

Is QCD ferromagnetic?

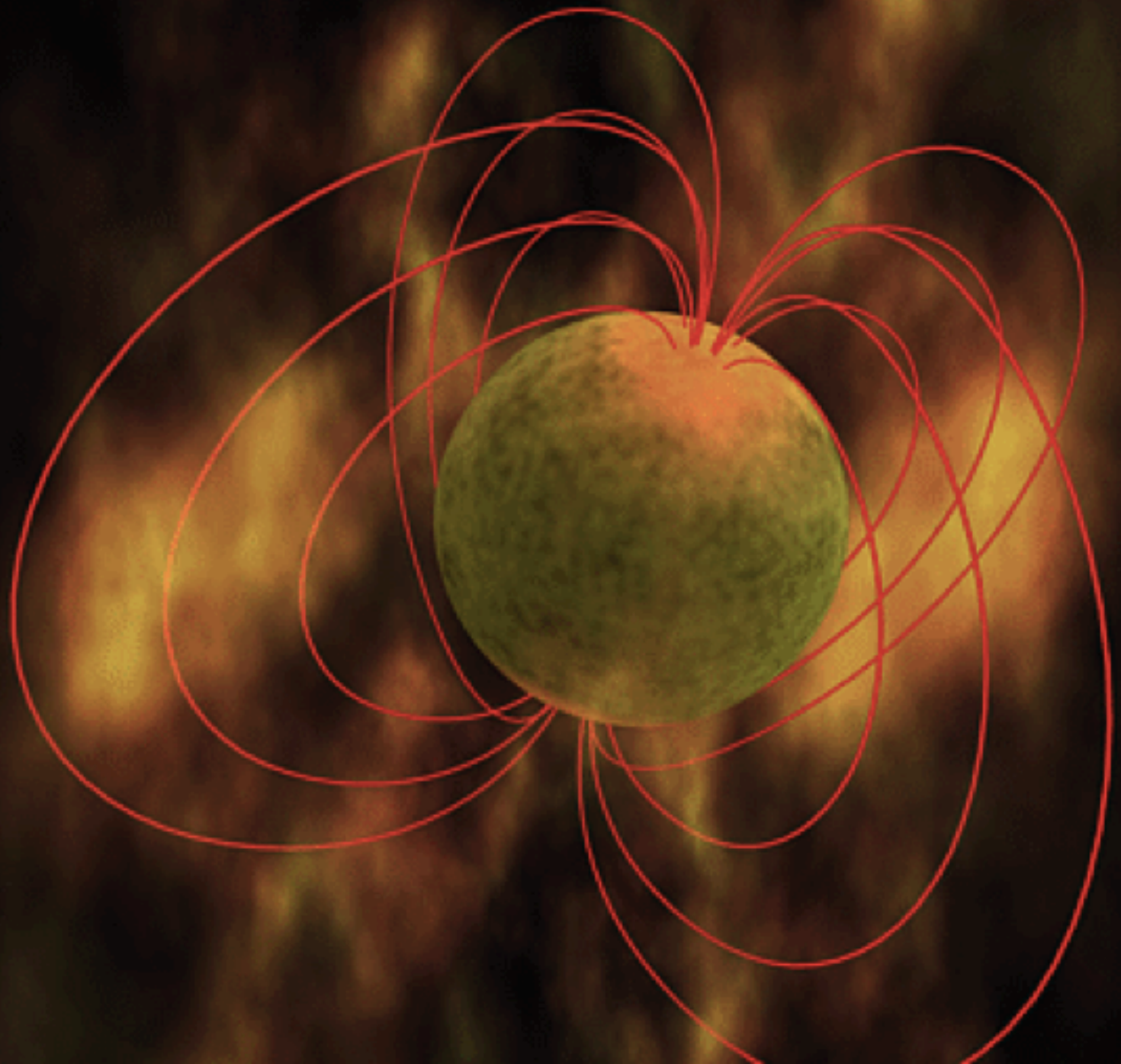


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@YKIS 2012, Yukawa institute

Holographic Magnetars

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(Osaka / RIKEN)

[arXiv/1209.4814\[hep-ph\]](https://arxiv.org/abs/1209.4814) + work in progress
w/ M. Eto (Yamagata), T.Hatsuda (RIKEN)



Problem

Origin of magnetars?

Cause

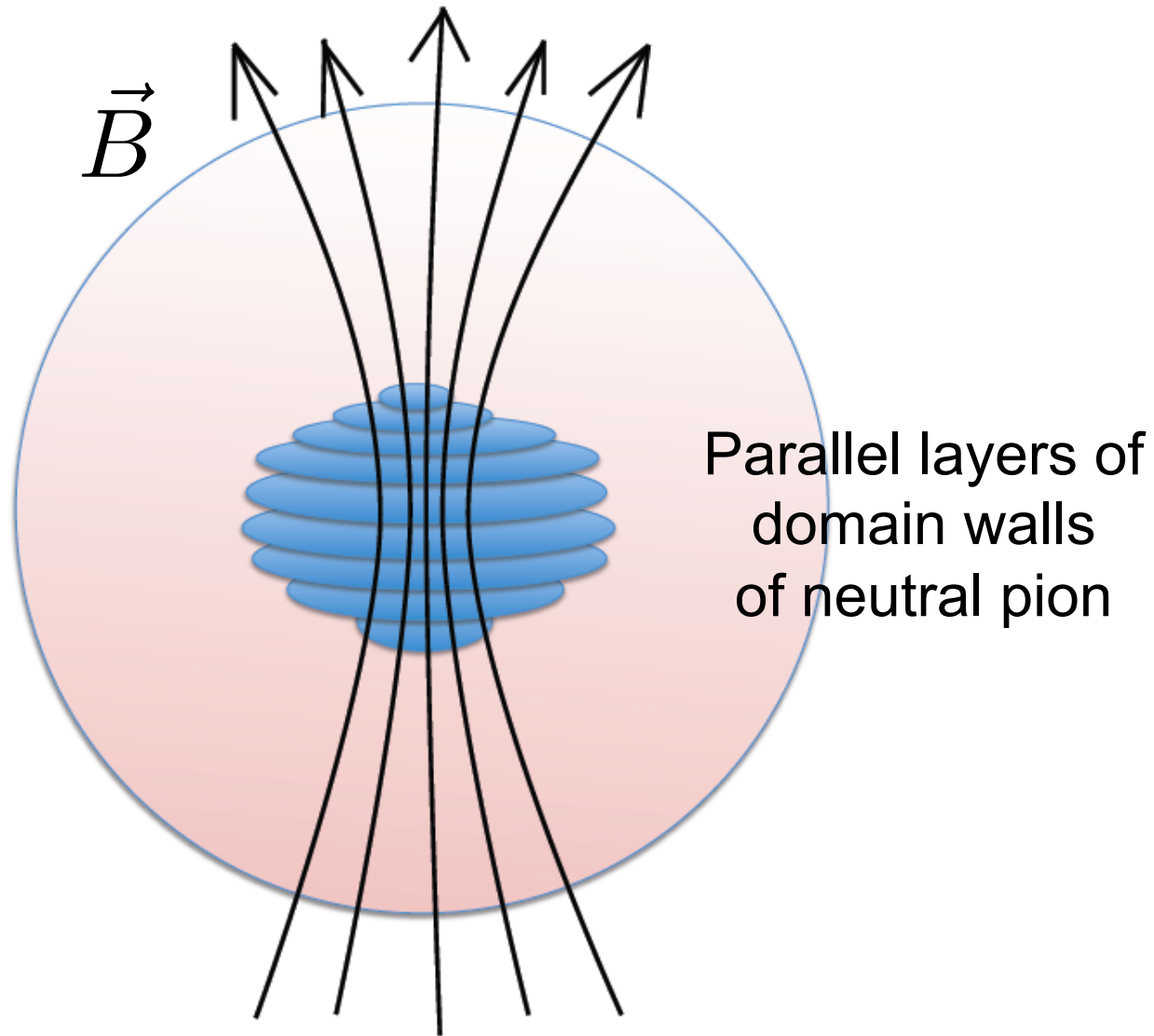
No magnetization at high density hadron phase?

Our solution

High density neutron matter

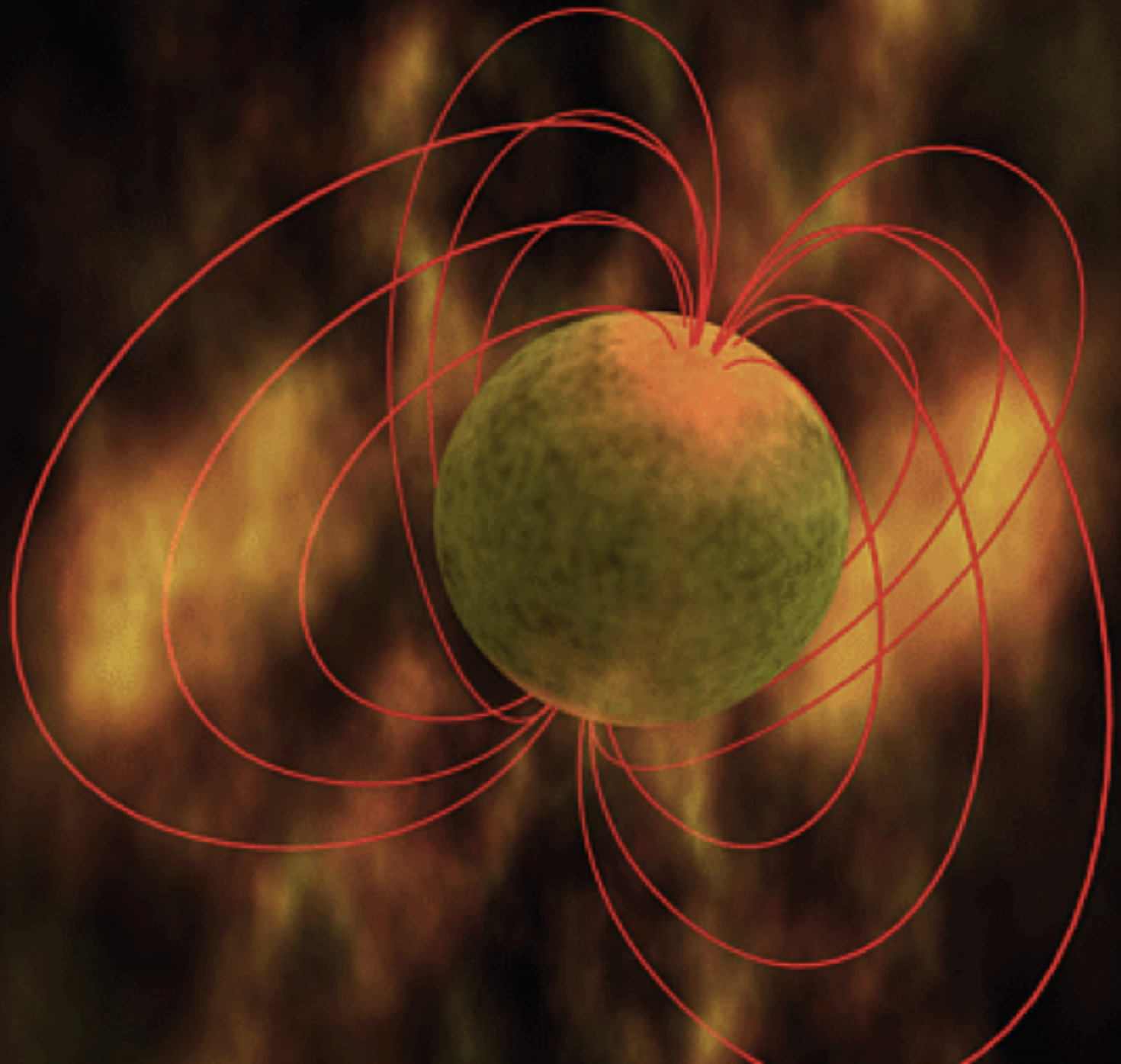
Axial anomaly + π^0 domain wall

New phase : QCD Ferromagnetism

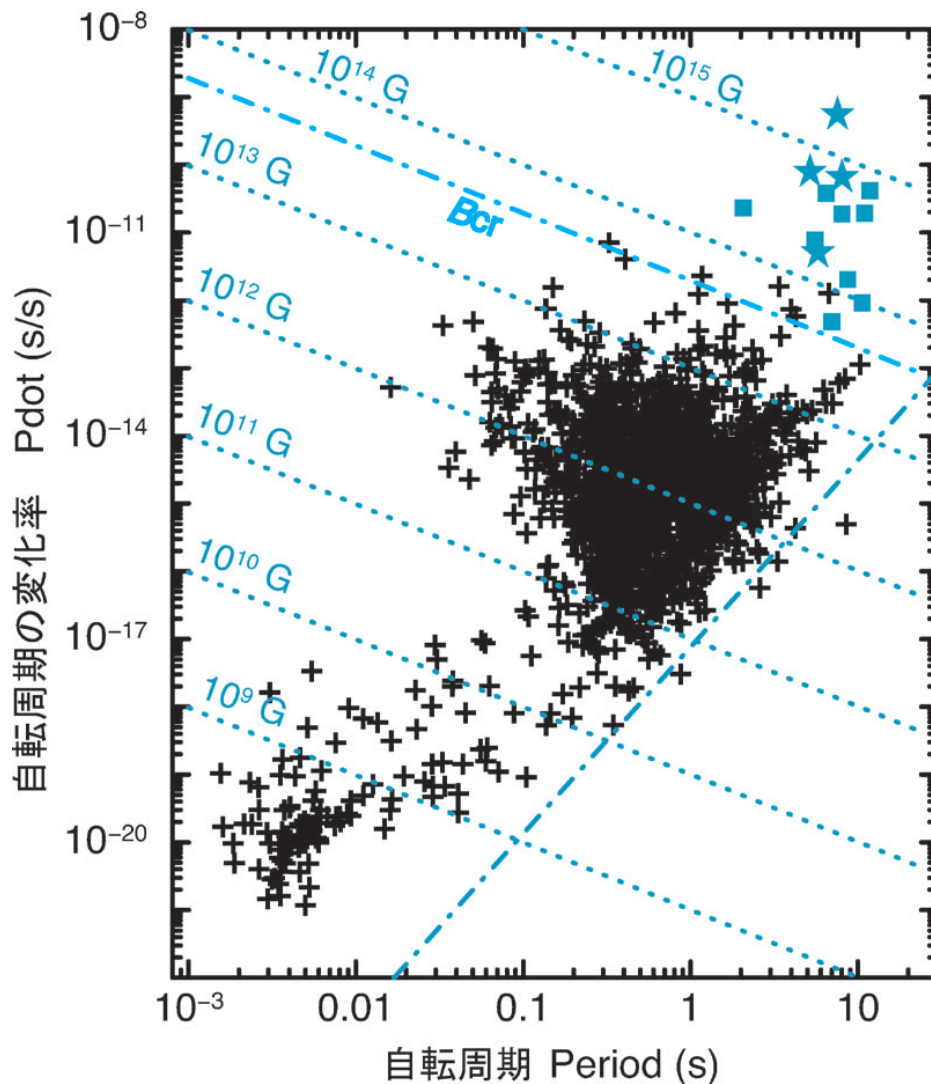


Plan

1. Magnetars. 3 pages
2. Pion domain wall. 3 pages
3. QCD Ferromagnetism. 3 pages
4. Holographization. 2 pages



Observation of Magnetars



[T.Enoto, 天文月報 2012.7]

Spindown luminosity

$$L = \partial_t(I\Omega^2) \propto \dot{P}/P^3$$

+

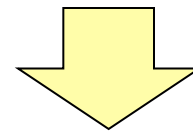
Magnetic dipole assumption

$$L = (2/3)\ddot{m}^2 \propto B^2/P^4$$

||

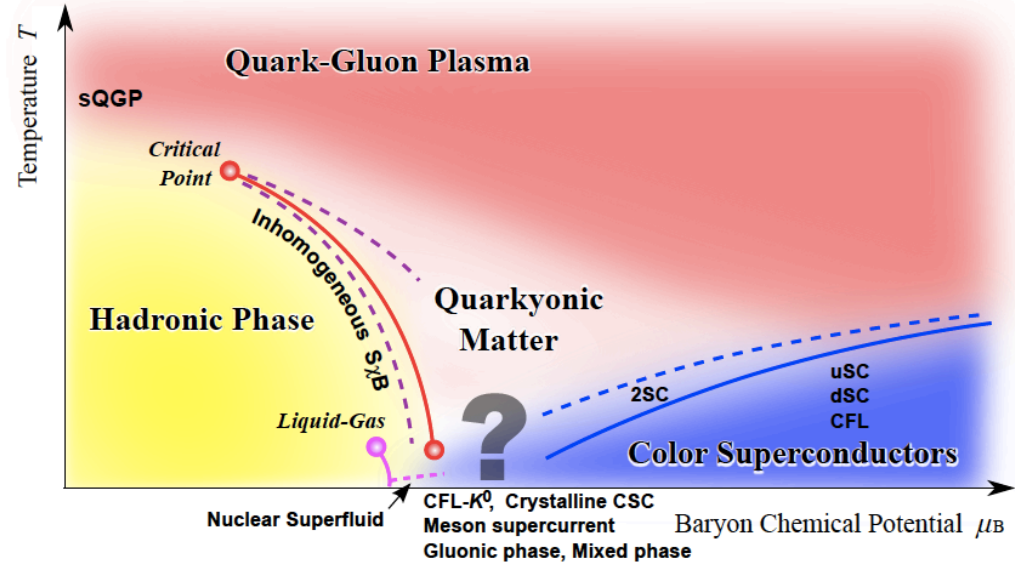
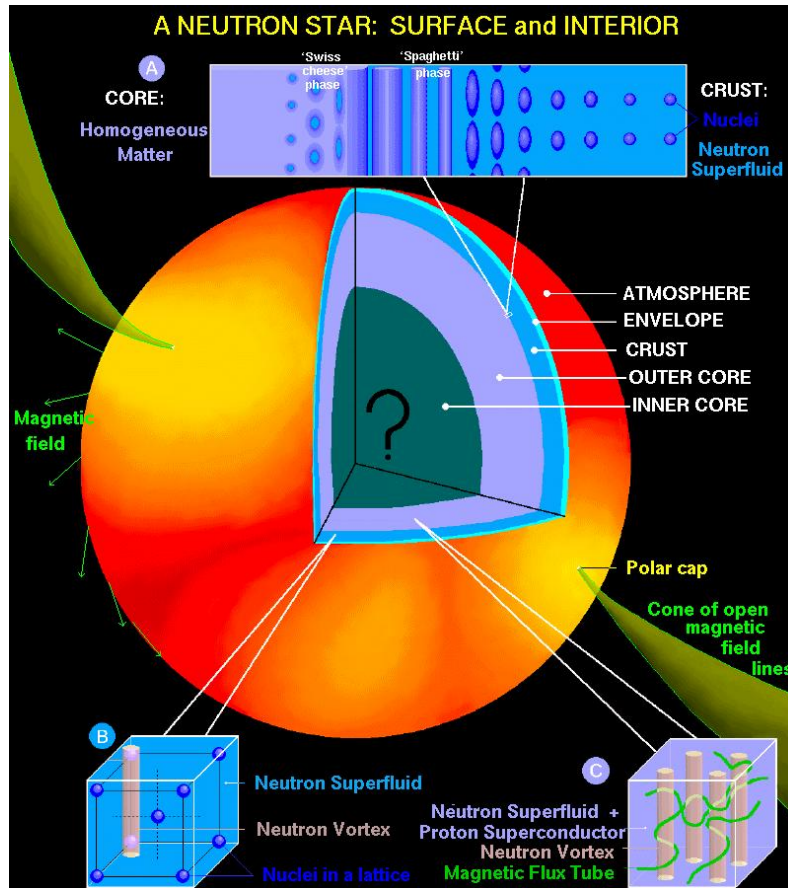
Surface magnetic field

$$B = 3.2 \times 10^{19} \sqrt{P\dot{P}} \text{ [G]}$$



Magnetar hypothesis

Origin of the strong magnetic field?



[Fukushima, Hatsuda 1005.4814]

[D.Page]

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Describing high density

Chiral lagrangian for neutron and neutral pion

$$\mathcal{L} = \bar{n}(i\gamma^\mu \partial_\mu + \mu\gamma_0 - M_n)n - \frac{1}{2f_\pi} (\bar{n}\gamma_\mu\gamma_5 n)\partial^\mu\theta + \frac{1}{2}(\partial_\mu\theta)^2 + f_\pi^2 m_\pi^2 (\cos\theta - 1)$$

Chemical potential giving high density neutron matter

Weinberg's action : nucleons and all pions

$$\mathcal{L} = \bar{\psi}(iD_\mu\gamma^\mu - M_N)\psi - g\bar{\psi}\gamma_5\gamma^\mu\boldsymbol{\tau}\psi \cdot D_\mu\boldsymbol{\phi} + \frac{1}{2}D_\mu\boldsymbol{\phi} \cdot D^\mu\boldsymbol{\phi} - \frac{1}{2}m_\pi^2(1 + \boldsymbol{\phi}^2/4f_\pi^2)^{-1}\boldsymbol{\phi}^2$$

$$D_\mu\boldsymbol{\phi} \equiv (1 + \boldsymbol{\phi}^2/4f_\pi^2)^{-1}\partial_\mu\boldsymbol{\phi}$$

$$D_\mu\psi \equiv \left[\partial_\mu + \frac{i}{4f_\pi^2}(1 + \boldsymbol{\phi}^2/4f_\pi^2)^{-1}\boldsymbol{\tau} \cdot (\boldsymbol{\phi} \times \partial_\mu\boldsymbol{\phi}) \right] \psi$$

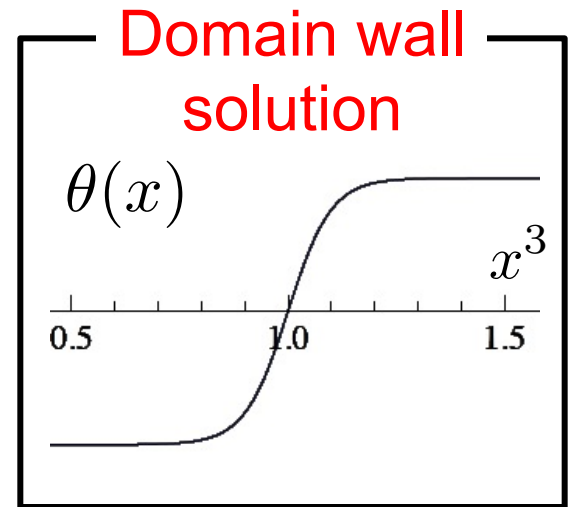
Domain wall of neutral pion

Chiral lagrangian = Sine-Gordon model

$$\mathcal{L} = \frac{1}{2} f_\pi^2 (\partial_\mu \theta)^2 - m_\pi^2 f_\pi^2 \cos \theta$$

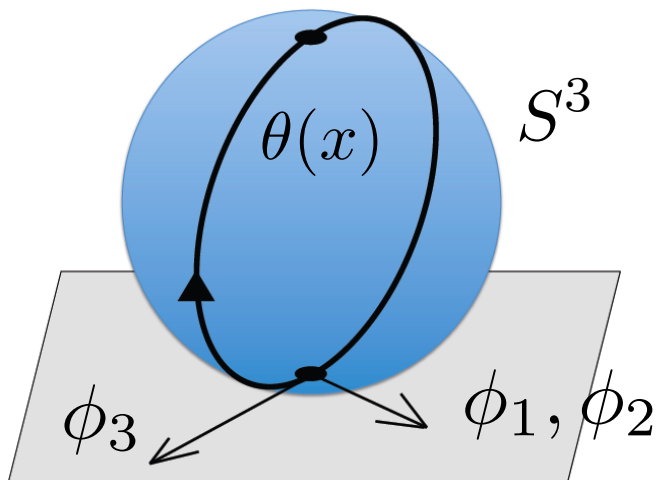
Classical solution connecting

$$\theta = 0 \text{ and } \theta = 2\pi$$



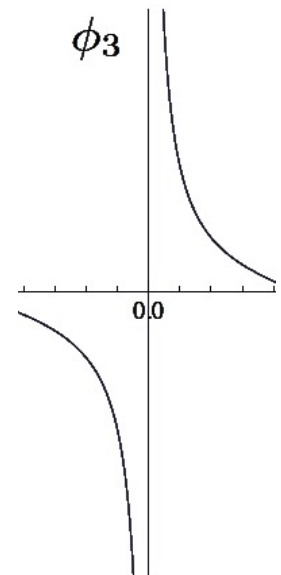
No topological obstruction against creation

[Hatsuda 1986]



$$\phi_3 = \frac{2f_\pi}{\sinh[m_\pi x^3]}$$

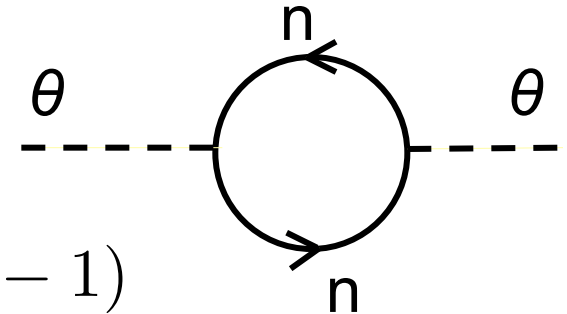
$$\phi_1 = \phi_2 = 0$$



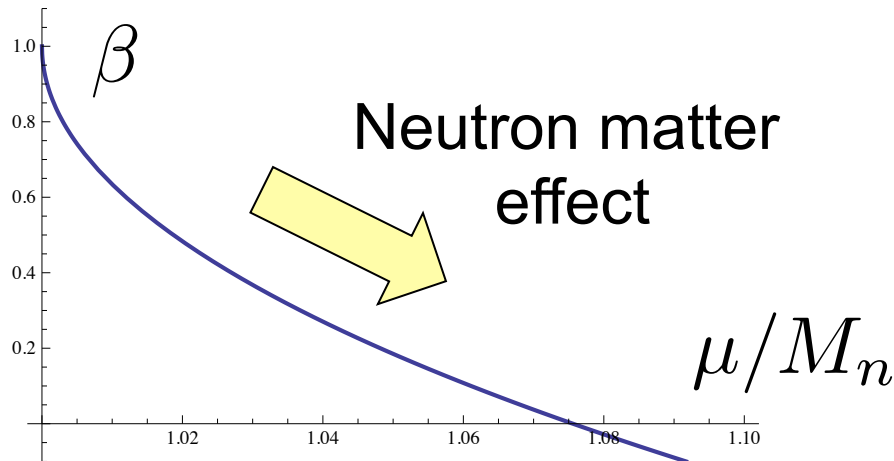
Light domain wall in neutron matter

[Hatsuda 1986]

Integration of neutron at 1-loop
→ pion effective action modified



$$\mathcal{L} = \frac{\alpha}{2} (\partial_0 \theta)^2 - \frac{\beta}{2} (\partial_i \theta)^2 + f_\pi^2 m_\pi^2 (\cos \theta - 1)$$



- β decreases at high density

- Light domain wall

$$\mathcal{E}/S = 8\sqrt{\beta} f_\pi^2 m_\pi$$

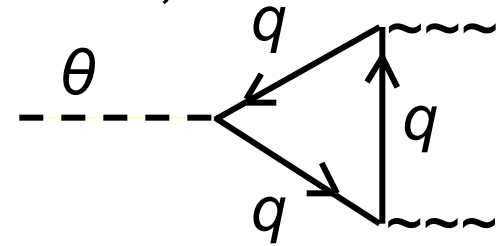
Plan

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Baryonically charged domain wall

QCD axial anomaly induces “ $\pi^0 \rightarrow 2\gamma$ ” term,

$$\mathcal{L}_{\text{anom}} = \frac{e}{4\pi^2} \mu_B B_3 \partial_3 \theta$$



Domain wall inducing a baryon charge,
under a magnetic field

$$N_B/S = eB_3/2\pi$$

[Son, Stephanov, 0710.1084]

Energy per baryon charge:

$$\mathcal{E}/N_B = \frac{16\pi f_\pi^2 m_\pi}{eB_i \hat{n}_i} \sqrt{\beta}$$

If smaller than neutron mass,
domain wall dominates!

$$B \sim \sqrt{\beta} \times 10^{19} \text{ [G]}$$

$$\left[\begin{aligned} L_{\text{WZW}} &= \frac{-ie}{16\pi^2} A_0^{(B)} B_3 \text{tr} [\tau_3 (U \partial_3 U^\dagger + \partial_3 U^\dagger U)] \\ U &\equiv (\sigma + i\boldsymbol{\tau} \cdot \boldsymbol{\pi})/f_\pi \end{aligned} \right]$$

Ferromagnetic QCD

Magnetization via wall $M = \frac{e\mu_B}{4\pi^2} \partial_3 \theta$ [Son, Zhitnitsky, hep-ph/0405216]

Wall layer with interval d : $M = \frac{e\mu_B}{2\pi d}$

Induced baryon density : $\rho = \frac{eB_3}{2\pi} \frac{1}{d}$

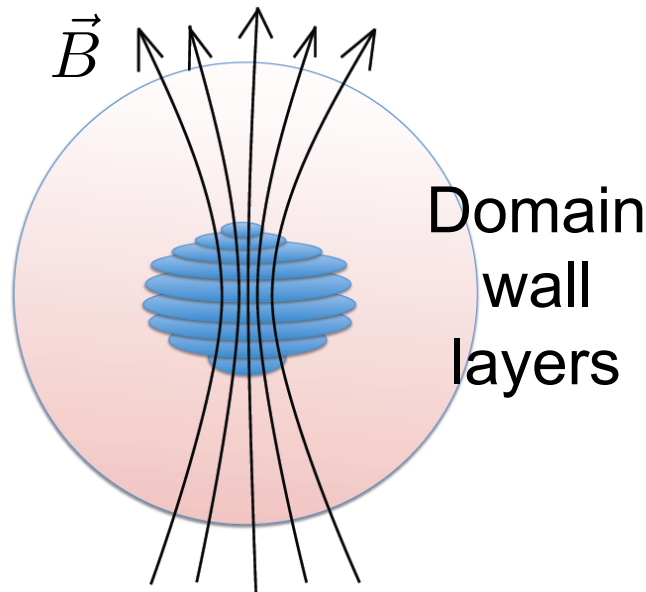
$$M = \frac{\mu_B \rho}{B_3}$$

Equilibrium

$$B_3 = M$$

$$B_3 \sim 10^{19} [\text{G}]$$

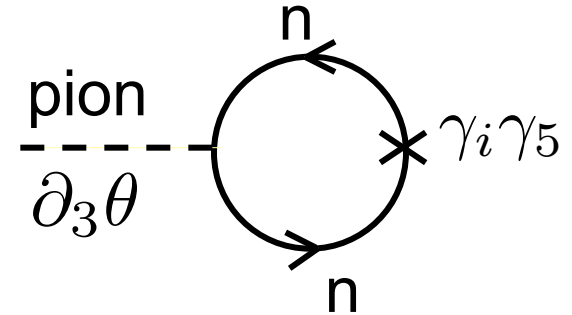
Spontaneous magnetization.
Very strong magnetic field,
QCD ferromagnetism



Spins align on the wall

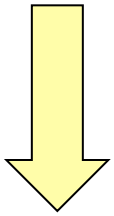
Neutron spins accumulate on the wall:

$$\mathcal{S}_i(x) = \langle \bar{n} \gamma_i \gamma_5 n \rangle = (\beta - 1) f_\pi^2 \partial_i \theta(x)$$



→ Domain wall carries neutron spins:

$$s_3/S = 2\pi(\beta - 1) f_\pi^2 \quad [\text{Hatsuda 1986}]$$



Neutrons have magnetic moment :

$$\mu = \frac{g e s_3}{2M_n} \quad \text{with} \quad g \sim -3.8$$

Domain walls are magnetized:

$$M = \frac{\pi |g| e f_\pi^2}{m_N} \frac{1}{d}$$

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--- **Anomaly magnetizes the core**
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Full meson theory (酒井杉本) giving the same

Pions are gauge field A_z

$$S_{D8} = -\frac{\lambda N_c}{216\pi^3} \int d^4x dz (1+z^2) F_{\mu z}^2$$

Quark mass effect : worldsheet instanton

$$S_{\text{mass}} = c \int d^4x \text{Ptr} \exp \left[-i \int_{-\infty}^{\infty} A_z dz \right] + \text{c.c.}$$

[Hirayama, Lin, Yee, KH 0803.4192]

[Aharony, Kutasov, 0803.3547]

EOM for $A_\mu(x)$ is solved as

$$A_\mu = 0, \quad A_z = \frac{1}{1+z^2} \tau_3 \theta(x^\mu)$$

Then the EOM for $A_z(x)$ reduces to sine-Gordon eq. of $\theta(x)$

Anomaly analysis and Magnetization

Anomaly = CS term of the D8-brane action

$$S_{\text{CS}} = \frac{N_c}{24\pi^2} \int d^4x dz A_0^{(\text{B})} B_3^{(\text{em})} \partial_3 A_z.$$

Consider a chiral limit for light domain walls,

$$A_z = \frac{1}{1+z^2} a x^3$$

Total energy with all D8 action + EM :

$$\begin{aligned} E &\sim \lambda N_c a^2 + N_c \mu_q B a + B^2 \\ &= \lambda N_c \left(a + \frac{\mu_q B}{\lambda} \right)^2 - \frac{N_c \mu_q^2}{\lambda} B^2 + B^2 \end{aligned}$$

$\mu_q > \sqrt{\lambda/N_c} \rightarrow$ Spontaneous magnetization

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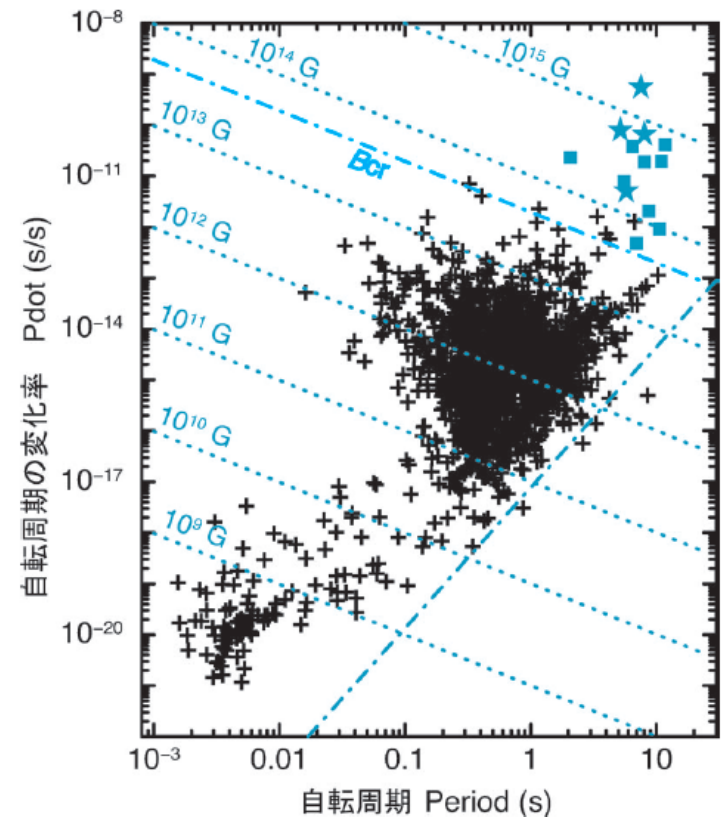
High density neutron matter

Axial anomaly + π^0 domain wall

New phase : QCD Ferromagnetism

More to come...

- Stability of the walls?
Charged pion condensates
- More realistic?
Nuclear forces, hyperons
- EOS via walls?
- Neutron superfluidity?
p-wave at high density,
pairing on the walls
- Gravitational waves from
isolated neutron stars?
- Classification of neutron stars?



[T.Enoto, 天文月報 (2012)]