

First-principles study of the lattice dielectric function in rutile TiO_2

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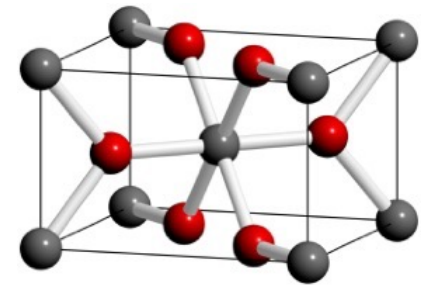
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- Lattice dielectric properties
- Anharmonic phonon theory
- Applied to strongly anharmonic rutile TiO_2



[TA et al. arXiv preprint arXiv:2210.15873 \(2022\).](#)

Dielectric function : 5G · 6G materials

5G:30-60GHz

6G:100-300GHz

- **dielectric loss** is important in **higher frequency region (~THz)**

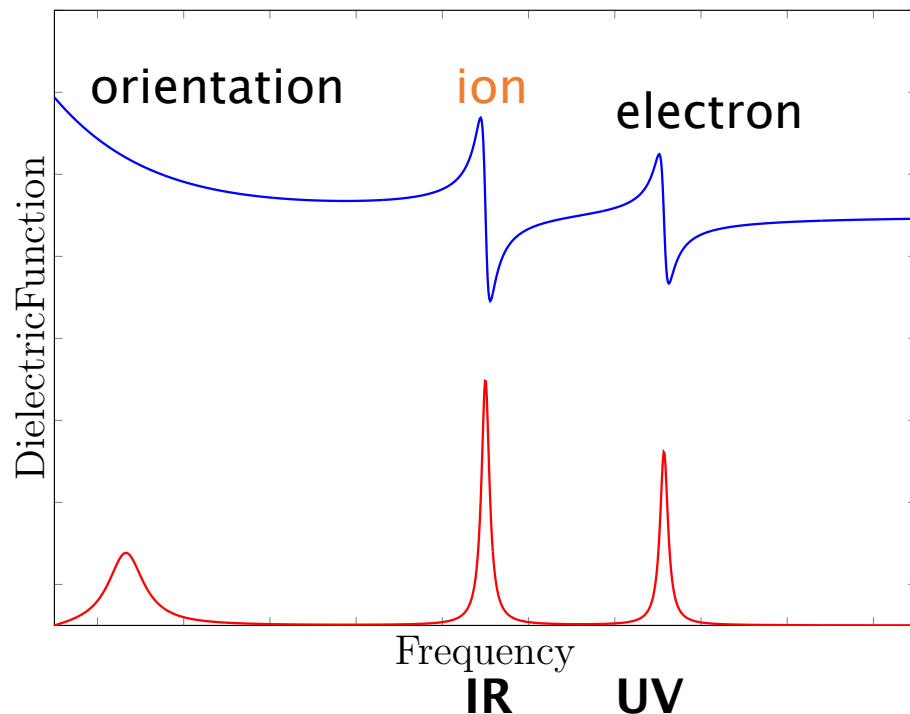
Dielectric loss

$$\tan \theta = \epsilon'' / \epsilon'$$

Dielectric function

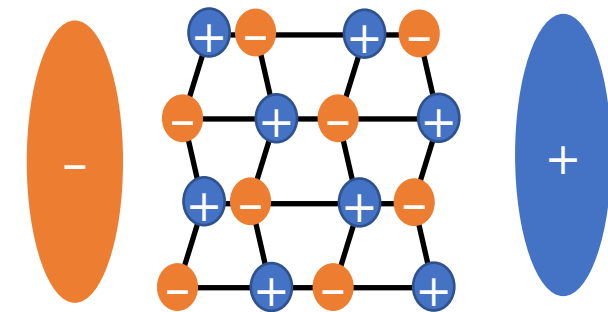
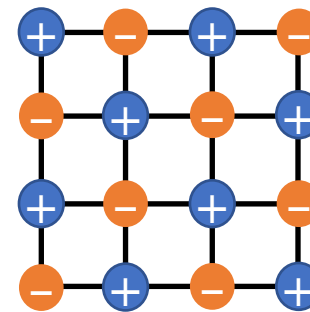
$$\epsilon(\omega) = \epsilon'(\omega) - i\epsilon''(\omega)$$

- **Ionic polarization** is dominant in THz region



Ionic polarization (IR-active phonon)

$$P = \sum_{\kappa} Z_{\kappa} r_{\kappa}$$



(+ and - move oppositely)

Lattice dielectric function by phonons

➤ Dielectric function of phonons^[1]

$$\epsilon_{\alpha\beta}(\omega) = \epsilon_{\alpha\beta}^{\infty} + \frac{1}{V_0} \sum_{(\mathbf{0},j)} \frac{S_{\alpha\beta}^j}{\omega_{0j}^2 - \omega^2 + 2\omega_{0j}\Sigma_{0j}(\omega)}$$

S :: mode oscillator strength

ω_0 :: phonon frequency

Σ :: phonon self energy
(real: frequency shift, imaginary: linewidth)

➤ Anharmonic phonon calculation

T. Tadano et al. 2014. *J. Phys. : Cond. Mat.* 26.22 (2014): 225402.

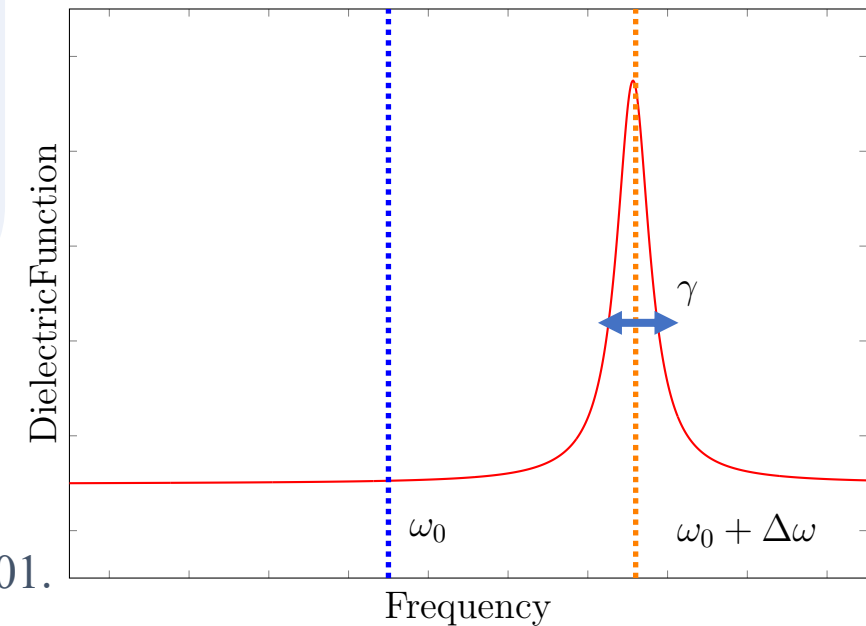
Hellman Olle et al. *Phys. Rev. B* 84.18 (2011): 180301.

Monacelli, Lorenzo, et al. *J. Phys.: Cond. Mat.* 33.36 (2021): 363001.

Phono anharmonic effect is essential for dielectric function

Lorentz model

$$\epsilon(\omega) = \epsilon^{\infty} + \sum_{(\mathbf{0},j)} \frac{\Delta\epsilon_{0j}\Omega_{0j}^2}{\Omega_{0j}^2 - \omega^2 + i\gamma_{0j}\omega}$$

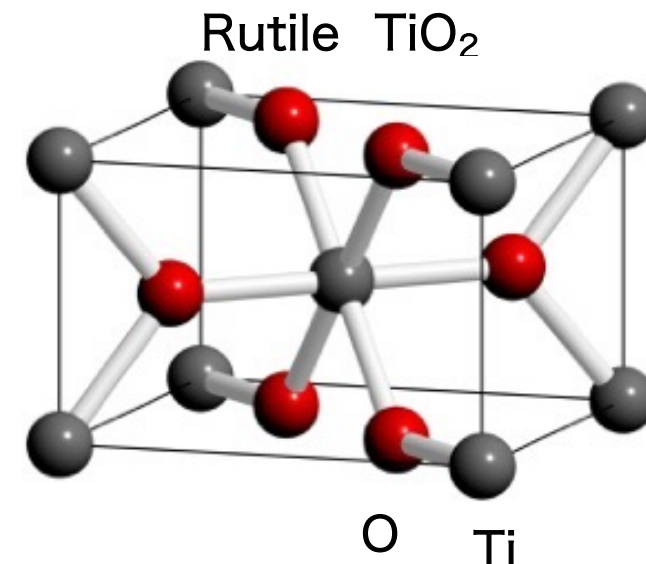


[1] R.A. Cowley. *Advances in Physics* 12(48):421–480, 10 1, 1963.

ab-initio calculation on Dielectric properties

DFT calculation on Dielectric properties from anharmonic-phonon theory

- anharmonic phonon calculations
 - Thermal Conductivity
 - Structure phase transition
 - **NOT** so many optical properties...
- Objective :: ab-initio dielectric function on rutile TiO_2
 - Self-consistent approach for **Strongly Anharmonic materials**
 - **Frequency-dependence** of self-energy



- Strong Anharmonicity
- gigantic dielectric constant (incipient ferroelectric)

[1] Fugallo, Giorgia, Benoit Rousseau, and Michele Lazzeri *Physical Review B* 98.18 (2018): 184307.

[2] Tong, Zhen, et al *Physical Review B* 101.12 (2020): 125416

Harmonic Phonon Theory

Lattice energy expansion (with atomic displacements)

$$\delta U = U_2 + \cancel{U_3} + \cancel{U_4} + \dots$$

Harmonic Hamiltonian (ω :: harmonic phonon frequency)

$$H_0 = \sum_{\mathbf{q}j} \omega_{\mathbf{q}j} A_{\mathbf{q}j}^\dagger A_{\mathbf{q}j}$$

Green's function

$$G_{\mathbf{q},jj'}(\omega) = \frac{2\omega_{\mathbf{q}j}\delta_{jj'}}{\omega_{\mathbf{q}j}^2 - \omega^2}$$

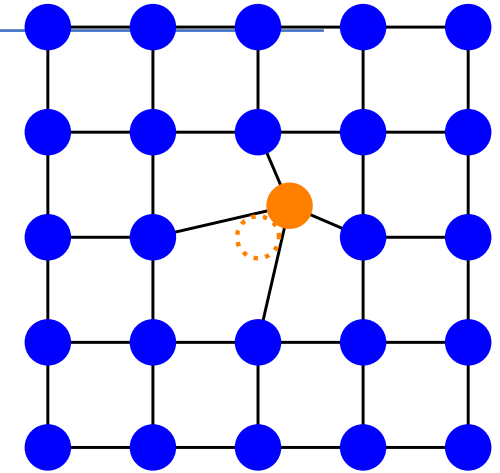
× NO phonon-phonon interaction

→ NO phonon linewidth

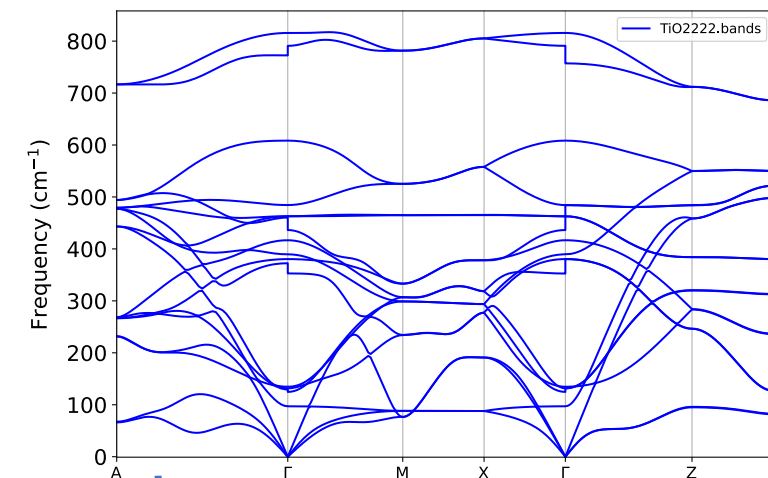
→ impossible to calculate dielectric properties

We need higher order terms to include phonon-phonon interaction

Atomic displacement



Phonon Band structure



Anharmonic Phonon Theory^[1]

Lattice energy expansion (with atomic displacement)

$$\delta U = U_2 + \underline{U_3 + U_4 + \dots}$$

anharmonic terms (3rd,4th,...)

Hamiltonian

$$H = H_0 + U_3 + U_4 + \dots$$

harmonic 3phonon 4phonon

Green's function

$$G_{\mathbf{q},jj'}(\omega) = \frac{2\omega_{qj}\delta_{jj'}}{\omega_{qj}^2 - \omega^2 - 2\omega_{qj}\Sigma_{\mathbf{q},jj'}(\omega)}$$

○ Include phonon-phonon interaction

① Frequency shift $\Delta\omega(q, \omega, T) = -\text{Re} \Sigma_q(\omega_q)$

② linewidth $\Gamma(q, \omega, T) = \text{Im} \Sigma_q(\omega_q)$

Anharmonic effect give rise to phonon frequency-shift & linewidth



[1] T. Tadano, Y. Gohda, & S. Tsuneyuki. (2014). *Journal of Physics: Condensed Matter*, 26(22), 225402

Estimate phonon properties from DFT^[1]

- Lattice energy expansion (with atomic displacement)

$$\begin{aligned}\delta U &= U_2 + U_3 + U_4 + \dots \\ &= \frac{1}{2} \sum_{i,j} \Psi_{i,j} u_i u_j + \frac{1}{3!} \sum_{i,j} \Psi_{i,j,k} u_i u_j u_k + \frac{1}{3!} \sum_{i,j} \Psi_{i,j,k,l} u_i u_j u_k u_l + \dots \\ &= \mathbf{b} \cdot \Psi\end{aligned}$$

- Force acting on atoms

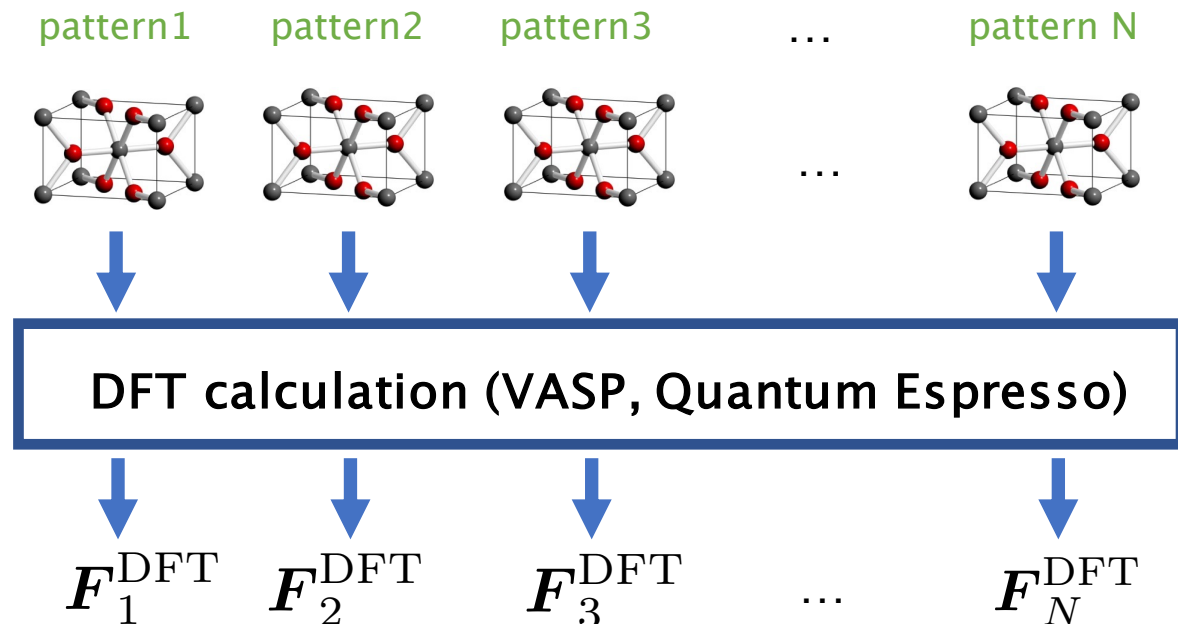
$$\mathbf{F} = -\frac{\partial U}{\partial \mathbf{u}} \equiv \mathbf{A} \Psi$$

- Least square fitting for IFCs

$$\tilde{\Psi} = \arg \min_{\Psi} |\mathbf{A} \Psi - \mathbf{F}^{\text{DFT}}|^2$$

Interatomic force constant
(parameter)

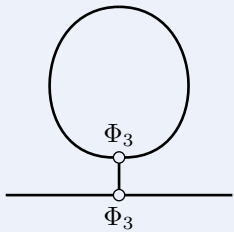
Atomic displacement
(descriptor)



[1] T. Tadano, Y. Gohda, & S. Tsuneyuki. (2014). *Journal of Physics: Condensed Matter*, 26(22), 225402

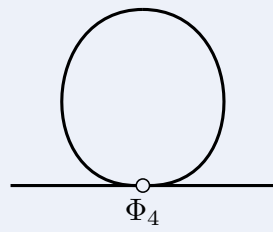
Self-energy of phonon

Tadpole Σ^T

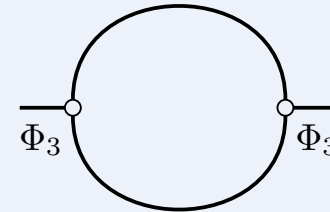


ω -independent, real

Loop Σ^L



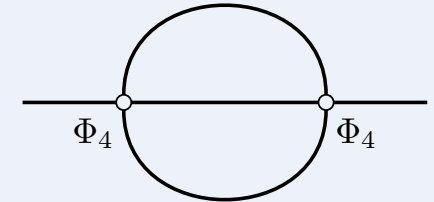
Bubble $\Sigma^B(\omega)$



2/3phonon-scattering

ω -dependent, complex

4ph $\Sigma^{4ph}(\omega)$



4phonon-scattering

Calculate phonon frequencies self-consistently, not perturbatively

	ω -dependent	Real / Complex	Real	Imaginary
Tadpole	×	Real	SCPH	-
Loop	×	Real	SCPH	-
Bubble	○	Complex	SCPH+B	One-shot
4ph	○	Complex	One-shot	One-shot

Self-Consistent-PHonon (SCPH)

Calculate renormalized phonon frequencies including Tadpole&Loop

SCPH equation [1] : Tadpole&Loop

- Self consistent equation for ω and G^S

$$[G_q^S(\omega)]^{-1} = [G_q^0(\omega)]^{-1} - \Sigma^T[G^S] - \Sigma^L[G^S]$$

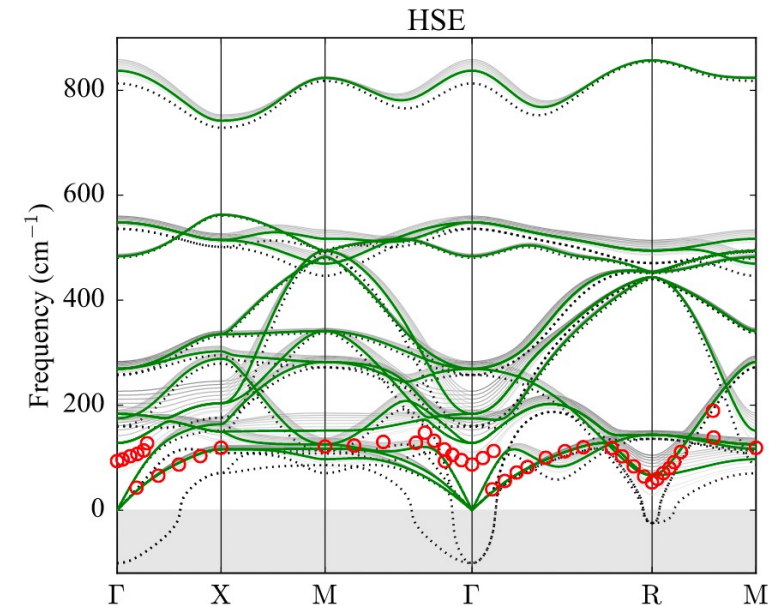
$$G_{\mathbf{q},jj'}^S(\omega) = \frac{2\omega_{qj}^S \delta_{jj'}}{(\omega_{qj}^S)^2 - \omega^2}$$

SCPH+bubble equation[2] : Bubble (real part)

- Self consistent equation for Ω . G^S is fixed during calculation

$$\Omega_q^2 = (\omega_q^S)^2 - 2\omega_q^S \text{Re} \Sigma_q^B[G^S, \Phi_3](\omega = \Omega_q)$$

$$\omega_{\text{harm}} \longrightarrow \omega^S \longrightarrow \Omega^{\text{SCPH+B}}$$



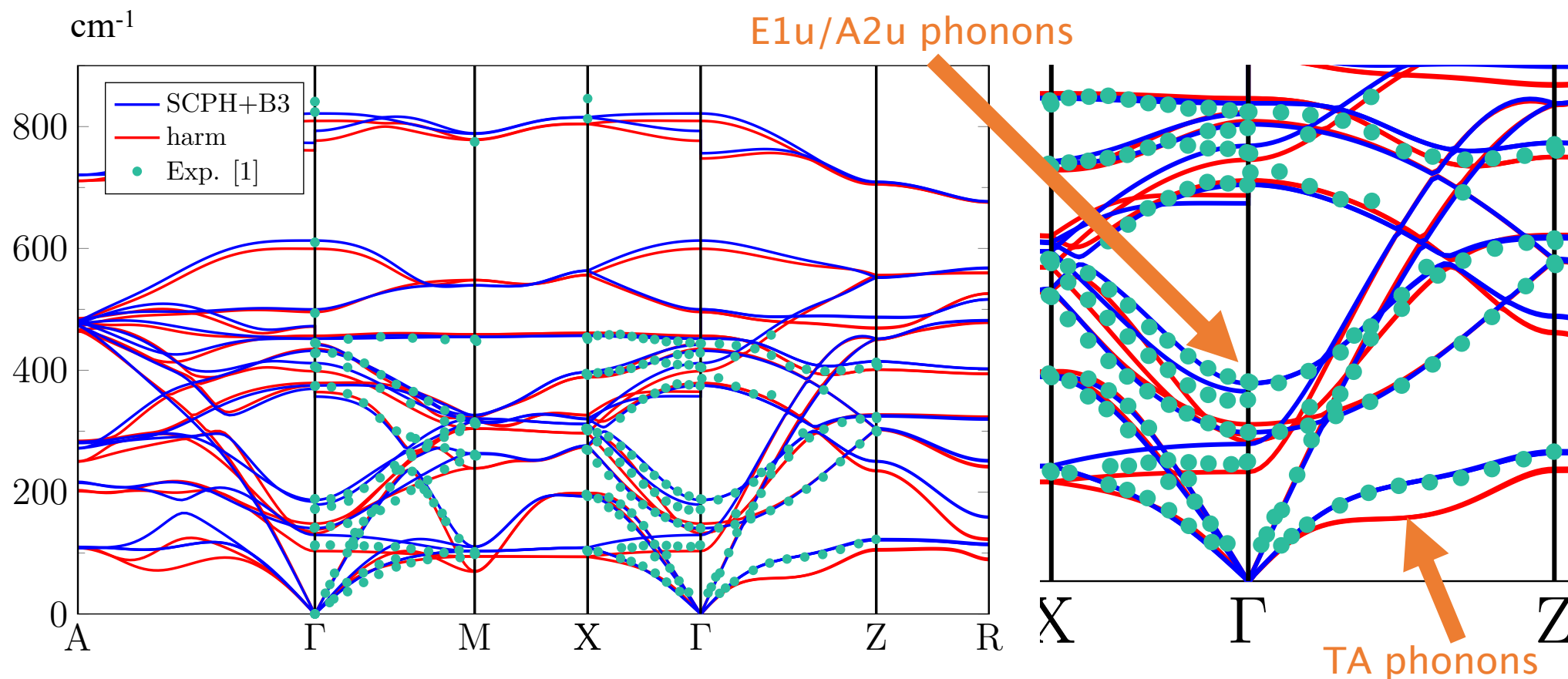
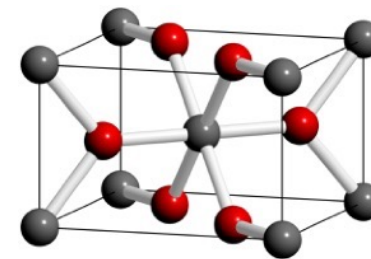
[1] T. Tadano, S. Tsuneyuki Phys. Rev. B .92.054301

[2] Tadano, Terumasa, and Wissam A. Saidi. *Physical Review Letters* 129.18 (2022): 185901.

[3] Tadano, Terumasa, and Shinji Tsuneyuki. *Journal of the Physical Society of Japan* 87.4 (2018): 041015.

Band : harmonic vs SCPH+B (r^2 SCAN)

SCPH+B well describes TA-phonons & Γ -point optical phonons.

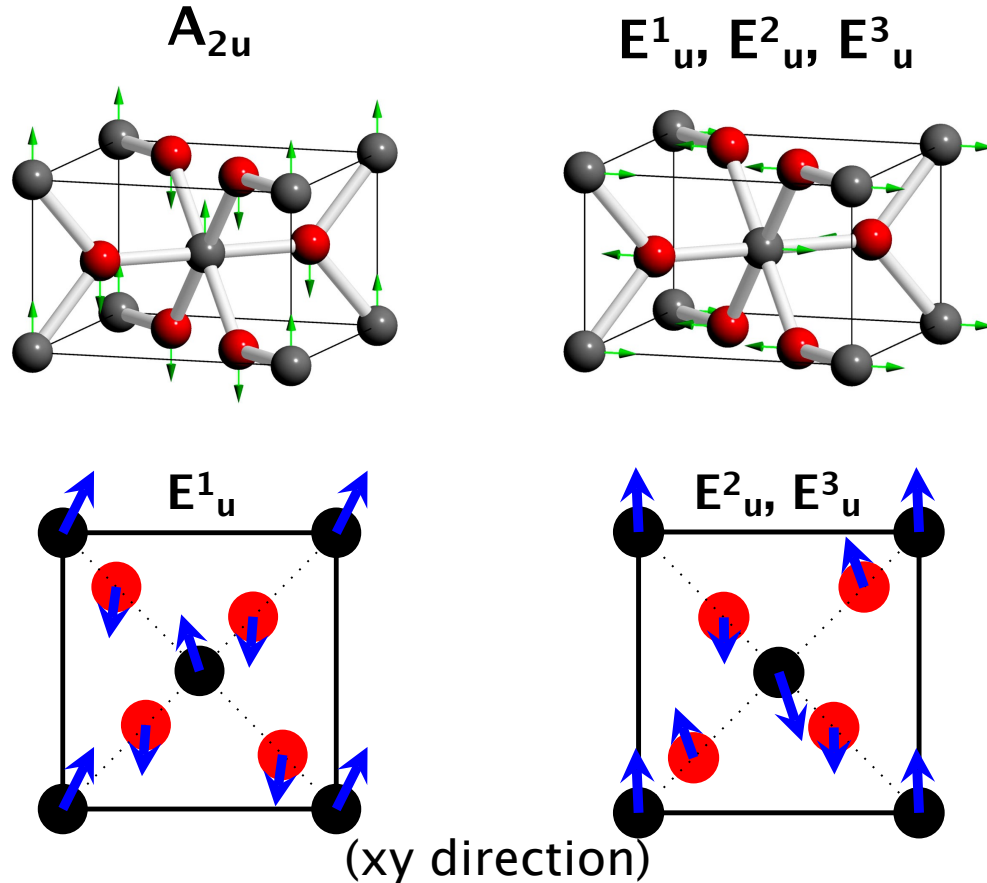


DFT	VASP
Supercell	2×2×4 supercell
Functional	r^2 SCAN
K-points	3×3×3
Phonon	ALAMODE
Temperature	300 [K]

[1]Traylor, Joseph Gibson, et al. *Physical Review B* 3.10 (1971): 3457.

IR active phonons : harmonic vs SCPH+B

SCPH+B well describes strongly anharmonic A_{2u} & E_{1u}



IR active phonon frequency

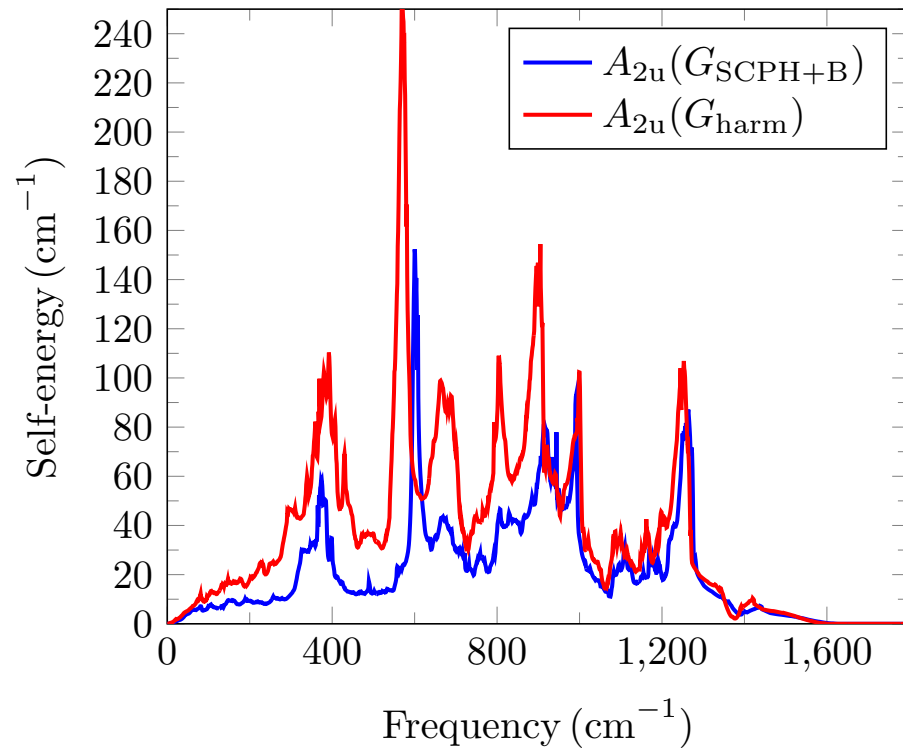
	R2SCAN		neutron [1]
	harm	SCPH+B3	
A_{2u}	138.5	179.4	172.6
E_{1u}^1	132.7	186.9	189
E_{2u}^2	378.8	374.9	374
E_{3u}^3	495.7	500.1	494

[1]Traylor, Joseph Gibson, et al. *Physical Review B* 3.10 (1971): 3457.

Linewidth (non-SC vs SCPH+B)

Perturbative calculation Overestimates linewidth

Imaginary part of Bubble self-energy



ω -independent linewidth

$$\gamma^{\text{SCPH+B}} = 2 \text{Im} \Gamma[G_{\text{SCPH+B}}](\omega = \omega_{\text{SCPH+B}})$$

$$\gamma^{\text{non-SC}} = 2 \text{Im} \Gamma[G_{\text{harm}}](\omega = \omega_{\text{harm}})$$

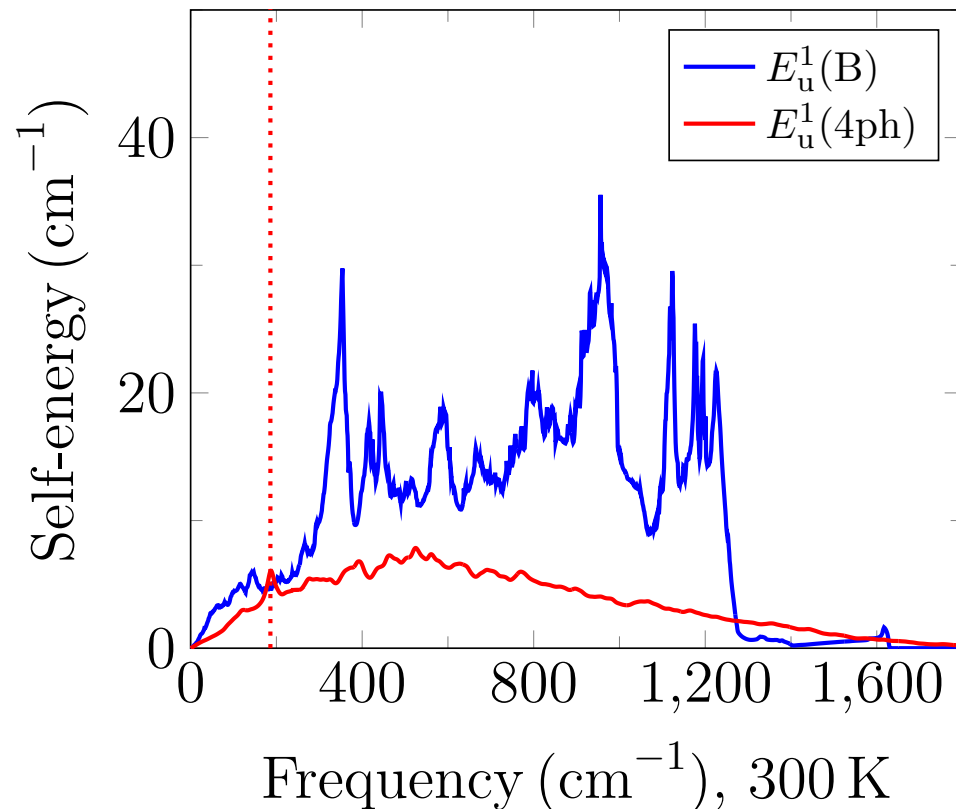
	non-SC			SCPH+B			FPSQ [1]
	bubble	4ph	total	bubble	4ph	total	
E_u^1	21	70	91	8.5	8.0	16	14.7
E_u^2	41	7.1	48	15	1.8	16	19.3
E_u^3	28	12	40	13	3.5	16	22.4
A_{2u}	29	52	81	16	9.6	26	20.0

[1] Schöche, S., et al. *Journal of Applied Physics* 113.16 (2013): 164102.

Linewidth (bubble vs 4ph)

4ph is as large as bubble linewidth

Self energy (E1u phonon)



ω -independent linewidth

$$\gamma^{\text{SCPH+B}} = 2 \text{Im} \Gamma[G_{\text{SCPH+B}}](\omega = \omega_{\text{SCPH+B}})$$

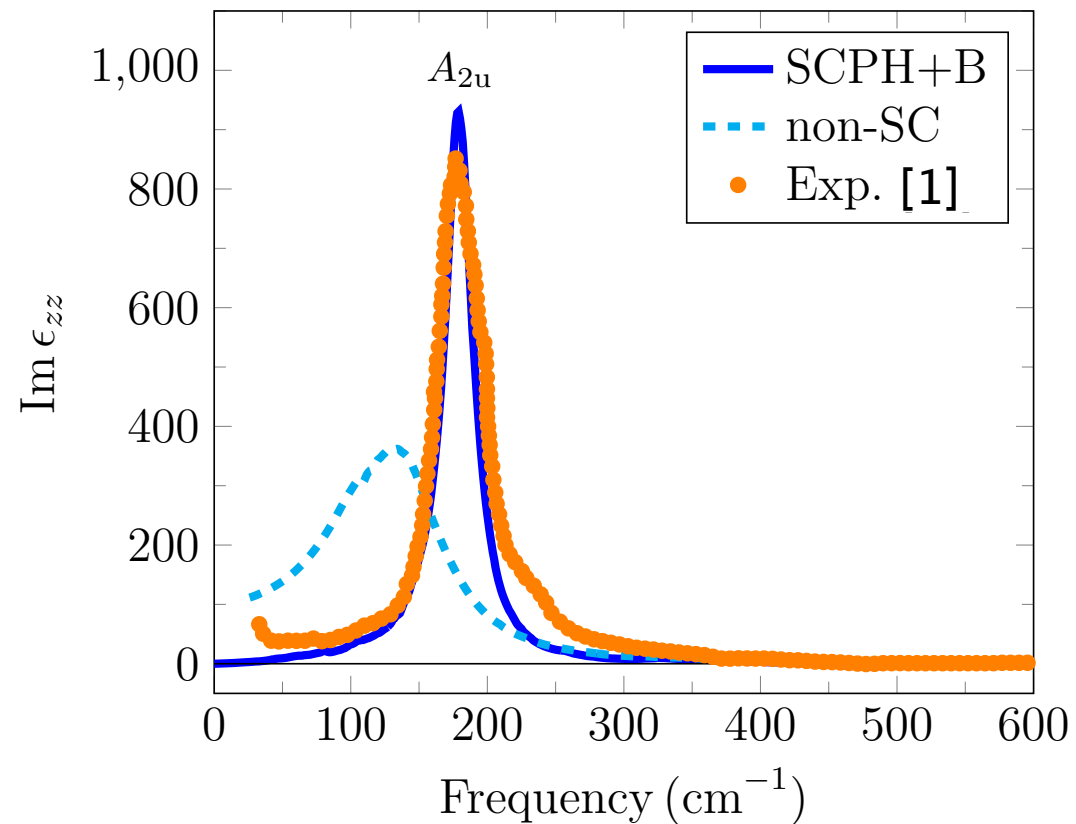
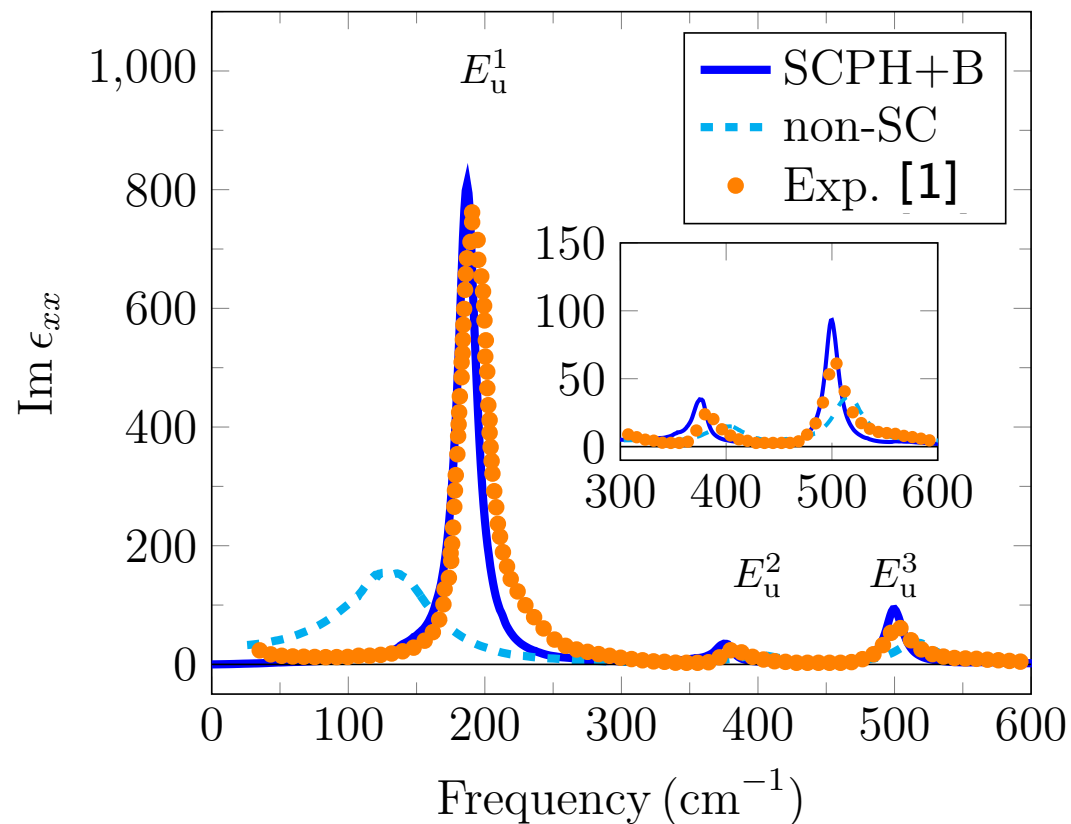
Calculated linewidth [cm-1]@300K

	SCPH+B3			Exp. [1]
	bubble	4 ph	total	
E1u	8.54	7.95	16.4	14.7 ± 8.0
E2u	14.5	1.81	16.2	19.3 ± 2.8
E3u	12.6	3.47	16.0	22.4 ± 3.1
A2u	16.3	9.57	25.7	20.0 ± 10.2

[1] Schöche, S., et al. *Journal of Applied Physics* 113.16 (2013): 164102.

Dielectric function

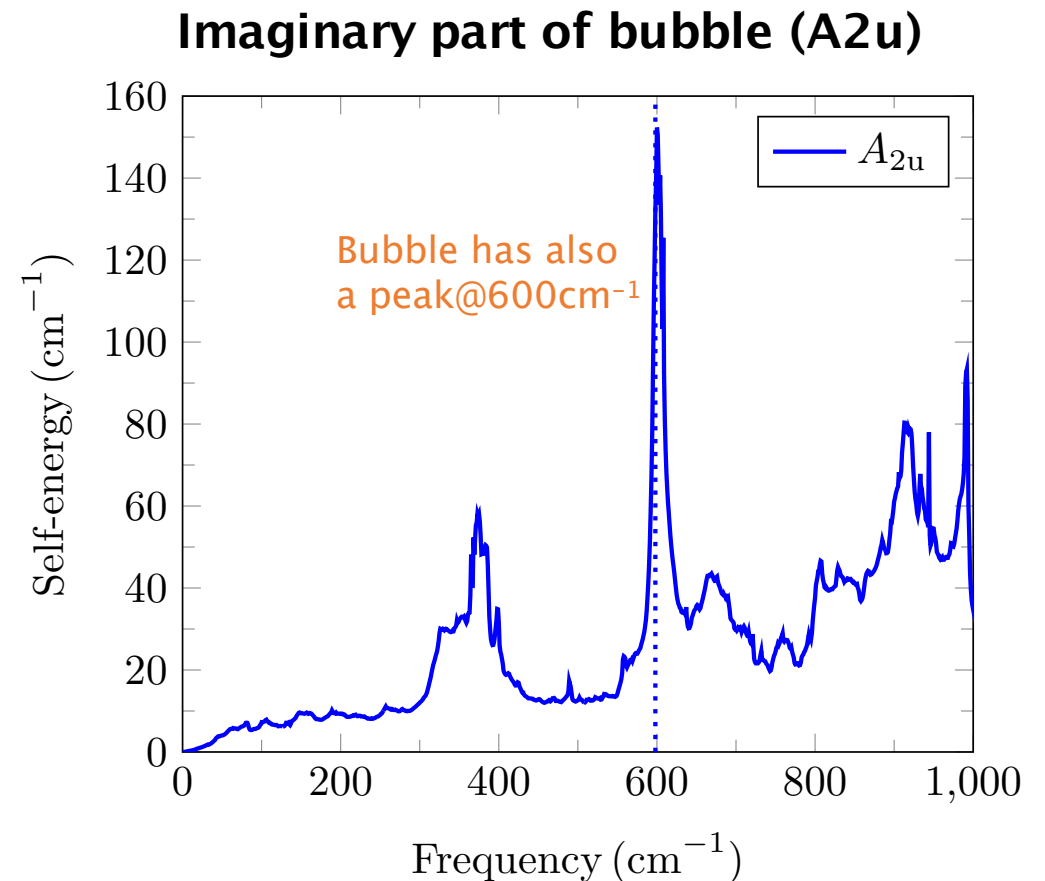
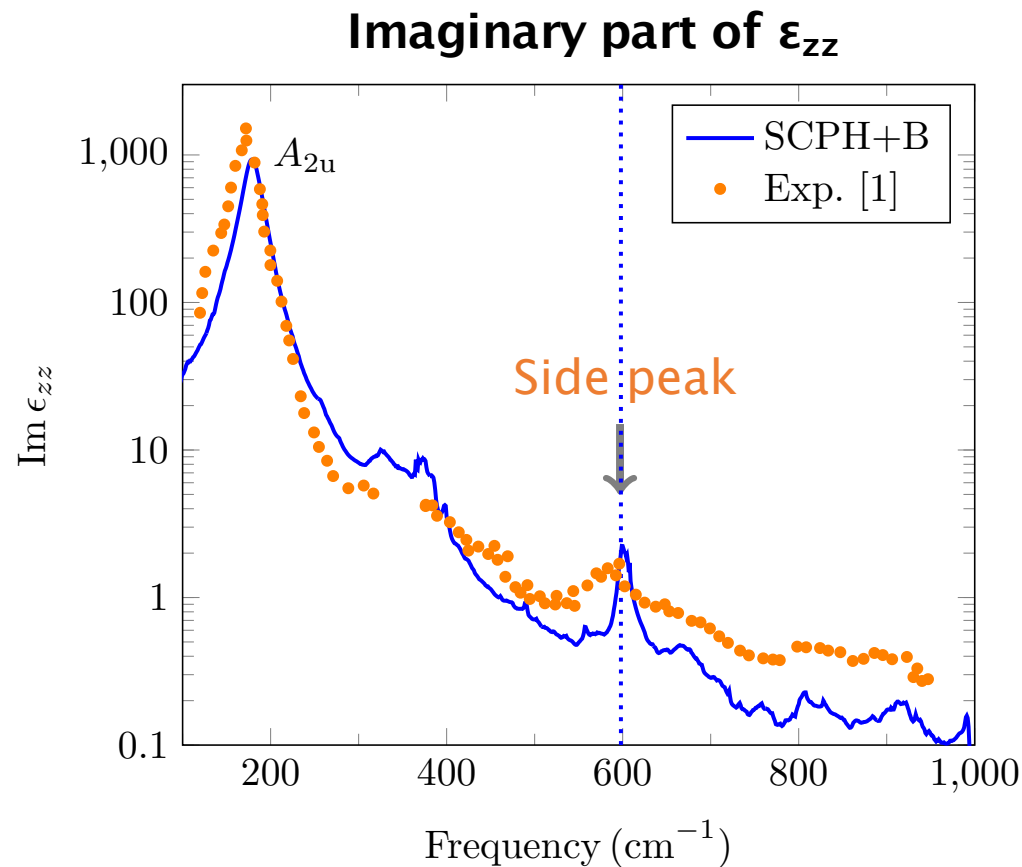
SCPH+B very well agree with experimental data



[1] Kanehara, Kazuki, et al. *Journal of the Ceramic Society of Japan* 123.1437 (2015): 303-306

Dielectric function : Side peak@600cm⁻¹

Bubble-selfenergy describes the side peak



[1] Schöche, S., et al. *Journal of Applied Physics* 113.16 (2013): 164102.

Conclusion

SCPH+B & frequency-dependent self-energy is important for dielectric property

- **Self-Consistent Phonon approach (SCPH+B) for strongly anharmonic material**
 - ※ SCPH+B well describe phonon frequencies of rutile TiO_2
 - ※ Phonon frequencies renormalization by SCPH+B is important for phonon linewidths
 - ※ 4ph linewidth is as large as bubble linewidth @ 300K
- **Cowley's theory to use frequency-dependent self-energy**
 - ※ Bubble self-energy can describe side-peak of dielectric function
 - ※ 4ph self-energy do not so frequency-dependent