

Electromagnetic radiation through transport process in quark-gluon plasma

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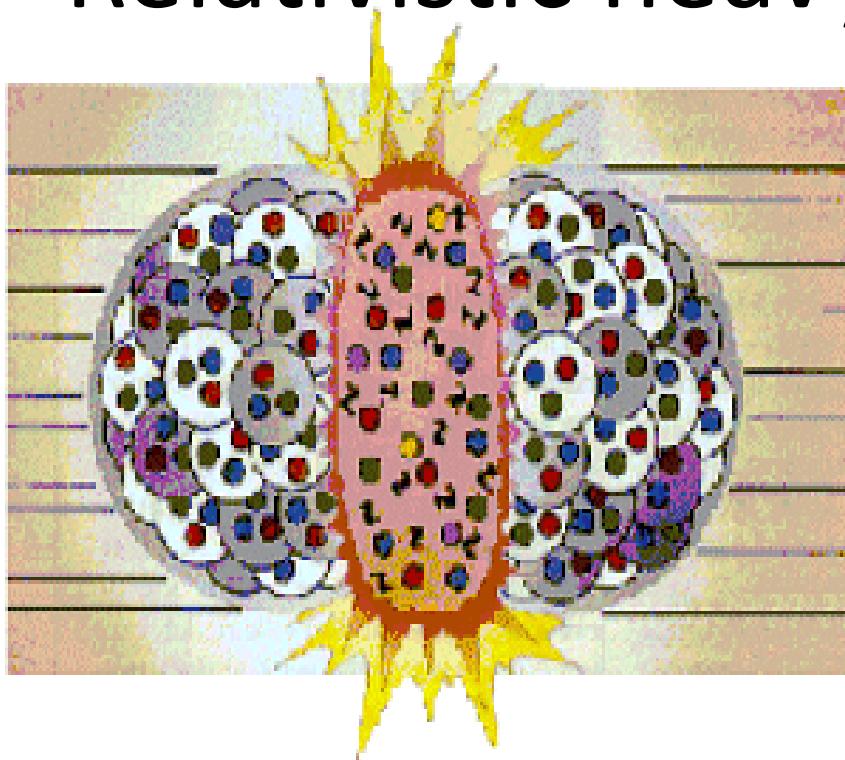
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1. Introduction

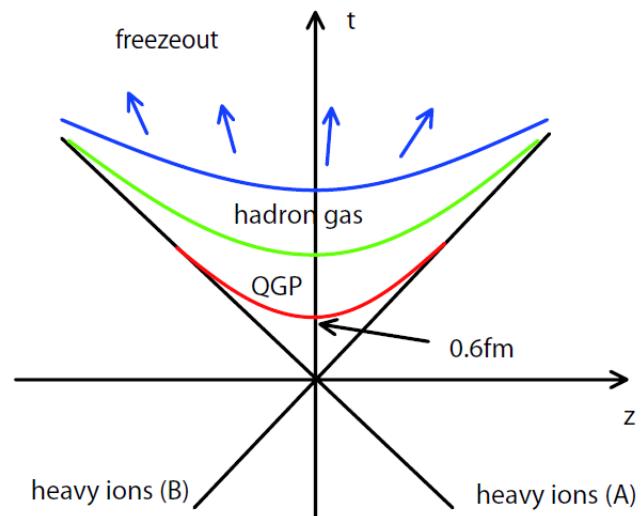
Relativistic heavy ion collisions



Hydrodynamic model gives
spacetime evolution profile
of the created matter.

RHIC
 $\text{Au+Au}, \sqrt{s_{\text{NN}}}=200\text{GeV}$

LHC
 $\text{Pb+Pb}, \sqrt{s_{\text{NN}}}=5.6\text{TeV}$
(2.76TeV at present)



Dilepton production at PHENIX

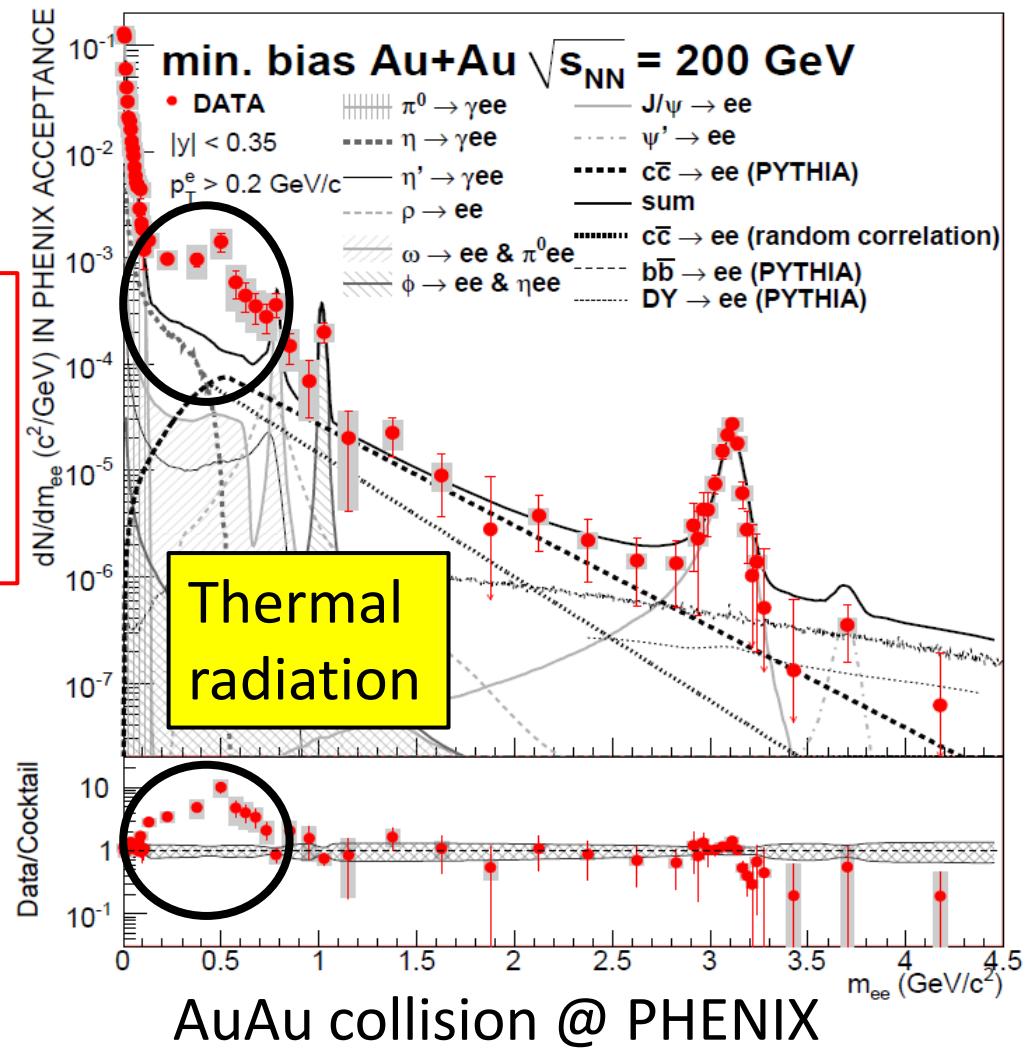
Low-mass ($m_{ee} < 0.6 \text{ GeV}$)
enhancement has been a
challenge to theorists:

Rapp ('01,'10)
Dusling, Zahed ('07,'09)
Bratkovaskaya, Cassing, Linnyk ('09,'10)
Ghosh, Sarkar, Alam ('10)
...

Hadron interaction, chiral symmetry
(Hadronic phase) / pQCD(QGP phase)

↓ Successful at SPS
but not at RHIC

Non-perturbative process
in QGP phase is important.



2. Spectral function from causal dissipative hydrodynamics

Formula for thermal radiation

Retarded correlator (or spectral function) of QCD-EM current

$$\frac{E_1 E_2 dN_{ee}}{d^3 p_1 d^3 p_2 d^4 x} = \frac{2e^4 L_{\mu\nu}(p_1, p_2)}{(2\pi)^6 q^4} \text{Im } G_R^{\mu\nu}(q; T) f_{BE}(q^0; T), \quad q^\mu \equiv p_1^\mu + p_2^\mu.$$
$$G_R^{\mu\nu}(\omega, k) \equiv \int d^4 x e^{iqx} i\theta(t) \langle [J^\mu(x), J^\nu(0)] \rangle_T$$

What's in the spectral function?

- Large ω : Towers of higher resonances/Quark pair annihilation
- Intermediate ω : Vector correlation (vector meson/qqbar)
- Small ω : Transport phenomena

Charge transport and dilepton

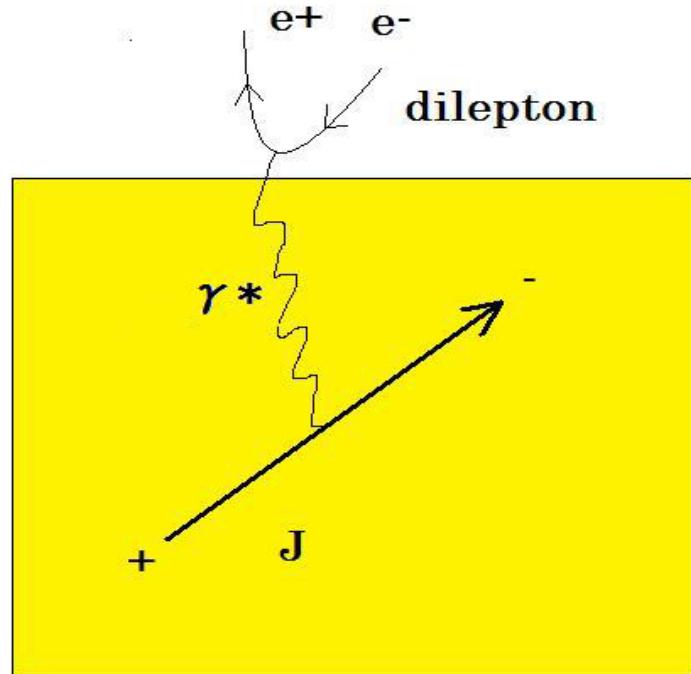
In each fluid element:

1. charge fluctuation → current
2. current → dileptons

*Remark:

net charge = 0 at RHIC/LHC

→ only induced current flows



Spectral function at small ω & k

Linear response theory:

$$\delta H(t) = \int d^3x J^\mu(x) \delta A_\mu(x)$$

$$\rightarrow \langle \delta J^\mu(q) \rangle_T = -G_R^{\mu\nu}(q;T) \delta A_\nu(q)$$

2nd order relativistic dissipative hydrodynamics Israel ('76)
 in external electromagnetic field: Israel, Stewart ('79)

$$J^\mu(x) = (\delta\rho(x), \vec{v}(x)), \quad \partial_\mu J^\mu = 0,$$

$$\vec{v}(x) = \sigma \vec{E}(x) - D \vec{\nabla} \delta\rho(x) - \tau_J \frac{\partial \vec{v}}{\partial t},$$

$$\sigma \equiv \chi D, \quad \chi \equiv \frac{\partial \rho}{\partial \mu}.$$

Ohm's law, Fick's law
 Memory effect



Linear response theory

$$\text{Im } G_R^{(L)}(q;T) = -\frac{\chi D \omega q^2}{\omega^2 + (\tau_J \omega^2 - D k^2)^2},$$

$$\text{Im } G_R^{(T)}(q;T) = -\frac{\chi D \omega}{\tau_J^2 \omega^2 + 1}.$$

Spectral shape & strength

Spectral shape:

$$\text{Im} G_R^{(L)}(q;T) = -\frac{\chi D \omega q^2}{\omega^2 + (\tau_J \omega^2 - D k^2)^2},$$

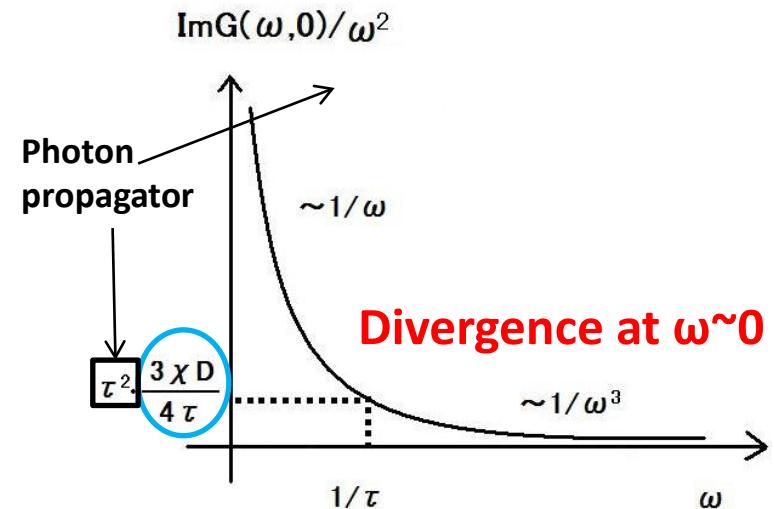
$$\text{Im} G_R^{(T)}(q;T) = -\frac{\chi D \omega}{\tau_J^2 \omega^2 + 1},$$

$$\text{Im} G_R^{\mu\mu} = 2 \text{Im} G_R^{(T)} + \text{Im} G_R^{(L)}$$

Stronger at larger χ , D and smaller τ :

- Susceptibility χ : how large the charge fluctuation is.
- Diffusion constant D: how effectively current is induced.
- Relaxation time τ_J : how swiftly current is induced.

induced current: $\vec{v}(x) = \sigma \vec{E}(x) - D \vec{\nabla} \delta \rho(x) - \tau_J \frac{\partial \vec{v}}{\partial t}$



Parameterization

$$D \propto 1/T, \tau_J \propto 1/T.$$

	D	τ_J	
pQCD	$4/T$	$15/T$	Arnold et al. ('00,'03), Hong, Teaney ('10)
AdS/CFT	$1/2\pi T$	$\ln 2/2\pi T$	Natsuume, Okamura ('08)

$$\chi(T) = 0.28T^2 \left[1 + \tanh\left(\frac{T - 0.155\text{GeV}}{0.023\text{GeV}}\right) \right].$$

$\chi(T)$: Using lattice result for quark number susceptibility

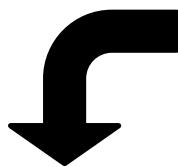
Allton et al. ('05)

3. Dilepton production at RHIC (PHENIX)

Hydro model setting

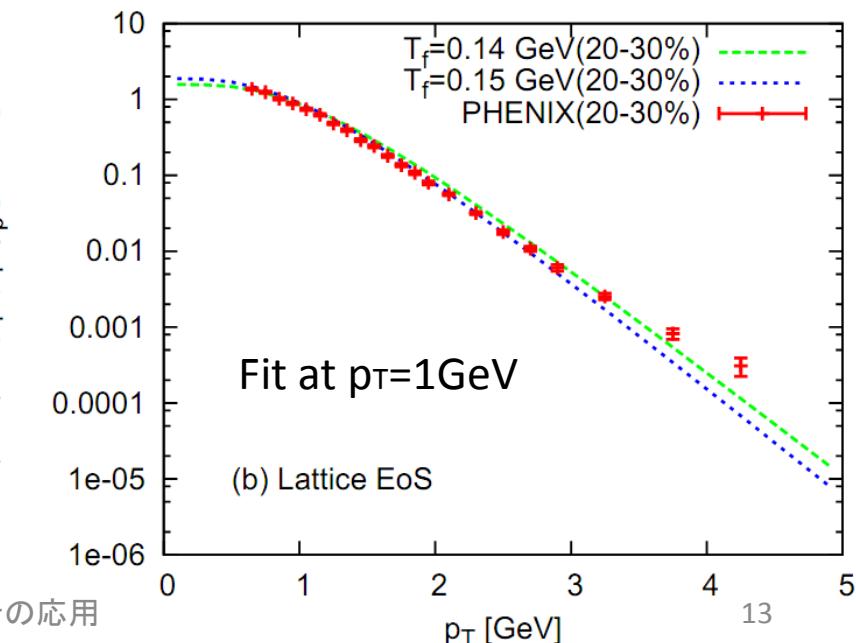
Viscosity	EoS	Initial	Hadronic phase
Perfect fluid	Lattice	Modified BGK	Chemical equilibrium

How far the matter
evolves under hydrodynamics



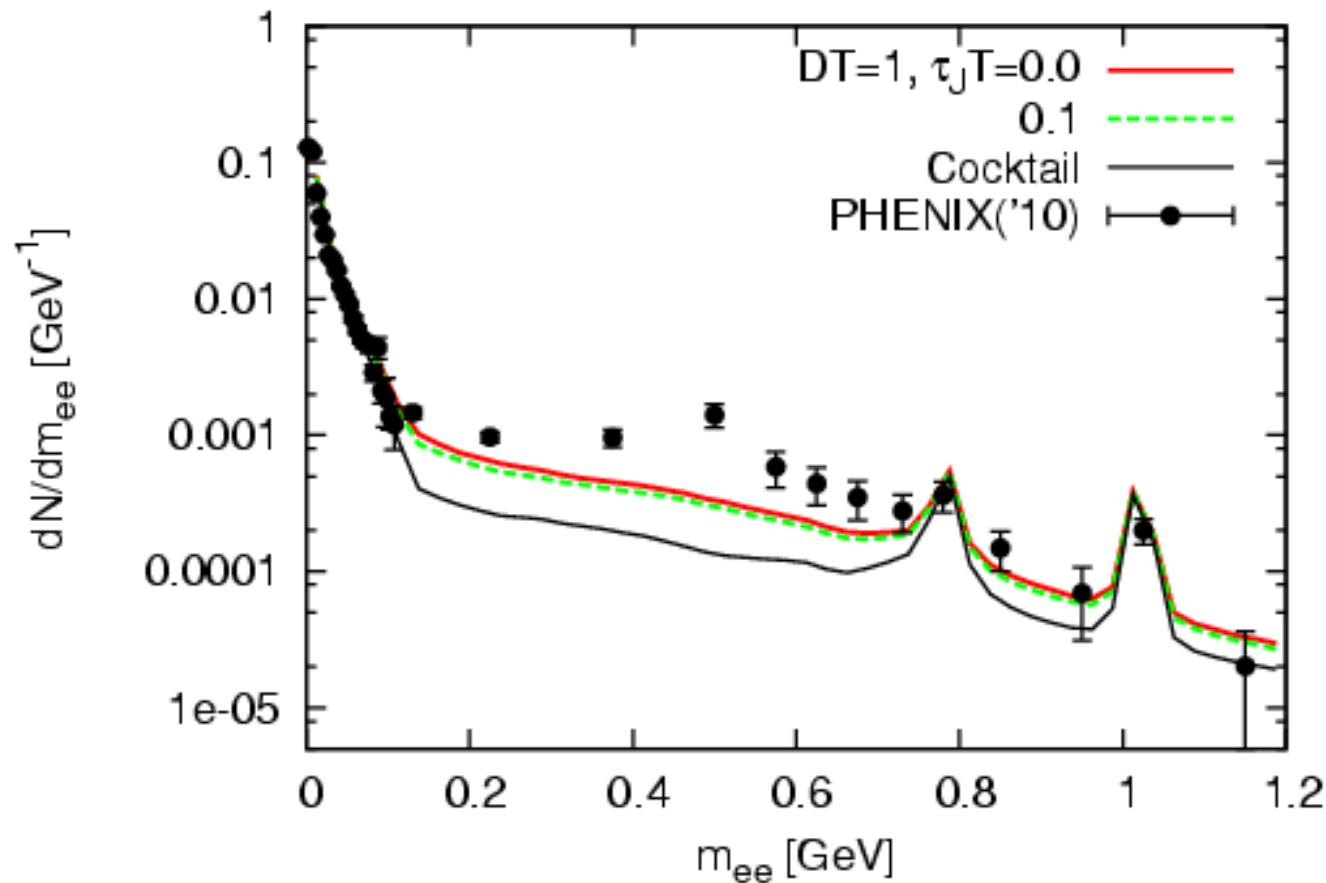
Until $T_f = 0.15 \text{ GeV}$

Proton spectra:
theory vs. experiment



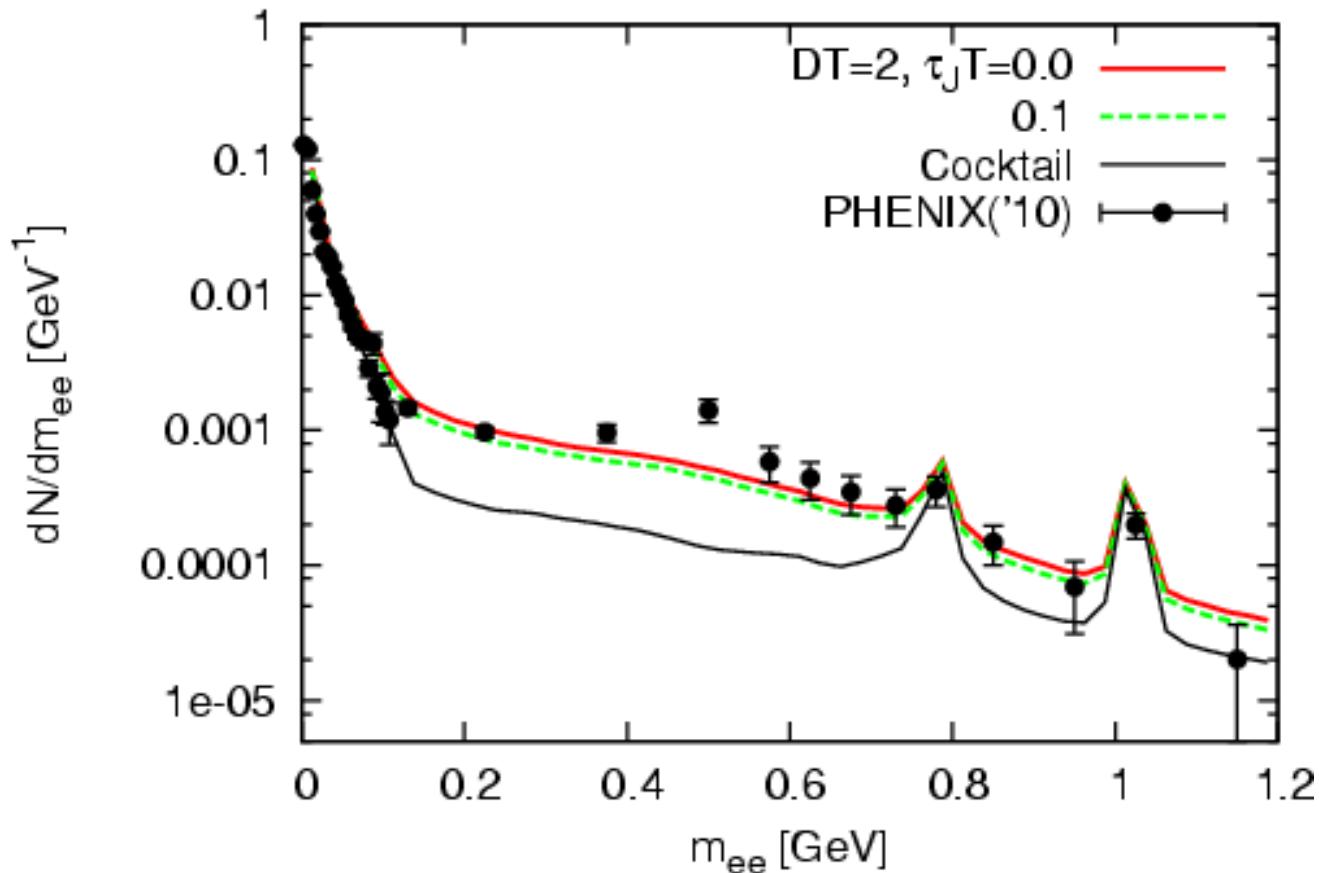
Invariant mass spectra

DT=1 : no solution



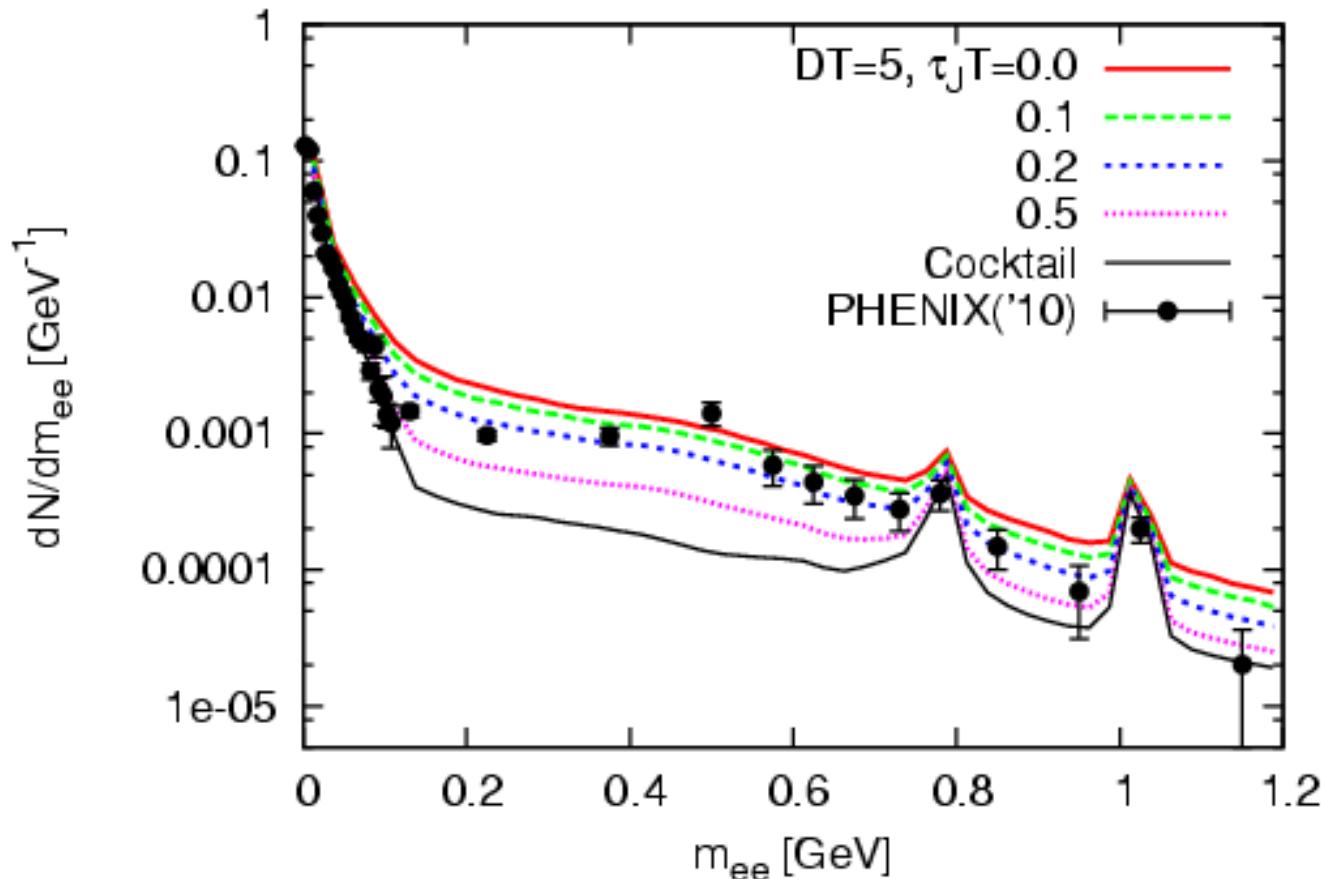
Invariant mass spectra

$(DT, \tau_J T) = (2, 0-0.1)$: good agreement



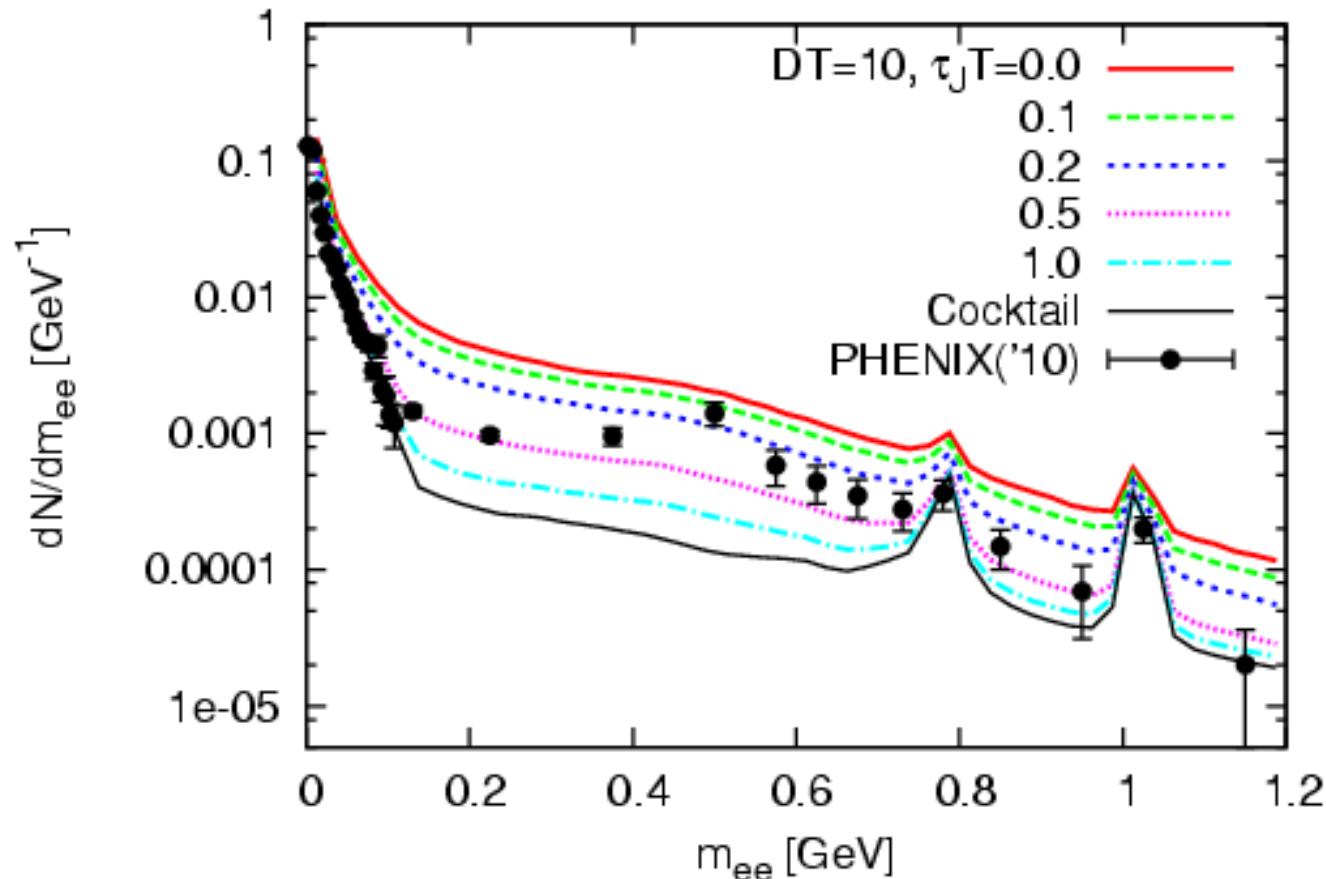
Invariant mass spectra

$(DT, \tau_J T) = (5, 0.2), (2, 0-0.1)$: good agreement



Invariant mass spectra

$(DT, \tau_J T) = (10, 0.5), (5, 0.2), (2, 0-0.1)$: good agreement
 $DT=1$: no solution



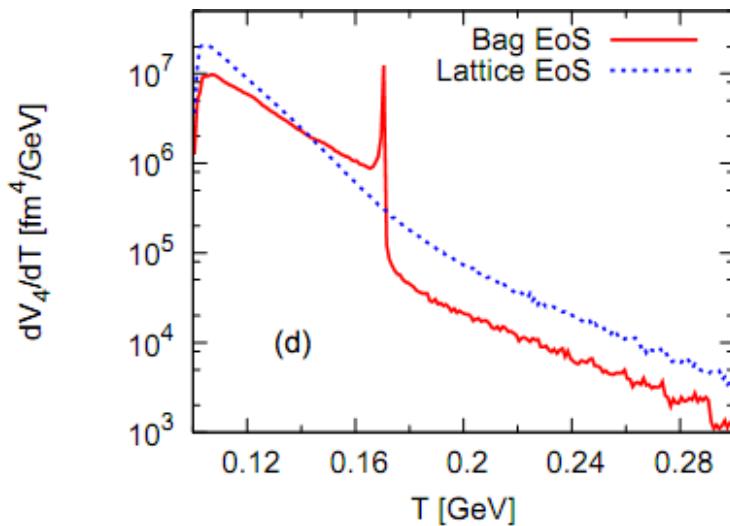
Main source: fluctuation or volume?

- Large charge fluctuation at **high-T**

Spectral strength $\propto \chi(T)$

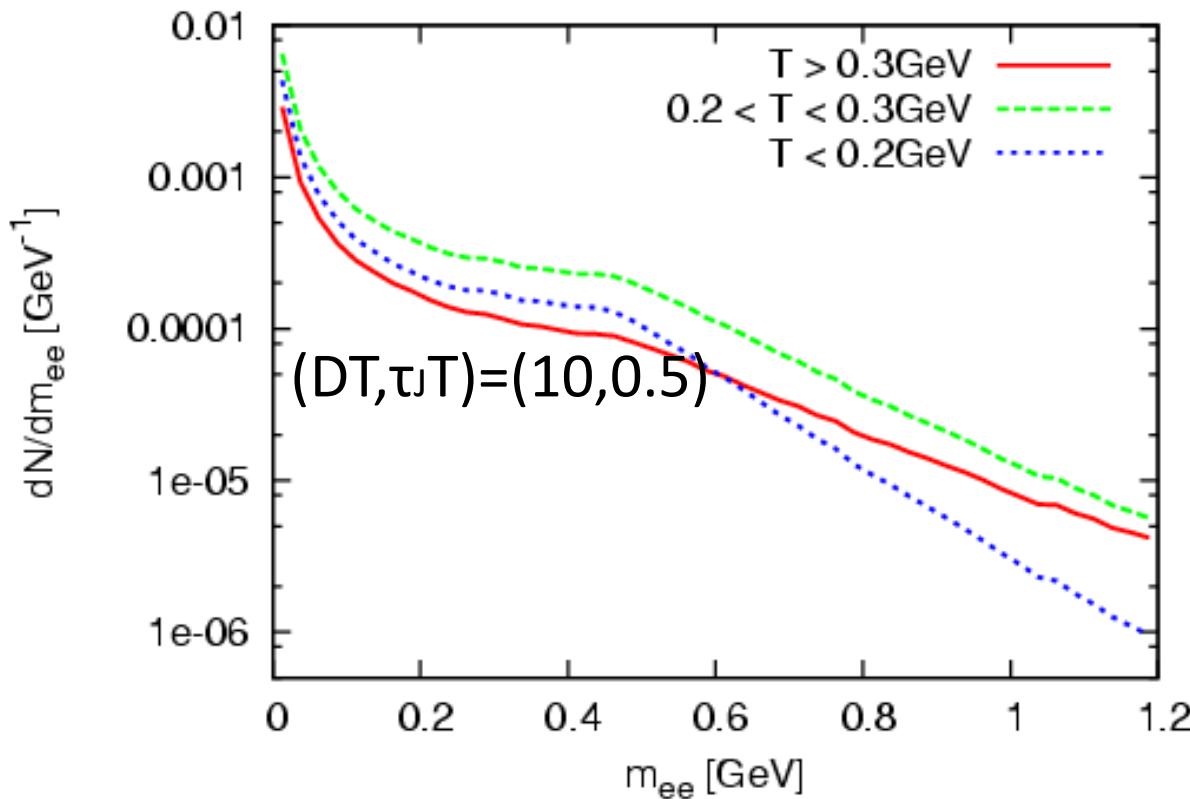
$$= 0.28T^2 \left[1 + \tanh\left(\frac{T - 0.155\text{GeV}}{0.023\text{GeV}}\right) \right]$$

- Large spacetime volume at **low-T**



Main source: volume or fluctuation?

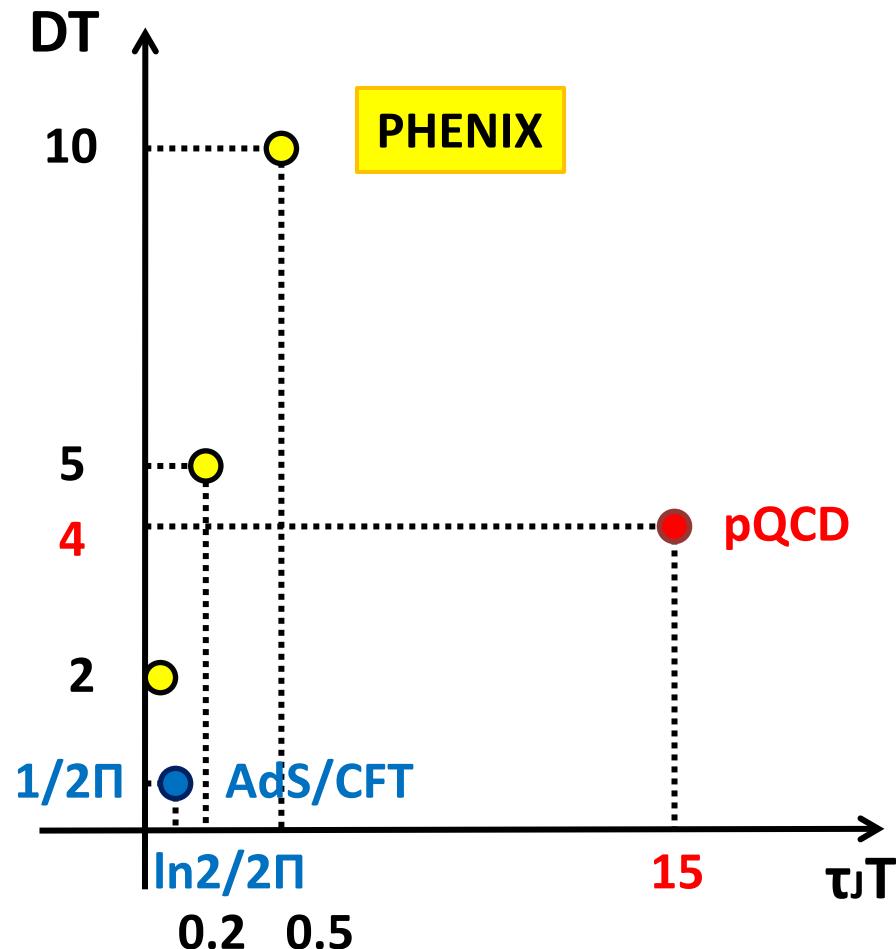
Large production rate due to **thermal fluctuation of charge** in hot region strongly compensates smallness of volume !



Transport coefficients

Assumptions made:

- Validity of transport-SPF
- Hydro with chem.eq.
- No radiation from pre-eq. stage



4. Summary

- Dilepton production via transport phenomena is studied. Transport spectral function is parameterized with diffusion constant D , relaxation time τ_J , and susceptibility χ .
- PHENIX Data requires $D \geq 2/T$. Solution (D, τ_J) is not uniquely obtained, but all the solutions are far from both pQCD and AdS/CFT results.
- Main source of thermal dilepton radiation is high- T QGP phase due to the large fluctuation.
- Application to dilepton/photon radiation at LHC.

Backup

5. Spectral function from causal dissipative hydrodynamics

Hydrodynamic equations

Hydrodynamic equations
in external field

$$\begin{aligned}\partial_\nu T^{\nu\mu} &= F^{\mu\nu} J_\nu, \\ \partial_\mu J^\mu &= 0,\end{aligned}$$

Tensor decomposition

$$T^{\mu\nu} = e u^\mu u^\nu - (P + \Pi) \Delta^{\mu\nu} + \pi^{\mu\nu},$$

$$J^\mu = \rho u^\mu + v^\mu,$$

$$\Delta^{\mu\nu} \equiv g^{\mu\nu} - u^\mu u^\nu,$$

Entropy current and constitutive equations

Entropy current

$$s^\mu = su^\mu - \frac{\mu}{T}\nu^\mu - \frac{1}{T}(\alpha_0\Pi\nu^\mu + \alpha_1\pi^{\mu\nu}\nu_\nu) - \frac{u^\mu}{2T}(\beta_0\Pi^2 - \beta_1\nu^\mu\nu_\mu + \beta_2\pi^{\rho\sigma}\pi_{\rho\sigma}),$$

Constitutive equations (to ensure 2nd law of thermodynamics)

$$-\Pi = \zeta(\partial_\mu u^\mu + \alpha_0\partial_\mu\nu^\mu + \beta_0\dot{\Pi}),$$

$$\pi_{\mu\nu} = 2\eta\langle\langle\partial_\mu u_\nu - \alpha_1\partial_\mu\nu_\nu - \beta_2\dot{\pi}_{\mu\nu}\rangle\rangle,$$

$$\nu^\mu = \sigma\Delta^{\mu\rho}\left[T\partial_\rho\left(\frac{\mu}{T}\right) + F_{\rho\sigma}u^\sigma + \alpha_0\partial_\rho\Pi + \alpha_1\partial_\sigma\pi_\rho^\sigma - \beta_1\dot{\nu}_\rho\right]$$

Linear response theory and spectral function

Linear response theory $\langle \delta J^\mu(q) \rangle_T = -G_R^{\mu\nu}(q; T)\delta A_\nu(q),$

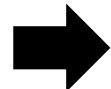
Linearized equation → spectral function

$$J^\mu(x) = (\delta\rho(x), \vec{\nu}(x)),$$

$$\vec{\nu}(x) = \sigma \vec{E} - D \vec{\nabla} \delta\rho - \tau_J \frac{\partial \vec{\nu}}{\partial t},$$

$$D \equiv \frac{\sigma}{\chi}, \quad \tau_J \equiv \beta_1 \sigma, \quad \chi \equiv \frac{\partial \rho}{\partial \mu}.$$

$$G_R^{00}(\omega, \vec{k}; T) = \frac{\sigma k^2}{-\tau_J \omega^2 - i\omega + Dk^2},$$

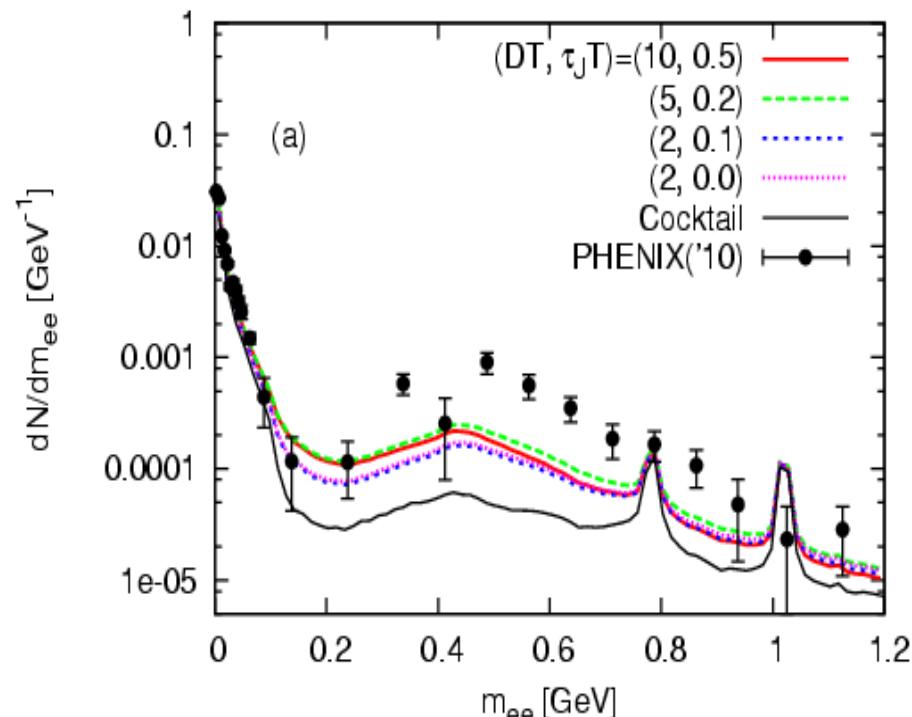


$$G_R^{0i}(\omega, \vec{k}; T) = G_R^{i0}(\omega, \vec{k}; T) = \frac{\sigma \omega k^i}{-\tau_J \omega^2 - i\omega + Dk^2},$$

$$G_R^{ij}(\omega, \vec{k}; T) = \frac{\sigma \omega}{-\tau_J \omega - i} \left(\delta^{ij} - \frac{k^i k^j}{k^2} \right) + \frac{\sigma \omega^2}{-\tau_J \omega^2 - i\omega + Dk^2} \frac{k^i k^j}{k^2},$$

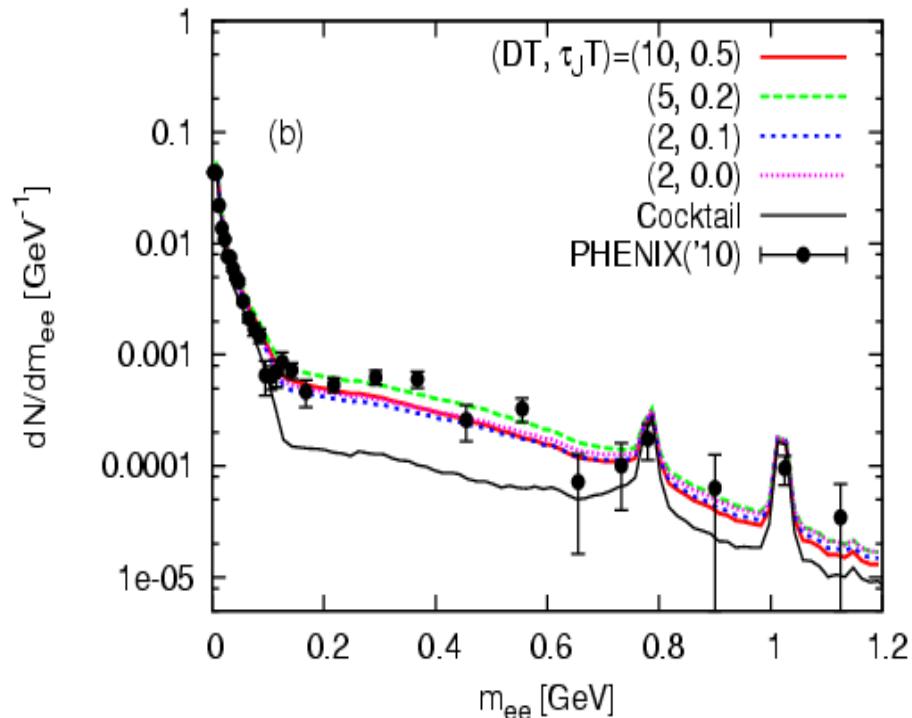
6. More results of dileptons and photons

Invariant mass spectra (p_T -window)



$0 < p_T < 0.5 \text{ GeV}$

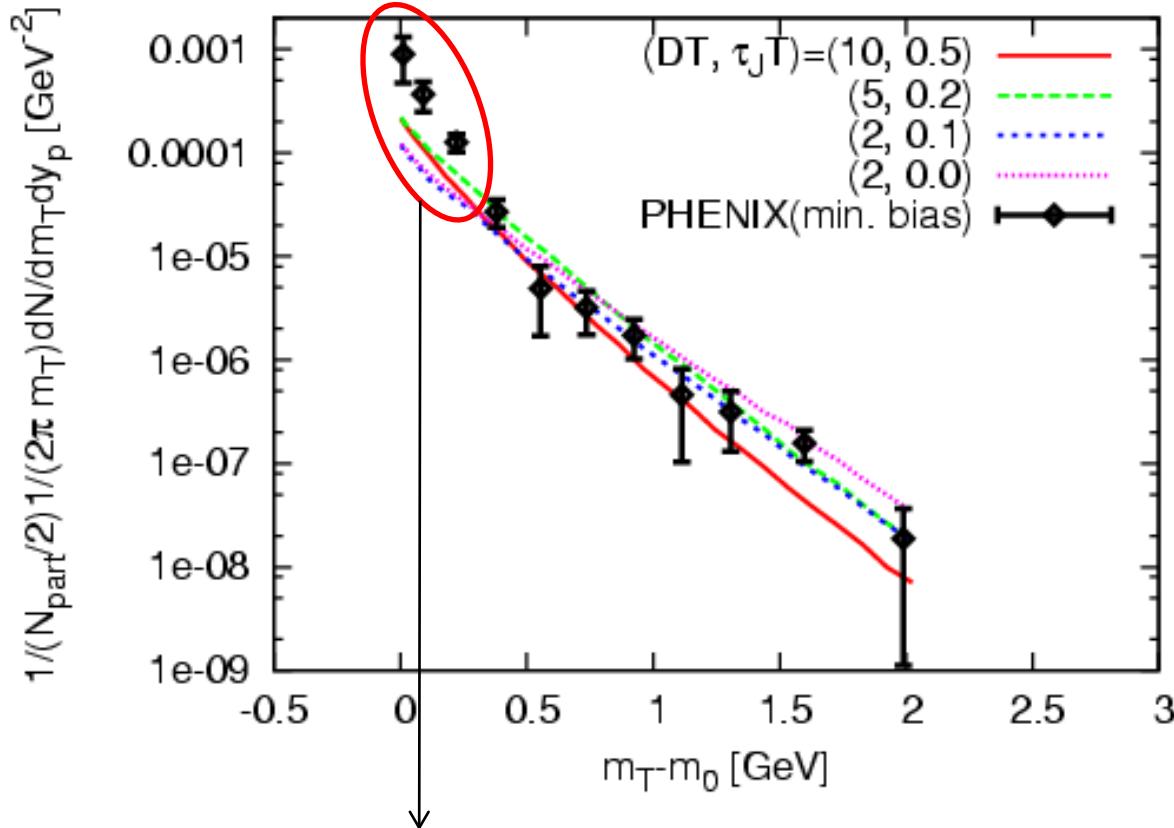
Cannot explain the spectrum
for the lowest p_T window



$0.5 < p_T < 1.0 \text{ GeV}$

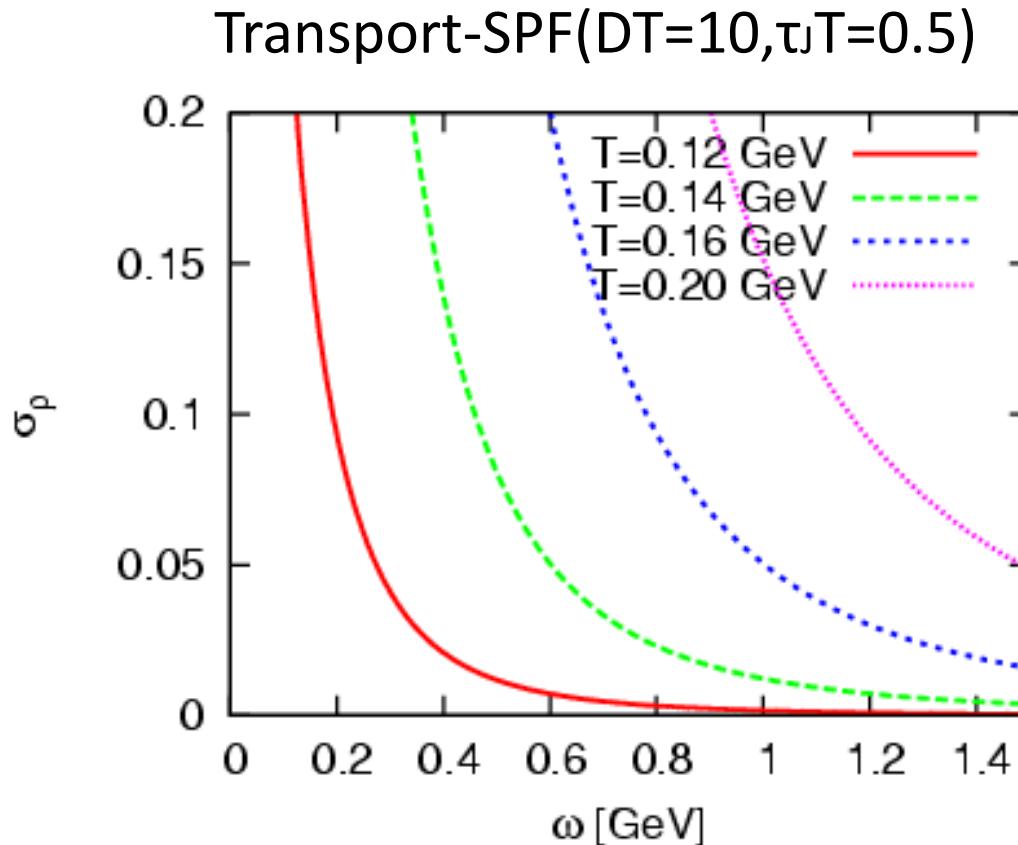
Can explain the spectrum
for this p_T window

m_T -slope



m_T -slope of the data is much steeper than those of theoretical ones.

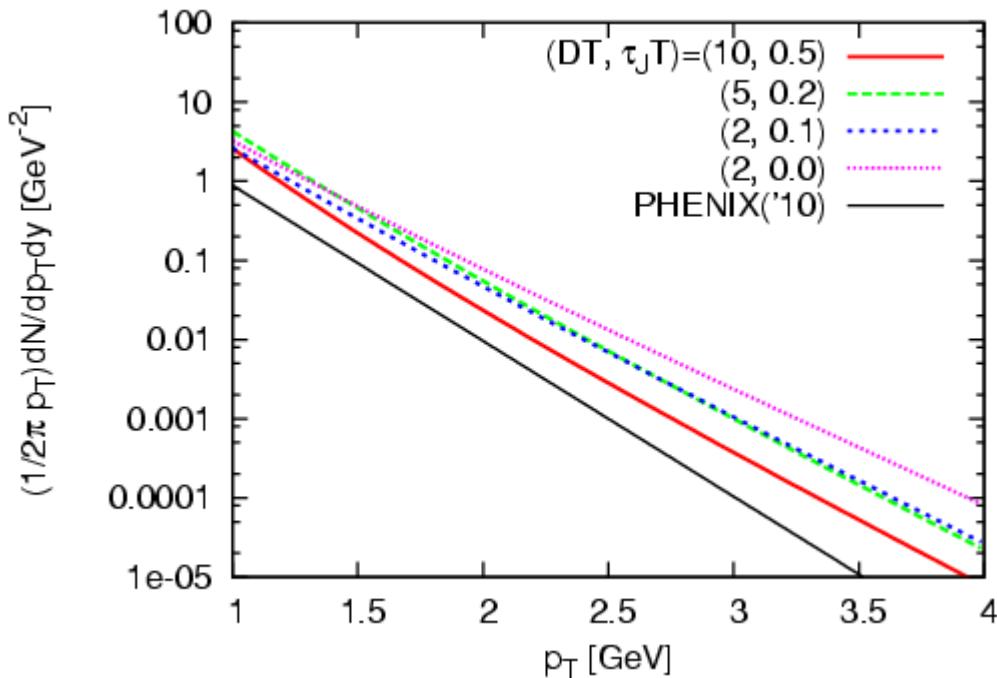
Transport-SPF



$$\sigma(\omega) \equiv -\text{Im} G_R^{\mu\mu}(\omega, 0) / 3\omega^2$$

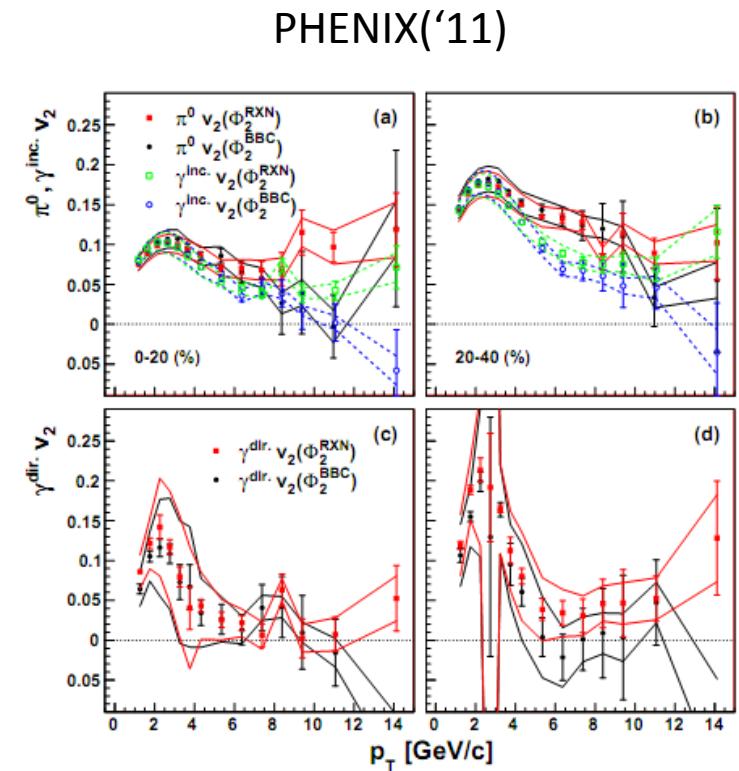
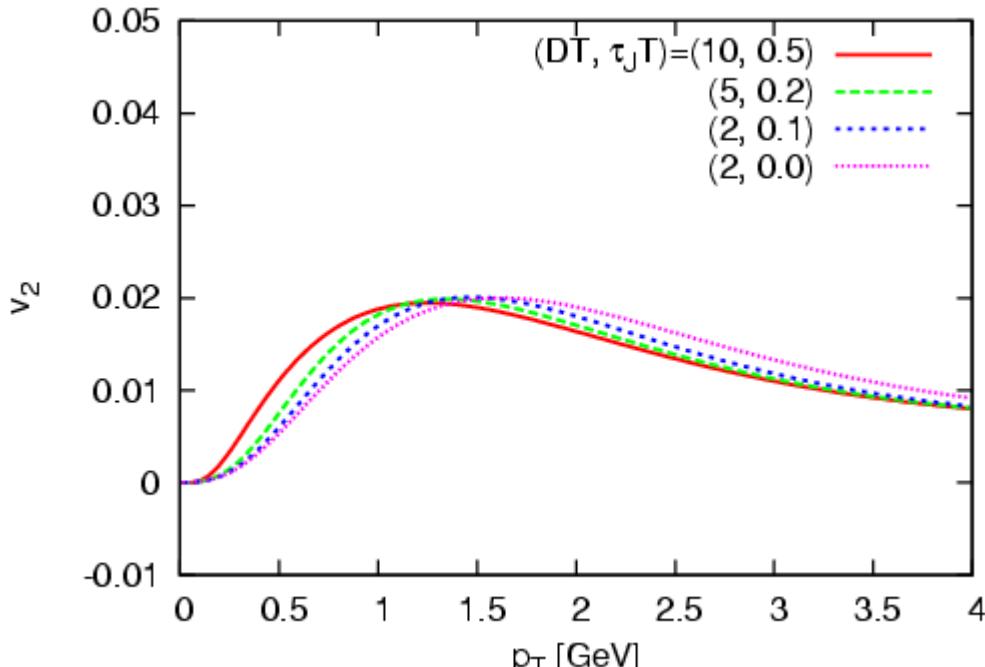
熱場の量子論とその応用

Photon : p_T-spectra



Parameters favored by dileptons overpredict the data.

Photon : V2



Parameters favored by dileptons do not explain data.
($v_2 \sim 0.1$ up to $p_T \sim 2\text{GeV} !!$)