

# QCD和則と最大エントロピー法を用いた有限温度 におけるクォークコニウムのスペクトル解析

(Spectral analysis of quarkonium from QCD sum rules  
and the maximum entropy method)

P. Gubler and M. Oka, Prog. Theor. Phys. **124**, 995 (2010).

P. Gubler, K. Morita and M. Oka, Phys. Rev. Lett. **107**, 092003 (2011).

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## 熱場の量子論とその応用

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Collaborators:

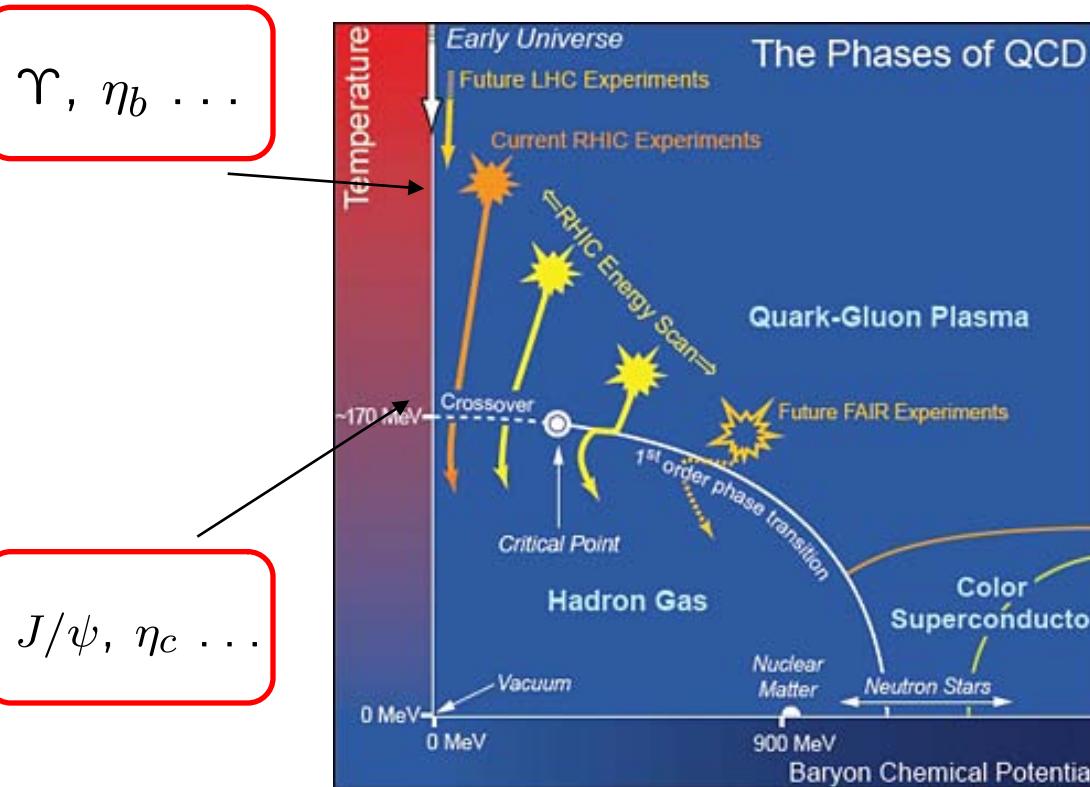
Makoto Oka (TokyoTech), Kenji Morita (YITP), Kei Suzuki (Tokyo Tech)

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- The method: QCD Sum Rules and the Maximum Entropy Method
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# Introduction: Quarkonia

General Motivation: Understanding the behavior of matter at high T.



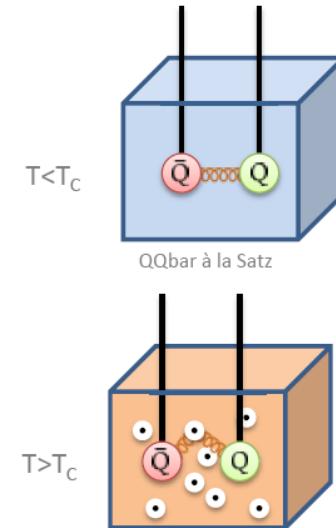
- Phase transition:  
 $QGP (T > T_c) \leftrightarrow \text{confining phase} (T < T_c)$
- Currently investigated  
at RHIC and LHC
- Heavy Quarkonium: clean probe  
for experiment

# Why are quarkonia useful?

Prediction of  $J/\psi$  suppression above  $T_c$   
due to Debye screening:

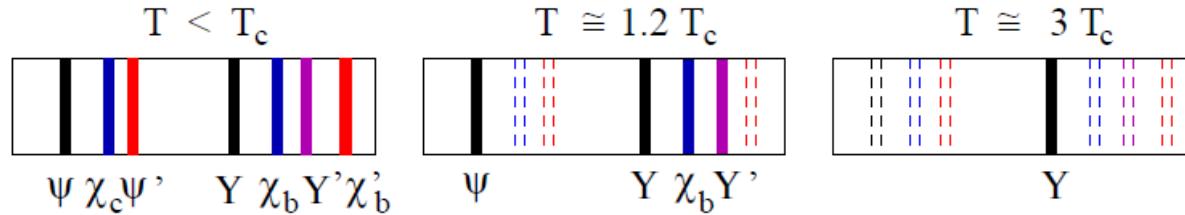
T. Matsui and H. Satz, Phys. Lett. B **178**, 416 (1986).

T. Hashimoto et al., Phys. Rev. Lett. **57**, 2123 (1986).



Lighter quarkonia melt at low  $T$ , while heavier ones melt at higher  $T$

→ Thermometer of the QGP



# Results from lattice QCD

- During the last 10 years, a picture has emerged from studies using lattice QCD (and MEM), where  $J/\psi$  survives above  $T_c$ , but dissolves below  $2 T_c$ .

M. Asakawa and T. Hatsuda, Phys. Rev. Lett. 92 012001 (2004).

S. Datta *et al*, Phys. Rev. D69, 094507 (2004).

T. Umeda *et al*, Eur. Phys. J. C39, 9 (2004).

A. Jakovác *et al*, Phys. Rev. D75, 014506 (2007).

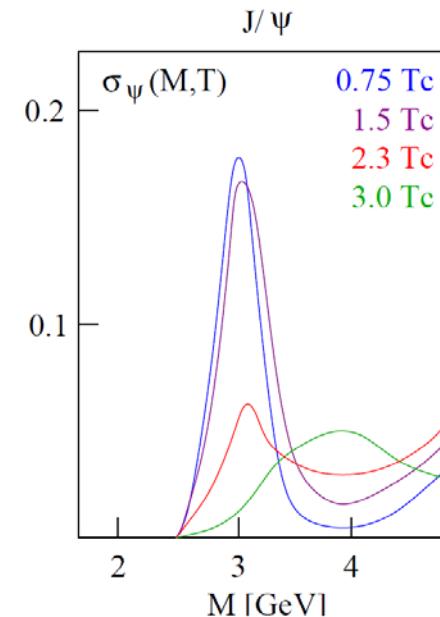
G. Aarts *et al*, Phys. Rev. D 76, 094513 (2007).

H.-T. Ding *et al*, PoS LAT2010, 180 (2010).

- However, there are also indications that  $J/\psi$  survives up to  $2 T_c$  or higher.

H. Iida *et al*, Phys. Rev. D 74, 074502 (2006).

H. Ohno *et al*, PoS LAT2008, 203 (2008).



Taken from  
H. Satz, Nucl.Part.Phys. **32**, 25 (2006).

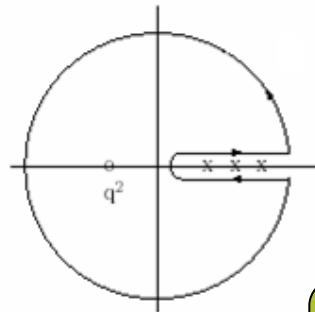
(schematic)

# QCD sum rules

M.A. Shifman, A.I. Vainshtein and V.I. Zakharov,  
Nucl. Phys. B147, 385 (1979); B147, 448 (1979).

In this method the properties of the two point correlation function is fully exploited:

$$\Pi(q^2) = i \int d^4x e^{iqx} \langle 0 | T\{\chi(x)\bar{\chi}(0)\} | 0 \rangle$$



$$\rightarrow \Pi(q^2) = \frac{1}{\pi} \int_0^\infty ds \frac{\text{Im} \Pi(s)}{s - q^2 - i\epsilon}$$

is calculated  
“perturbatively”,  
using OPE

spectral function  
of the operator  $\chi$

After the Borel transformation:

$$G_{OPE}(M) = \frac{1}{\pi} \int_0^\infty ds \frac{1}{M^2} e^{-\frac{s}{M^2}} \text{Im} \Pi(s)$$

# The Maximum Entropy Method

→ Bayes' Theorem

$$P[\rho|G, I] = \frac{P[G|\rho, I]P[\rho|I]}{P[G|I]}$$

likelihood function

$$P[G|\rho, I] = \frac{1}{Z_L} e^{-L[\rho]}$$

prior probability

$$P[\rho|I] = \frac{1}{Z_s} e^{\alpha S[\rho]}$$

$$L[\rho] = \frac{1}{2(M_{\max} - M_{\min})} \int_{M_{\min}}^{M_{\max}} dM \left[ \frac{G_{OPE}(M) - G_\rho(M)}{\sigma^2(M)} \right]^2$$

$$S[\rho] = \int_0^\infty d\omega [\rho(\omega) - m(\omega) - \rho(\omega) \log \left( \frac{\rho(\omega)}{m(\omega)} \right)]$$

Corresponds to ordinary  
 $\chi^2$ -fitting.

(Shannon-Jaynes entropy)  
“default model”

M. Jarrel and J.E. Gubernatis, Phys. Rep. 269, 133 (1996).

M. Asakawa, T. Hatsuda and Y. Nakahara, Prog. Part. Nucl. Phys. 46, 459 (2001).

# The charmonium sum rules at finite T

The application of QCD sum rules has been developed in:

A.I. Bochkarev and M.E. Shaposhnikov, Nucl. Phys. B 268, 220 (1986).

T.Hatsuda, Y.Koike and S.H. Lee, Nucl. Phys. B 394, 221 (1993).

$$M(\nu) = \int_0^\infty e^{-\nu t} \rho(4m_c^2 t) dt \quad (\nu \equiv \frac{M^2}{4m_c^2})$$

$$M(\nu) = A(\nu) \left[ 1 + a(\nu) \alpha_s(\nu) + b(\nu) \frac{\langle \frac{\alpha_s}{\pi} G^2 \rangle_T}{m_c^4} + c(\nu) \frac{\langle \frac{\alpha_s}{\pi} G^2 \rangle_{T,2}}{m_c^4} + d(\nu) \frac{\langle g^3 G^3 \rangle_T}{m_c^6} \right]$$

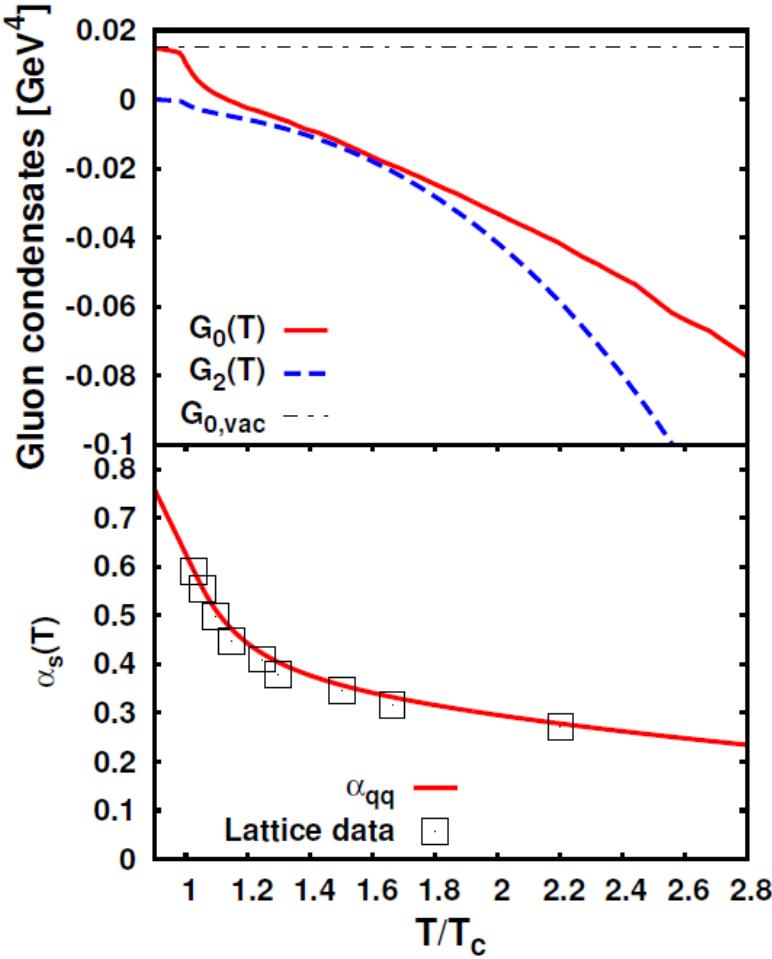
depend on T

Compared to lattice:

No reduction of data points that can be used for the analysis,  
allowing a direct comparison of T=0 and T $\neq$ 0 spectral functions.

## The T-dependence of the condensates

K. Morita and S.H. Lee, Phys. Rev. Lett. 100, 022301 (2008).



taken from:

K. Morita and S.H. Lee, Phys. Rev. D82, 054008 (2010).

Considering the trace and the traceless part of the energy momentum tensor, one can show that in the quenched approximation, the T-dependent parts of the gluon condensates by thermodynamic quantities such as energy density  $\epsilon(T)$  and pressure  $p(T)$ .

$$\langle \frac{\alpha_s}{\pi} G^2 \rangle_T = \langle \frac{\alpha_s}{\pi} G^2 \rangle_{\text{vac.}} - \frac{8}{11}(\epsilon - 3p)$$

$$\langle \frac{\alpha_s}{\pi} G^2 \rangle_{T,2} = -\frac{\alpha_s(T)}{\pi}(\epsilon + p)$$

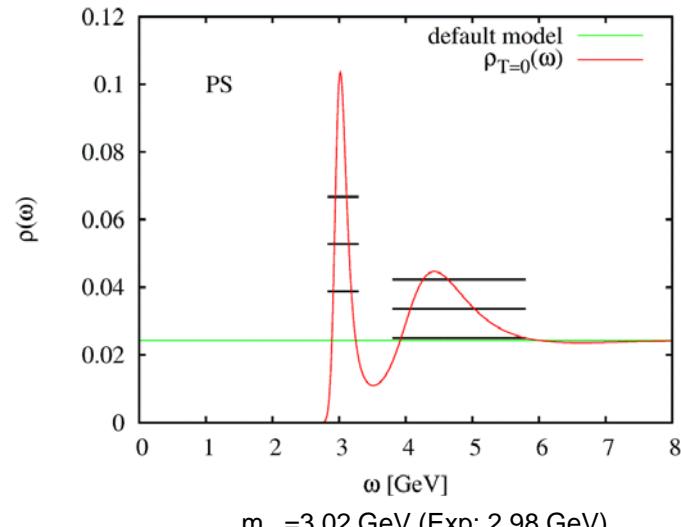
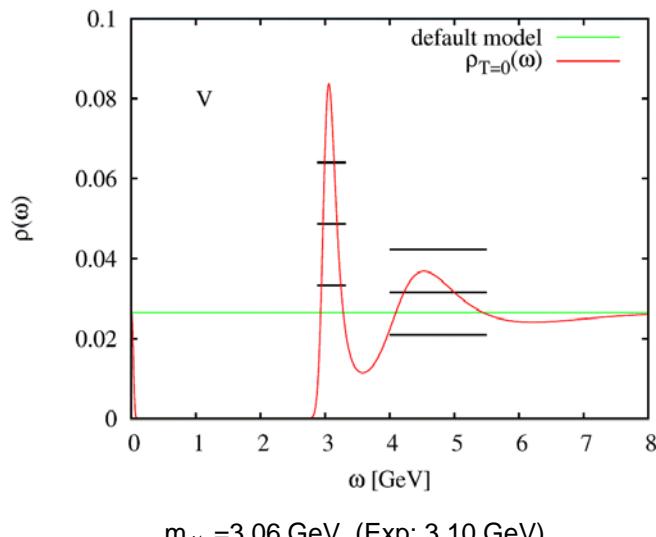
The values of  $\epsilon(T)$  and  $p(T)$  are obtained from quenched lattice calculations:

G. Boyd *et al*, Nucl. Phys. B 469, 419 (1996).

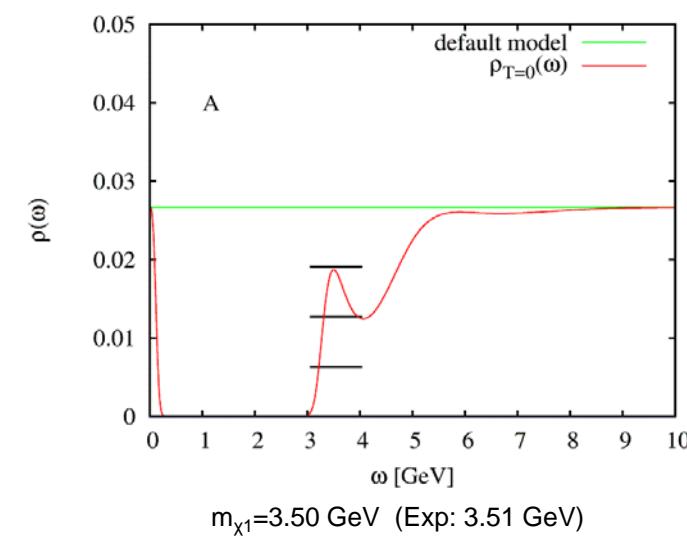
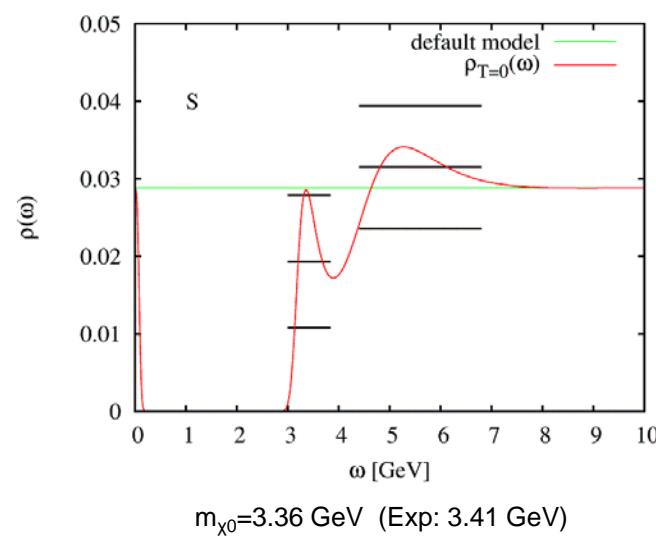
O. Kaczmarek *et al*, Phys. Rev. D 70, 074505 (2004).

## MEM Analysis at T=0

S-wave



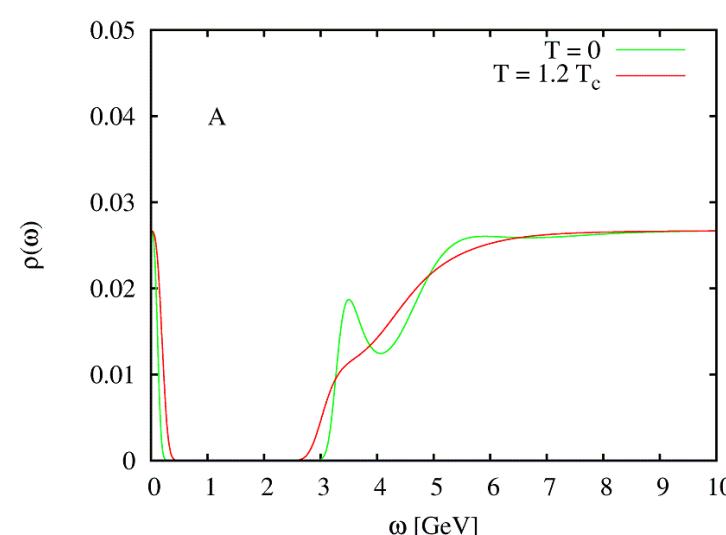
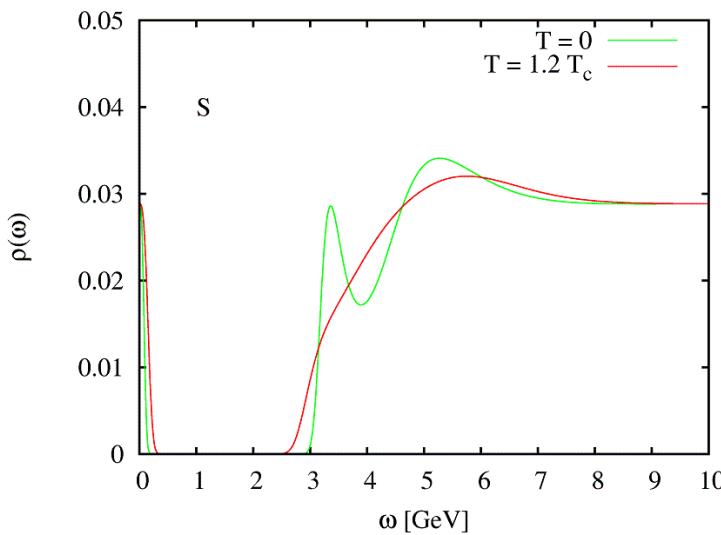
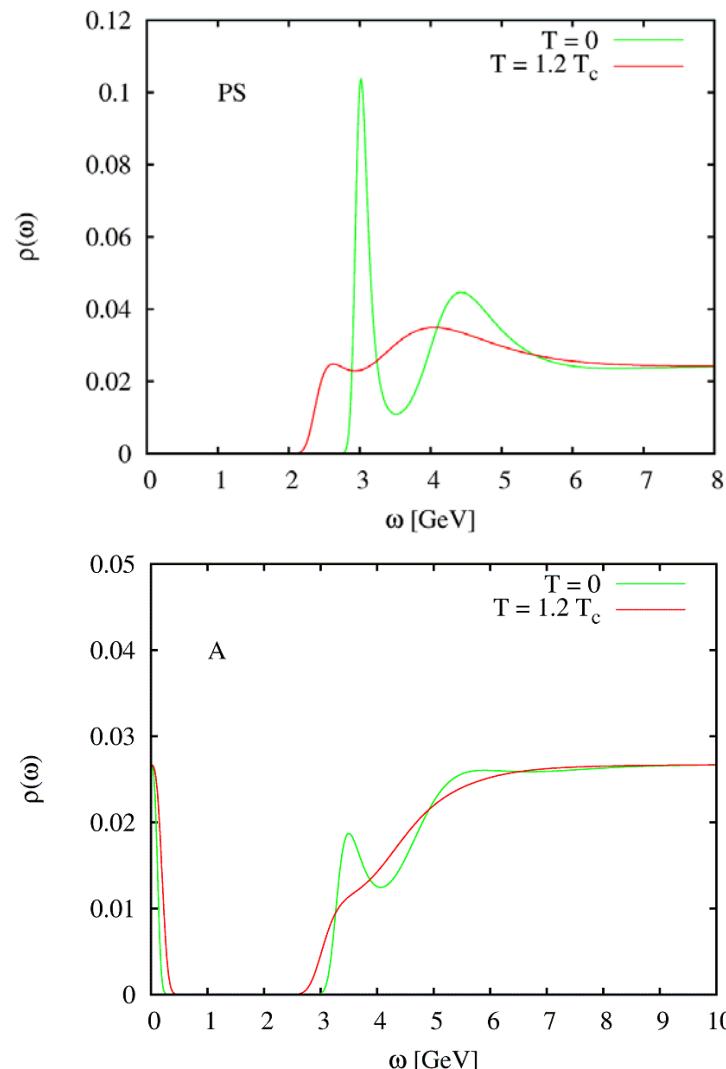
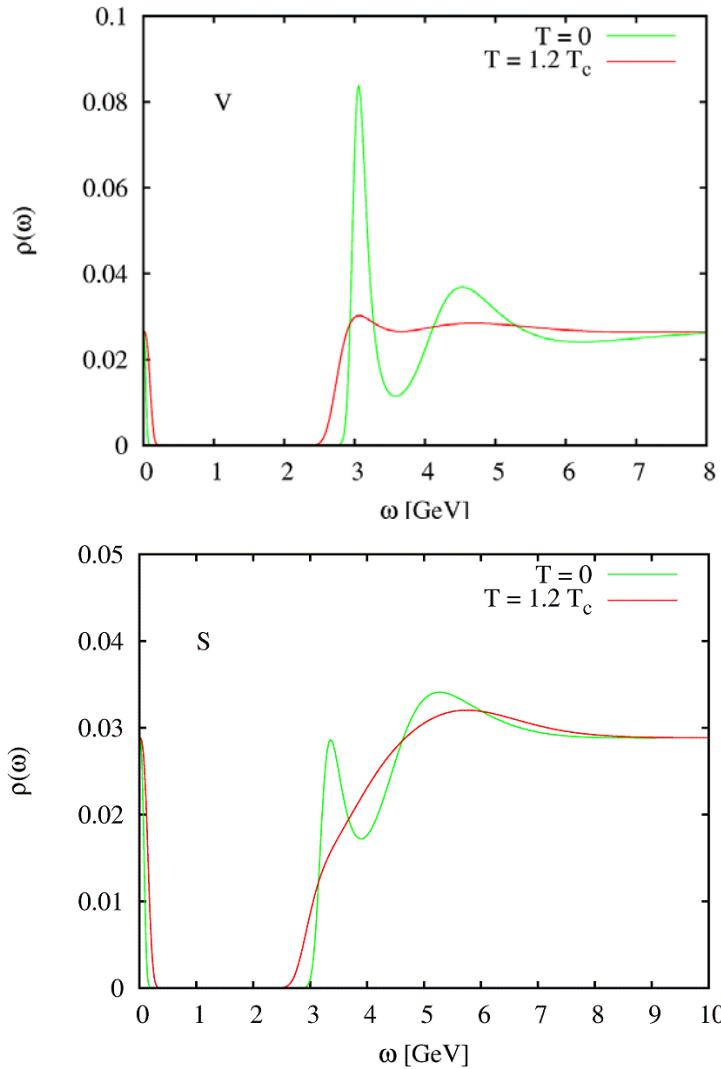
P-wave



$$\langle \frac{\alpha_s}{\pi} G^2 \rangle = 0.012 \pm 0.0036 \text{ GeV}^4$$

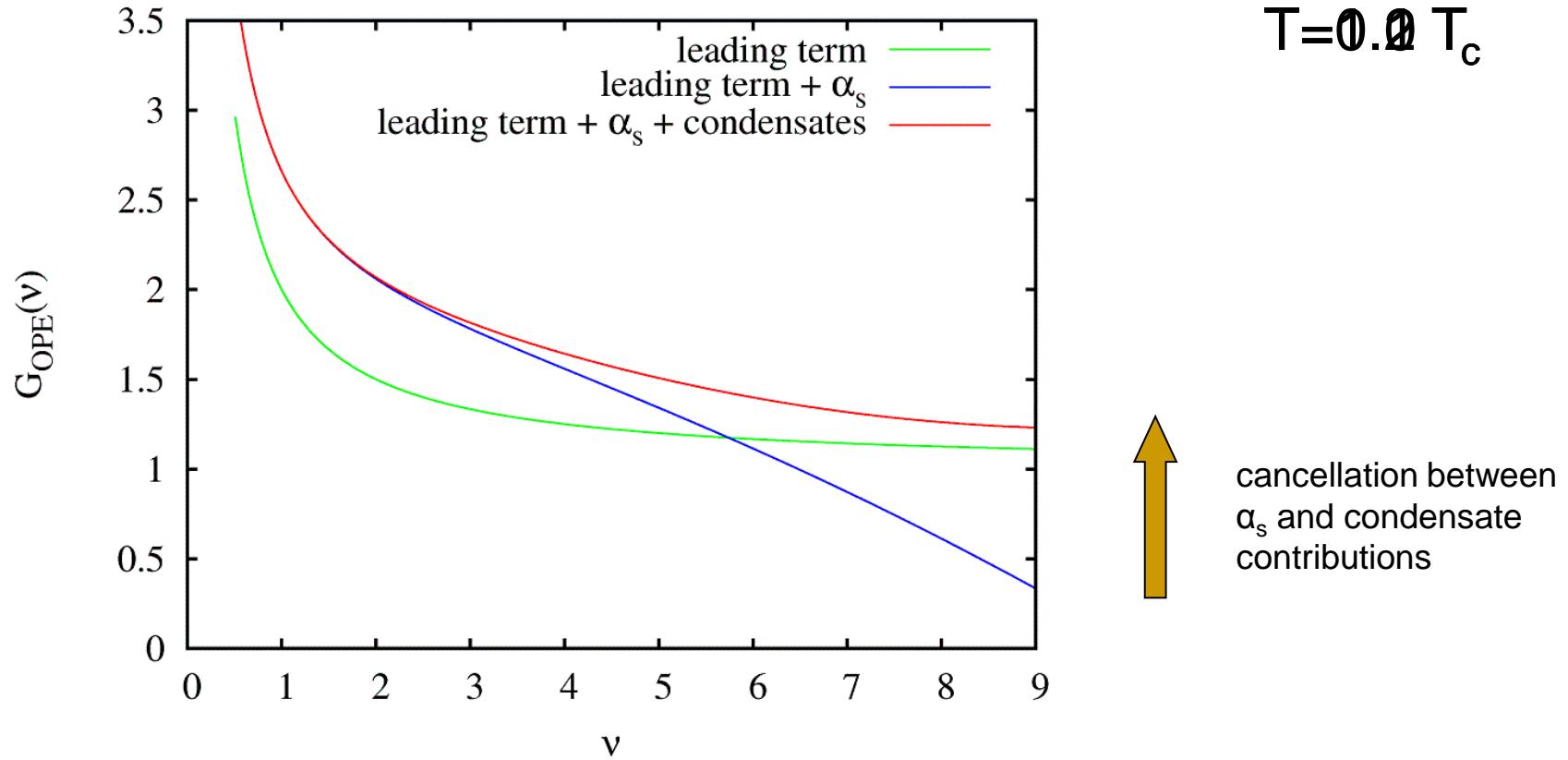
$$m_c = 1.277 \pm 0.026 \text{ GeV}$$

# The charmonium spectral function at finite T



# What is going on behind the scenes ?

The OPE data in the Vector channel at various T:



# Conclusions

- We have shown that MEM can be applied to QCD sum rules
- We could observe the melting of the S-wave and P-wave charmonia using finite temperature QCD sum rules and MEM
- Both  $J/\psi$ ,  $\eta_c$ ,  $X_{c0}$ ,  $X_{c1}$  melt between  $T \sim 1.0 T_c$  and  $T \sim 1.2 T_c$ , which is below the values obtained in lattice QCD

# Outlook

- Bottomium (see poster of K. Suzuki)
- Including higher orders ( $\alpha_s$ , twist)
- Extending the method to investigations of other particles (D, ...)

# Backup slides

# The basic problem to be solved

$$G_{OPE}(M) = \frac{1}{M^2} \int_0^\infty ds e^{-\frac{s}{M^2}} \rho(s)$$

↑  
given  
(but only incomplete and  
with error)

↑  
“Kernel”

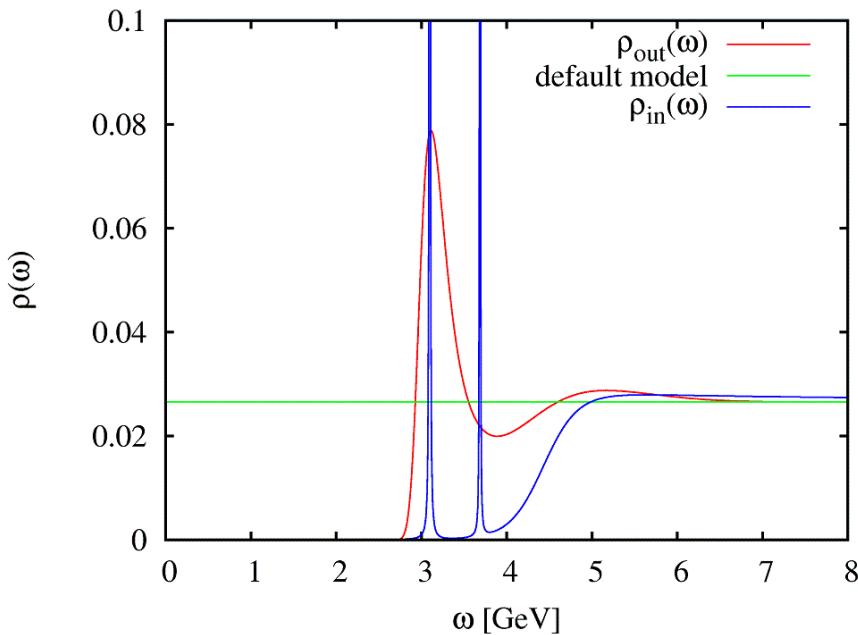
? ↗

This is an ill-posed problem.

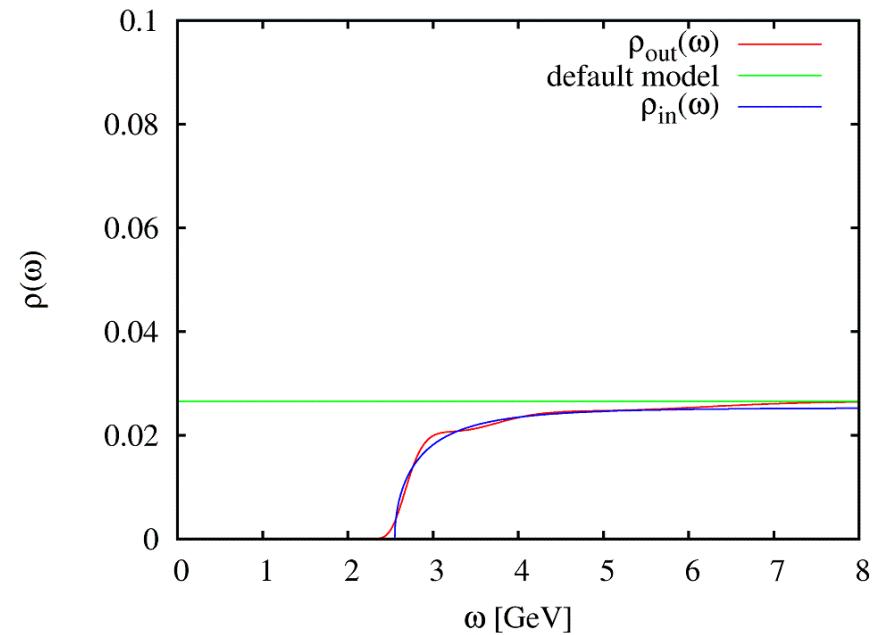
But, one may have additional information on  $\rho(\omega)$ , which can help to constrain the problem:

- Positivity:  $\rho(\omega) \geq 0$
- Asymptotic values:  $\rho(\omega = 0), \rho(\omega = \infty)$

# A first test: mock data analysis



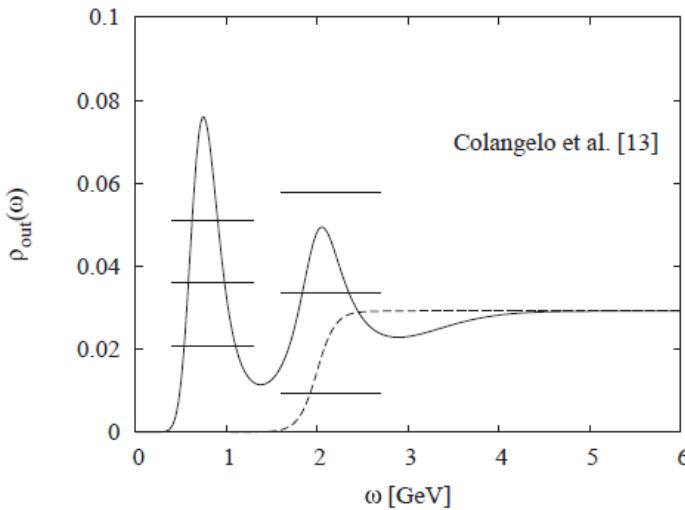
Both  $J/\psi$  and  $\psi'$  are included into the mock data, but we can only reproduce  $J/\psi$ .



When only free c-quarks contribute to the spectral function, this should be reproduced in the MEM analysis.

# First applications in the light quark sector

## p-meson channel



$$m_\rho = 0.76 \pm 0.07 \text{ GeV}$$

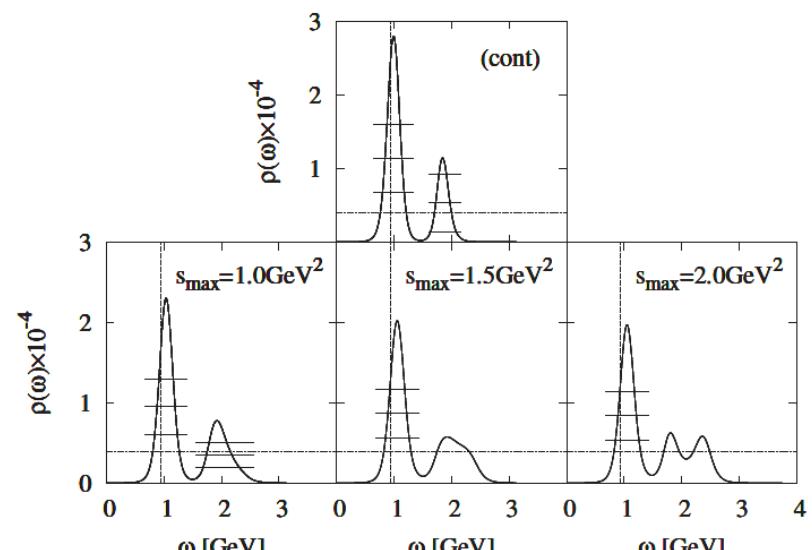
$$F_\rho = 0.174 \pm 0.039 \text{ GeV}$$

Experiment:

$$m_\rho = 0.77 \text{ GeV}$$

$$F_\rho = 0.141 \text{ GeV}$$

## Nucleon channel



$$m_N = 0.99 \pm 0.13 \text{ GeV}$$

Experiment:

$$m_N = 0.94 \text{ GeV}$$