Equation of state and magnetic monopoles in hot SU(2) gluodynamics

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- Thermodynamics of Yang-Mills theory
- Models of color confinement at T < Tc
 - Abelian monopoles
- In deconfinement (gluon plasma at T > Tc)
 - Are they (still) alive as real object?
- Contribution to (trace of) energy-momentum tensor from Abelian monopoles

Thermodynamics

Free Energy (T is temperature and V is spatial volume) F = -T log Z(T, V)
Pressure $p = \frac{T}{V} \frac{\partial \log Z(T, V)}{\partial \log V} = -\frac{F}{V} = \frac{T}{V} \log Z(T, V)$

• Energy density $\varepsilon = \frac{T}{V} \frac{\partial \log Z(T, V)}{\partial \log T}$

• Entropy density $s(T) = \frac{\varepsilon + p}{T} = \frac{\partial p(T)}{\partial T}$

Thermodynamics: Trace Anomaly

• Trace anomaly of the energy-momentum tensor $T_{\mu\nu}$ $\theta(T) = \langle T^{\mu}_{\mu} \rangle \equiv \varepsilon - 3p = T^5 \frac{\partial}{\partial T} \frac{p(T)}{T^4}$ Pressure via trace anomaly $p(T) = T^4 \int^T \frac{dT_1}{T_1} \frac{\theta(T_1)}{T_1^4}$ Energy density via trace anomaly $\varepsilon(T) = 3T^4 \int^T \frac{dT_1}{T_1} \frac{\theta(T_1)}{T_1^4} + \theta(T)$ Trace anomaly is a key quantity

Trace Anomaly for SU(2) pure gluons

Partition Function

$$Z(T,V) = \int DU \exp\{-\beta \sum_{P} S_{P}[U]\}, \ \underline{S_{P}[U]} = (1 - \frac{1}{2}Tr\underline{U_{P}})$$

Plaquette action Plaquette

Asymmetric N_s³ × N_t lattice:
T = 1/(N_ta), V = (N_sa)³
Trace anomaly on the lattice
$$\frac{\theta(T)}{T^4} = 6 N_t^4 \left(\frac{\partial \beta(a)}{\partial \log a}\right) \cdot (\langle S_P \rangle_T - \langle S_P \rangle_0)$$

Mechanisms of color confinement

Dual superconductor picture

['t Hooft, Mandelstam, Nambu, '74-'76]

- Based on existence of special gluonic configurations, called ``magnetic monopoles''
- □ Monopoles are classified with respect to the Cartan subgroup $[U(1)]^{N-1}$ of the SU(N) gauge group
- Confinement is due to monopole condensation
- □ Monopole dominance for various quantities
 - String tension
 String tension
 Polyakov loop behaviors
 - Critical exponents

Magnetic monopoles play an important role for color confinement.

Confinement (T<Tc) and plasma (T>Tc)

- The monopoles are percolating and condensed in confining vacuum
- The percolating monopole cluster disappears and monopole condensate vanishes in deconfinement phase
- <u>Suggestion</u>: the monopoles must emerge as a real (thermal) component of deconfinement plasma similar to electrically neutral electron-positron plasma:
 - □ individual particles exist at high temperatures in a heat bath
 - □ annihilate at low temperatures, but still present in the vacuum

[V.I.Zakharov, M.N.Chernodub,'07] [Liao and Shuryak, '06-'07]

 <u>Check</u>: if the suggestion is true, then the monopoles must contribute to the equation of state of the gluon plasma

Gauge fixing (MA gauge): maximize

$$R = \sum_{s,\mu} Tr[\sigma_3 U_{\mu}(s)\sigma_3 U_{\mu}^{\dagger}(s)]$$

[A.S.Kronfeld, M.L.Laursen, G.Schierholz, U.J.Wiese '87]

 Define particular singular gluon objects (monopoles)
 [T.A.DeGrand, D.Toussaint '80]

 $k_{\mu}(s) = \epsilon_{\mu\nu\rho\sigma}\partial_{\nu}n_{\rho\sigma}(s+\hat{\mu})/2$

- Extract the plaquettes around the monopole
- Decompose the trace anomaly into two parts
 The contribution around the monopole and the rest [for center vortex ; M.N.Chernodub, A.Nakamura, V.I.Zakharov '08]

Action density:

$$\begin{split} \langle S_P \rangle &= \langle S_P \rangle^{\text{mon}} + \langle S_P \rangle^{\text{rest}} \\ &= \frac{1}{6N_s^3 N_t} \left[\langle \sum_P \rho_P S_P \rangle + \langle \sum_P (1 - \rho_P) S_P \rangle \right] \\ \rho_P &= 1 \ (P \in \Sigma) \text{ or } 0 \ (P \notin \Sigma) \\ \Sigma \text{ : plaquettes around monopoles} \end{split}$$

$$\begin{aligned} & \text{Trace anomaly (naive regularization):} \\ & \frac{\theta_{\text{naive}}^{\text{mon}}}{T^4} = 6N_t^4 \left(\frac{\partial \beta}{\partial \log a} \right) \left[\langle S_P \rangle_T^{\text{mon}} - \rho(T) \langle S_P \rangle_0 \right] \\ & \frac{\theta_{\text{naive}}^{\text{rest}}}{T^4} = 6N_t^4 \left(\frac{\partial \beta}{\partial \log a} \right) \left[\langle S_P \rangle_T^{\text{rest}} - (1 - \rho(T)) \langle S_P \rangle_0 \right] \end{split}$$

Trace anomaly (naive regularization):



T = 0: 16^4 lattice, 1000 conf. T > 0: $16^3 \times 4$ lattice, 5000 conf.

Calculated by RICC at RIKEN SX8 at RCNP

Specific action density:

$$\langle s_P \rangle^{\mathsf{mon}} = \frac{\langle S_P \rangle^{\mathsf{mon}}}{\rho} = \frac{\langle \Sigma_P \rho_P S_P \rangle^{\mathsf{mon}}}{\langle \Sigma_P \rho_P \rangle} \\ \langle s_P \rangle^{\mathsf{rest}} = \frac{\langle S_P \rangle^{\mathsf{rest}}}{1 - \rho} = \frac{\langle \Sigma_P (1 - \rho_P) S_P \rangle^{\mathsf{rest}}}{\langle \Sigma_P (1 - \rho_P) \rangle}$$

(action density per an elementary plaquette)

Regularized trace anomaly :

$$\frac{\theta_{\text{reg}}^{\text{mon}}}{T^4} = 6N_t^4 \left(\frac{\partial\beta}{\partial\log a}\right)\rho(T) \left[\langle s_P \rangle_T^{\text{mon}} - \langle s_P \rangle_0^{\text{mon}}\right]$$

Specific action density :



The difference between $\langle s_P
angle^{
m mon}$ and $\langle s_P
angle^{
m rest}$ is seen clearly.





•Sensitive to the phase transition.

•The behavior is similar to the case of center vortex. [M.N.Chernodub et.al. '08]

Summary and future works

Summary

Found: strong contributions from the plaquettes around Abelian monopoles to the trace anomaly, and, consequently, to the pressure and to the energy density of the gluon plasma.

□ Gluonic configurations around the Abelian monopoles are similar to the worldsheets of the center vortex.

Future works

- Check of scaling for trace anomaly (finite volume effect , Nt-dependence)
- Study the contribution to the electric (magnetic) part of the trace anomaly from the Abelian monopoles.