

Cast of Play

3d CS-VM.

$U(N)$ or $SU(N)$ gauge group.

CS $A_\mu + \square$ matter field

boson ϕ or fermion ψ .

non-SUSY.

$$S = kCS(A) + \int |D_\mu \phi|^2 + \frac{f(x)}{N^2} (\phi^2)^3$$

$$\text{or } + \int \bar{\psi} D^\mu \psi.$$

Now consider N_f flavors of \square scalars AND fermions.

- not (yet) supersymmetric.

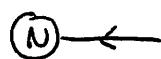
SUSY vector models obtained by dbl + triple trace deformations + possibly gauging flavor currents.

e.g. $N=2$ CS-VM w/ \square chiral

$$S = CS(A) + \int D\sigma + \bar{X}X + \text{(minimal coupling of } N=2 \text{ chiral matter)}$$

$$\text{from } D, \sigma \Rightarrow \frac{1}{k} (\bar{\Phi}\Phi)(\bar{\Psi}\Psi), \frac{1}{k^2} (\bar{\Phi}\Phi)^3$$

$$\text{from } \bar{X}, X \Rightarrow \frac{1}{k} (\bar{\Phi}\bar{\Psi})(\bar{\Psi}\Phi)$$



$$N=3. \quad \text{Diagram: } (N) \xrightarrow{\square} \bar{Q} \quad + W = \frac{1}{k} (Q\bar{Q})^2.$$

$$N=4 \quad \text{Diagram: } (N) \xrightarrow{\square} (M) \quad \Leftrightarrow \quad W=0.$$

$$N=6 \quad \text{Diagram: } (N) \xrightarrow{\square} (M) \quad \Leftrightarrow \quad W = \frac{1}{k} \text{Tr}(Q_i \tilde{Q}_i Q_j \tilde{Q}_j) - \frac{1}{k} \text{Tr}(\tilde{Q}_i Q_i \tilde{Q}_j Q_j)$$

4d Vasiliev theory.

x^μ, Y, Z .

$A(x|Y, Z), B(x|Y, Z)$.

$$dA + A \star A = f_*(B \star K) dz^2 + \text{c.c.}$$

$$dB + A \star B - B \star \pi(A) = 0.$$

$$f(x) = X e^{i\Theta(x)}, \quad \Theta(x) = \Theta_0 + \Theta_2 x^2$$

real, even. + ...

AdS₄ vacuum

$$A = W + S, \quad W = W_0 = \omega_0 + \epsilon_0$$

$$S = 0, B = 0.$$

Perturbation theory around AdS₄.

⇒ interacting HS fields.

SUSY generalization. "n-extended Vasiliev" (a fake)

$$\psi_i, i=1, \dots, n. \quad \Gamma = i^{\frac{n(n+1)}{2}} \psi_1 \dots \psi_n.$$

$$A(x|Y, Z, \psi), B(x|Y, Z, \psi)$$

$$R = K \bar{K} \Gamma, [R, W] = \{R, S\} = [R, B] = 0. \quad = [R, E]$$

$$\text{reality condition: } U(A)^* = -A,$$

$$U(B)^* = \Gamma K \star B \star K = \bar{K} \star B \star \bar{K} \Gamma.$$

Involution on \star -algebra, reverses order.

$$dA + A \star A = f_*(B \star K) dz^2 + \bar{f}_*(B \star \bar{K} \Gamma) d\bar{z}^2.$$

Most susyles \star all HS sym

will be broken by bdry condition for generic Θ_0 .



Spectrum and correlators.

VM.

single trace operators.

$$\bar{\Phi}_i \Phi_i, \quad \bar{\Phi}_i \partial^S \Phi_i;$$

$$\bar{\Phi}_i \Psi_i, \quad \bar{\Psi}_0 \Psi_0, \quad \bar{\Psi} \gamma_\mu \partial_{\mu_1} \dots \partial_{\mu_S} \Psi,$$

$$- \Delta = S+1 + \mathcal{O}(1/N) \text{ @ FINITE } \lambda.$$

HS GT in AdS₄.

minimal bosons.

$$S=0, 2, 4, \dots$$

$$\text{run-min } S=0, 1, 2, 3, \dots$$

susy. include half-integer spins.

Correlators.

Consider CS-scalar - VM as example.

$$\langle J_S J_{S'} \rangle = \tilde{N} \delta_{SS'}$$

$$\langle J_{S_1} J_{S_2} J_{S_3} \rangle = \tilde{N} \left[\frac{1}{1+\tilde{\lambda}^2} (\text{FB}) + \frac{\tilde{\lambda}^2}{1+\tilde{\lambda}^2} (\text{FF}) \right. \\ \text{for } S_i \neq 0, \quad \left. + \frac{\tilde{\lambda}}{1+\tilde{\lambda}^2} (\text{odd}) \right] + \mathcal{O}(N^0)$$

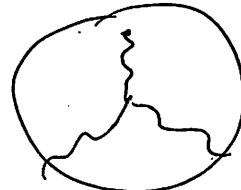
 $\tilde{N}, \tilde{\lambda}$ related to N, λ .

$$\langle J_0 J_S J_{S_2} \rangle = \tilde{N} \left[\frac{1}{\sqrt{1+\tilde{\lambda}^2}} (\text{FB}) + \frac{\tilde{\lambda}}{\sqrt{1+\tilde{\lambda}^2}} (\text{odd}) \right]$$

HS sym breaking:

$$\partial \cdot J = \frac{f(\lambda)}{N} \sum J J + \frac{g(\lambda)}{N^2} \sum J J J.$$

given boundary source $J_{i_1 \dots i_N}$
 solve bulk spin- s field $\varphi_{i_1 \dots i_N}$
 $\varphi_{i_1 \dots i_N} \rightarrow z^{S+1} \delta(\vec{z} \cdot \vec{x}_0) \varepsilon_{i_1 \dots i_N}$

translate into A_μ, B master fieldsThen solve $\tilde{D}_\mu B^{(2)} = -A^{(1)} \star B^{(1)} + B^{(1)} \star A^{(1)}$
extract 2nd order φ from Weyl tensorbdry value \propto 3-pt fn. $B^{(2)}(x|Y, z=0)$ 

$$\langle J_0 J_S J_{S_2} \rangle$$

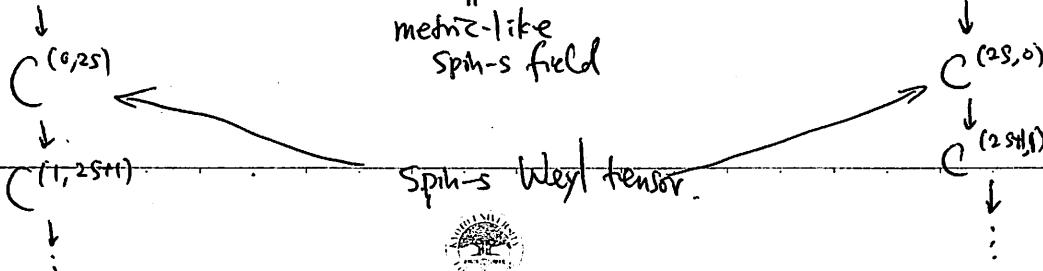
$$= \cos \theta_0 (\text{FB}) + \sin \theta_0 (\text{odd})$$

$$\boxed{\tilde{\lambda} = \tan \theta_0}$$

structure of master fields.

$$A = W + S, \quad W = W_0 + \hat{W}, \quad \text{set } \Omega^{(x|Y)} = \hat{W}|_{z=0}, \quad C = B|_{z=0}.$$

$$\Omega^{(0,2S)} \leftarrow \Omega^{(S-1,S)} \leftarrow \Omega^{(S-1,S-1)} \rightarrow \Omega^{(S,S-2)} \rightarrow \dots \Omega^{(2S-2,0)}$$

metric-like
spin- s field

Symmetry & breaking.

Symmetries of AdS₄ sol'n: $\epsilon(x|Y, z, \psi)$ that obey.

$$d_z \epsilon = 0 \Rightarrow \epsilon = \epsilon(x|Y, \psi).$$

$$D_0 \epsilon(x|Y, \psi) = 0, \quad W_0 = L^{-1} * dL.$$

$$\Rightarrow \epsilon = L^{-1} * \epsilon(Y, \psi) * L.$$

e.g. SUSY generators.

$$\epsilon(x|Y, \psi) = \Lambda^\alpha(x|\psi) y_\alpha + \bar{\Lambda}^\alpha(x|\psi) \bar{y}_\alpha$$

$$\Lambda(x|\psi) = z^{-\frac{1}{2}} (\Lambda_0(\psi) + \vec{x} \cdot \vec{\sigma} \sigma^z \Lambda_-(\psi)) + z^{\frac{1}{2}} \Lambda_+(\psi).$$

$\uparrow \quad \quad \quad \downarrow \quad \quad \quad \uparrow$
Q S

$$\delta_\epsilon B = -\epsilon * B + B * \pi(\epsilon).$$

\uparrow mixes fields of different spins.



Consider fields sourced by $J^{(s)}$ in AdS₄,
perform spin-s' ~~gauge trans~~ global symm transf.

$$\delta_\epsilon \circledast B^{(s)} \neq 0.$$

\uparrow scalar ~~$\vec{M}_z + \vec{p}_z^2 \vec{z}$~~

$$= \cos \theta_0 (--) z + \sin \theta_0 (---) z^2 + \dots$$

Preserved by: $\Delta=1$ b.c. if $\theta_0=0$.

$\Delta=2$ b.c. if $\theta_0=\frac{\pi}{2}$.

broken for e.i.s.c. b.c. if θ_0 generic.

general n -extended Vasiliev
 $U(2^{\frac{n}{2}-1}) \times U(2^{\frac{n}{2}-1})$
bosonic spin-1 gauge fields,

$N=3$ b.c. for $n=4$ Vasiliev theory.

$$\psi_1, \dots, \psi_4.$$

SUSY generator.

$$E \leftrightarrow A_0(\psi) = \gamma \psi_1, \gamma \psi_2, \gamma \psi_3.$$

out of $2^3 = 8$ fermion sym.
(monomials in ψ_i)

b.c.

scalar.

$$B^{(0)}(x, z | \psi) = (e^{i\gamma} + \Gamma e^{-i\gamma}) \tilde{f}_1(\psi) z + (e^{i\gamma} - \Gamma e^{-i\gamma}) \tilde{f}_2(\psi) z^2 + \mathcal{O}(z^3)$$

fermion

$$B^{(\frac{1}{2})}(x, z | Y, \psi) = z^{\frac{3}{2}} [e^{i\alpha} (xy) - \Gamma e^{-i\alpha} (\bar{x}\bar{y})] + \dots$$

$$X = \sigma^z \bar{X}$$

spin-1 gauge field

$$B^{(1)}(x, z | Y, \psi) = z^2 [e^{i\beta} (yFy) + \Gamma e^{-i\beta} (\bar{y}\bar{F}\bar{y})] + \dots$$

$$F = -\sigma^z \bar{F} \sigma^z.$$

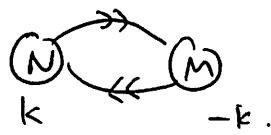
Characterized by: γ, α, β . - linear operators on fn of ψ .

$N=3$ b.c. - $\beta = \theta_0$, $\alpha = \theta_0(1 - P_{\psi_1 \psi_2 \psi_3}) - \theta_0 P_{\psi_1 \psi_2 \psi_3}$, $\gamma = \theta_0$.

$$P_{\psi_1 \psi_2 \psi_3} \tilde{f}_{1,2} = \tilde{f}_{1,2}.$$

vs $N=4$ b.c. $\beta = \theta_0(1 - P_F)$, $\alpha = \theta_0 P_{\psi_1}$, $\gamma = \theta_0 P_F$, $P_F \tilde{f}_{1,2} = 0$.

ABJ triality



↔

IIA

on
 $AdS_4 \times \mathbb{CP}^3$

$$\text{w/ } \int_{\mathbb{CP}^1} B = \frac{N-M}{k} + \frac{1}{2}$$



$n=6$ Vasiliev theory. $U(M)$ CP factor

$$\text{w/ } \theta_0 = \frac{\pi}{2} \lambda.$$

and b.c.

$$\beta = \theta_0 (1 - P_T) - \theta_0 P_T,$$

$$\alpha = \theta_0 (1 - P_{\psi_i T}) - \theta_0 P_{\psi_i T},$$

$i=1, \dots, 6$.

$$\gamma = \theta_0 P_{\psi_i \psi_j T}$$

$$P_{\psi_i \psi_j T} \hat{f}_{1,2} = 0.$$

- preserve global SUSY gen. by

$$\Lambda_0 = \eta \psi_i, \quad i=1, \dots, 6.$$

Comments

- HS/VM duality, as we understand it, involves only local gauge invariant operators on the CFT side. Putting VM on 3-manifolds of nontrivial topology requires extension of Vasiliev's system
(Also clear in AdS_3/CFT_2 HS - W_n-min model duality)
- Holography is not mysterious, bulk locality is.
- Embedding into string theory?

2 constructions.

• ABJ $U(N)_k \times U(M)_{-k}$ \longrightarrow II A on $AdS_4 \times \mathbb{C}P^3$.

$\downarrow N \gg M$.

$n=6$ $U(M)$ - Vasiliev $\leftarrow ?$ single string
w/ $N=6$ b.c. \nearrow
multi-HS-particles.

- Giotto - Jaffers

