\frac{D_T}{D} \Delta T\right)$$



$$S_T = D_T / D : \text{Soret coefficient}$$

# The mechanism of localization (2/4)

$$c_{PEG}(r) = c_0 \exp\left[-S_T^{PEG}(T(r) - T_0)\right]$$



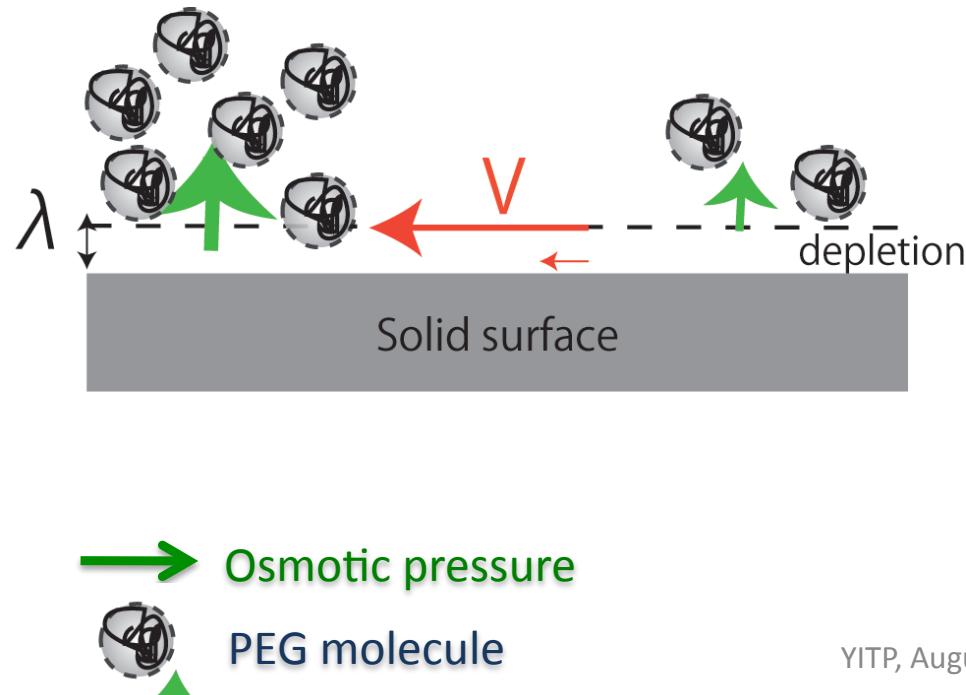
JL Anderson  
*Ann. Rev. Fluid Mech.* **21**, 61 (1989)

H-R Jiang, H Wada, N Yoshinaga, M Sano  
*Physical Review Letters* **102**, 208301 (2009)

# The mechanism of localization (3/4)

Depletion of PEG from the surface of a solute

$$c(x,z) \approx c_0(x) \exp\left[-\frac{U(z)}{k_B T}\right]$$



- Balance of osmotic pressure and hydrostatic Pressure (No net force on a fluid in the bulk)

$$p(x,z) - k_B T c(x,z) = p_0 - k_B T c_0(x) = \text{constant}$$

$$\partial_x p(x,z) = k_B T \partial_x [c(x,z) - c_0(x)]$$

- The diffuse-osmotic flow is balanced by viscous shear stress at equilibrium

$$\partial_x p(x,z) = \eta \partial_x^2 v_x$$

- The fluid velocity increase in the bulk is

$$V \approx \int_0^z dz' (e^{-U(z')/kT} - 1) \frac{dc_0}{dx} \approx -\frac{k_B T}{\eta} \lambda^2 \nabla c_0$$

Drop of solute

# The mechanism of localization (4/4)

$$J = -D\nabla c^b - c^b D_T \nabla T + c^b u,$$

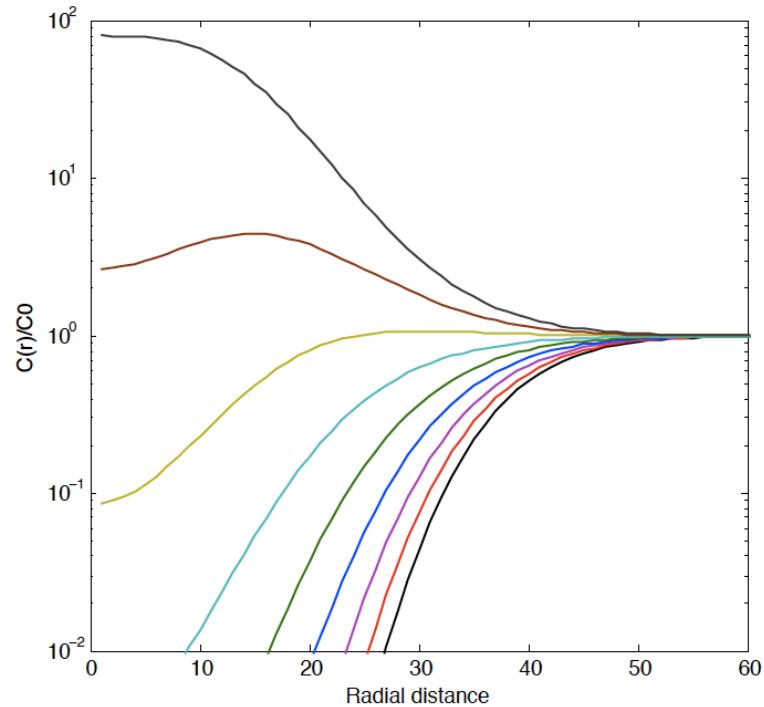
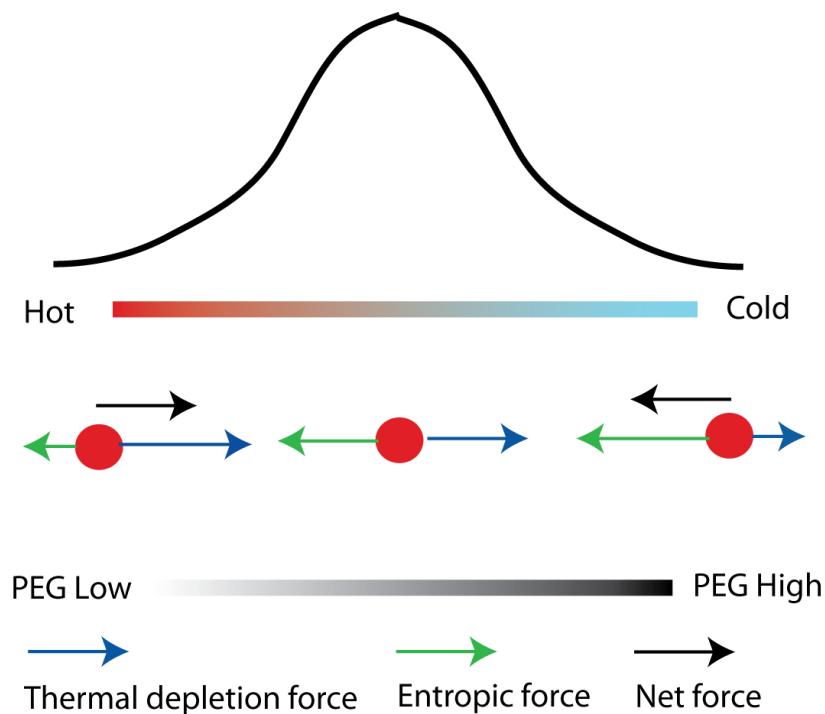
### No slip condition

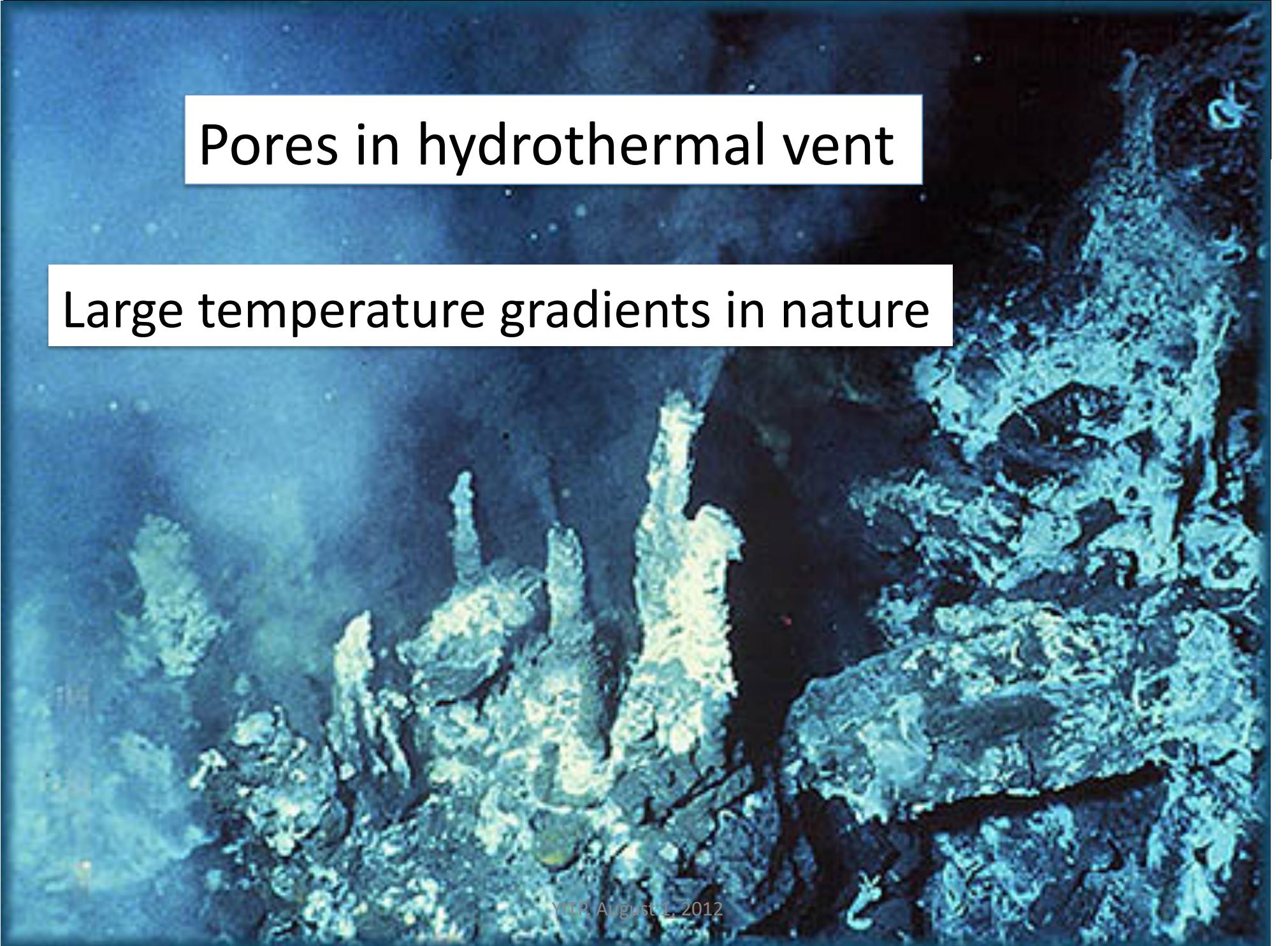
$$u = \frac{k_B T}{3\eta} \left( S_T^{PEG} - \frac{1}{T} \right) \lambda^2 c^{PEG}(r) \nabla T(r)$$

$$c^b(r) = c_0^b \exp \left[ -S_T^b \Delta T(r) + \left( c_0^{PEG} - c^{PEG}(r) \right) V \right]$$

Soret effect      Entropic force

$$V = \pi \lambda^2 d$$





Pores in hydrothermal vent

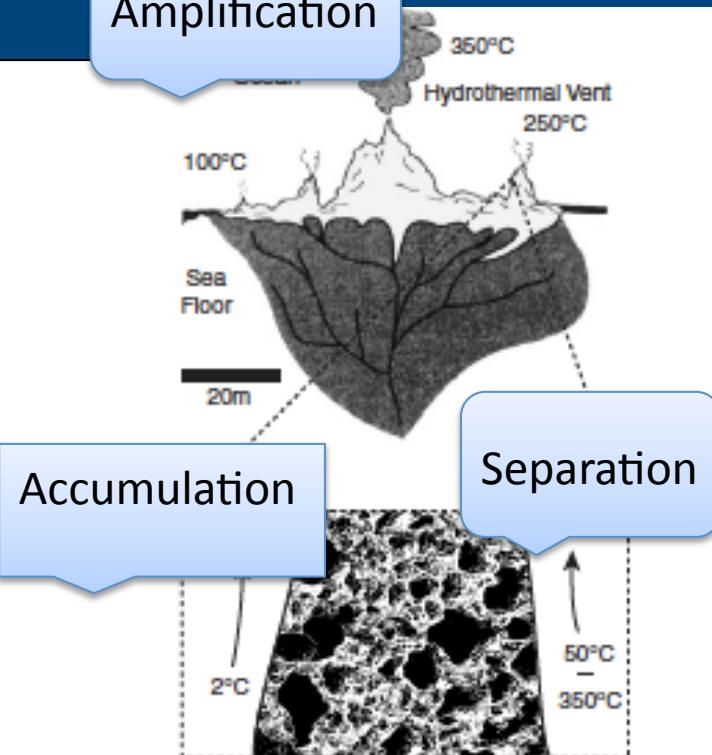
Large temperature gradients in nature

Amplification

4000,000,000 BC

Accumulation

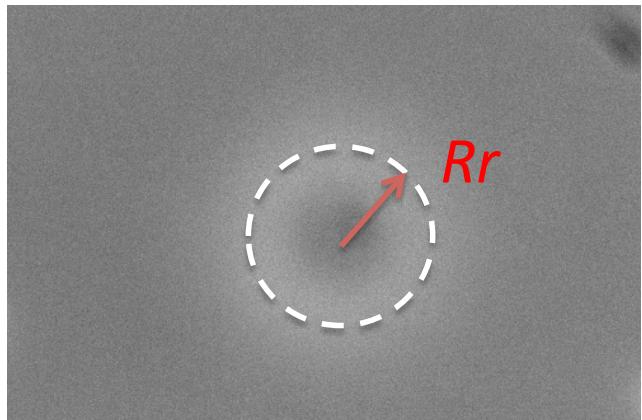
Separation



2,000 AD



## 2. The effect of molecular folding on thermophoresis



非公開

Below 5.6 kbp

Ring shrinks as DNA becomes longer

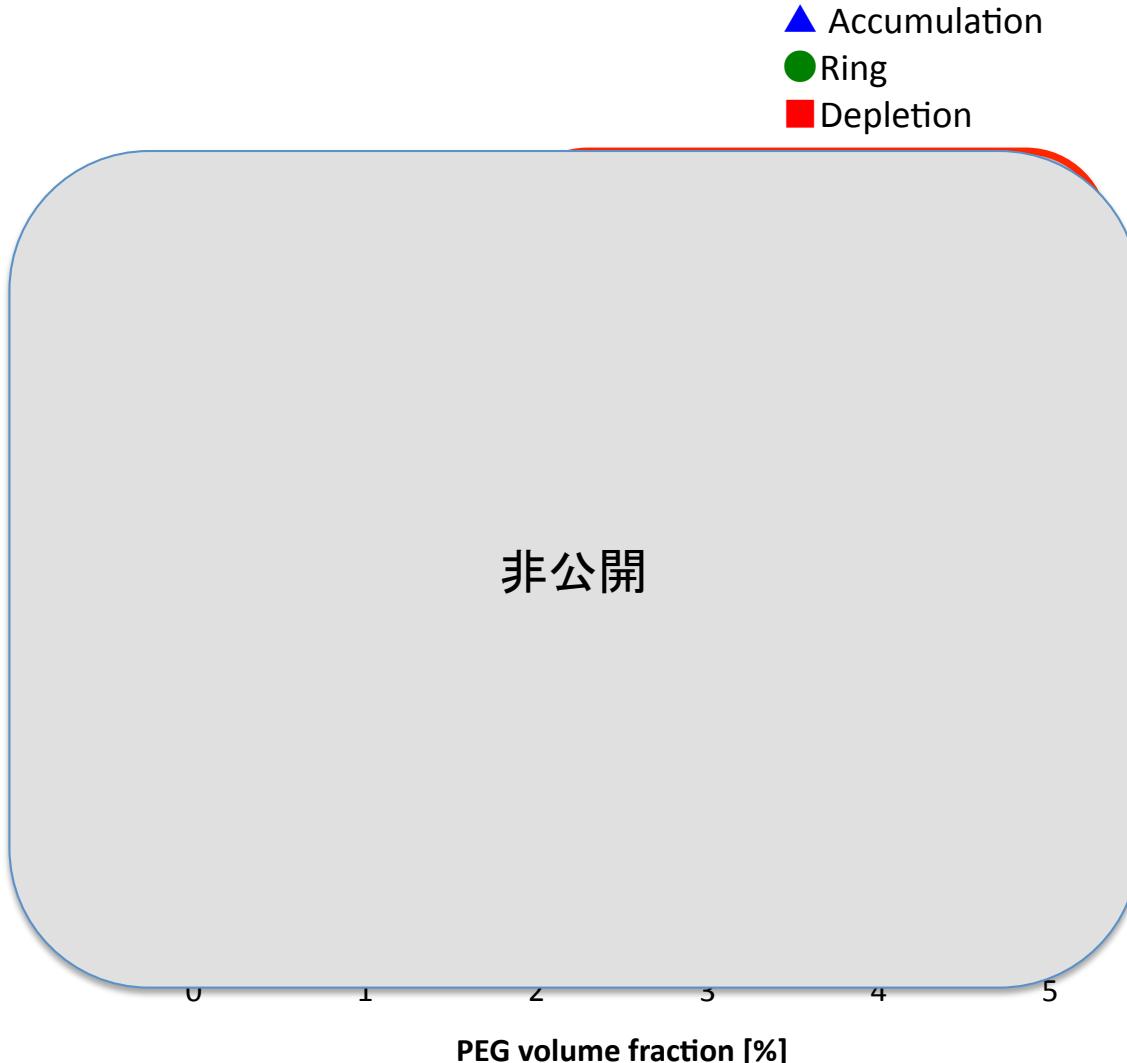
Above 5.6 kbp

Ring expands as DNA becomes longer

Behavior is reversed at around 5.6 kbp

DNA size [kbp]

# Reversed size-dependence above 5.6kbp



150 bp = 1 persistence length  
5 kbp > 20 persistence lengths

Is DNA molecule expanded coil?

# Coil-globule transition occurs in a PEG

非公開

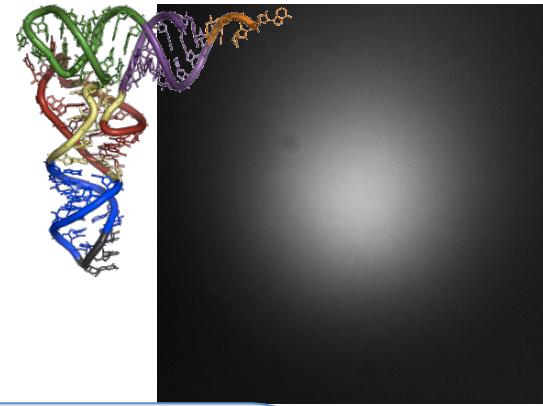
# Single molecule observation

非公開

# The effect of stem-loop on RNA transport

Next question: Why the force from a PEG gradient can act on small RNA?

tRNA ( $\sim 5$  nm) can be accumulated in a PEG ( $\sim 3$  nm) solution.



非公開

# Folding dependent accumulation

非公開

- The accumulation of small RNA and DNA is folding-dependent.

# Rigid part is essential

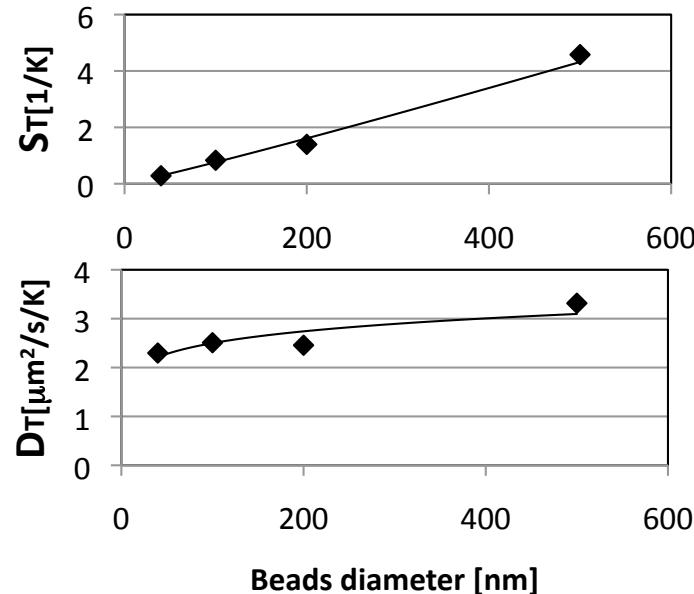
非公開

- The stem (double-stranded, rigid) of RNA and DNA strongly enhances accumulation while the loop (single-stranded, flexible) has few effect.

### 3. The origin of entropic force: diffusiophoresis

Thermophoresis

$$V_T = -D_T \nabla T(r)$$



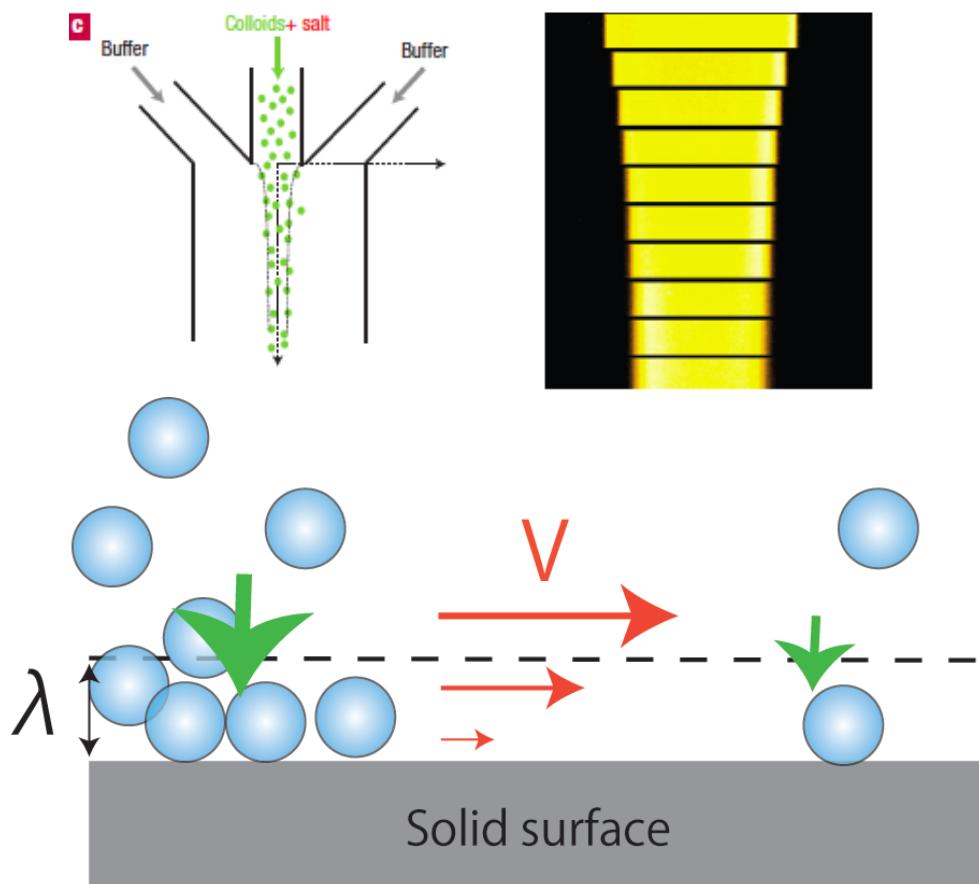
Entropic force (Diffusiophoresis)

$$V = ??$$

- $D_T$  has little size-dependence and it cannot account for observed separation.
- Size-dependence should present in entropic force.

$$S_T \sim a^{1.1} \quad D_T \sim a^{0.1}$$

# Salt and migrate: Diffusiophoresis



→ Osmotic pressure



Sakt

- The fluid velocity increase in the bulk is

$$V = -\frac{k_B T}{\eta} \Gamma L \nabla_x c_0$$

Excess of solute  $\Gamma = \int_0^\infty \left( e^{-\frac{U(z)}{k_B T}} - 1 \right) dz$

The range of attraction  $L = \Gamma^{-1} \int_0^\infty z \left( e^{-\frac{U(z)}{k_B T}} - 1 \right) dz$

$$V \approx -\frac{k_B T}{\eta} \lambda^2 \nabla_x c_0 \approx -D_{DP} \nabla_x \log c_0$$

$$\lambda^2 = \varepsilon k T / 2q^2 c_0 \quad D_{DP} = \varepsilon k^2 T^2 / 2q^2 \eta$$

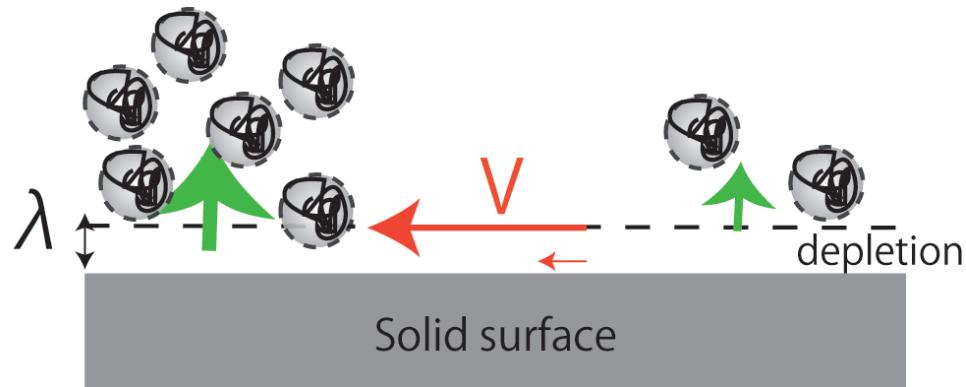
B Abecassis, et al  
*Nature Materials* 7, 785 (2008)

# PEG and Diffusiophoresis

Repulsion of PEG in the surface of a solute

$$c(x,z) \approx c_0(x) \exp\left[-\frac{U(z)}{k_B T}\right]$$

$$U(z) > 0$$



- The fluid velocity increase in the bulk is

$$V \approx -\frac{k_B T}{3\eta} \lambda^2 \nabla c_{PEG}$$

Diffusiophoresis velocity in no-slip condition is independent of the particle size.

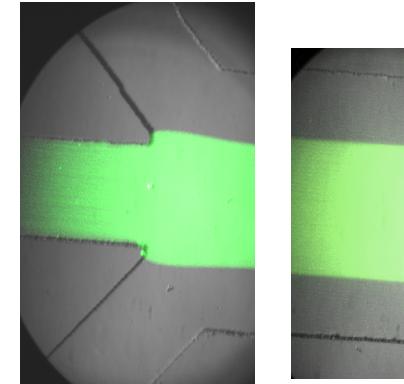
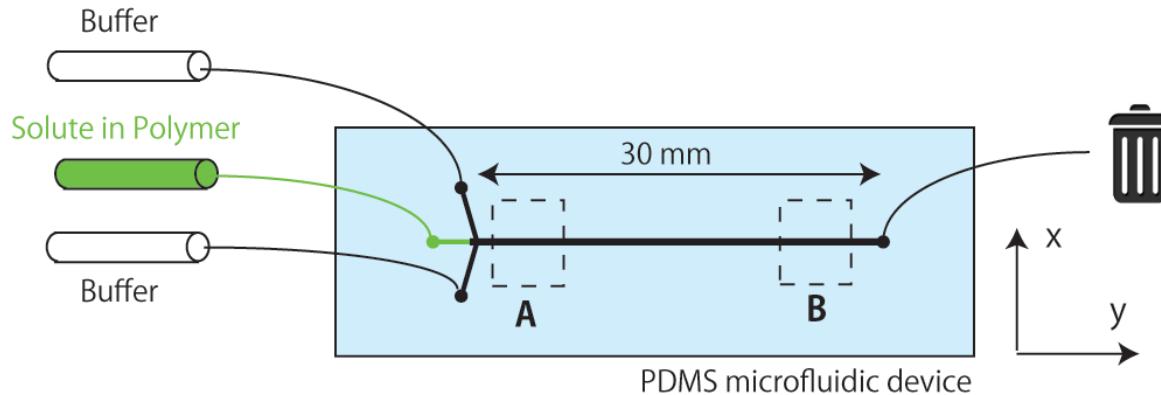
We check if this relation holds or not, in the absence of temperature gradient.

→ Osmotic pressure



PEG molecule

# Experimental setup: microfluidics



- Laminar flows in PDMS microfluidics
- Microchannel: Width 2 mm, Height 50  $\mu\text{m}$
- Polymer in a solution is PEG1000

$$c_{PEG}(x,y) = \frac{c_{PEG}^0}{2} \operatorname{erf}\left(\frac{x - h/4}{\tau}\right) - \frac{c_{PEG}^0}{2} \operatorname{erf}\left(\frac{x - 3h/4}{\tau}\right)$$

$$\operatorname{erf}(x) = \frac{2}{\sqrt{\pi}} \int_0^x dx e^{-x^2}$$

$$\tau = \sqrt{\frac{y}{U}}$$

$$D_{PEG} = 1.35 \times 10^2 \text{ } \mu\text{m}^2/\text{s}$$

$$U = 1.16 \text{ mm/s}$$

$$\nabla c_{PEG} \approx c_{PEG}^0 / \sqrt{D_{PEG} y / U} \approx 0.5 \%/\mu\text{m}$$

# PEG1,000 gradient-driven transport

Colloid (fluorescent)

$0.1 \mu\text{m}$ ,  $2.0 \times 10^{-3} \%$



$y=0 \text{ mm}$

PEG1,000, 20%

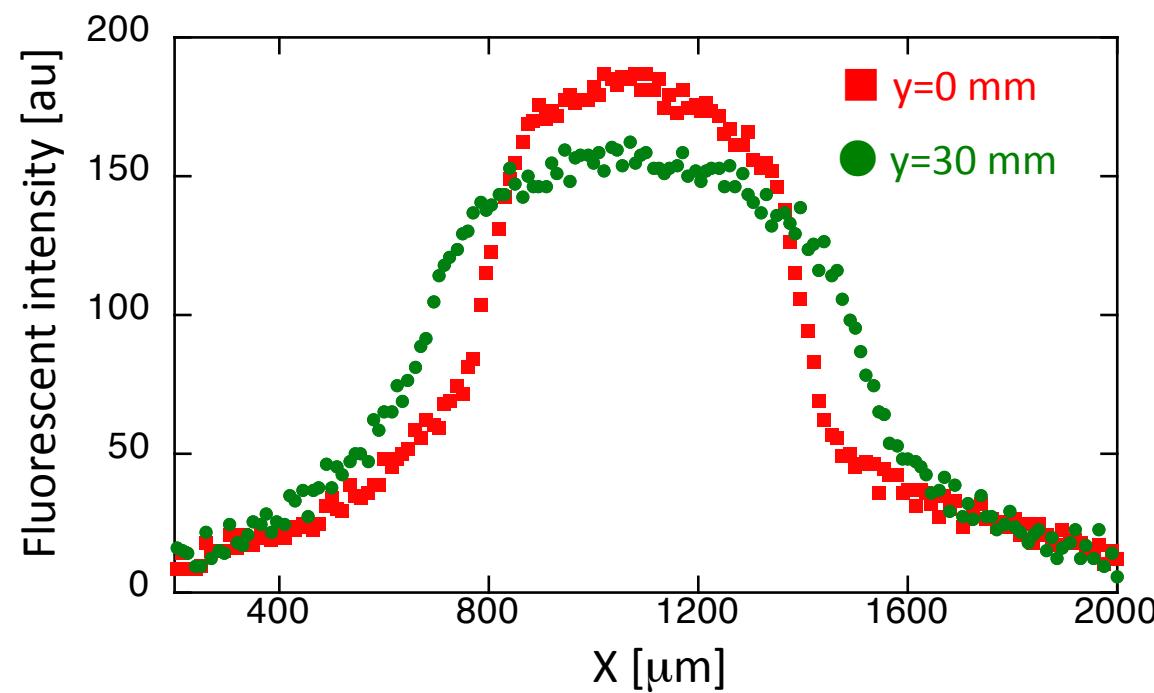
(Not PEG10,000 !)



$10 \text{ mm}$

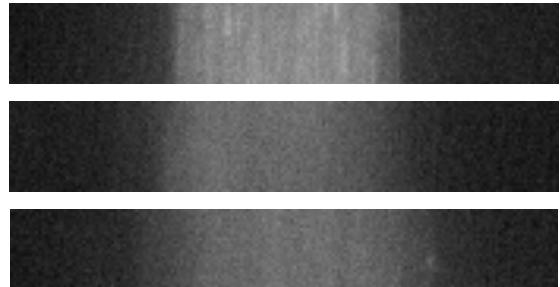


$30 \text{ mm}$

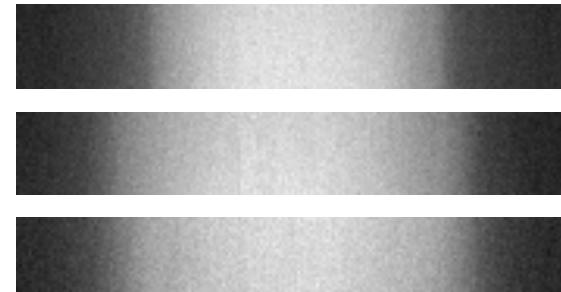


# Different sized-colloids in a gradient of small PEG1,000

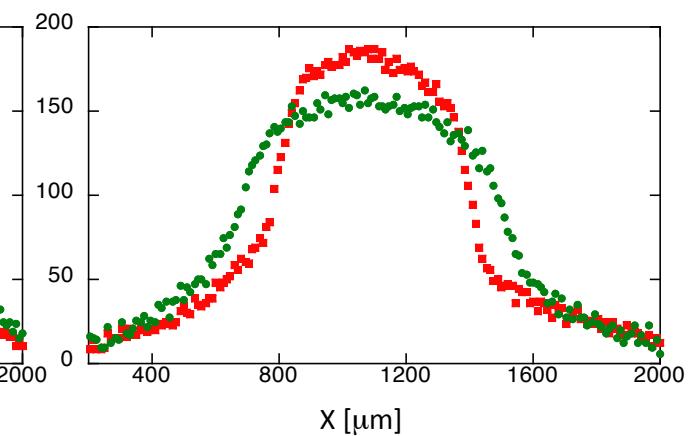
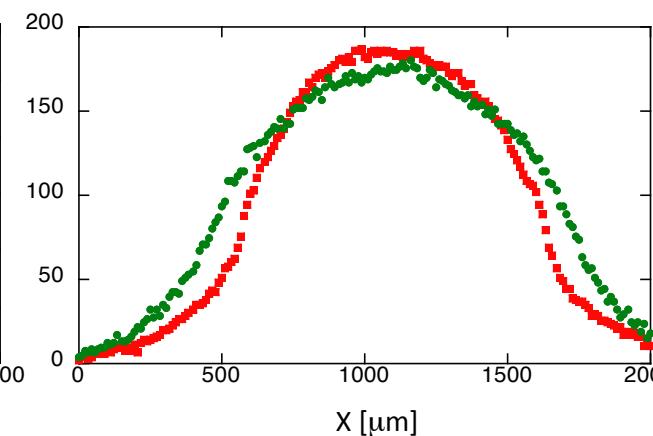
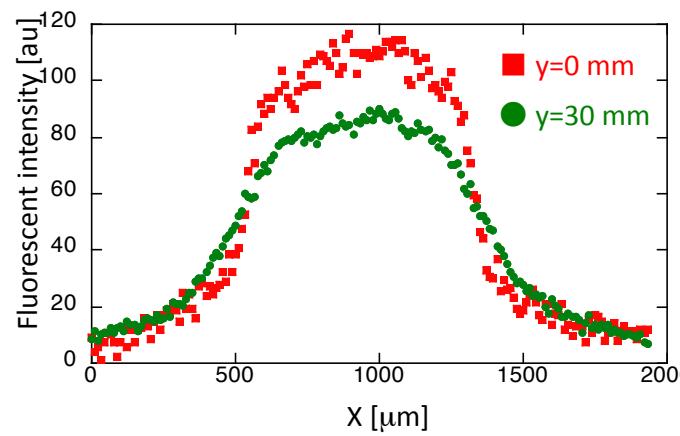
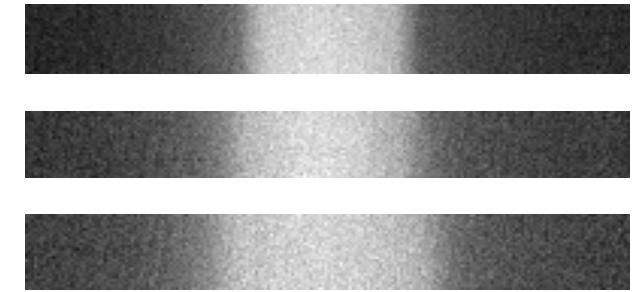
Colloid 0.5  $\mu\text{m}$



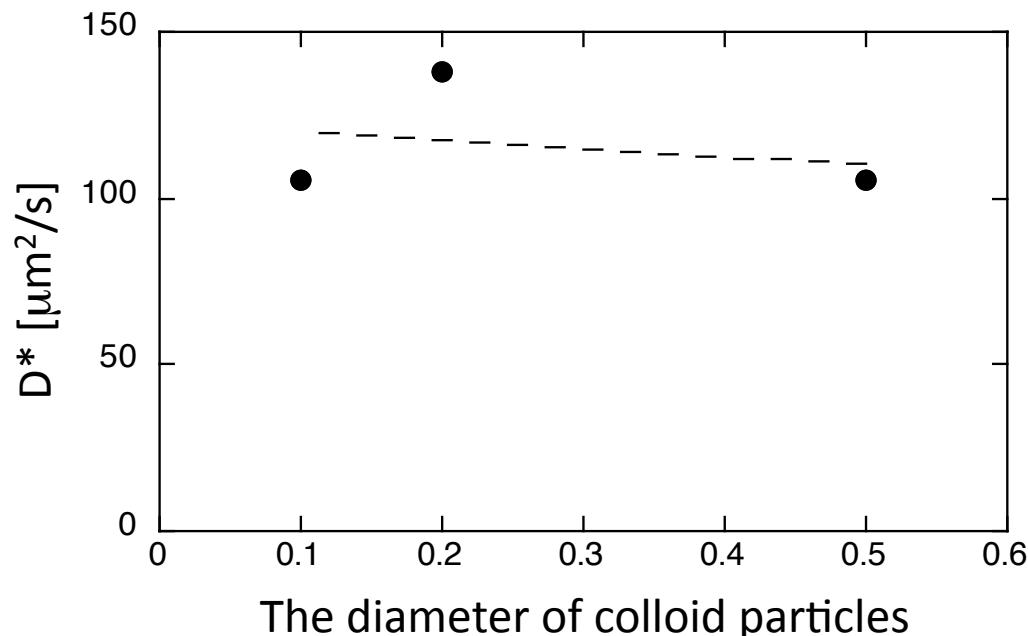
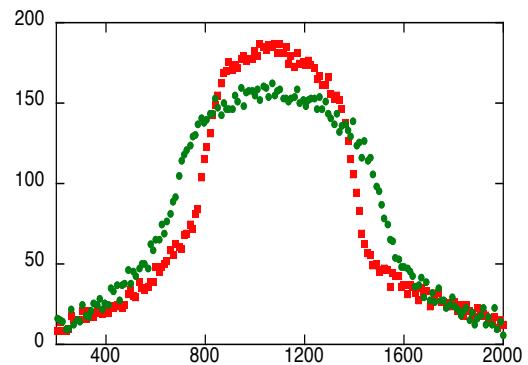
Colloid 0.2  $\mu\text{m}$



Colloid 0.1  $\mu\text{m}$



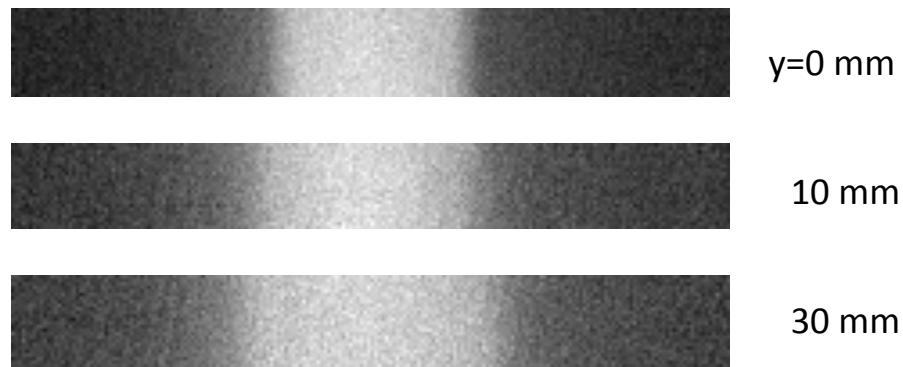
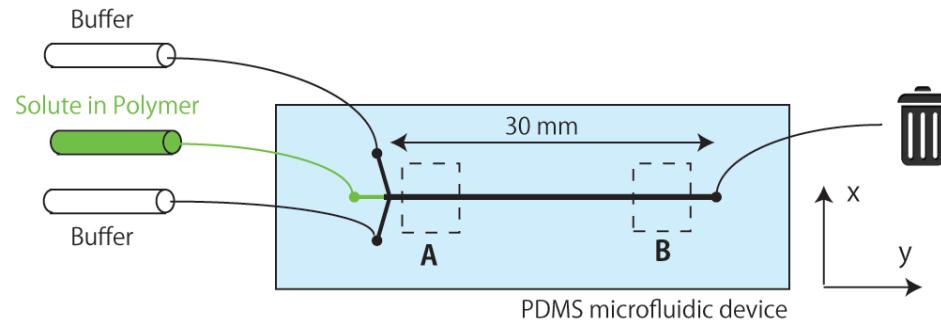
$$w(y) - w_0 \approx \sqrt{D^* \cdot U / 2y}$$



- In small PEG such as PEG1,000, the velocity of diffusiophoresis is not strongly dependent on particle size.
- It is because 10% PEG1,000 is still in a dilute regime, which is included in a model.

# Future works

- The mechanism of size-dependent separation



Diffusiophoresis through the repulsive interactions of the solute (colloid, DNA) and another polymer (PEG)

Viscosity?

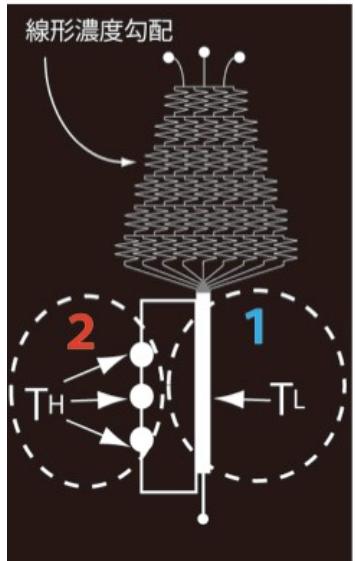
Many-body interactions?

K Odagiri, K Seki, K Kudo  
*Soft Matter* 8, 2775 (2012)

- The mechanism of trapping small RNA by comparable PEG

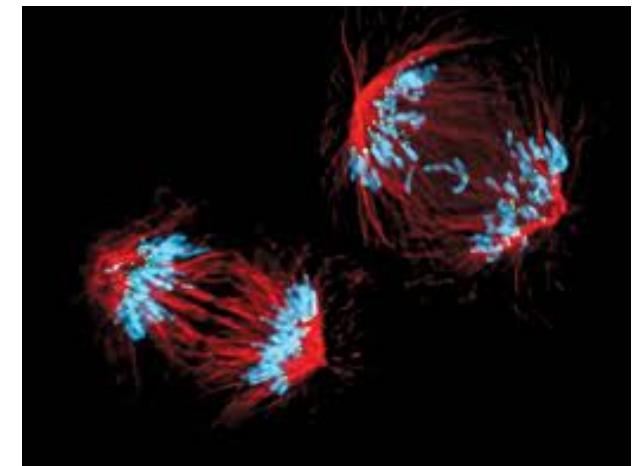
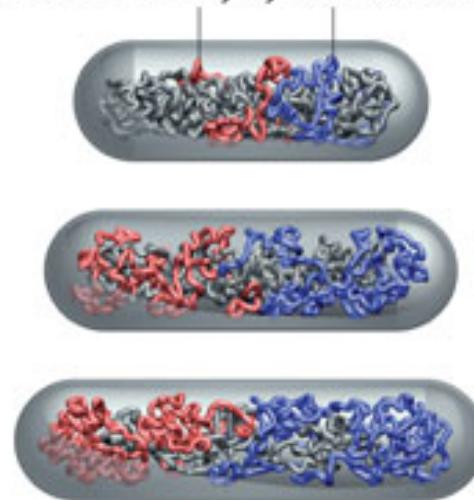
# (Putative) Relevance of biological systems

## 1. Selection and evolution of small RNA



## 2. Transport and segregation of genome

Blobs of newly synthesized DNA



# Conclusions

- DNA and RNA of various sizes show ring-like localization or accumulation under a temperature gradient in PEG.
- The entropic force gradient generated by the PEG gradient on the surface of solute, called diffusiophoresis, can push solute to the hot region.
- Simple microfluidic device allows us to study the transport by repulsive PEG. In a dilute regime, the clear size-dependence is not observed.
- Size-dependence of DNA localization is reversed for a long DNA >5.6kbp. PEG osmotic pressure induces coil-globule transition and decreases its effective size.
- The attraction by PEG concentration gradient is folding-dependent: Small RNA having stem-loop can be accumulated while poly-U of random chain cannot.

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