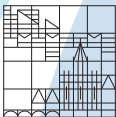




Active microrheology: Brownian motion, strong forces and viscoelastic media

Matthias Fuchs

Fachbereich Physik, Universität Konstanz



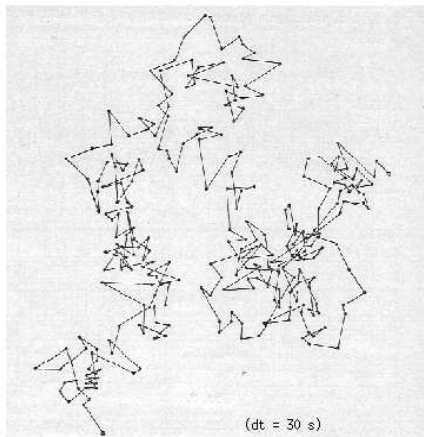
New Frontiers in Non-equilibrium Statistical Physics 2015

- Intro to dense colloidal dispersions
- Active microrheology: Theory
- Delocalization transition
- Length Scales & Populations
- Steady state motion
- Transient dynamics



Intro to dense colloidal dispersions

Brownian motion



The trajectory of a
0.53 micron particle
J. Perrin "Atoms" 1916



Jean Perrin

<http://ndburl.prize.org/physics/laureates/1926>

- Stokes–Einstein–Sutherland:

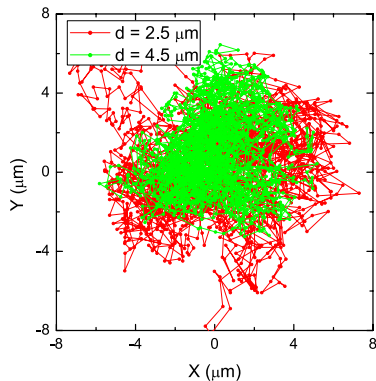
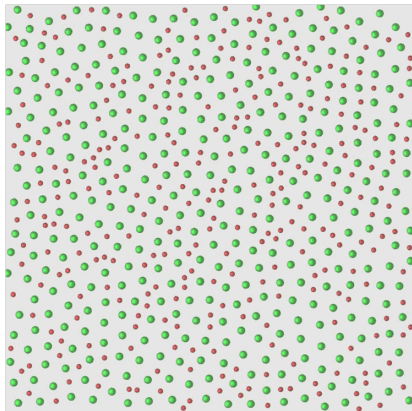
$$D_0 = \frac{k_B T}{\zeta_0}$$

- mean squared displacement

$$\langle (\mathbf{r}(t) - \mathbf{r}(0))^2 \rangle = 6D_0 t$$

Displacement fluctuations in glass

Binary glass in $d = 2$



finite displacements

$$\langle (\mathbf{r}_i(t) - \bar{\mathbf{r}}_i)^2 \rangle < \infty$$

with center of trajectory:

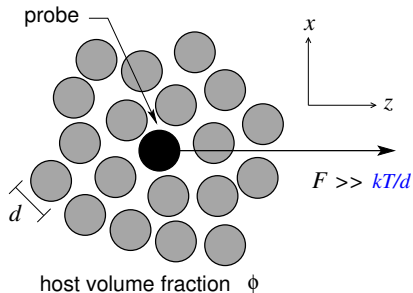
$$\bar{\mathbf{r}}_i = \frac{1}{\Delta t} \int_0^{\Delta t} dt \mathbf{r}_i(t)$$



Active microrheology: Theory

Active microrheology

- Hard-sphere suspension
- Fluid and glass states
- Probe displacement under strong forces
- Tails, heterogeneities etc.
- Parameters: ϕ and F



Microscopic (statistical) description

Model: Interacting Brownian particles, **no HI** ($D_0 = \frac{k_B T}{\zeta_0}$)

Smoluchowski equation:

$$\partial_t \Psi = \Omega \Psi$$

Micro rheology

$$\Omega = \sum_i^N D_0 \partial_i \cdot (\partial_i - \beta \mathbf{F}_i) + D_s \partial_s \cdot (\partial_s - \beta \mathbf{F}_s) - D_s \partial_s \cdot \mathbf{F}^{\text{ex}}$$

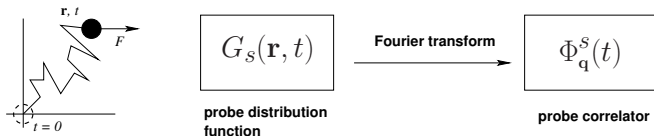
ITT force-velocity relation* (nonlinear, exact)

$$\zeta \langle \mathbf{v} \rangle_{t \rightarrow \infty} = F^{\text{ex}} \quad , \quad \zeta = \zeta_0 + \frac{1}{3k_B T} \int_0^\infty dt \langle \mathbf{F}_s e^{\Omega_{\text{irr}}^\dagger t} \mathbf{F}_s \rangle^{(e)}$$

[* Gazuz, Puertas, Voigtmann, MF, Phys. Rev. Lett. **102**, 248302 (2009)]

Microscopic (statistical) description (II)

Mode coupling approximation:



Evolution equation:

$$\partial_t \Phi_q^S(t) + \frac{1}{\tau_q(F)} \Phi_q^S(t) + \int_0^t m_q(t-t') \partial_{t'} \Phi_q^S(t') dt' = 0$$

Parallel relaxation channels*:

$$m_q(t) = \mathcal{F}\{\Phi_q^S(t), \Phi_q^{\text{host}}(t), F\}$$

Input (full MCT):

Equilibrium structure of the host at ϕ

[* Lang et al., PRL **105**, 125701 (2010)]

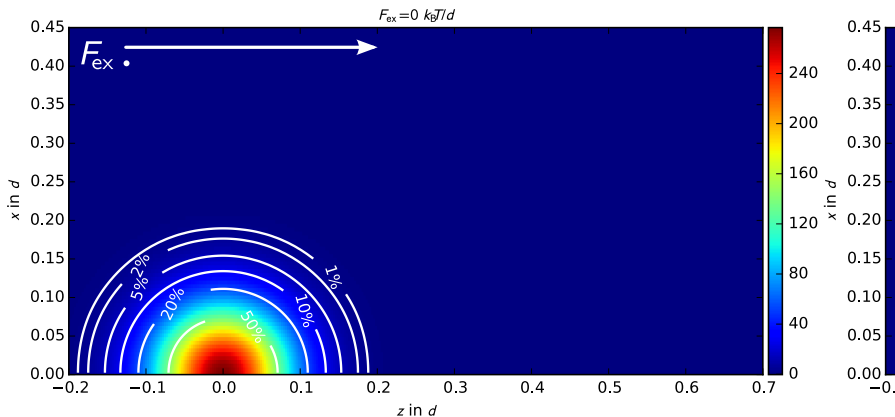


Delocalization transition

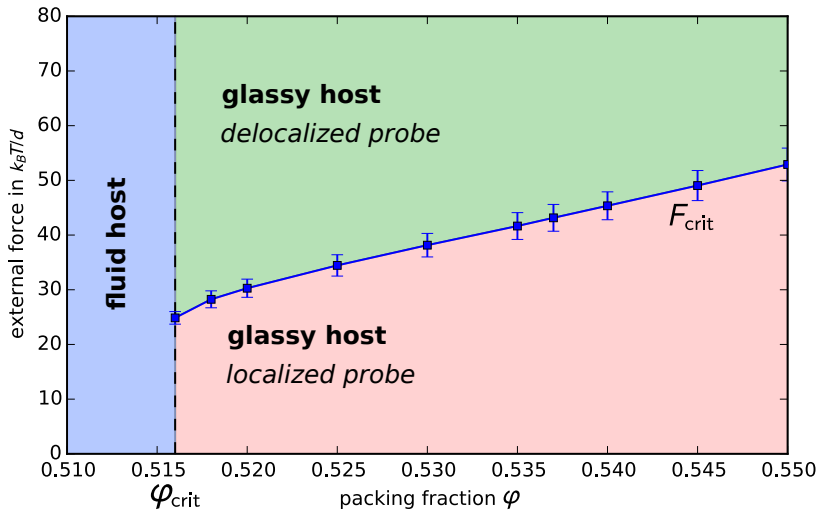
Tracer Nonergodicity Parameter in Real Space

$$f^s(\mathbf{r}, t) = \lim_{t \rightarrow \infty} \rho^s(\mathbf{r}, t)$$

$$\varphi = .537$$

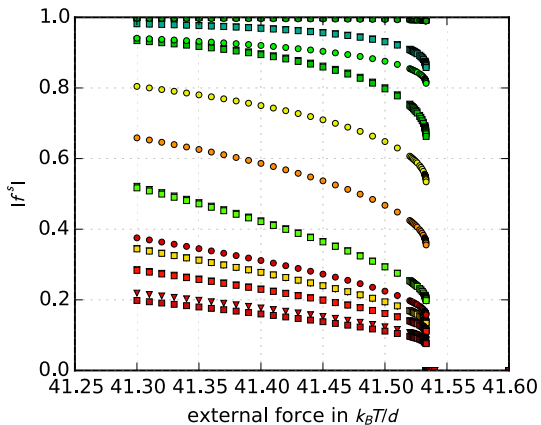


Delocalization transition: phase diagram (MCT)



Delocalization transition

$$\phi = 0.537$$



● type B transition



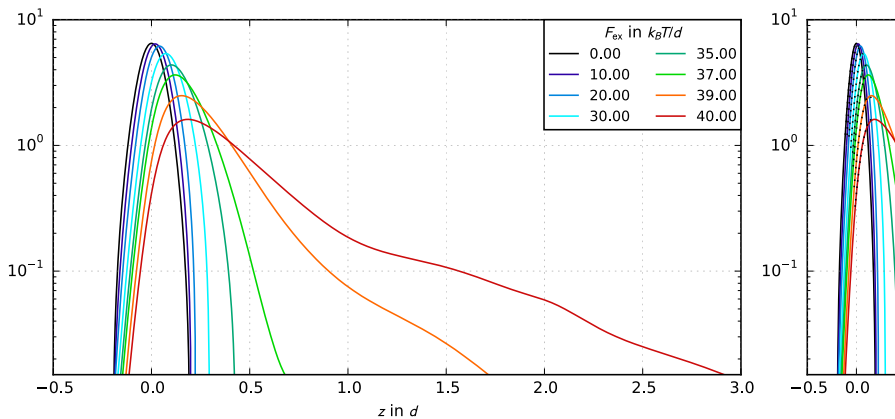
Length Scales & Populations

Exponential tails

Glassy host · $F \lesssim \text{threshold}$ · **localized probe**

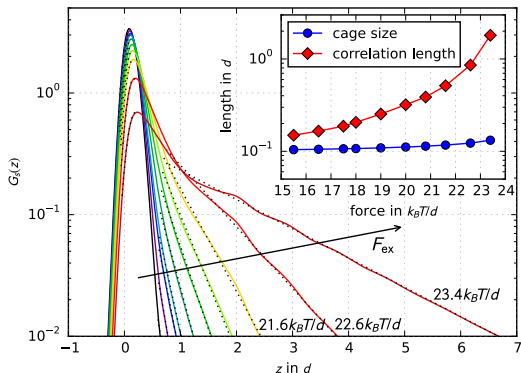
MCT, $\phi = 0.516$

Marginal probability distribution *parallel* to the force



Probe distribution function $G_s(z, t \rightarrow \infty)$

Localized probe



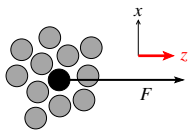
$$G_s^{\text{tail}} \propto \exp \left[-\frac{z}{\xi(F)} \right]$$

Growing correlation length ξ



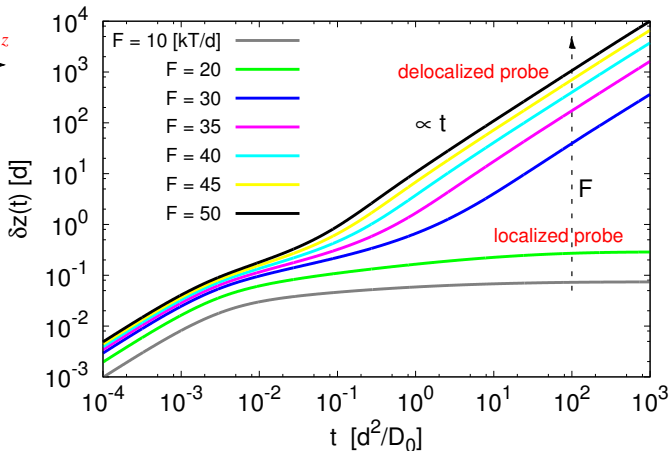
Steady state motion

Probe motion in force direction: average drift $\delta z(t)$



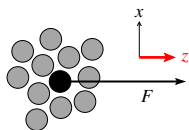
**GLASSY
HOST**

MCT
 $\phi = 0.516$



- Delocalization transition

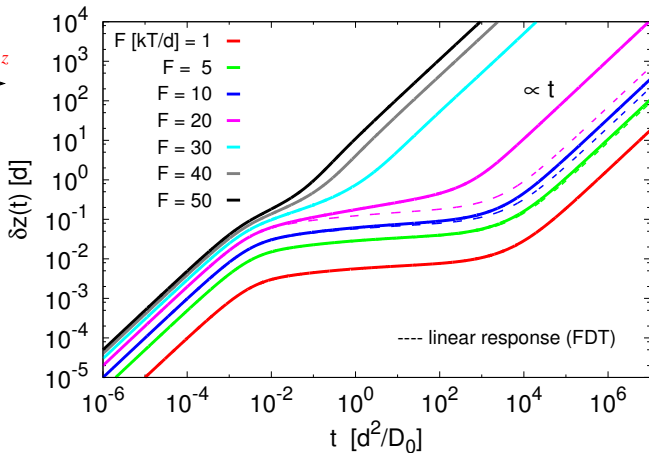
Probe motion in force direction: average drift $\delta z(t)$



FLUID
HOST

MCT

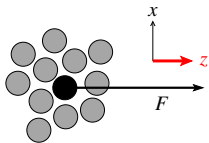
$\phi = 0.515$



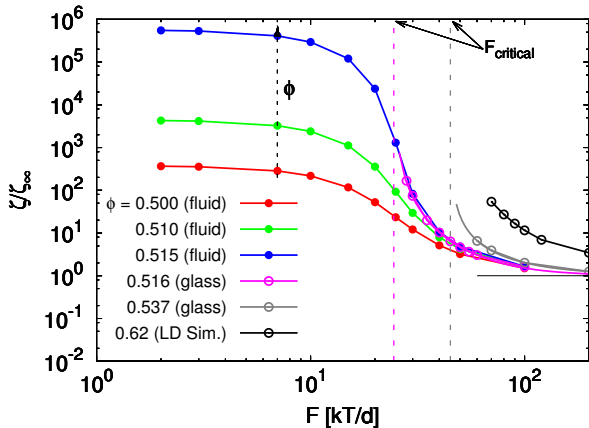
- Steady probe velocity for long times \rightarrow friction coefficient

Nonlinear friction coefficient: $v = [\zeta(F)]^{-1}F$

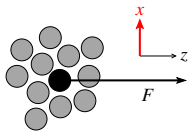
Force-thinning



$$\zeta_{\infty} = \zeta(F \rightarrow \infty)$$

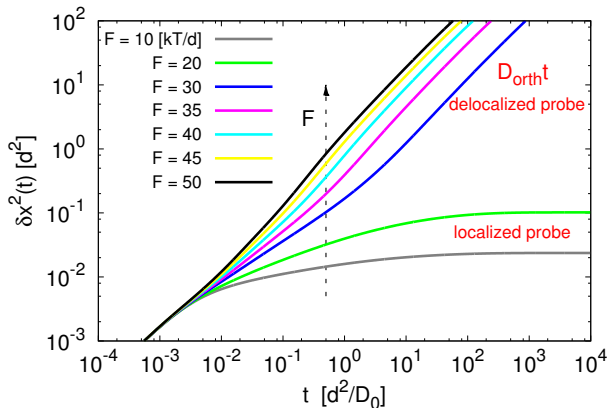


Probe motion in perpendicular direction: $\delta x^2(t)$



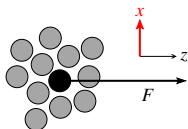
**GLASSY
HOST**

MCT
 $\phi = 0.516$

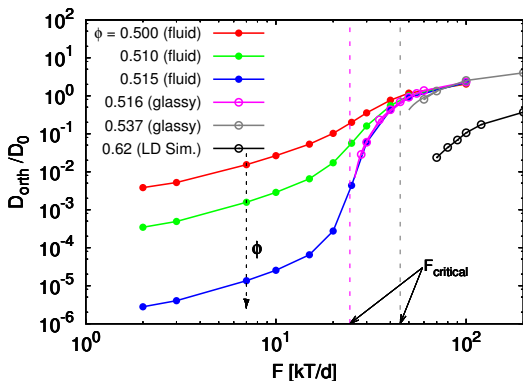


- Delocalization transition
- Diffusive motion for long times

Nonlinear diffusion coefficient - perpendicular direction



$$\delta x^2 \sim D_{\text{orth}} t$$



- Orthogonal diffusion enhanced by increasing F

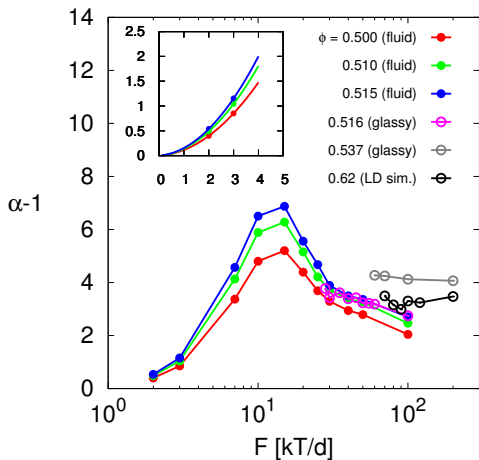
Stokes-Einstein Relation

Equilibrium:

$$\frac{D(\varphi)\zeta(\varphi)}{kT} = 1$$

Nonequilibrium:

$$\frac{D_{orth}(\varphi, F)\zeta(\varphi, F)}{kT} = \alpha(\varphi, F)$$





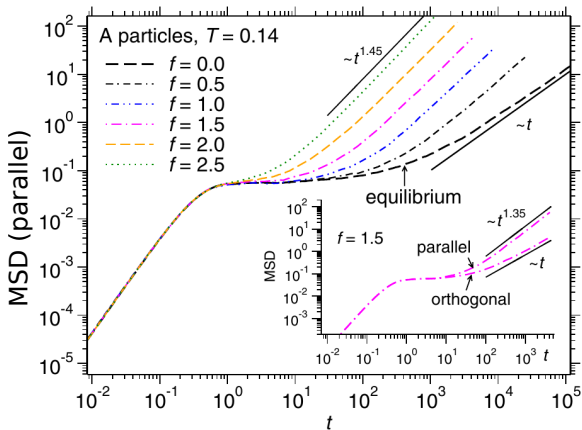
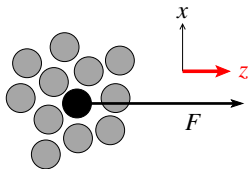
Transient dynamics

Probe motion in force direction: **superdiffusion**

MD Simulations - Yukawa mixture
Winter *et al.*, PRL **108** (2012)

MSD

$$\delta z^2(t) - [\delta z(t)]^2$$



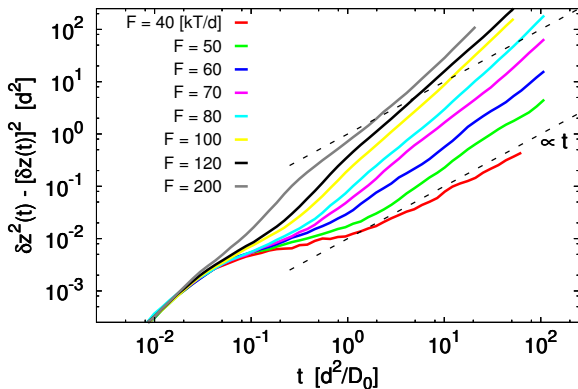
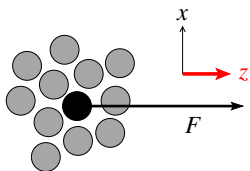
Probe motion in force direction: **superdiffusion**

Langevin Dynamics simulations

$$\phi = 0.62$$

MSD

$$\delta z^2(t) - [\delta z(t)]^2$$



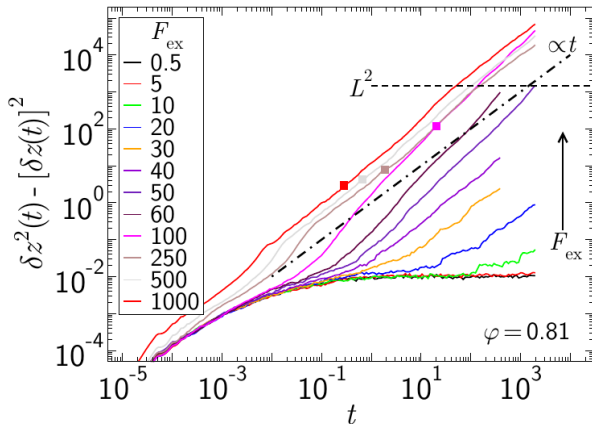
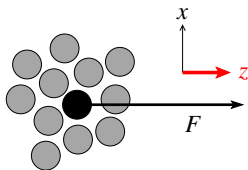
[A. Puertas (U. Almeria), unpublished]

Probe motion in force direction: **superdiffusion**

Brownian Dynamics simulations (2D)

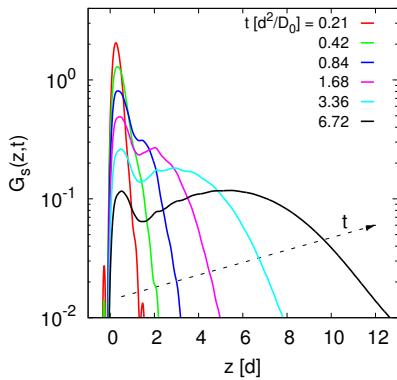
MSD

$$\delta z^2(t) - [\delta z(t)]^2$$



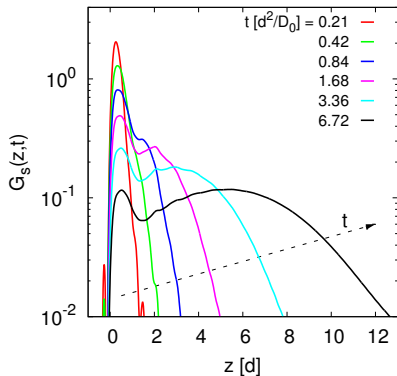
Evolution of the probe distribution function $G_s(z, t)$

Delocalized probe

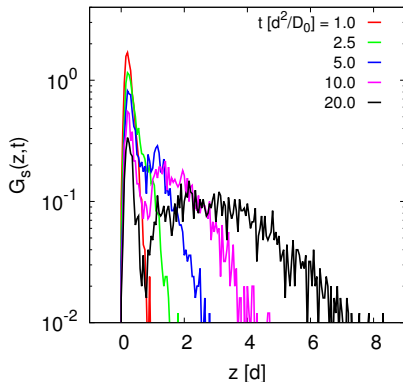


Evolution of $G_s(z, t)$: Delocalized probe

MCT, $\phi = 0.537$, $F = 50kT/d$



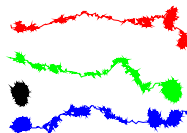
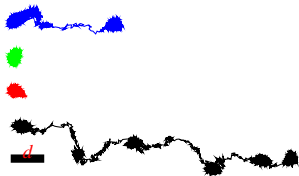
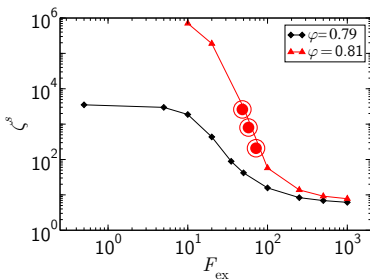
LD Simulations, $\phi = 0.62$,
 $F = 100kT/d$



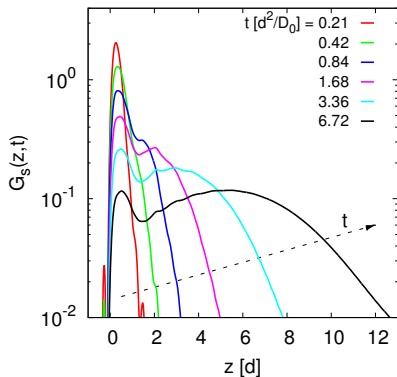
[A. Puertas (U. Almeria), unpublished]

Trajectories

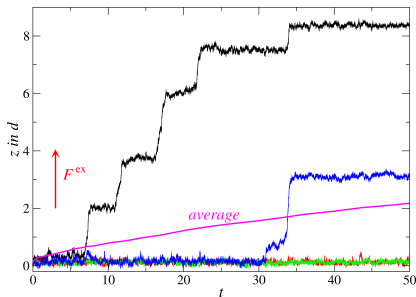
- Brownian dynamics simulations
- individual trajectories in the **glassy host**
- intermittent dynamics



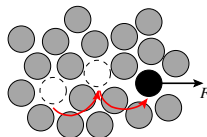
Evolution of $G_s(z, t)$: Delocalized probe



Probe trajectories
Brownian dynamics simulations



- Active and inactive probes



Conclusions

- micro-rheology delocalization threshold $F_c^{\text{ex}} \gg k_B T / \sigma$
- force-induced diffusion \perp force
- “two population” behavior close to depinning
- rare large excursions on diverging length scale

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Thank you for your attention