Spin Glass Approach to Restricted Isometry Constant

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Outline

- Background: Compressed sensing
- Problem setup: Restricted isometry constant (RIC)
- Spin glass approach
 - Replica symmetric analysis
 - Improvement by replica symmetry breaking
- Summary

Compressed sensing

Reconstruct original signal **X** from its undersampled measurement **Y**.





Compressed sensing



when **X** is sufficiently "**sparse**".

Practical relevance

- The problem is related to various technologies of modern signal processing
- Many application domains
 - Refraction seismic survey (mine examination)
 - Tomography(X-ray CT, MRI)
 - Single pixel camera
 - Noise removal of image
 - Data streaming computing
 - Group testing
 - etc.

Simulation of tomography



LT: Original(Logan-Shepp Phantom) #512x512

- RT: Sampling 512 points of 2D FT from 22 directions.
- LB: Recovery of pseudo-inverse (standard approach)

RB: Recovery utilizing the "sparseness" of spatial variations. "Original" is perfectly recovered.

Perfect recovery is realized by only 2% samples of what Nyquist-Shannon's theory requires. →Breaking of the conventional limit!

EJ Candes J Romberg and T. Tao, IEEE Trans. IT Vol. 52, 489-502 (2006)

Two major reconstruction methods

l_0 reconstruction

$$\hat{\mathbf{x}} = \arg\min_{\mathbf{x}} \|\mathbf{x}\|_{0}$$
, subject to $\mathbf{Y} = \mathbf{A}\mathbf{x}$

 \hat{x}_{1} reconstruction $\hat{x} = \arg \min_{X} ||x||_{1}$, subject to Y = Ax

Sufficient conditions for the reconstruction are given by <u>Restricted Isometriy Constant (RIC)</u>.

Restricted Isometry Constant (RIC)

Definition (Candès and Tao (2006)).

 $\sum_{\mu=1}^{M} A_{\mu i}^2 = 1 \text{ for } \forall i$

Let *A* be a <u>column-wisely normalized</u> $M \times N$ matrix.

Let $x \in \mathbb{R}^N$ be an arbitrary *S*-sparse vector, whose number of non-zero components is smaller than *S*.

Then, if there exists such a constant $d_S = \max\{d_S^{\min}, d_S^{\max}\}$ that satisfies

$$(1 - \delta_{S}^{\min}) ||\mathbf{x}||_{F}^{2} \leq ||A\mathbf{x}||_{F}^{2} \leq (1 + \delta_{S}^{\max}) ||\mathbf{x}||_{F}^{2},$$

A is said to satisfy the *S*-restricted isometry property (RIP) with the restricted isometry constant (RIC) d_S .

Intuitively, d_s quantifies how A deviates from orthogonal transforms in terms of Frobenius (L₂) norm for S-sparse vectors.

Sufficient conditions for l_0 , l_1 reconstruction

- **1.** l_0 reconstruction gives the unique *S*-sparse solution when $d_{2S} < 1$.
- **2.** l_1 reconstruction gives the same unique *S*-sparse solution as l_0 reconstruction when $d_{2S} < 2^{1/2} - 1$.

[Candes and Tao, IEEE Trans. Inform. Theory (2005)]

3. (Improvement of **2.**)

 l_1 reconstruction gives the same unique *S*-sparse solution as l_0 reconstruction when $(4\sqrt{2}-3)\delta_{2S}^{\min} + \delta_{2S}^{\max} < 4(\sqrt{2}-1)$.

[Foucart and Lai, Appl. Comput. Harmon. Anal. (2009)]

Computational difficulty

- RIC plays a key role in theoretical analysis and performance guarantee of compressed sensing.
- On top of this, it is purely of great interest as a fundamental problem of linear algebra.
- Unfortunately, "there are no known large (systematic) matrices with bounded RICs (and computing RICs is strongly NP-hard),...".

– Wikipedia

- On the other hand, many random matrices have been shown to remain bounded. Therefore, much effort has been paid for <u>improving the bounds</u>.
 - Our study is along this line.

Why computationally difficult?

Suppose a situation where non-zero components are fixed.



$$\lambda_{\min}(\mathbf{A}_{T}^{\mathrm{T}}\mathbf{A}_{T})\|\mathbf{x}_{T}\|_{F}^{2} \leq \|\mathbf{A}_{T}\mathbf{x}_{T}\|_{F}^{2} \leq \lambda_{\max}(\mathbf{A}_{T}^{\mathrm{T}}\mathbf{A}_{T})\|\mathbf{x}_{T}\|_{F}^{2}$$

The change of norm can be <u>easily characterized</u> by eigenvalues of "sub-matrix" A_{T} .

Why computationally difficult?

- In evaluation of RIC, the inequalities must hold for <u>all possible</u> <u>combinations of the positions of non-zeros.</u>
- This causes a combinatorial difficulty, yielding an exact expression of RIC as

$$\delta_{S} = \max\{1 - \min_{T:|T|=S, T \subseteq V} \lambda_{\min}(\mathbf{A}_{T}^{T}\mathbf{A}_{T}), \max_{T:|T|=S, T \subseteq V} \lambda_{\max}(\mathbf{A}_{T}^{T}\mathbf{A}_{T}) - 1\}.$$

All possible column choices
=> Combinatorial difficulty

Structures of the problem and the difficulty are <u>analogous to those of spin glass problems.</u>

Analogy to spin glasses

The choice of *T* are represented by binary vector $\{c_i\} \in \{0, 1\}^N$.

 $\lambda_{\min}(\mathbf{A}_T^{\mathsf{T}}\mathbf{A}_T)$ and $\lambda_{\max}(\mathbf{A}_T^{\mathsf{T}}\mathbf{A}_T)$ can be regarded as `energy functions' of *c* given **A**.

Analogy to spin glasses

$$\delta_{S} = \max \{1 - \min_{\mathbf{c}: \sum_{i=1}^{N} c_{i}=S} \lambda_{\min}(\mathbf{c} \mid \mathbf{A}), \max_{\mathbf{c}: \sum_{i=1}^{N} c_{i}=S} \lambda_{\max}(\mathbf{c} \mid \mathbf{A}) - 1\}$$

RIC is given by minimum and maximum 'energy' of *c*.

• Energy of '0-1 spin' *c*

$$\Lambda_{+}(\boldsymbol{c} \mid \boldsymbol{A}) = \lambda_{\min}(\boldsymbol{c}, \boldsymbol{A})$$
$$\Lambda_{-}(\boldsymbol{c} \mid \boldsymbol{A}) = \lambda_{\max}(\boldsymbol{c}, \boldsymbol{A})$$

• 'Canonical distribution' of
$$c$$

 $P(c \mid A; \mu) \propto \exp(-\mu N \Lambda_{\operatorname{sgn}(\mu)}(c \mid A)) \delta\left(\sum_{i=1}^{N} c_i - S\right)$

• 'Quenched randomness'

$$P(A) = (2\pi M^{-1})^{-MN/2} \exp\left(-\frac{M}{2} \sum_{\mu,i} A_{\mu i}^{2}\right)$$

Spin glass approach

- Free entropy (free energy) $Z(\mu \mid A)$: Partition function $\phi(\mu \mid A) = \frac{1}{N} \log \left[\sum_{c \in \{0,1\}^N} \exp(-\mu N \Lambda_{\operatorname{sgn}(\mu)}(c \mid A)) \delta\left(\sum_{i=1}^N c_i - S\right) \right]$
- Typical properties can be assessed by the replica method.

Possible minimum and maximum eigenvalues

$$\delta_{S} = \max\{1 - \lambda_{\min}^{*}, \lambda_{\max}^{*} - 1\}$$

Replica symmetric analysis

$$\begin{split} \phi(\mu) &= -\frac{\alpha}{2} \log\{\alpha + \chi + \mu(1-q)\} + \frac{\alpha}{2} \log(\alpha + \chi) \\ &- \frac{\alpha \mu q}{2\{\alpha + \chi + \mu(1-q)\}} + \frac{\hat{Q}}{2} - \frac{\hat{q}_1}{2} \left(1 + \frac{\chi}{\mu}\right) + \frac{\hat{q}_0 q}{2} + K\rho \\ &+ \int Dz \log\left(1 + e^{-\kappa} \int Dy \exp\left(\frac{(\sqrt{\hat{q}_1 - \hat{q}_0} y + \sqrt{\hat{q}_0} z)^2}{2\hat{Q}}\right)\right) \end{split}$$

- $\alpha = M / N$ (compression rate)
- $\rho = S / N$ (fraction of non-zero components)
- $\{q, \chi, \hat{Q}, \hat{q}_1, \hat{q}_0, K\}$ are determined by saddle point equations.

Replica symmetric entropy

Replica symmetric RIC

• d_s at $\alpha = 0.5$

$$\approx \alpha = \frac{M}{N}, \rho = \frac{S}{N}$$

l_0, l_1 reconstruction limit

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Improvement by RSB

- The RS RIC estimate is lower than <u>any existing upper</u> <u>bounds</u>, being consistent with a known lower bound.
- On the other hand, there are non-negligible deviations from the experimental data.
- In fact, detailed analysis shows that the replica symmetry is broken for the left and right edges of the entropy curves.
- However, physical interpretation of RSB indicates that the RS estimates still serve as upper bounds of RIC, and the bounds are improved as the higher RSBs are taken into account.

Replica symmetric entropy (again)

Physical interpretation of RSB

Complexity = entropy of pure states

$$\Sigma(\lambda,s) = \frac{1}{N} \log(\# \text{pure states specified by } (\lambda,s))$$

Non-negativity constraint: Must be non-negative.

Total entropy

$$\omega(\lambda) = \frac{1}{N} \log \left(\int ds \exp \left(N \left(s + \Sigma(\lambda, s) \right) \right) \right) \simeq \max_{s} \left\{ s + \Sigma(\lambda, s) \right\}$$

Cf) Montanari and Ricci-Tersenghi (2003), Krzakala et al (2007), ...

RS evaluation: the complexity constraint is ignored.

$$\omega_{\rm RS}(\lambda) = \max_{s} \left\{ s + \Sigma(\lambda, s) \right\}$$

1RSB evaluation: the complexity constraint is taken into account.

$$\omega_{1\text{RSB}}(\lambda) = \max_{s} \left\{ s + \Sigma(\lambda, s) \middle| \Sigma(\lambda, s) \ge 0 \right\}$$

The non-negativity constraint of the 1RSB evaluation means

$$\omega_{\rm RS}(\lambda) \ge \omega_{\rm 1RSB}(\lambda).$$

Physical interpretation of RSB

If necessary, a similar argument may be applied for the dominant cluster of 1RSB, which yields

$$\omega_{\rm RS}(\lambda) \ge \omega_{\rm 1RSB}(\lambda) \ge \omega_{\rm 2RSB}(\lambda).$$

Applying the similar argument repeatedly concludes

$$\omega_{\rm RS}(\lambda) \ge \omega_{\rm 1RSB}(\lambda) \ge \omega_{\rm 2RSB}(\lambda) \ge \dots$$

Monotonic improvement by RSB

The series of inequalities

$$\omega_{\rm RS}(\lambda) \ge \omega_{\rm 1RSB}(\lambda) \ge \omega_{\rm 2RSB}(\lambda) \ge \dots$$

indicates that bounds are monotonically improved by incorporating the higher RSB.

1RSB results

- The estimates are actually improved by the 1RSB solution!
- Two scenarios for the RSB transition
 - λ_{\min} : Random first order transition (RFOT)
 - λ_{\max} : de Almeida-Thouless instability (full RSB)

Summary

- Evaluation of the restricted isometry constant (RIC) can be formulated as a spin glass problem.
- We provided a replica based-framework for accurate evaluation of RICs.
 - Replica evaluation provides the current best accuracy.
 - Although the RS solution is not thermodynamically stable, the physical interpretation of RSB implies that the RS estimates still have the meaning of "upper-bounds".
 - The bounds are monotonically improved by taking the higher RSB into account.
- Future work
 - Application to other matrix ensembles
 - Mathematical justification

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