

C and P violations in the PNJL model

H. Kouno (Saga Univ.)

Y. Sakai, T. Sasaki, K. Kashiwa,
M. Yahiro (Kyushu Univ.)

Contents

- ① Deconfinement is a remnant of C violation
- ② Chiral restoration is a remnant of P restoration
- ③ D and C happen almost simultaneously in the real world
 $\Leftrightarrow$ C and P may happen simultaneously in the θ world

Deconfinement as remnant of C violation

- **Deconfinement** at zero density can be regarded as a remnant of **C violation** at $\theta=\pi/3$, where $\mu_i = i\theta T$ is the imaginary chemical potential.

- **PNJL model analysis**

H.K., Yuji Sakai , Kouji Kashiwa, Masanobu
Yahiro

Journal of Physics G 36 (2009)

115010-1-115010-17

PNJL model

- K. Fukushima, Phys. Lett. B591, 277(2004)
Polyakov loop+NJL
- Chiral symmetry chiral restoration
extended Z_3 symmetry deconfinement
- at imaginary chemical potential,
there is **Roberge-Weiss transition**
as C violation

RW periodicity and extended Z_3 symmetry (EZ_3)

$$\theta = \mu_I/T$$

$$\Omega = -2N_f \int \frac{d^3\mathbf{p}}{(2\pi)^3} \left[3E(\mathbf{p}) + \frac{1}{\beta} \ln \left[1 + 3(\Phi + \Phi^* e^{-\beta E^-(\mathbf{p})}) e^{-\beta E^-(\mathbf{p})} + e^{-3\beta E^-(\mathbf{p})} \right] \right. \\ \left. + \frac{1}{\beta} \ln \left[1 + 3(\Phi^* + \Phi e^{-\beta E^+(\mathbf{p})}) e^{-\beta E^+(\mathbf{p})} + e^{-3\beta E^+(\mathbf{p})} \right] \right] + U_M + \mathcal{U}$$

$$E(\mathbf{p}) = \sqrt{\mathbf{p}^2 + M^2}, \quad \sigma = \langle \bar{q}q \rangle, \quad \Sigma_s = -2G_s \sigma, \quad U_M = G_s \sigma^2, \quad M = m_0 + \Sigma_s.$$

extended Z_3 trans.

$$\theta \rightarrow \theta + 2\pi k/3, \quad \Phi(\theta) \rightarrow \Phi(\theta) e^{-i2\pi k/3} \quad \theta = \mu_I/T$$

修正版Polyakovループ $\Psi \equiv e^{i\theta} \Phi$

$$\theta \rightarrow \theta + 2\pi k/3, \quad \Psi(\theta) \rightarrow \Psi(\theta), \quad \Psi(\theta)^* \rightarrow \Psi(\theta)^*$$

Thermodynamic potential

$$\theta = \mu_I/T$$

$$\Omega = -2N_f \int \frac{d^3p}{(2\pi)^3} \left[3E(p) + \frac{1}{\beta} \ln \left[1 + \cancel{3\Psi} e^{-\beta E(p)} + \cancel{3\Psi^*} e^{-2\beta E(p)} \cancel{e^{3i\theta}} + e^{-3\beta E(p)} \cancel{e^{3i\theta}} \right] \right. \\ \left. + \frac{1}{\beta} \ln \left[1 + \cancel{3\Psi^*} e^{-\beta E(p)} + \cancel{3\Psi} e^{-2\beta E(p)} \cancel{e^{-3i\theta}} + e^{-3\beta E(p)} \cancel{e^{-3i\theta}} \right] \right] + U_M + \mathcal{U}$$

$$E(p) = \sqrt{\mathbf{p}^2 + M^2}, \quad \sigma = \langle \bar{q}q \rangle, \quad \Sigma_s = -2G_s \sigma, \quad U_M = G_s \sigma^2, \quad M = m_0 + \Sigma_s.$$

extended Z_3 trans.

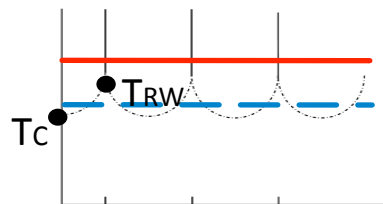
$$\theta \rightarrow \theta + 2\pi k/3, \quad \Phi(\theta) \rightarrow \Phi(\theta) e^{-i2\pi k/3}$$

修正版Polyakovループ $\Psi \equiv e^{i\theta} \Phi$

$$\theta \rightarrow \theta + 2\pi k/3, \quad \Psi(\theta) \rightarrow \Psi(\theta)$$

$$\Omega(\theta) = \Omega(\Psi(\theta), \Psi(\theta)^*, e^{3i\theta})$$

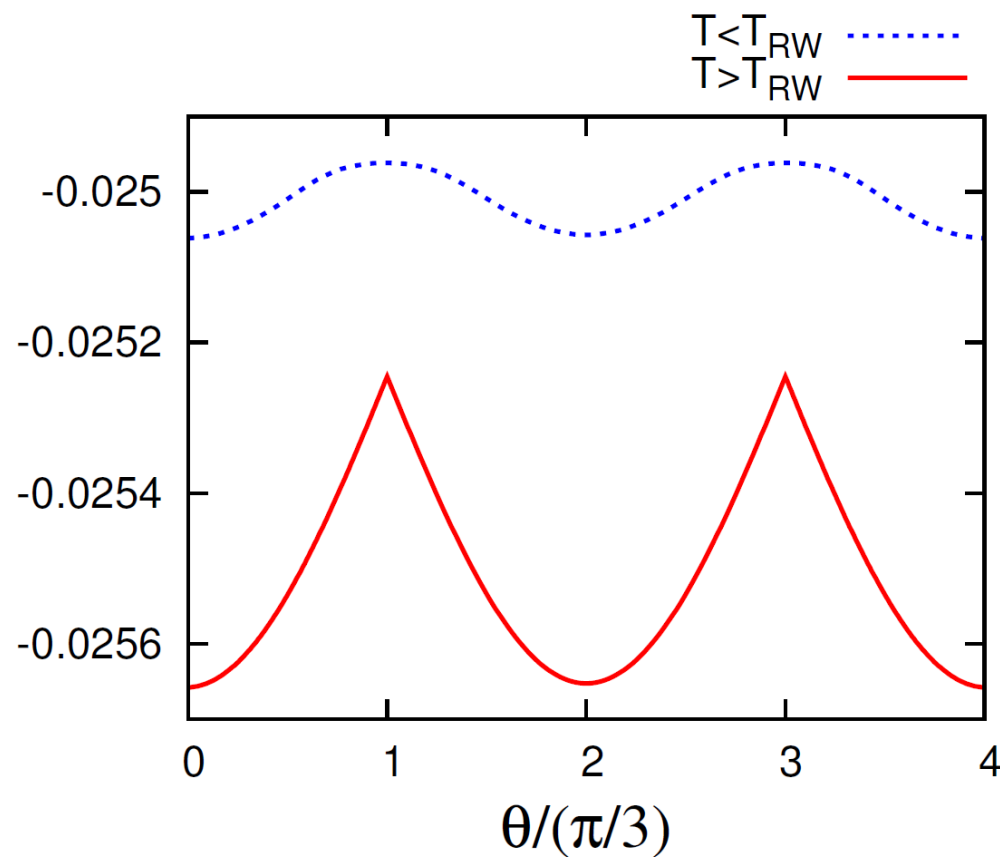
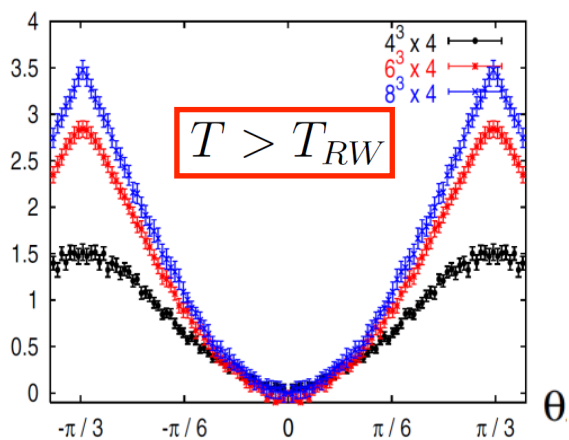
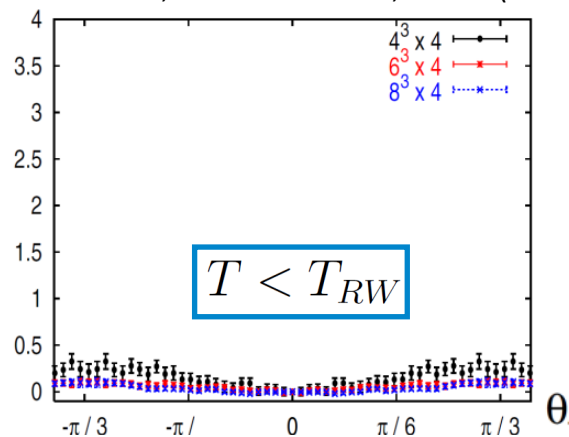
Extended Z_3 inv.



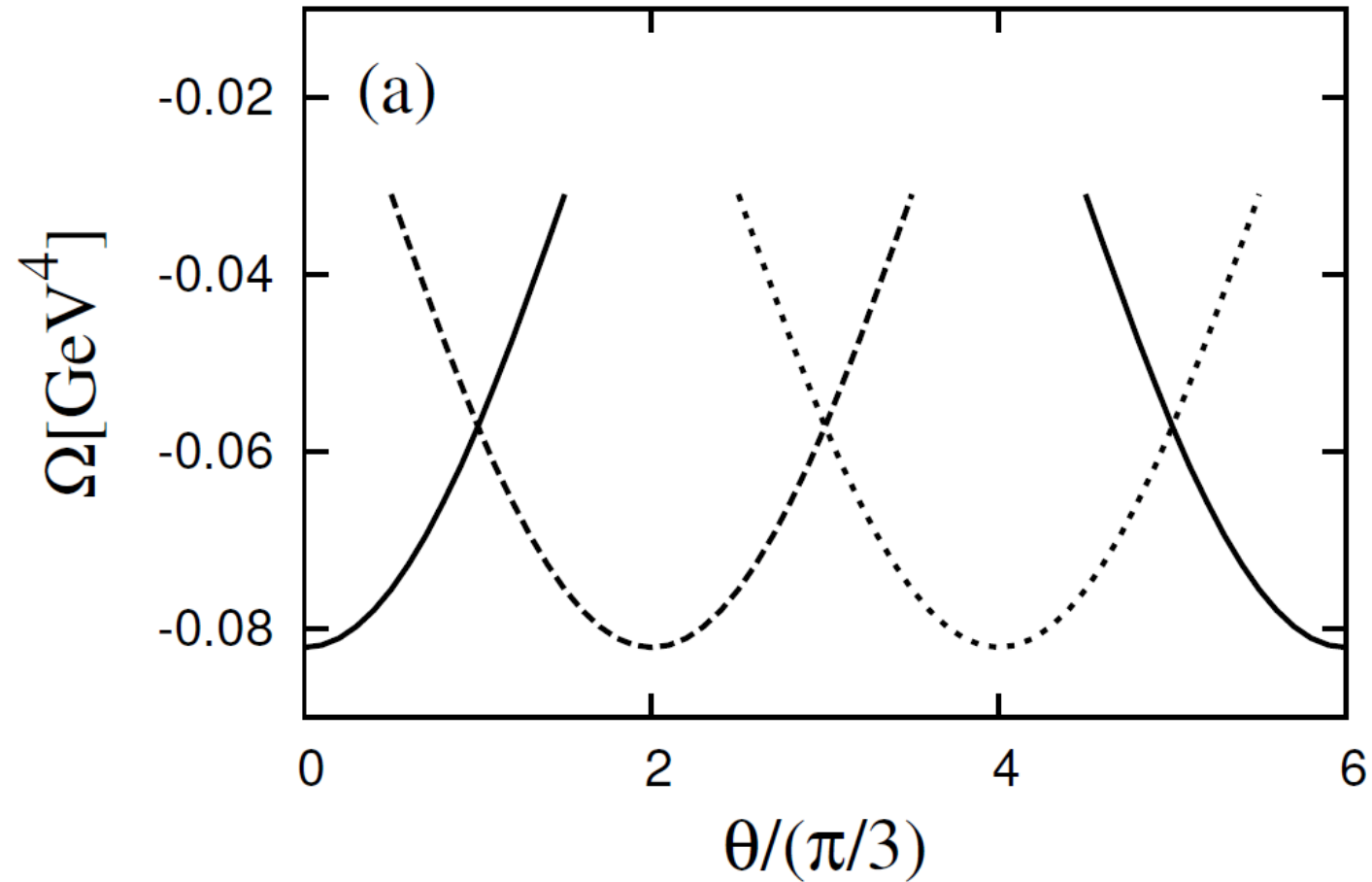
Roberg-Weiss periodicity and transition

$$\Omega(\theta) = \Omega(\theta + 2\pi k/3) = \Omega(-\theta).$$

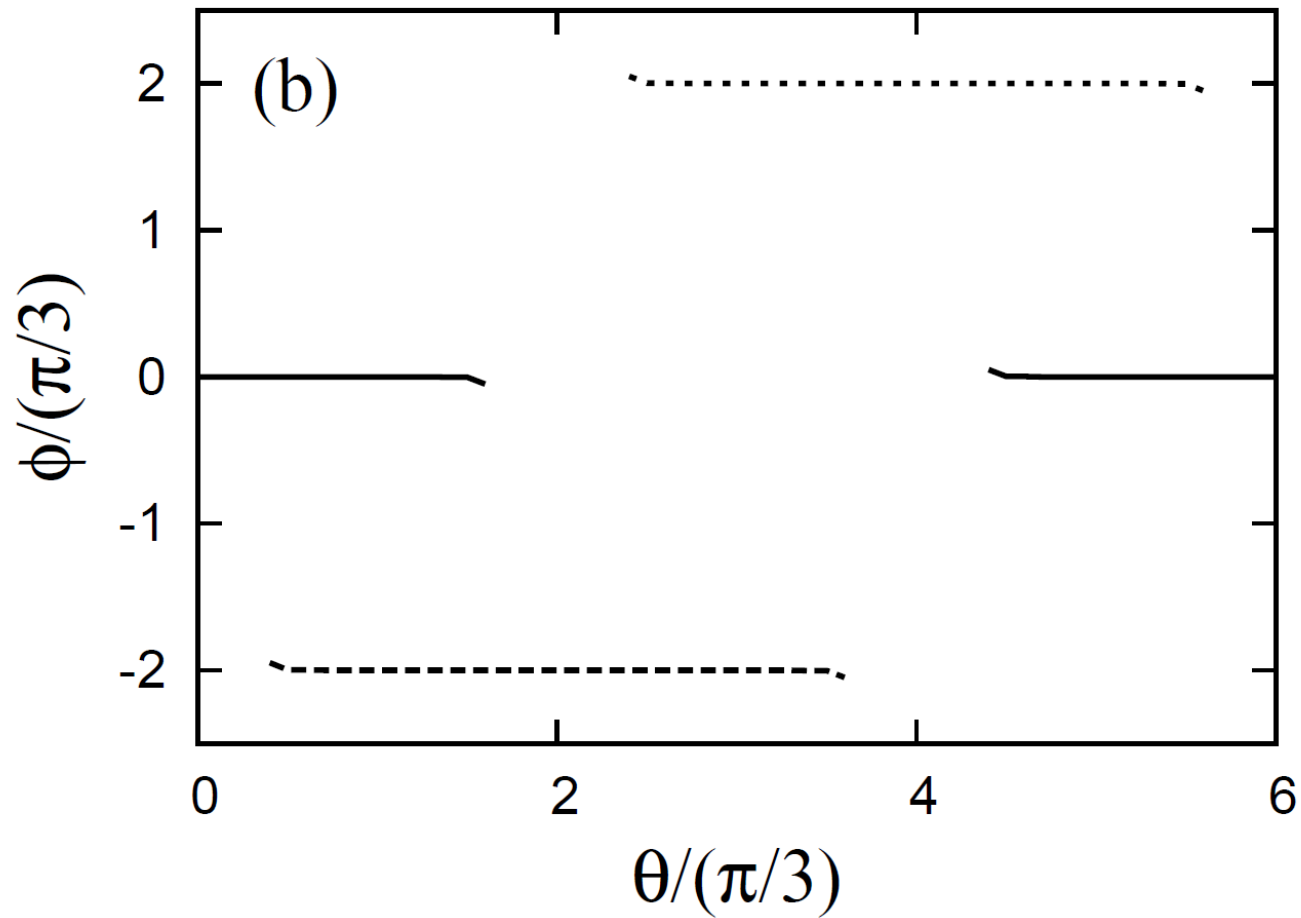
Kratochvila, Forcrand PRD**73**,114512(2006)



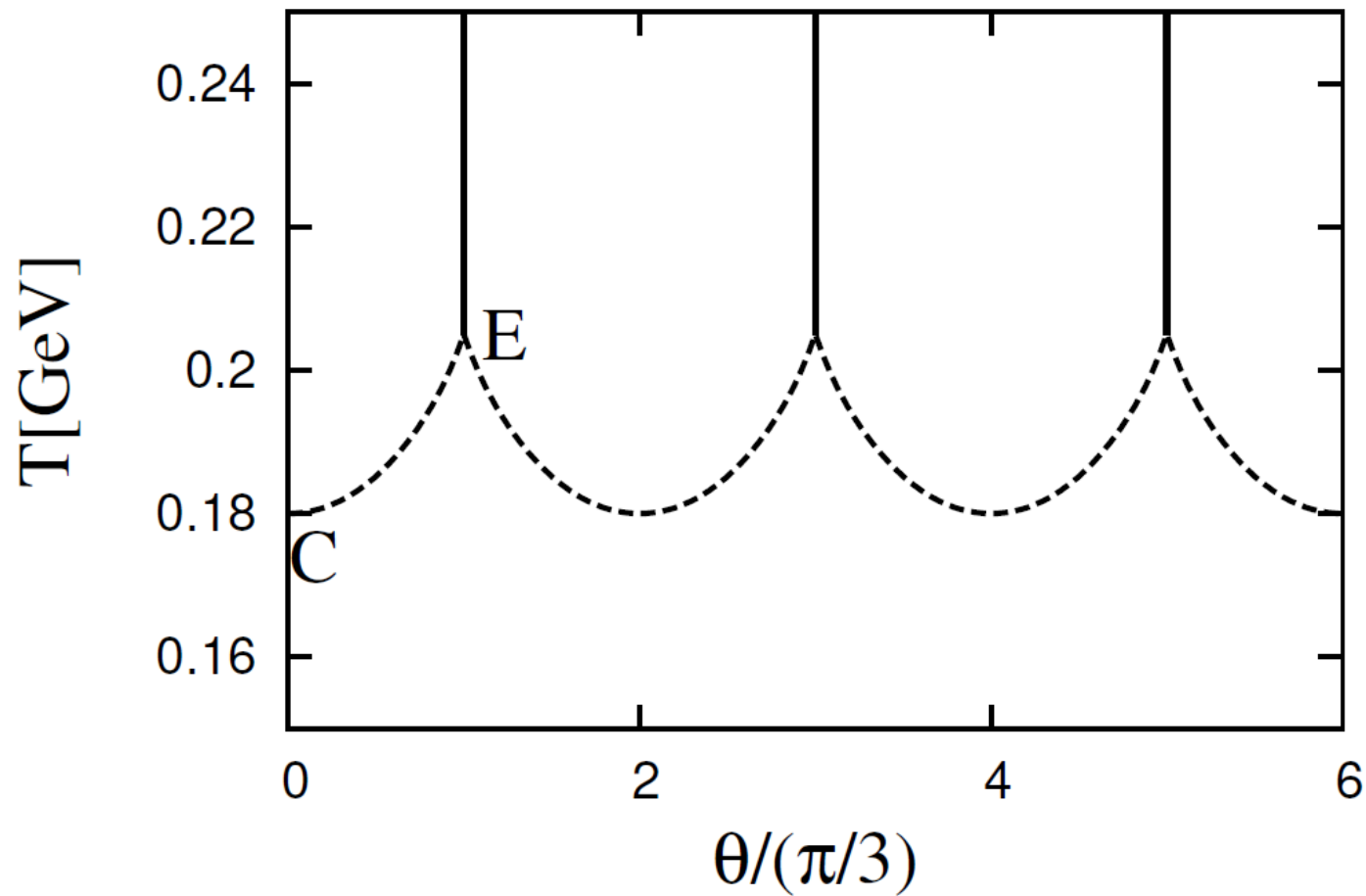
Thermodynamic potential



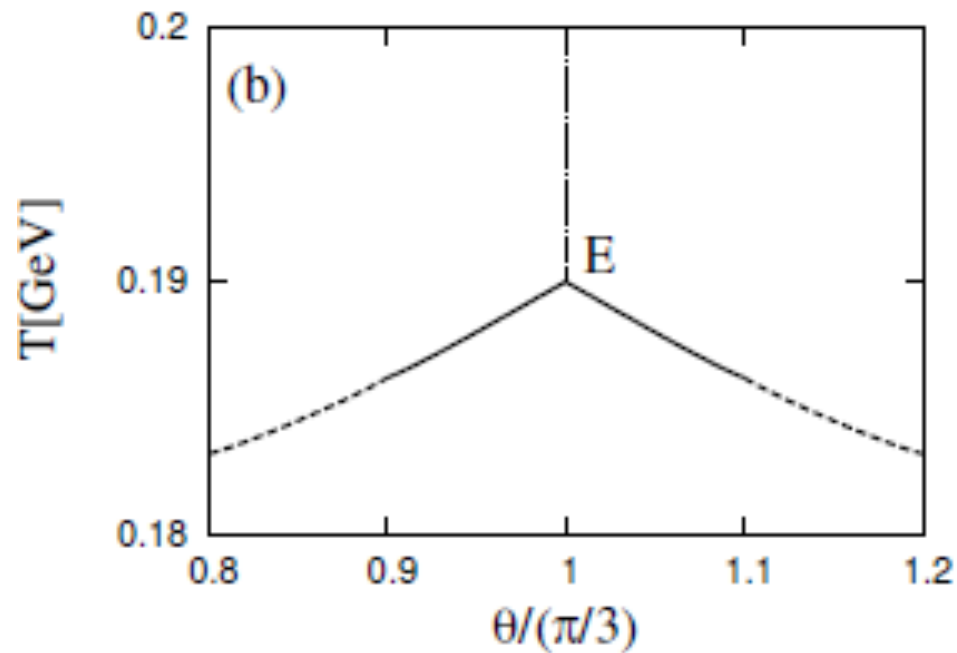
Phase of Polyakov loop



Phase diagram



Nontrivial CEP of deconfinement



Y. Sakai, H. K. and M. Yahiro
J. Phys. G: Nucl. Part. Phys. 37 (2010) 105007

RW transition

- θ -even quantities have cusp when $T > T_{RW}$, $\theta = (2k+1)\pi/3$ (k integer), while θ -odd quantities are discontinuous there.
- At $\theta = (2k+1)\pi/3$ and $T = T_{RW}$, Spontaneous breaking of C symmetry happens as T increases.

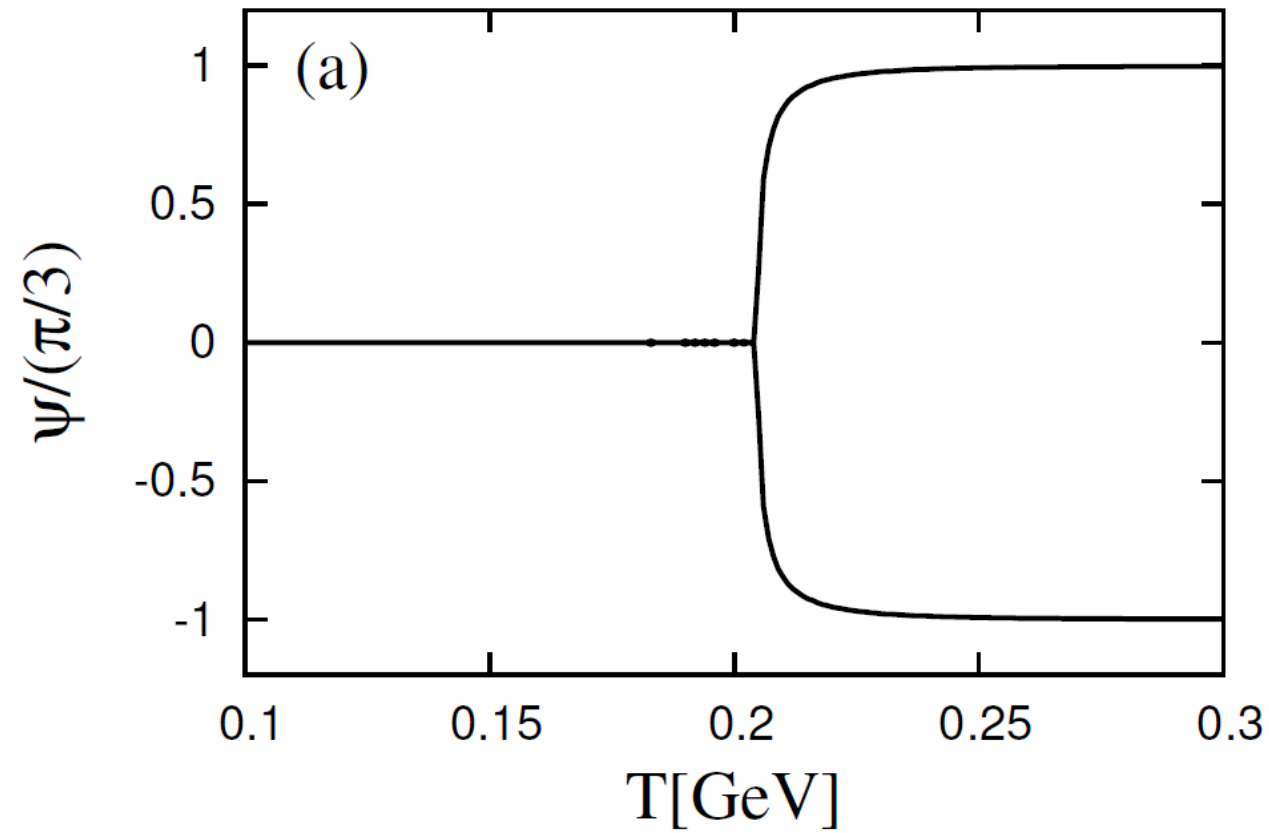
θ -odd quantities are order parameters.

Imaginary part or phase of modified Polyakov loop Ψ

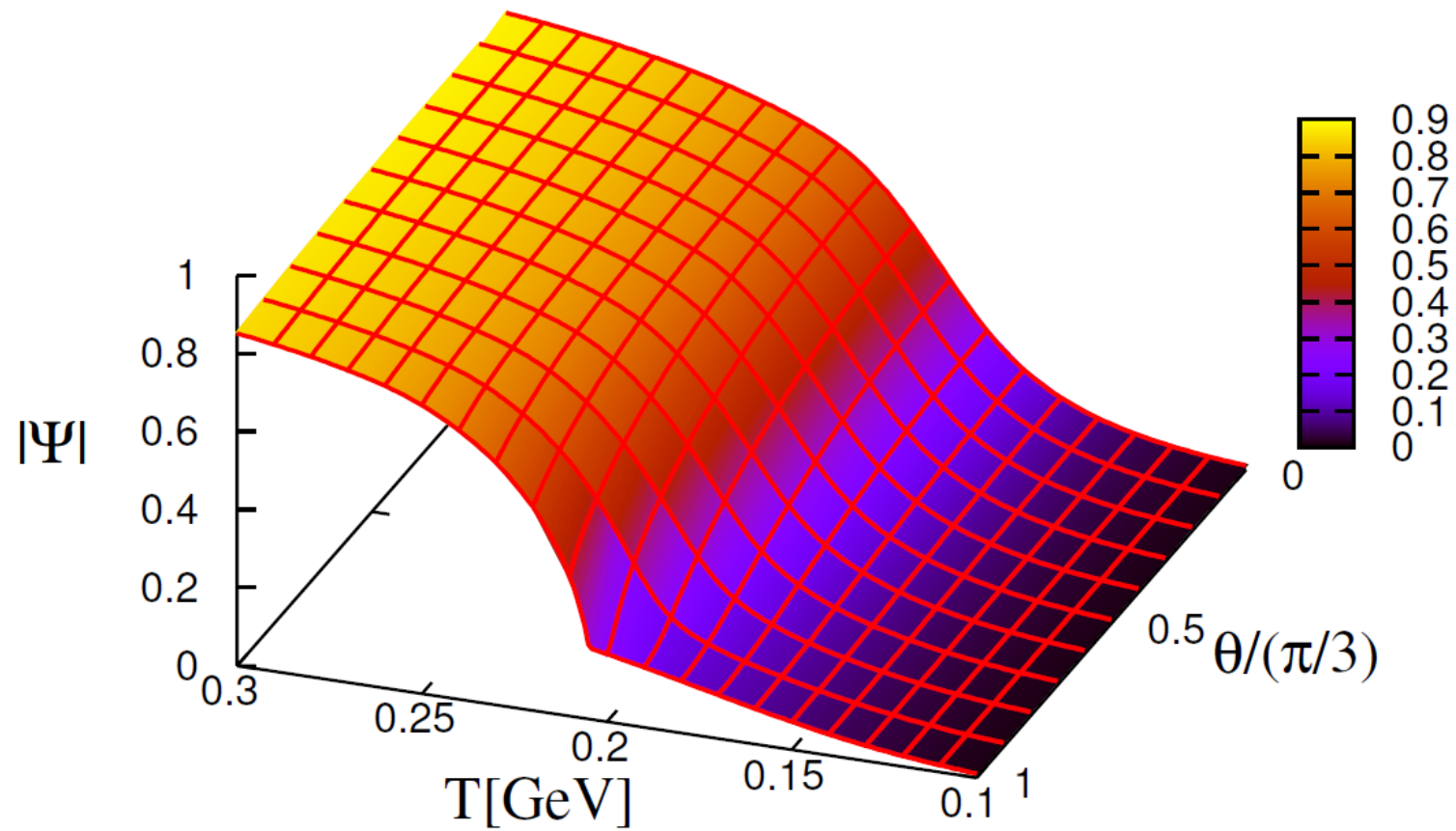
phase ψ of modified Polyakov loop Ψ

baryon number density (imaginary)

C violation at $\theta = \pi/3$



Deconfinement as remnant of C violation



Dashen mechanism

- CP violation in θ vacuum.

Dashen, Phys. Rev. D3, 1879(1971)

NJL version

Boer and Boomsma,

Phys. Rev. D78,054027 (2008)

Θ_v -term in two-flavor NJL

$$L_{NJL} = \bar{\psi} (i\gamma^\mu \partial_\mu - m_0 + \gamma_0 \mu) \psi + L_4 + L_\theta$$

$$L_4 = (1 - c)G[(\bar{\psi}\tau_a\psi)^2 + (\bar{\psi}\tau_a i\gamma_5\psi)^2]$$

$$L_\theta = ce^{i\theta_v} \det(\bar{\Psi}_R \Psi_L) + H.C.$$

PNJL model

- Replace derivative by D_μ
- Treat $A_4=iA_0$ as a back ground field.
- Add Polyakov potential $U(\Phi)$

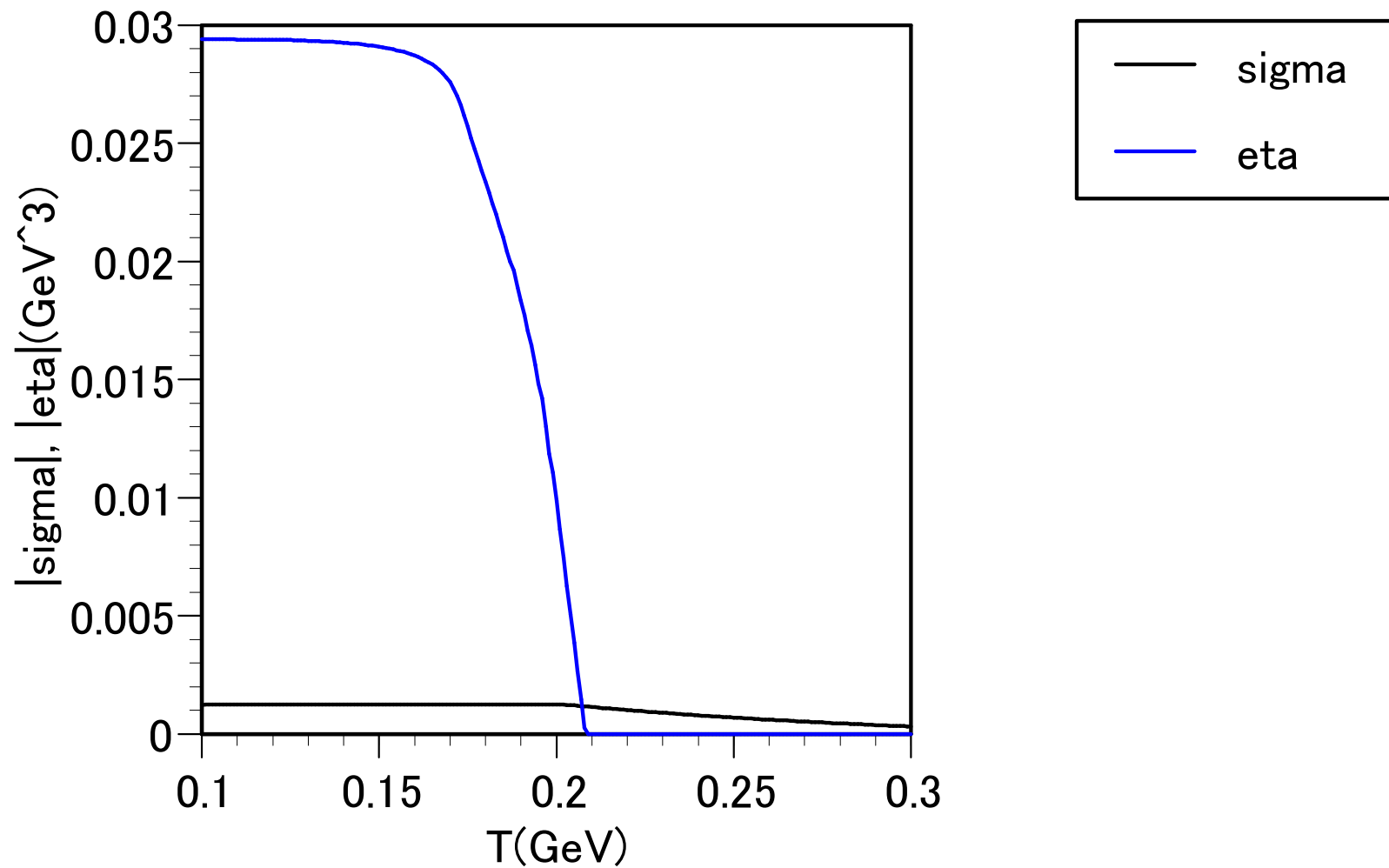
P violation and restoration at $\theta_v = \pi$

- For large c ($c > c_{\text{criti}}$), P symmetry is spontaneously broken at low temperature.
- P is restored at high T.
- P-odd quantities such as η are order parameters.

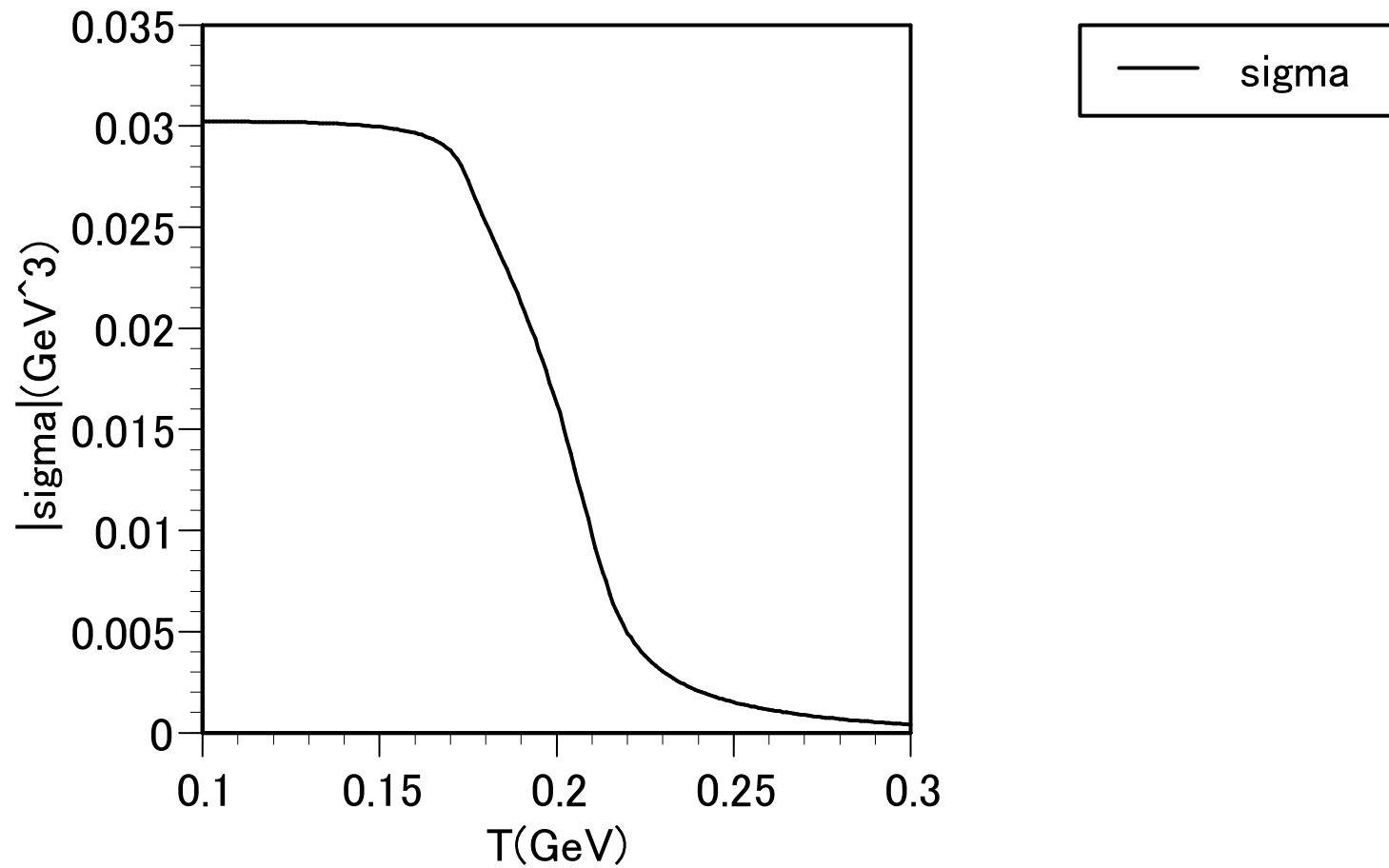
In broken phase,

$$\eta = \langle \bar{\psi} i \gamma_5 \psi \rangle \neq 0$$

P violation and restoration at $\theta_v = \pi$



Chiral breaking and restoration at $\theta_v=0$



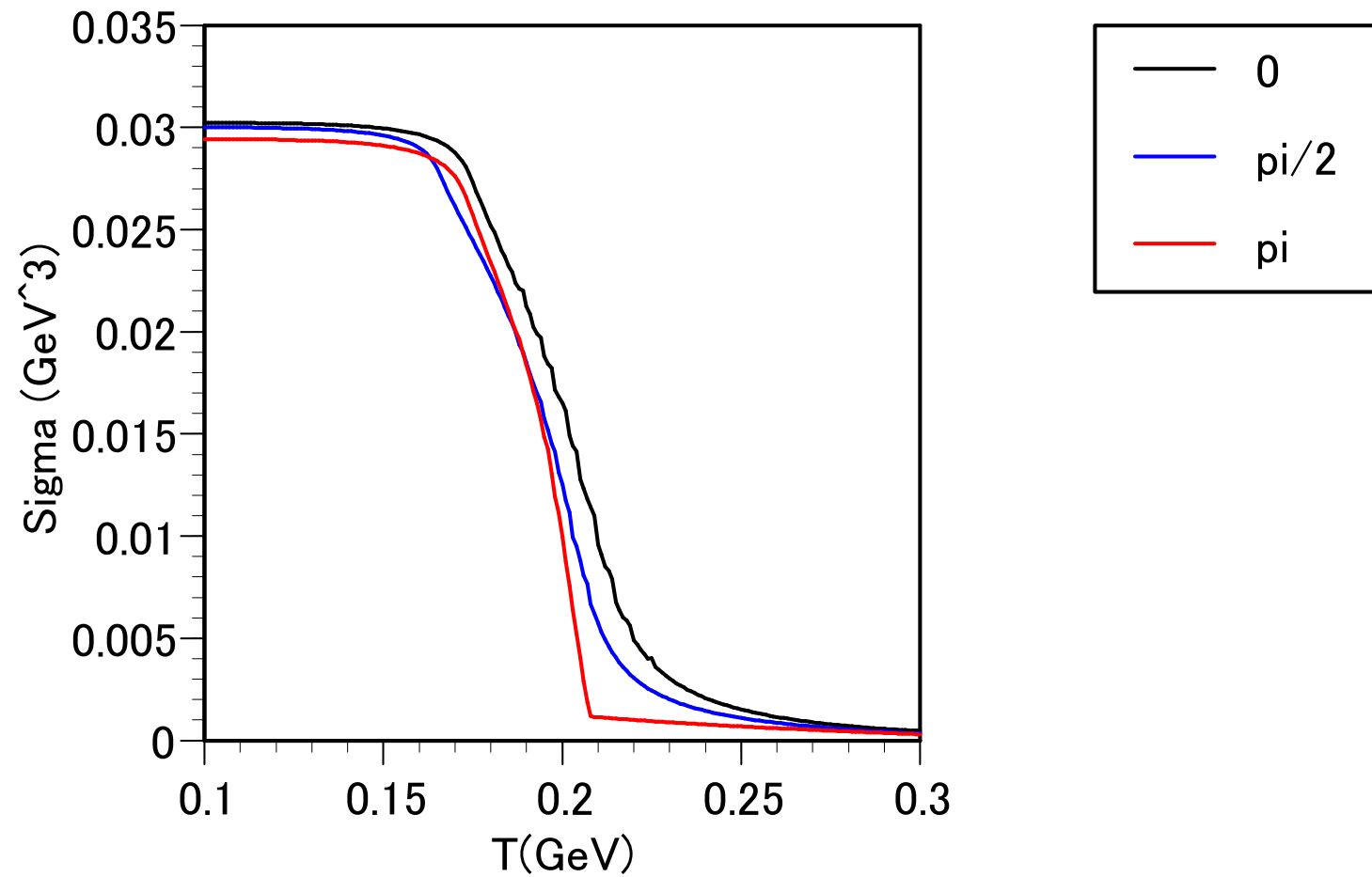
Chiral restoration as remnant of P restoration

Define $\Sigma = (\sigma^2 + \eta^2)^{1/2}$

Chiral restoration at $\theta_v = 0$ is a remnant of P restoration at $\theta_v = \pi$.

PNJL model can treat both C and P at the same time!! ($\theta_v = \pi$ and $\theta = \pi/3$)

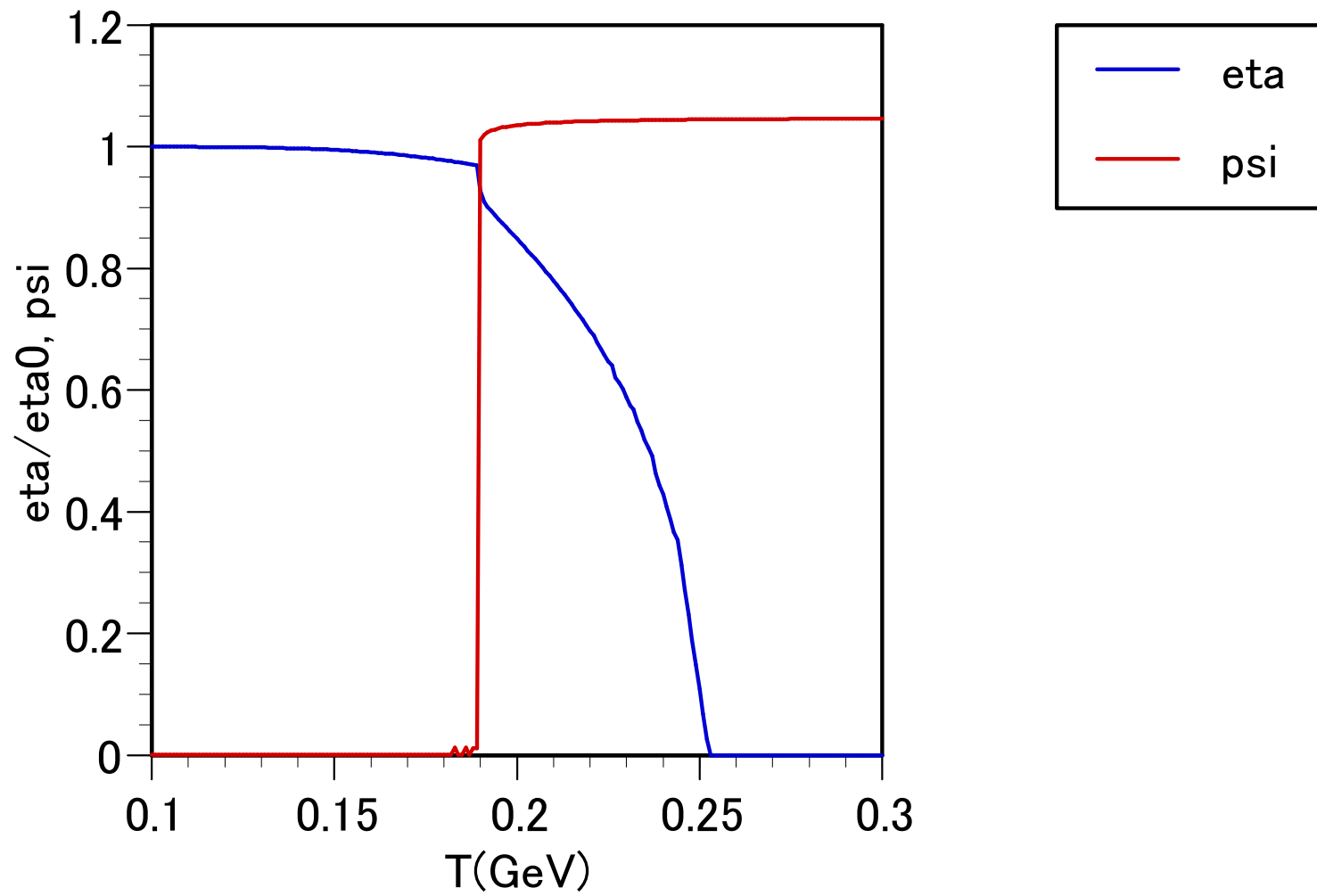
Chiral restoration as remnant of P restoration



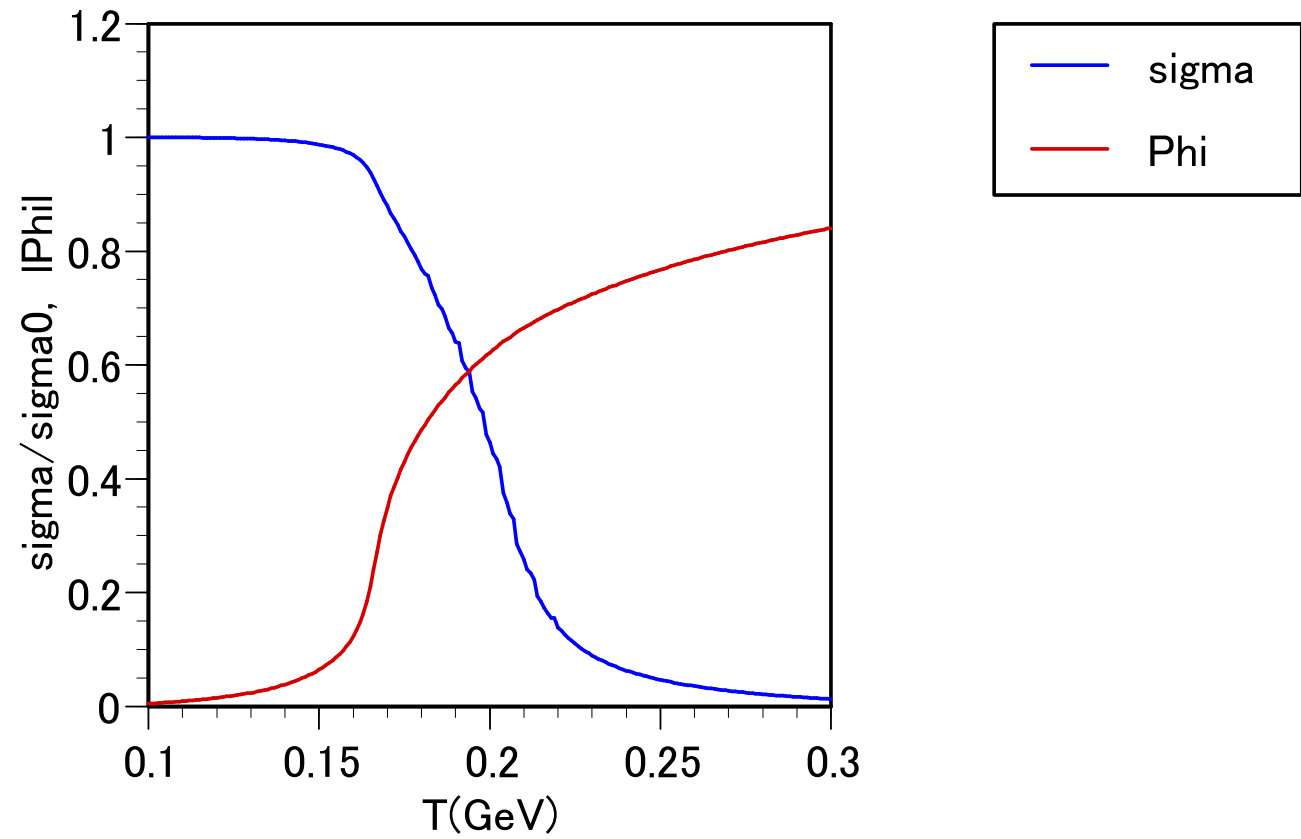
Dashen vs. RW

	Dashen	RW
Symmetry	P	C
Condition	At $\theta_v = \pi$	At $\theta = \pi/3$
Order parameter	η	ψ, n_q
Ω	has a cusp	has a cusp
Broken phase	low T	high T
Origin	anomaly instanton	term with $e^{3i\theta}$

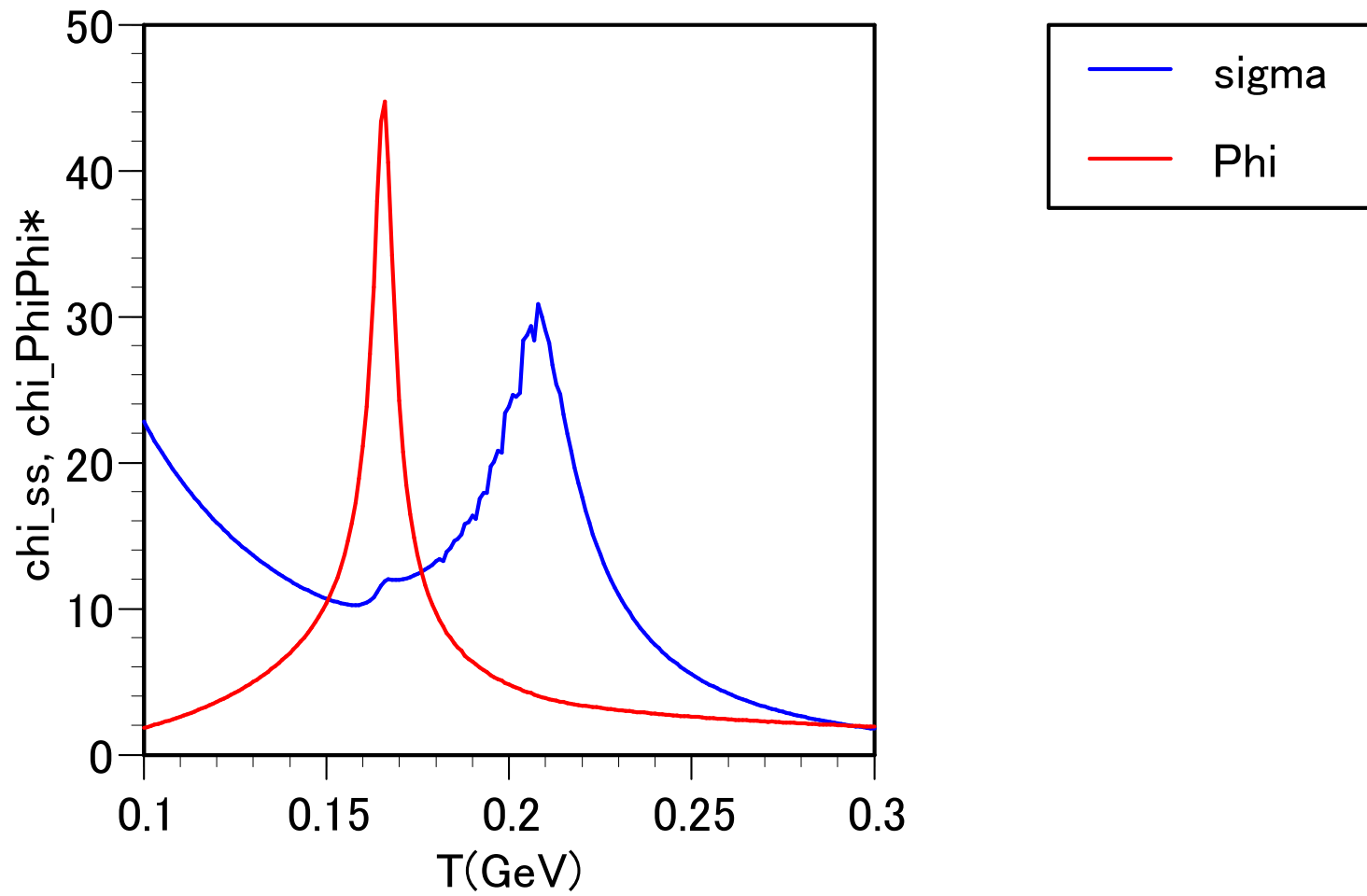
P restoration and C violation at $\theta=\pi/3$ and $\theta_v=\pi$ (θ world)



P restoration and C violation at $\theta=0$ and $\theta_v=0$ (Real World)



susceptibilities at $\theta=0$ and $\theta_v=0$ (Real World)



PC and CP

Real World

PC

Polyakov

Chiral

PC and CP

Real World

PC

Polyakov

Chiral



θ world

CP

C violation

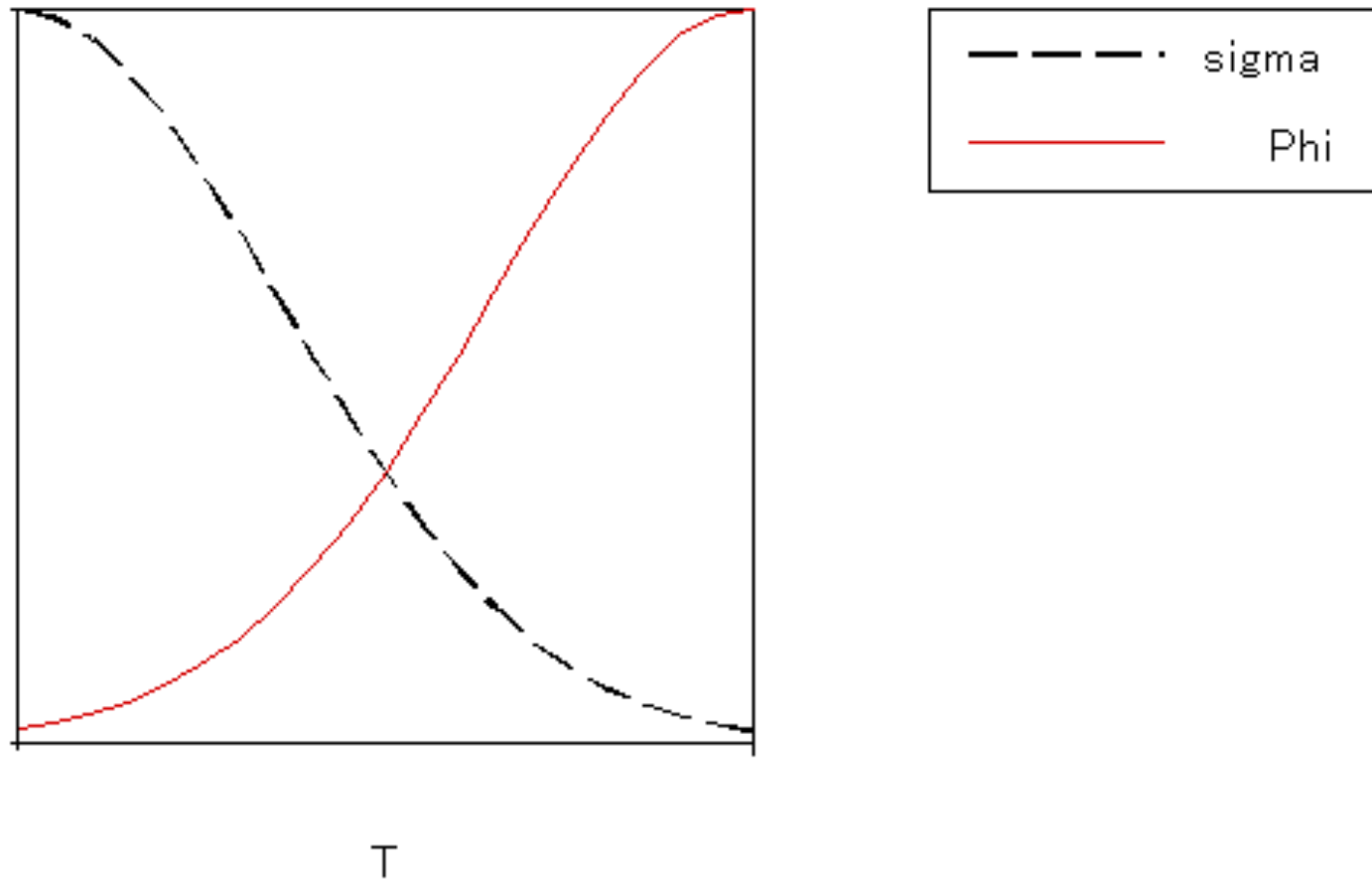
P restoration

Prediction

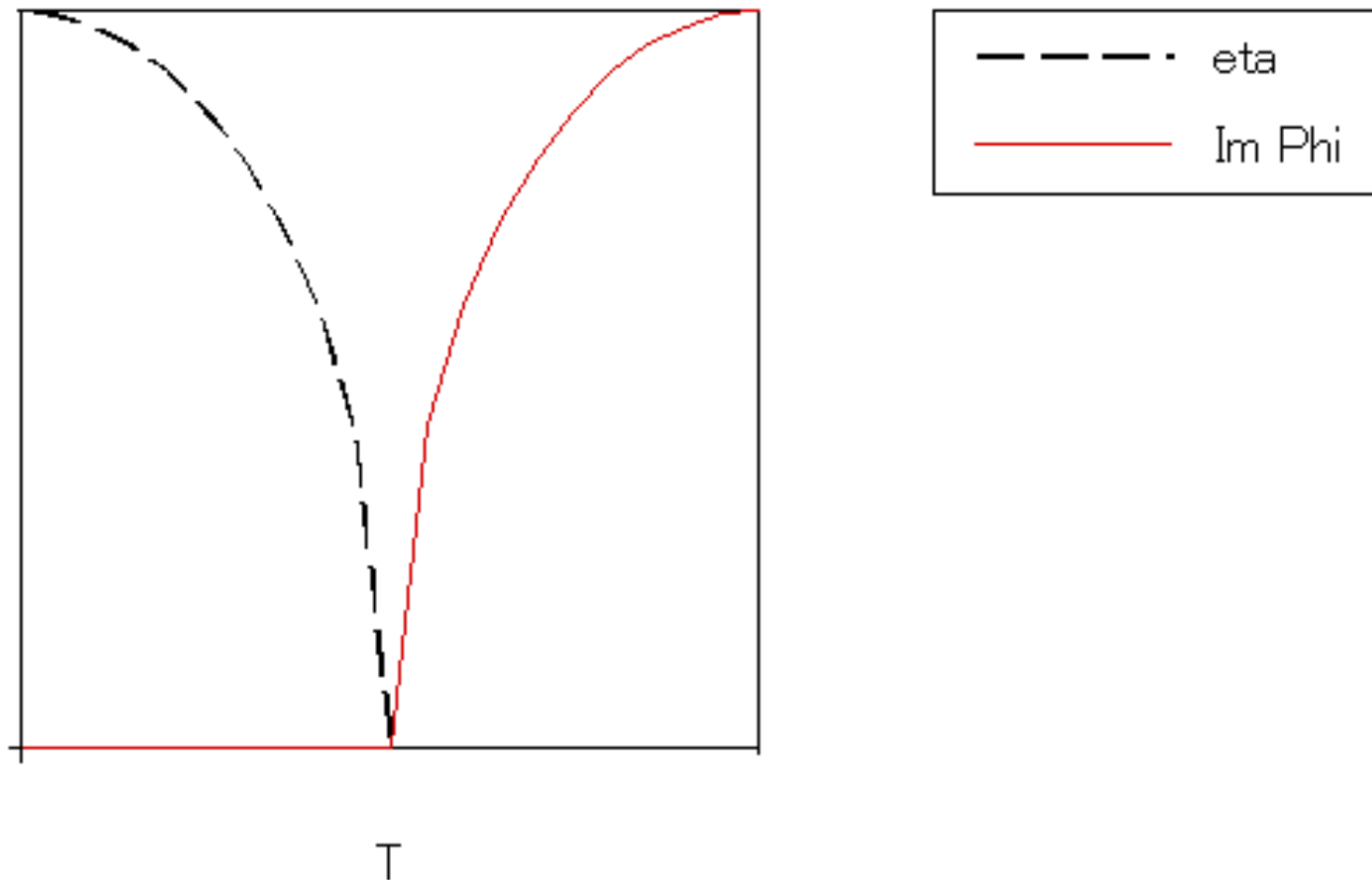
- If P and C happen almost simultaneously at real world,

C and P happen (almost) simultaneously at θ world ! ! ?

Real World

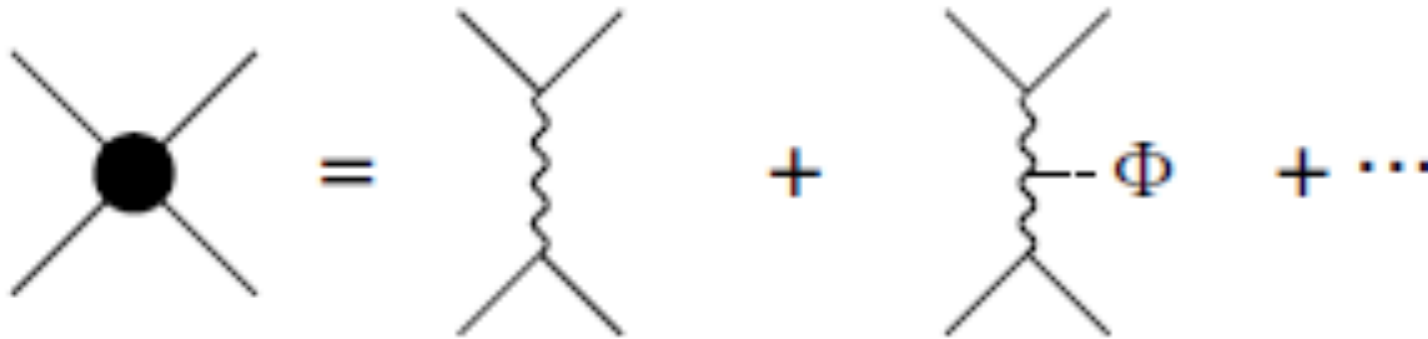


θ World

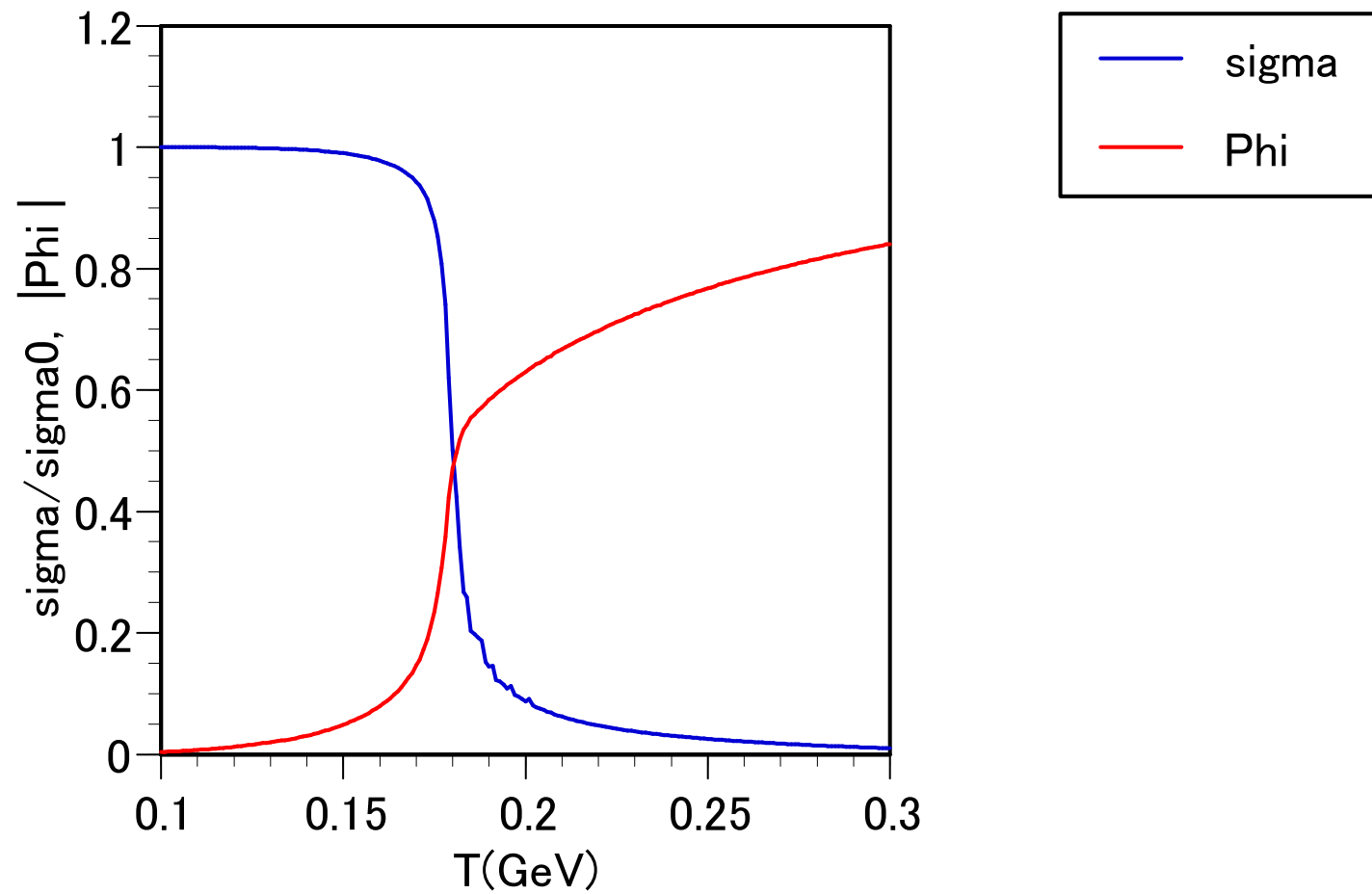


Entanglement PNJL model

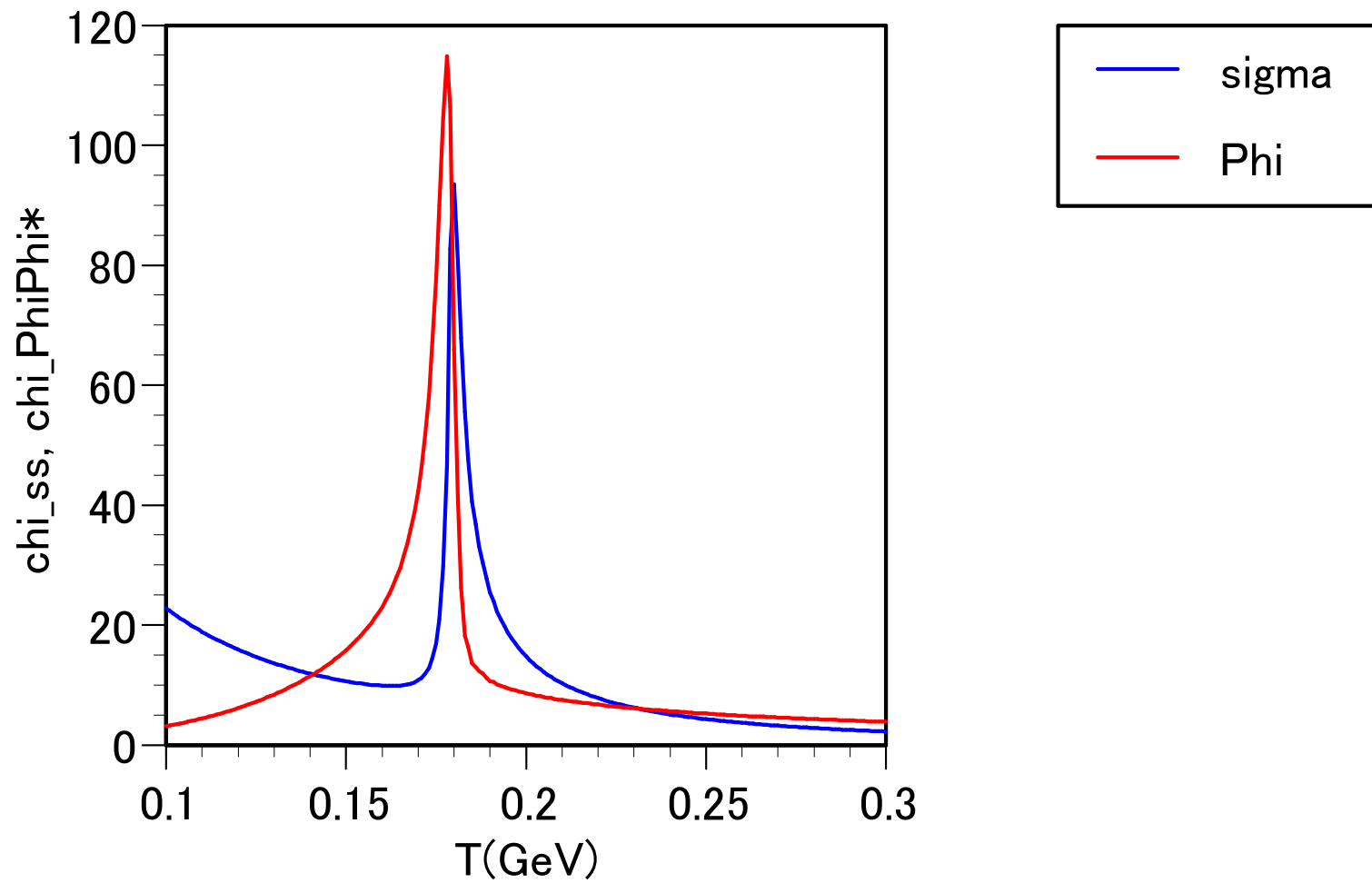
- $G_s \Rightarrow G_s[1 - \alpha \Phi \Phi^* - \beta(\Phi^3 + \Phi^{*3})]$
- Y. Sakai et al., arXiv:1006.3648[hep-ph]



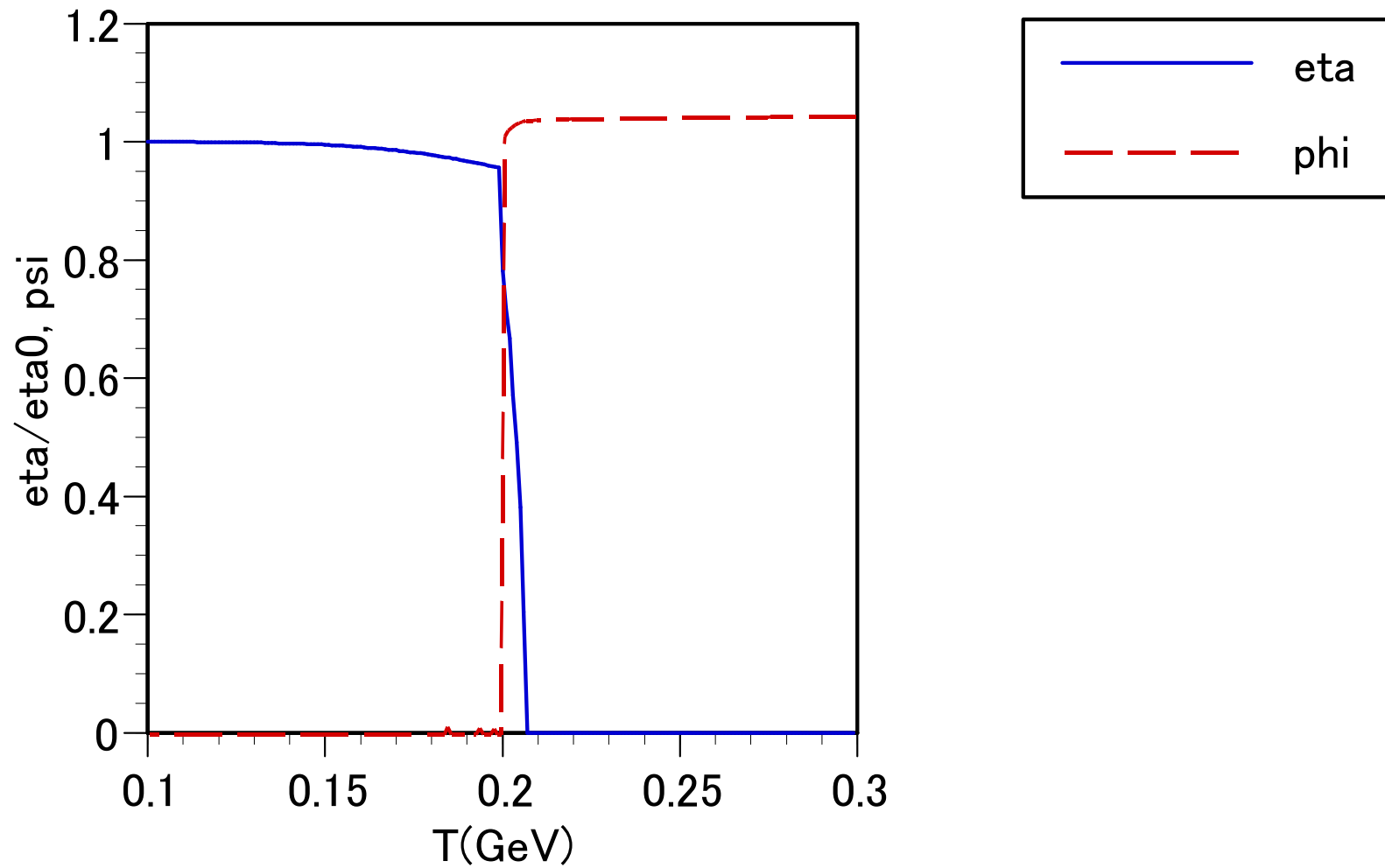
EPNJL at $\theta=0$ and $\theta=0$ (Real World)



Susceptibilities at $\theta=0$ and $\theta=0$ (Real World)



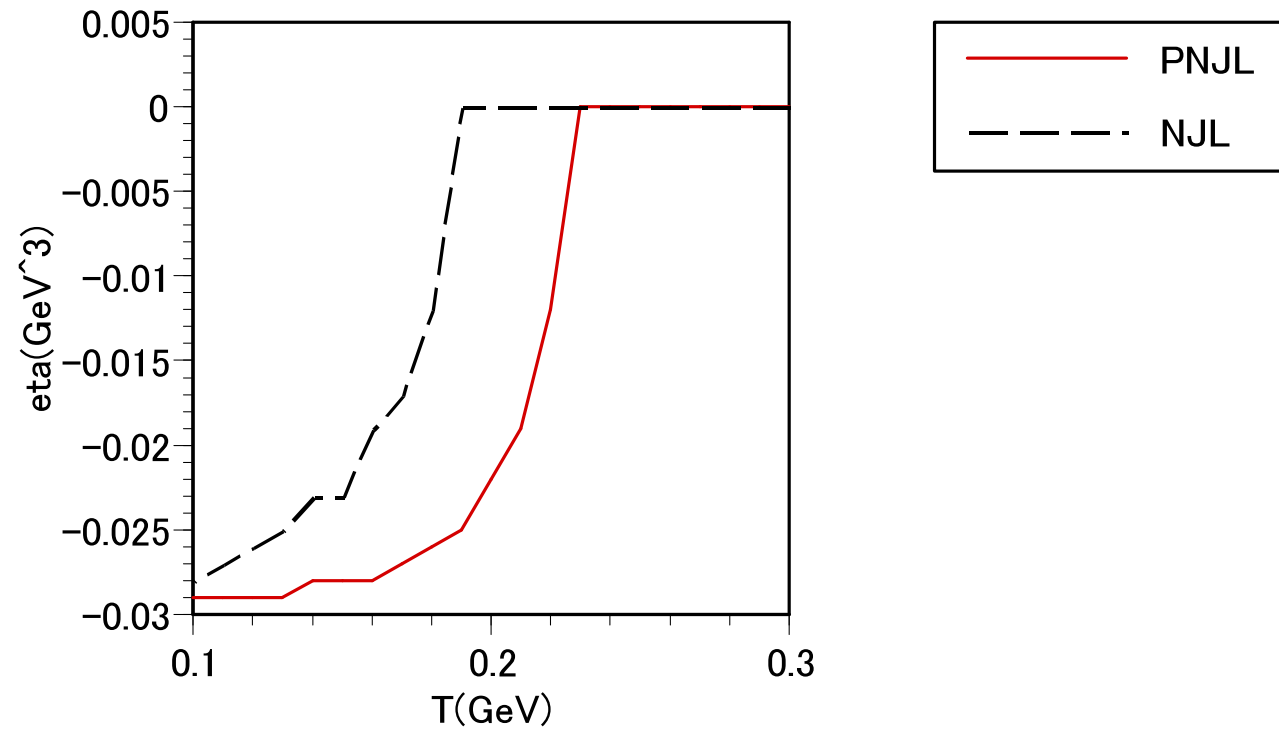
EPNJL at $\theta=\pi$ and $\Theta=\pi/3$



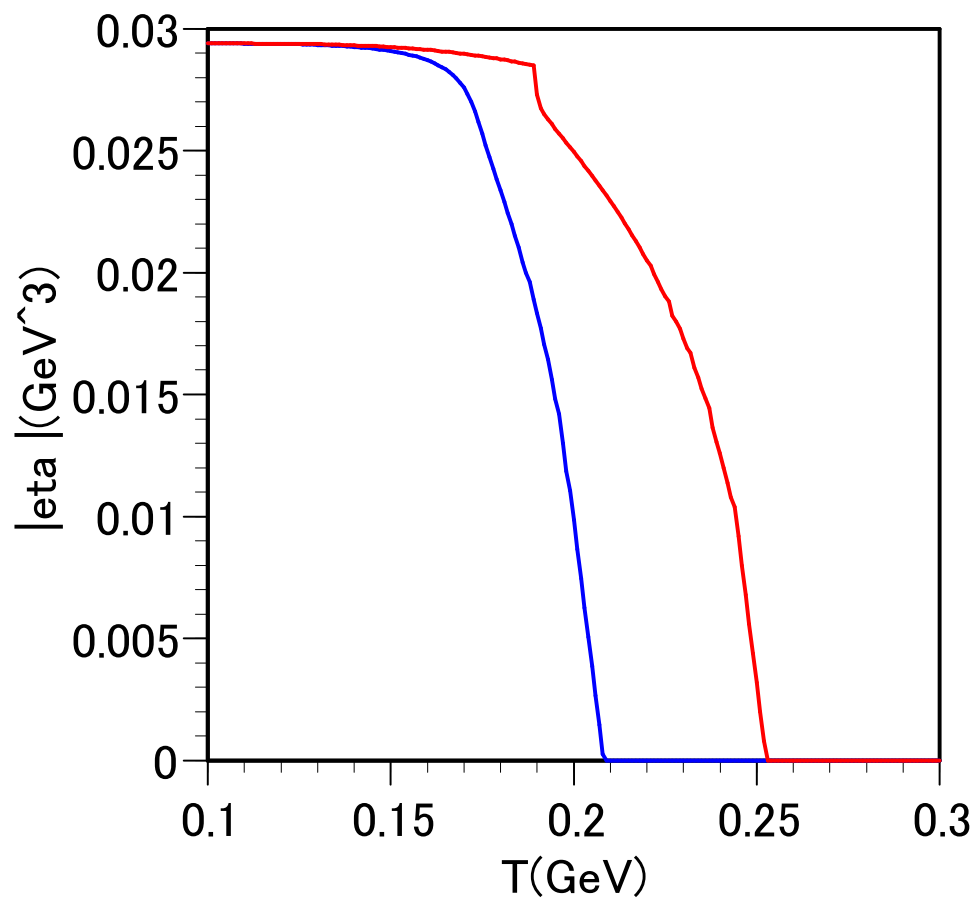
Summary

- ① Deconfinement is a remnant of C violation
- ② Chiral restoration is a remnant of P restoration
- ③ P and C happen almost simultaneously in the real world
 $\Leftrightarrow$ C and P may happen (almost) simultaneously in the θ world

NJL vs PNJL ($\theta=0, \theta_v=\pi$)



θ -dependence ($\theta_v=\pi$)



θ_v -dependence ($\theta=\pi/3$)

