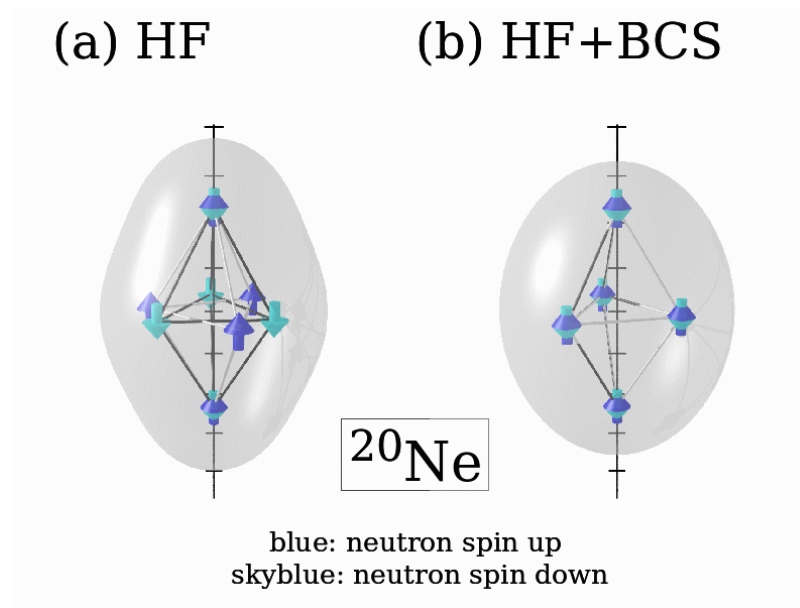


Visualization of nuclear many-body correlations with the most probable nucleon arrangement

Yusuke Tanimura and Moemi Matsumoto (Tohoku University)



← Most probable arrangements of neutrons in ^{20}Ne

Many-body theory

- Solution of Schrödinger eq. $H\Psi = E\Psi$

$$\rightarrow \Psi(\mathbf{r}_1\sigma_1, \mathbf{r}_2\sigma_2, \dots, \mathbf{r}_A\sigma_A)$$

- Structure, observables, ...

- ✓ $\langle \Psi | O^{(1)} | \Psi \rangle = \text{Tr}[\rho^{(1)} O^{(1)}]$ -- ex. size, EM moment, EM transitions, ...

- ✓ $\langle \Psi | O^{(2)} | \Psi \rangle = \text{Tr}[\rho^{(2)} O^{(2)}]$ -- ex. binding energy, 2-body correlation, ...

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- **Method to visualize many-body correlations based on the full information of wave function**
- **Structure of light nuclei**
 - ✓ **Hartree-Fock**
 - ✓ **Hartree-Fock + BCS**

How to extract information out of a many-body wave function?

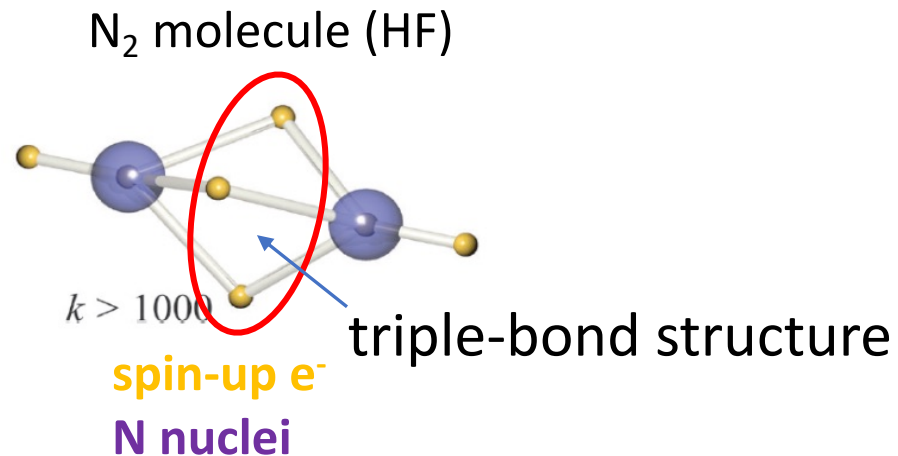
In quantum chemistry:

“Likely” arrangement of electrons in molecule
(N-body correlation in N-body system)

Liu et al., Phys. Chem. Chem. Phys. **18**, 13385 ('16).

Liu et al., Nat. Commun. **11**, 1210 ('20).

...



$$\langle \mathbf{X} \rangle = \int d\mathbf{X} \mathbf{X} |\Psi(\mathbf{X})|^2$$

$$\mathbf{X} = (\mathbf{r}_1\sigma_1, \dots, \mathbf{r}_N\sigma_N)$$

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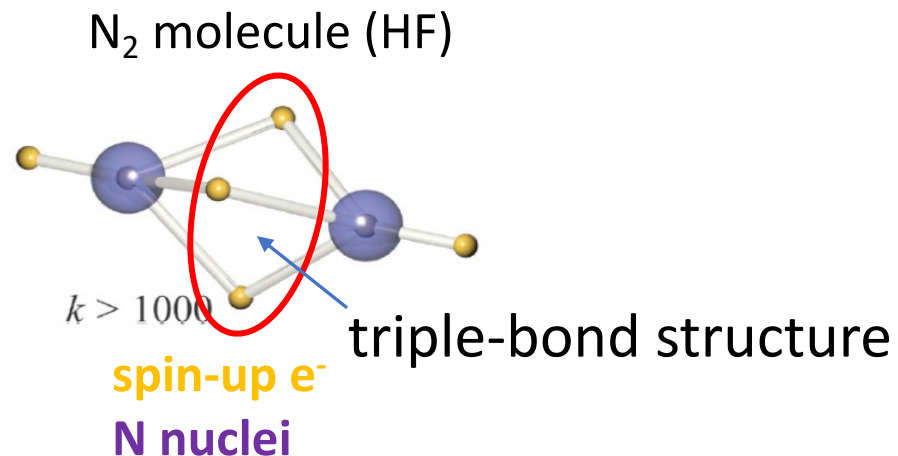
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Similar analysis can be made for nuclear wave functions?
→ Simple $|\Psi|^2$ maximization as the 1st step

$|\Psi|^2$ maximization

[M. Matsumoto and YT, PRC106, 014317 \('22\).](#)

System of N (= even) identical Fermions

Probability density of finding the particles simultaneously at $(\mathbf{r}_1\sigma_1, \dots, \mathbf{r}_N\sigma_N)$:

$$\rho^{(N)}(\mathbf{r}_1\sigma_1, \mathbf{r}_2\sigma_2, \dots, \mathbf{r}_N\sigma_N) = |\Psi(\mathbf{r}_1\sigma_1, \mathbf{r}_2\sigma_2, \dots, \mathbf{r}_N\sigma_N)|^2$$

to be maximized

$|\Psi|^2$ maximization

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to be maximized

$$\rho_0^{(N)}(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N) = |\Psi(\mathbf{r}_1 \uparrow, \dots, \mathbf{r}_{N/2} \uparrow, \mathbf{r}_{N/2+1} \downarrow, \dots, \mathbf{r}_N \downarrow)|^2$$

Spins fixed
#up = #down



$|\Psi|^2$ maximization

[M. Matsumoto and YT, PRC106, 014317 \('22\).](#)

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Spins fixed
#up = #down

Short-range & attractive nature of nuclear force

$\rightarrow \max(\rho_0^{(N)})$ is the global maximum of $|\Psi|^2$

Set up

α -cluster (and other possible) correlations in N=Z light nuclei

- Hartree-Fock calculation
 - ✓ SLy4 interaction
 - ✓ Time-reversal symmetry
 - ✓ Axial and reflection symmetries
 - ✓ HF+BCS: constant-gap approx.
- Maximization
 - ✓ $\rho^{(N)}_0(\mathbf{r}_1, \dots, \mathbf{r}_A)$ is maximized with conjugate-gradient method
 - ✓ initial $(\mathbf{r}_1, \dots, \mathbf{r}_A)$ is generated randomly

E. Chabanat et al., Nucl. Phys. **A635**, 231 (1998).
Vautherin, Phys. Rev. **C7**, 296 (1973).

$$\rho_0^{(N)}(\mathbf{r}_1, \mathbf{r}_2, \dots, \mathbf{r}_N) = |\Psi(\mathbf{r}_1 \uparrow, \dots, \mathbf{r}_{N/2} \uparrow, \mathbf{r}_{N/2+1} \downarrow, \dots, \mathbf{r}_N \downarrow)|^2$$

Remarks: symmetries

- Due to axial symmetry, $|\Psi|^2$ is invariant under rotation about z axis.

- Hartree-Fock wave function

$$\langle x_1, \dots, x_A | \Psi \rangle = \det B_n \det B_p = (\text{neutron determinant}) * (\text{proton determinant})$$

$$(B_q)_{ij} = \langle x_i | \varphi_j^{(q)} \rangle \quad q = p \text{ or } n$$

→ No explicit correlation between neutrons and protons

→ $|\Psi|^2$ does not depend on relative angle between n and p about z axis.

- The same applies to reflection symmetry

Remarks: correlations in Hartree-Fock

Mean-field (Hartree-Fock) theory

$$H = \sum_{i=1}^A \frac{\mathbf{p}_i^2}{2m} + \sum_{i>j}^A \underline{v(\mathbf{r}_i, \mathbf{r}_j)} \quad \rightarrow \quad \sum_{i=1}^A h(i) = \sum_{i=1}^A \left[\frac{\mathbf{p}_i^2}{2m} + v_{\text{MF}}(\mathbf{r}_i) \right]$$

2-body int. **"Mean field"**

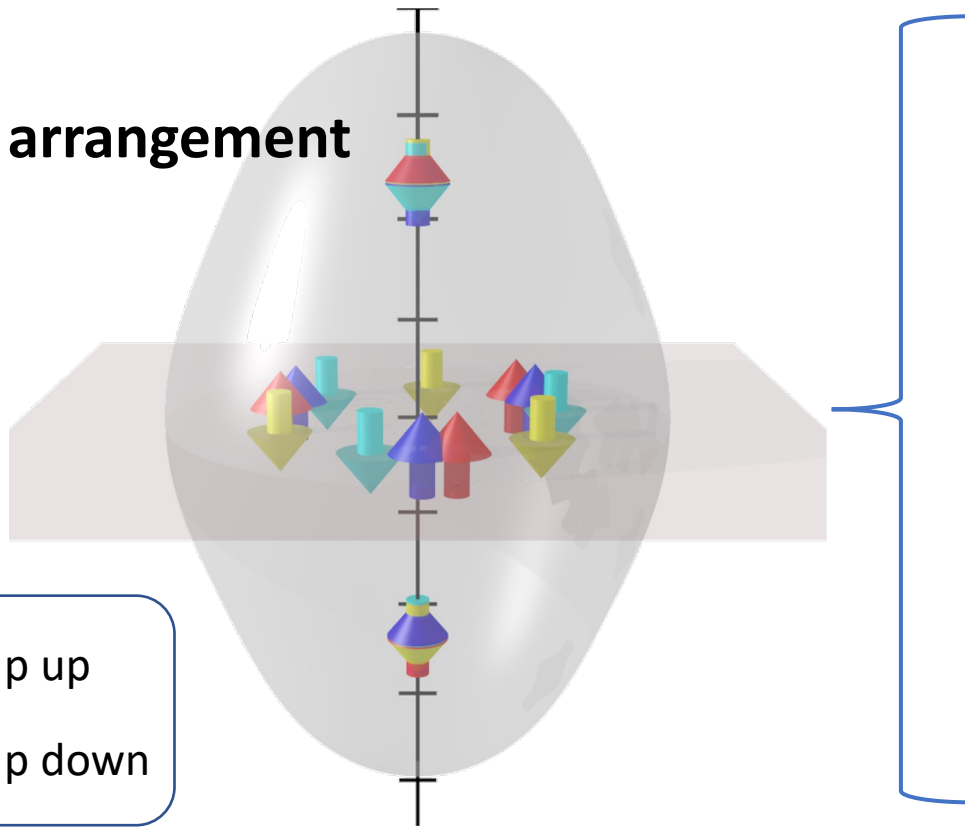
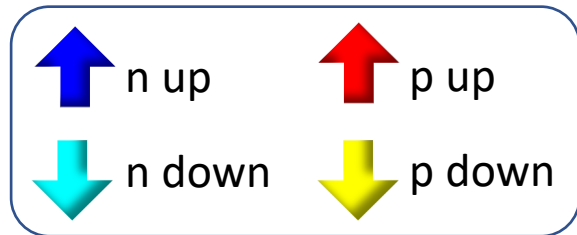
$$v_{\text{MF}}(\mathbf{r}) \sim \int d^3 r' v(\mathbf{r}, \mathbf{r}') \rho(\mathbf{r}')$$

Correlations in HF

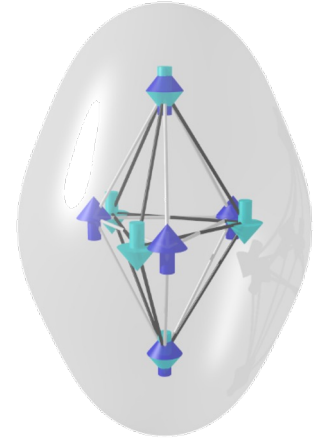
- Explicit: Pauli principle
- **Implicit/explicit: interaction through the mean field ("long-range correlation")**
 - size
 - deformation (symmetry breaking)

Result: HF ground state of ^{20}Ne

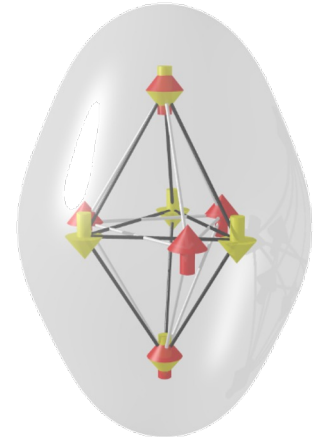
most probable arrangement
of all nucleons



neutrons →



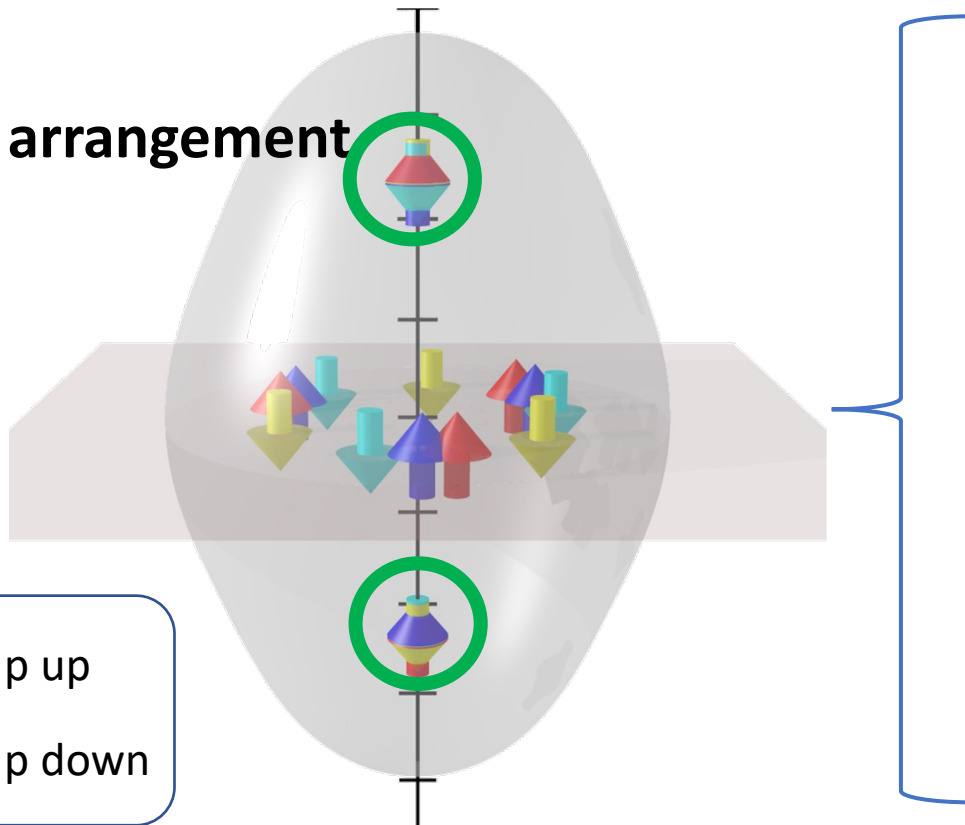
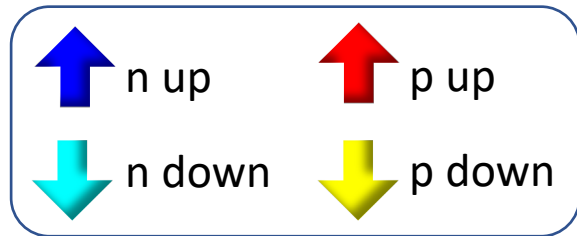
protons →



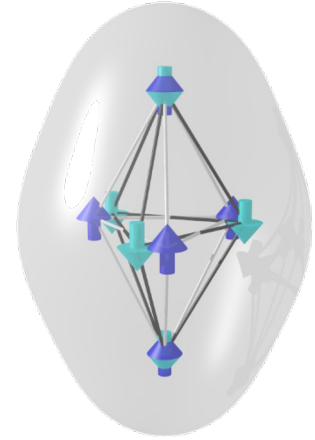
protons \approx neutrons

Result: HF ground state of ^{20}Ne

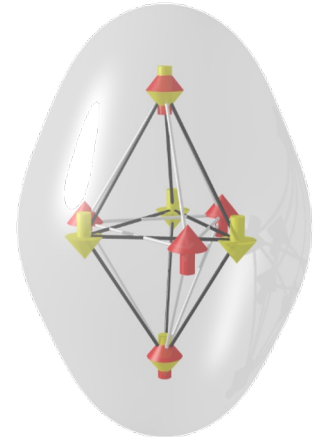
most probable arrangement
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neutrons \rightarrow



protons \rightarrow

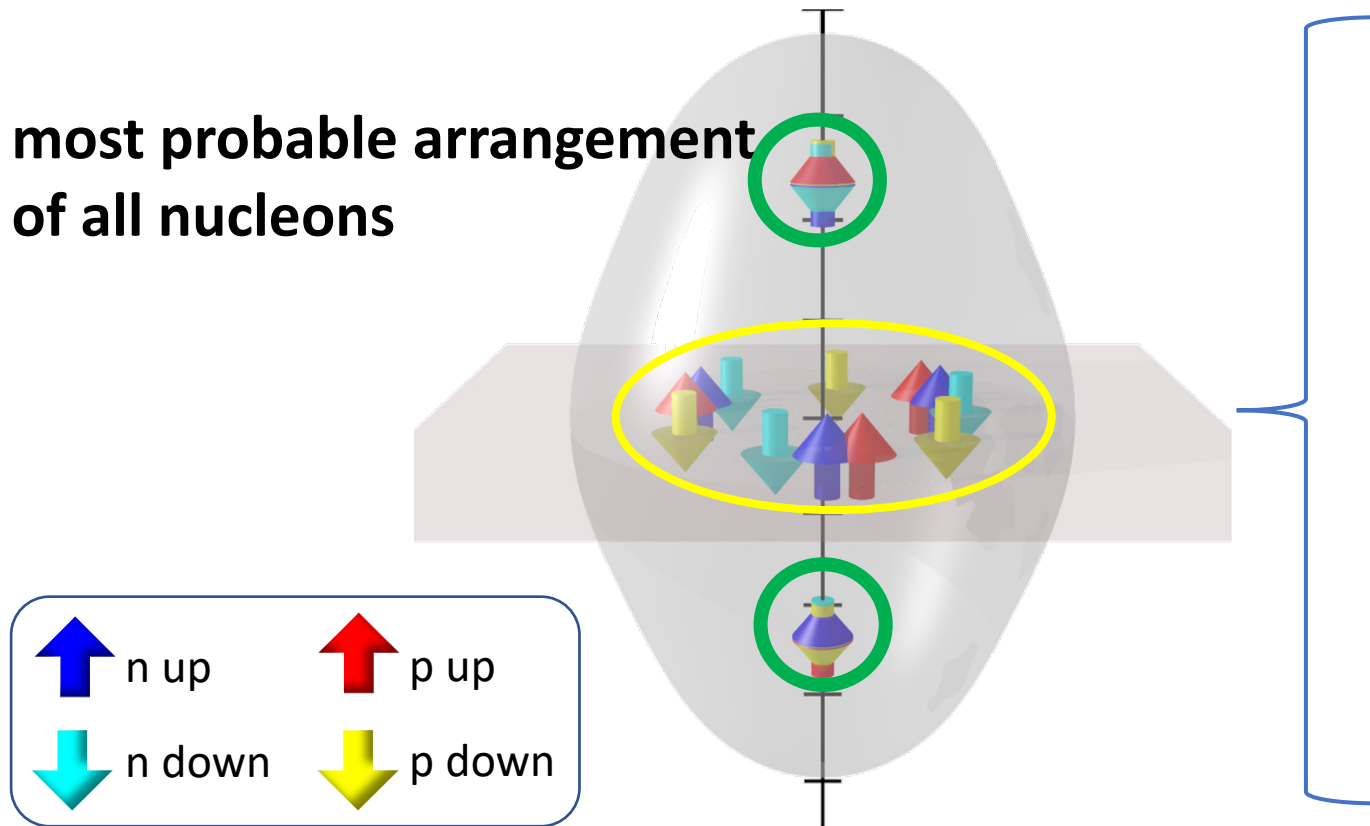


- α clusters at top/bottom

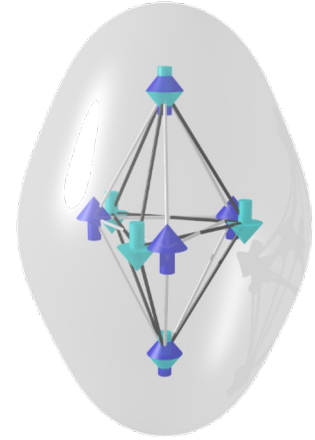
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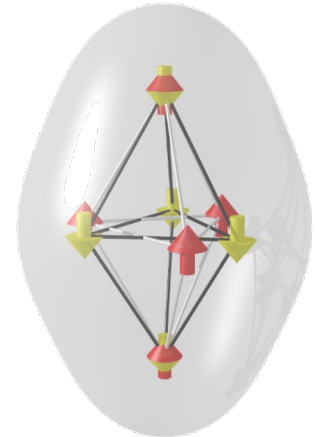
most probable arrangement
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neutrons →



protons →



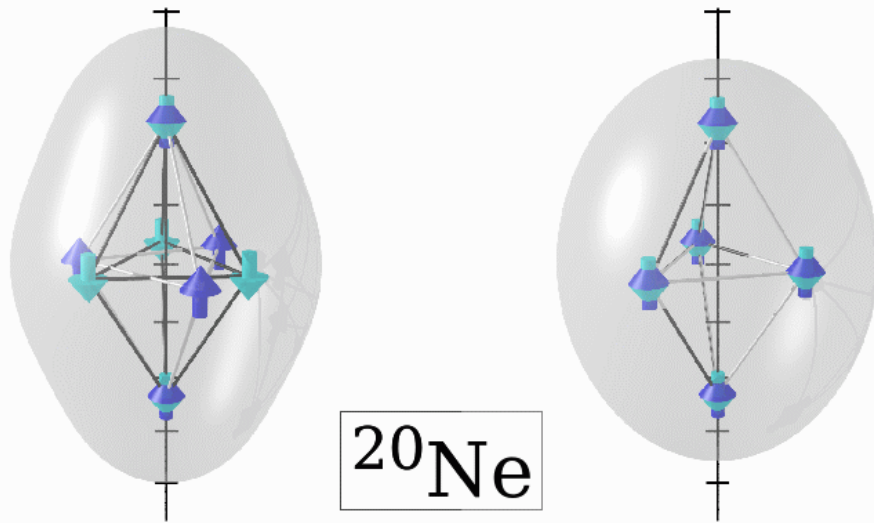
protons \approx neutrons

- α clusters at top/bottom
- 12 nucleons on xy plane

Results: HF+BCS g.s. of ^{20}Ne

(a) HF

(b) HF+BCS



blue: neutron spin up

skyblue: neutron spin down

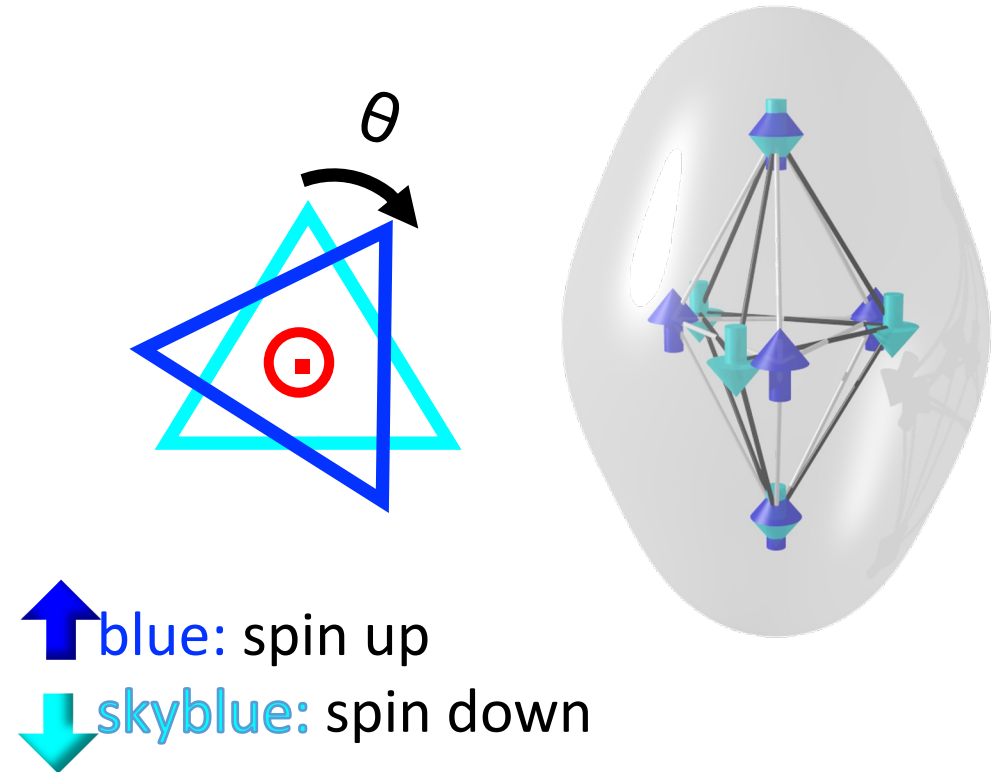
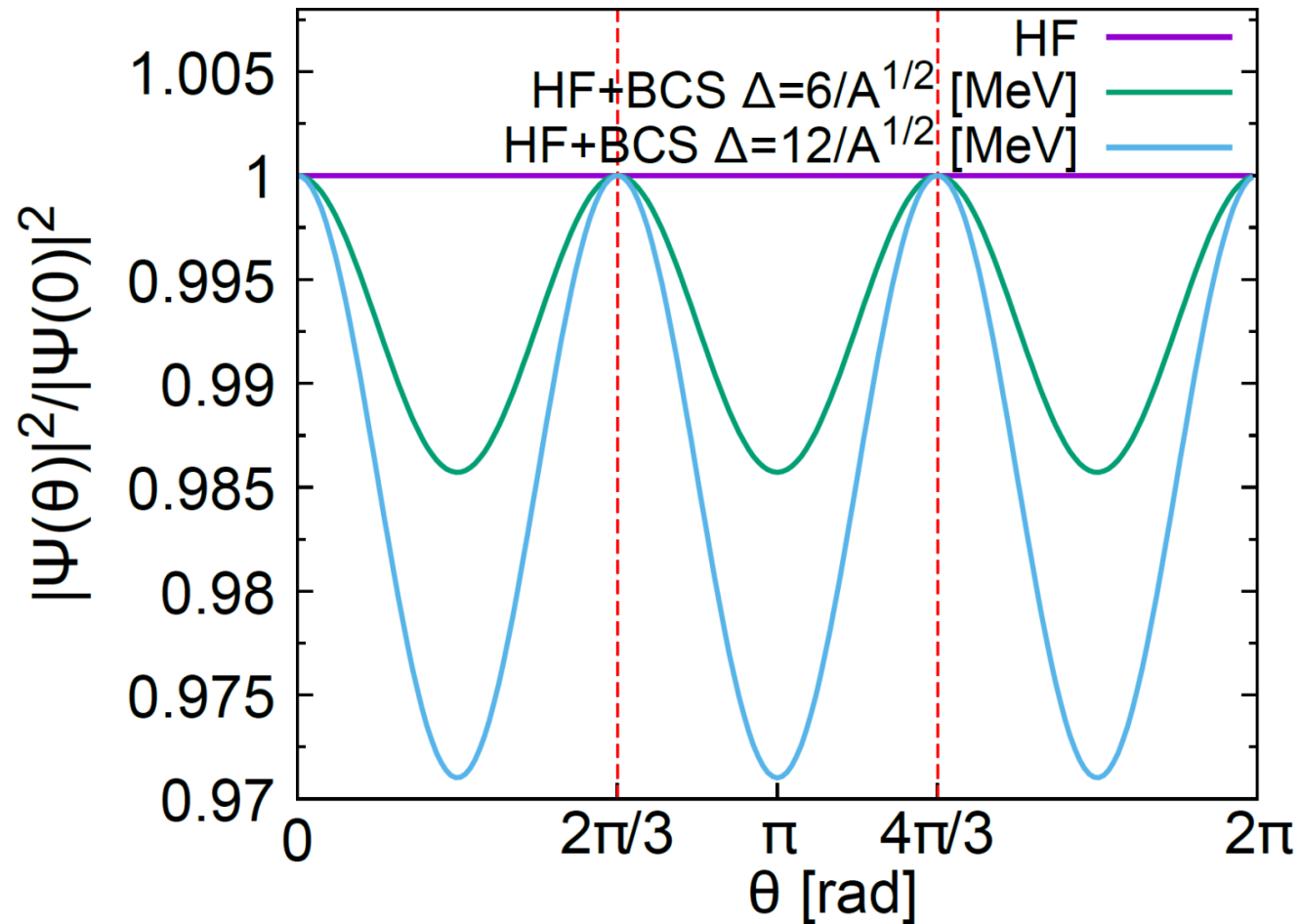
Most probable arrangement of neutrons

Protons \approx neutrons

Attractive correlation between
spin-up and -down

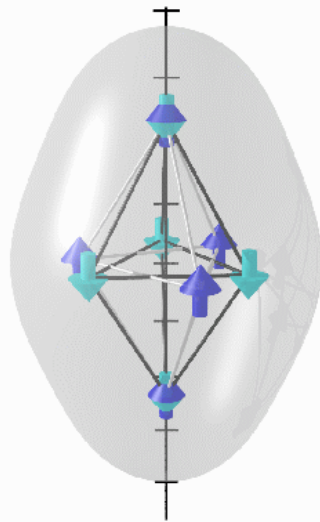
Result (1)²⁰Ne

Probability variation as a function of the relative angle (θ) around the z axis between spin-up and -down neutrons

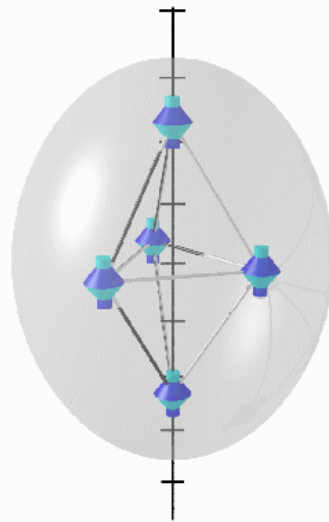


Results: HF+BCS g.s. of ^{20}Ne

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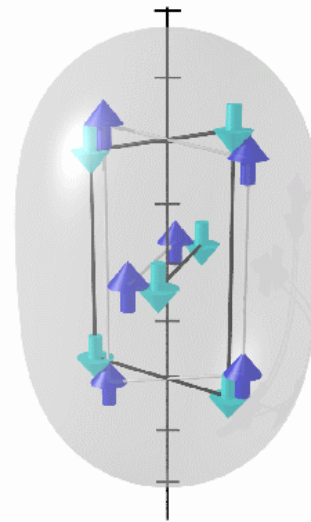
^{20}Ne

blue: neutron spin up
skyblue: neutron spin down

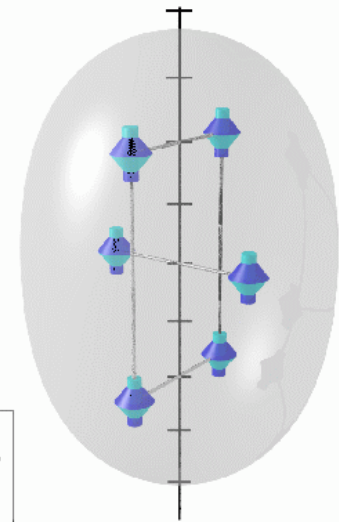
Most probable arrangement of neutrons

Protons \approx neutrons

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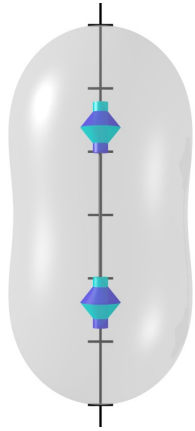


^{24}Mg

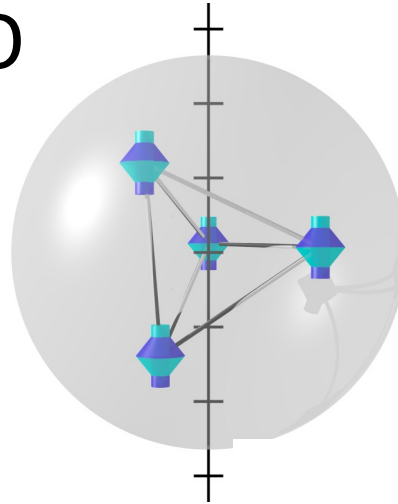
blue: neutron spin up
skyblue: neutron spin down

Ground states of ^8Be , ^{12}C , ^{16}O , ^{20}Ne , ^{24}Mg , and ^{28}Si

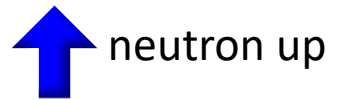
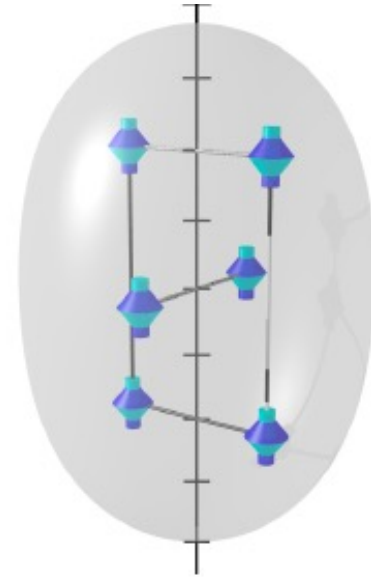
^8Be



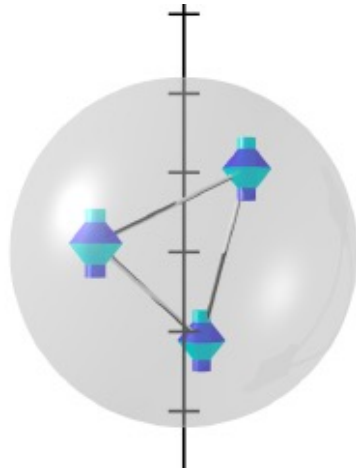
^{16}O



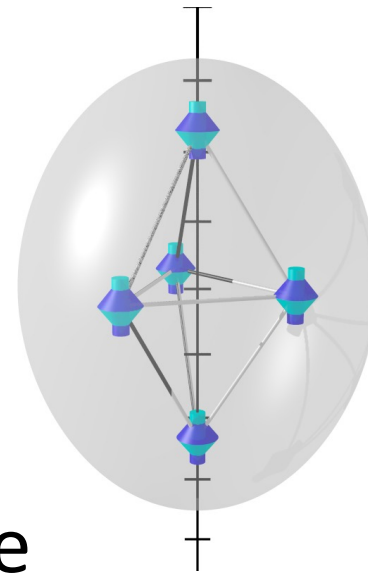
^{24}Mg



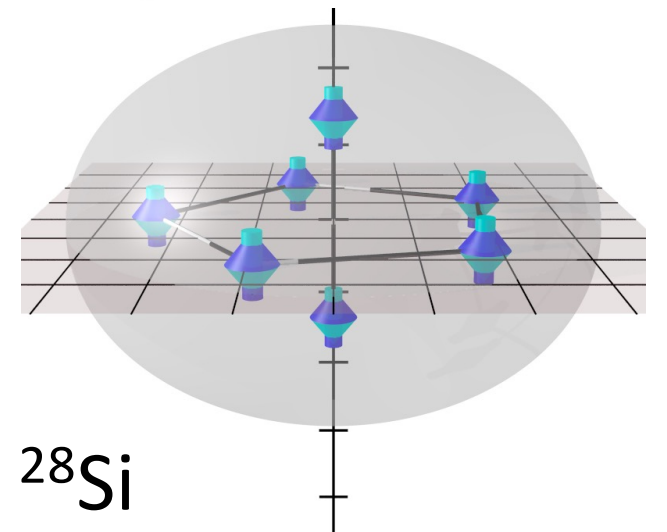
^{12}C



^{20}Ne

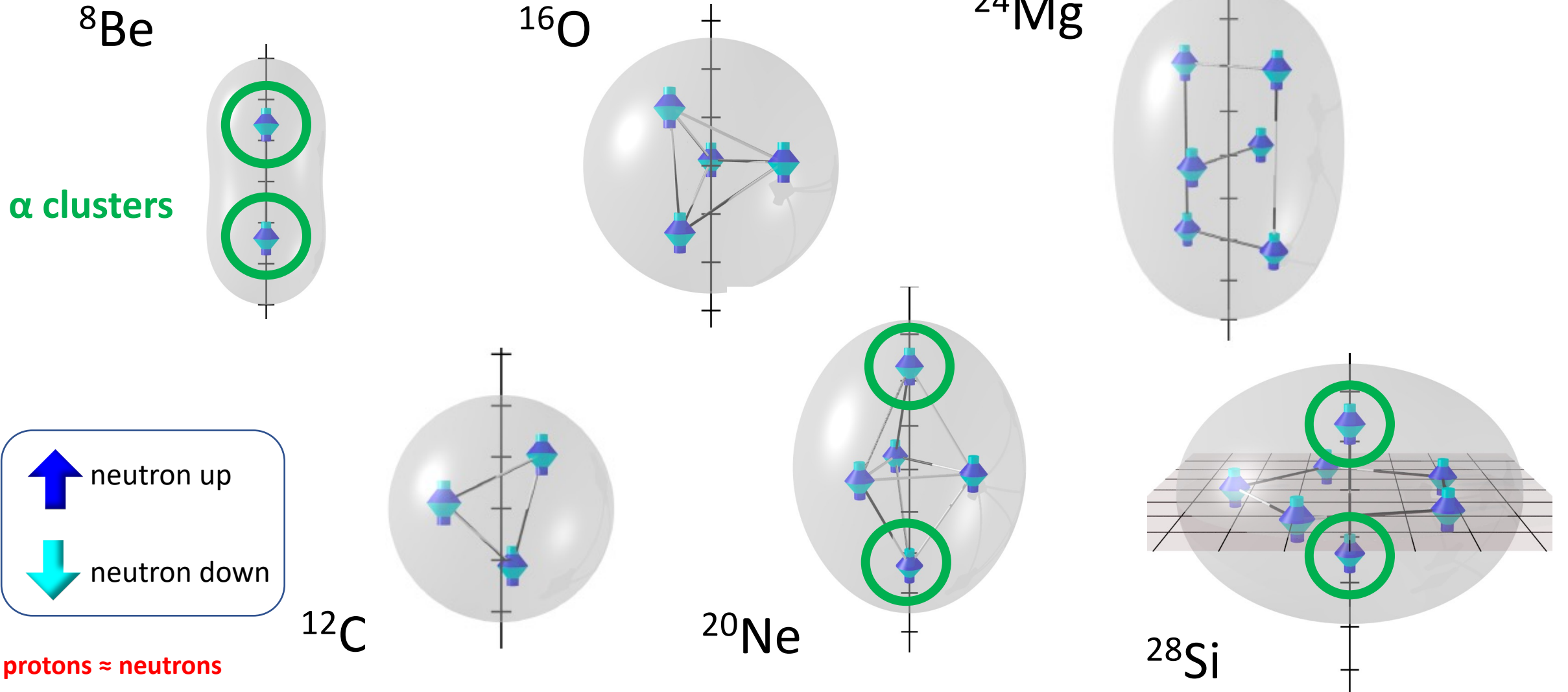


^{28}Si



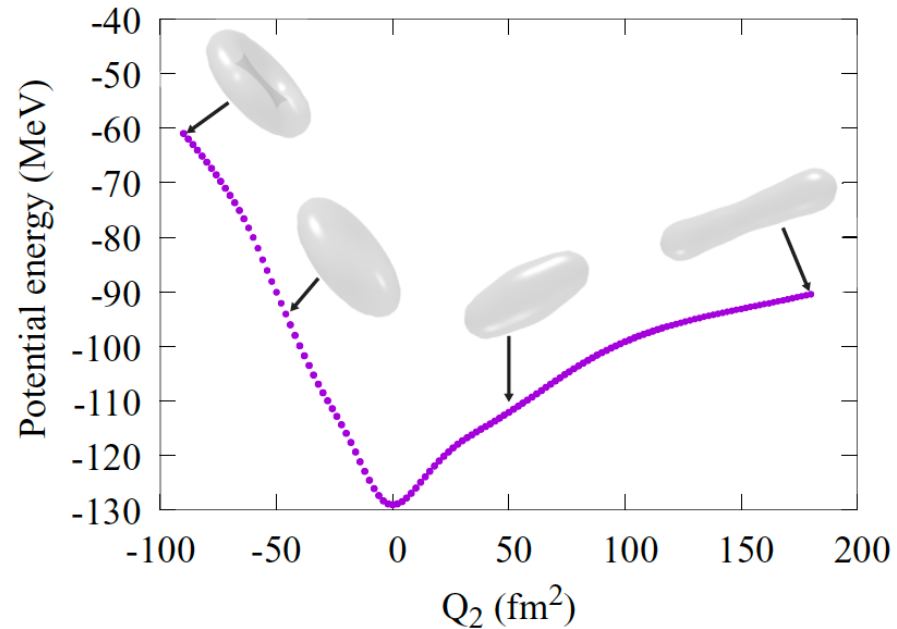
protons \approx neutrons

Ground states of ^8Be , ^{12}C , ^{16}O , ^{20}Ne , ^{24}Mg , and ^{28}Si



Results:

Deformed states of ^{16}O



← Potential energies obtained with Q_2 -constrained HF + BCS

$$Q_2 = \int d^3r r^2 Y_{20}(\hat{\mathbf{r}}) \rho(\mathbf{r})$$

FIG. 9. The HF+BCS potential-energy curve of the ^{16}O nucleus as a function of the quadrupole moment Q_2 . The HF + BCS result

Results:

Deformed states of ^{16}O

proton arrangement \approx neutron arrangement

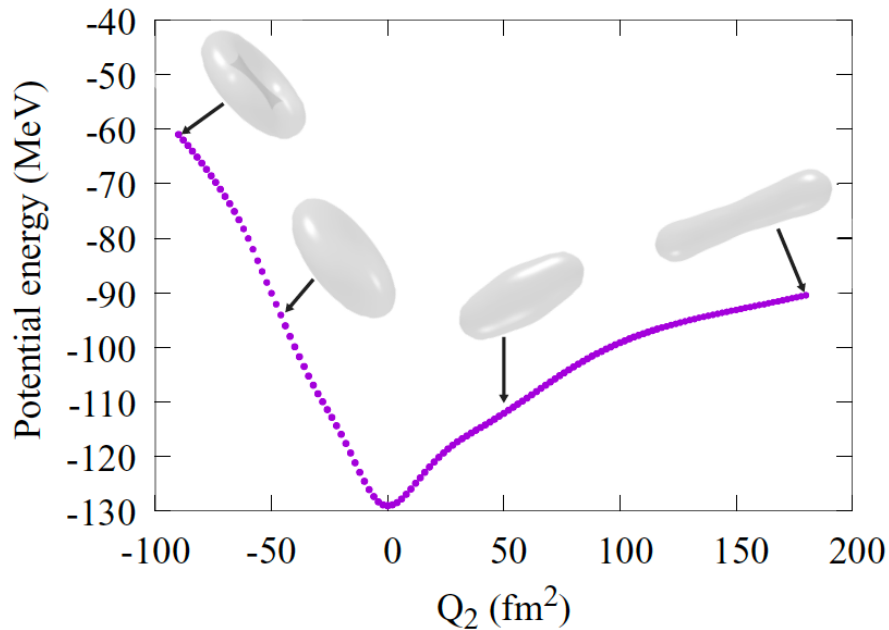
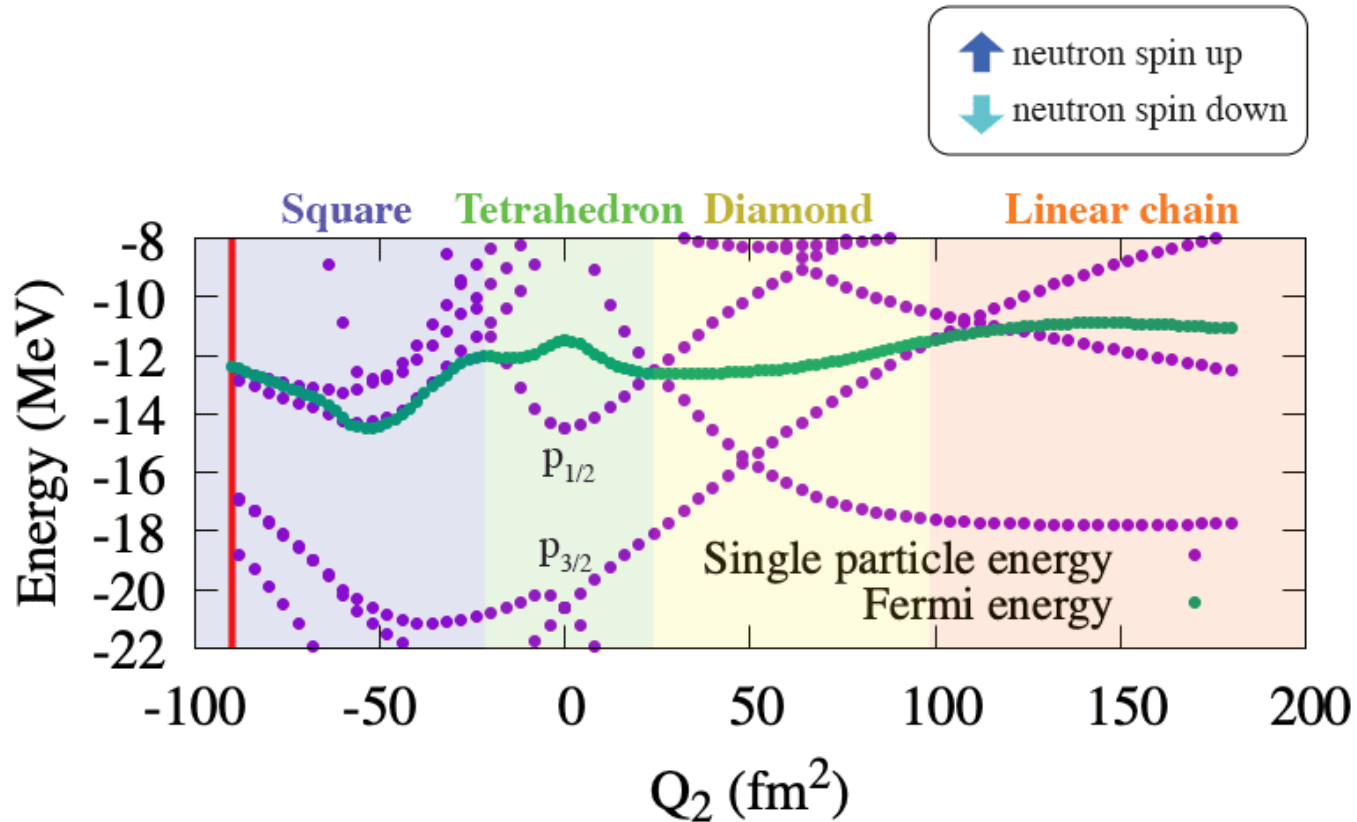
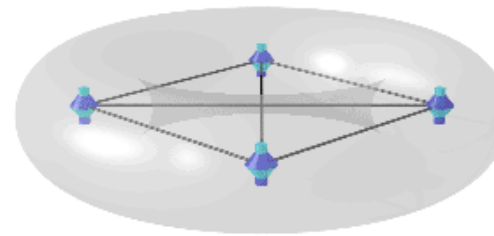
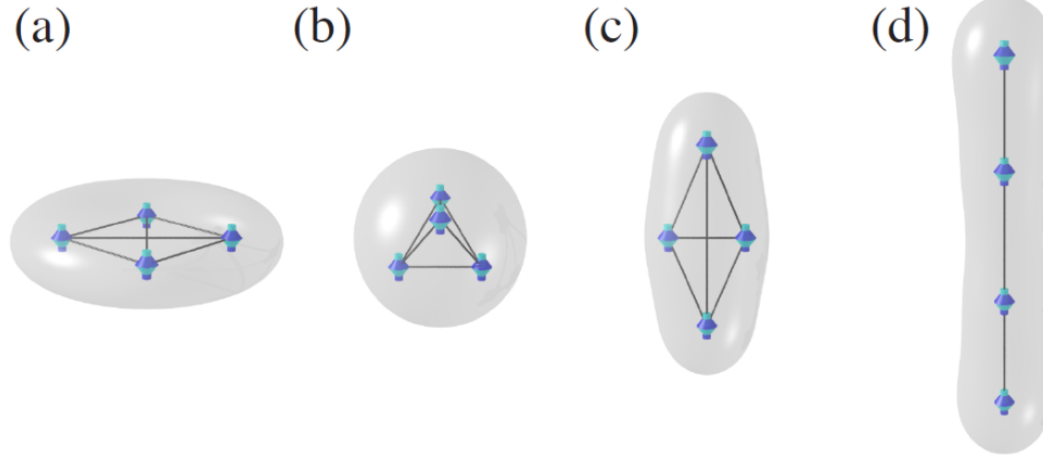
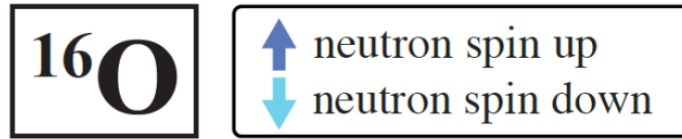


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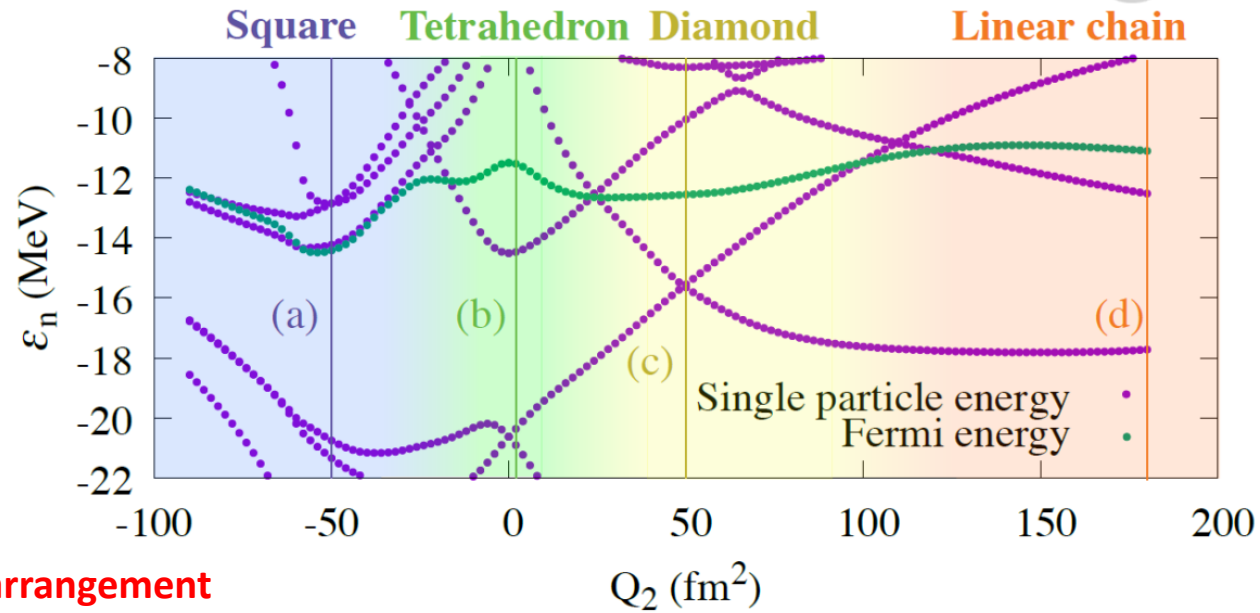
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Results: Deformed states of ^{16}O



Structure changes at alternations of the Fermi level



proton arrangement \approx neutron arrangement

Summary and perspectives

- $|\psi|^2$ -maximization method
 - ✓ cluster correlations in mean-field wave functions in light nuclei
 - ✓ Intuitive picture on correlations embedded in many-body state
 - ✓ **New viewpoint to a nuclear many-body wave function**

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- $|\psi|^2$ -maximization method
 - ✓ cluster correlations in mean-field wave functions in light nuclei
 - ✓ Intuitive picture on correlations embedded in many-body state:
New viewpoint to a nuclear many-body wave function
- Global behaviors of $|\psi|^2$ with Markov-chain Monte-Carlo sampling?
 - ✓ fluctuation around the maximum
 - ✓ local maxima
- More correlated states
 - n-p mixing (explicit n-p correlation)
 - ✓ (semi-)magic core/cluster appears?
 - RPA, GCM (collective motion/fluctuation)
- Other phenomena
 - ✓ molecular-bond structure
 - ✓ valence neutrons in n-rich nuclei
 - ✓ nucleon motions in reactions