

Quantum Black Holes as Holograms

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ICREA + UBarcelona

Extreme Universe Online Colloquium

20 December 2024



Early ideas

Takahiro Tanaka

gr-qc/0203082

RE+Alessandro Fabbri+Nemanja Kaloper

hep-th/0206155

Recent work with

Antonia Frassino

Raimon Luna

Juan Pedraza

Ryotaku Suzuki

Andy Svesko

Marija Tomašević

Manus Visser

Benson Way

arXiv

2007.15999

2207.03302

2301.02587

Half a century of Black Hole radiance

1974

The beginning of the era of the Quantum Era of Black Holes

Half a century of Black Hole radiance

Gravity should be the classical limit of a quantum theory

Black holes should be the classical limit of a quantum object

What is a Quantum Black Hole?

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Hawking did not have a quantum theory of gravity

He studied quantum fields in a fixed classical background

The black hole emits radiation

What is a Quantum Black Hole?

Hawking did not have a quantum theory of gravity

He studied quantum fields in a fixed classical background

The black hole emits radiation

Radiation comes out of the black hole

but the *black hole remains unaffected and classical*

What is a Quantum Black Hole?

Black hole remains unaffected and classical

- Good enough if radiation emission has a small effect on the black hole
- Fails if effects are large - because they accumulate over time, or because single quantum emission is significant

Backreaction problem

What is a Quantum Black Hole?

Hawking did not have a quantum theory of gravity

Nowadays, *we do* have good quantum theories of gravity
where black holes appear in the semiclassical limit

AdS/CFT

Then, what is a Quantum Black Hole?

AdS/CFT:

Black hole = a large-N matrix, with fastly scrambling entries

But too poorly understood to deal with backreaction
problem

Then, what is a Quantum Black Hole?

Simpler goal

Classical geometry of black hole modified (possibly a lot)
by quantum effects

Quantum-backreacted black hole

Much insight gained this way

Quantum backreaction

$$G_{\mu\nu}(g_{\alpha\beta}) = 8\pi G \langle T_{\mu\nu}(g_{\alpha\beta}) \rangle$$

classical Einstein tensor & metric

quantum matter renorm stress tensor
(many fields)

Coupled system: metric + \langle QFT \rangle

Very hard to solve simultaneously

Perturbative backreaction: limited insight

Quantum backreaction

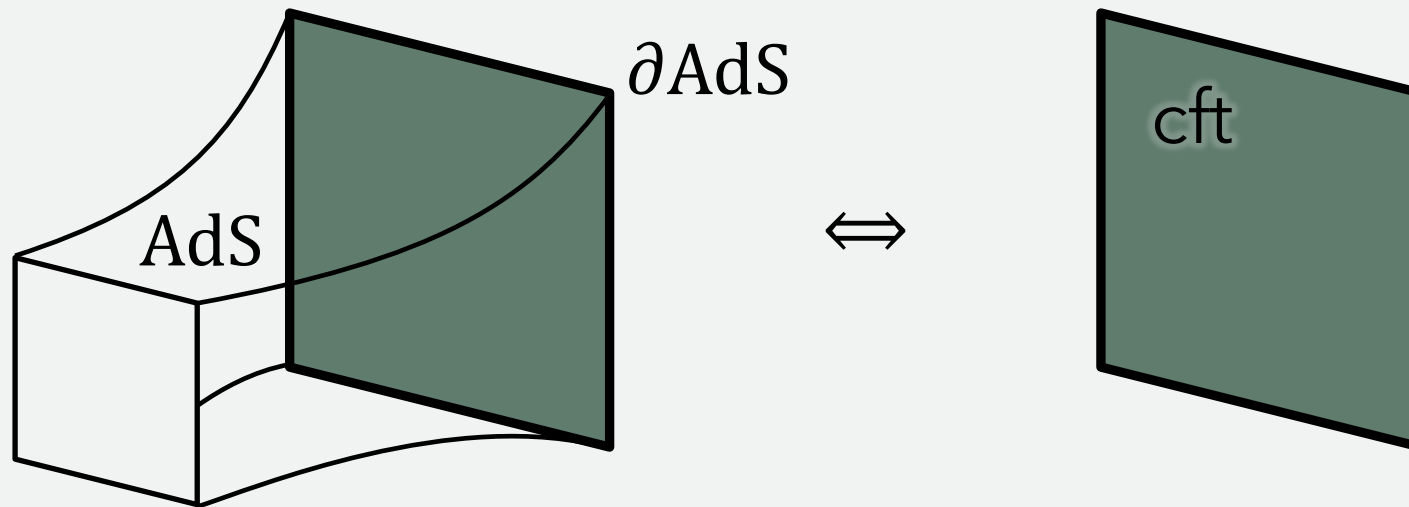
Exact backreaction:

2D models: CGHS/RST, JT+CFT

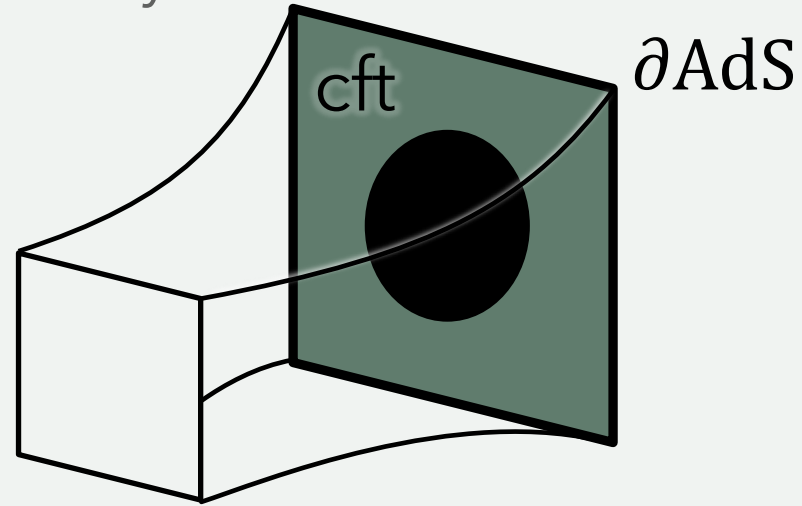
Holographic reformulation

Solving QFT with Holography

- AdS/CFT maps quantum field dynamics into classical gravitational bulk dynamics
- The AdS boundary geometry, where the CFT lives, is often flat, or spherical



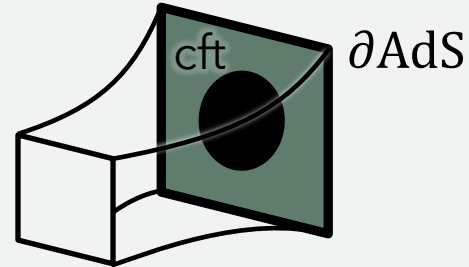
- But it can be a black hole geometry



- In this case, holographic CFT lives on black hole background (fixed)
- CFT state in the presence of black hole is dual to some classical bulk
- to be found

CFT on black hole background – holographically

solve bulk \equiv solve cft in black hole



- Quantum field on a black hole background \Rightarrow Hawking radiation
- Quantum Hawking radiation of CFT is *dual to some classical bulk dynamics* – not trivial, but possibly easier

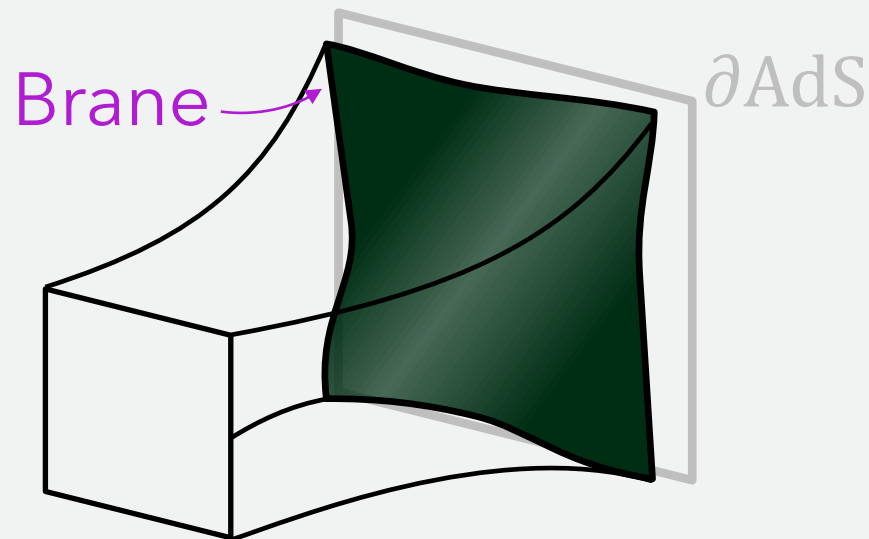
review: Holographic thermal field theory on curved spacetimes

Marolf+Rangamani+Wiseman 2013

CFT on *dynamical* geometry – holographically

Make *gravity on boundary dynamical*: *Braneworld holography*

- Gravitating brane is *wall boundary* that can fluctuate
- Gravity localized near AdS brane (Randall-Sundrum)



CFT on *dynamical* geometry – holographically

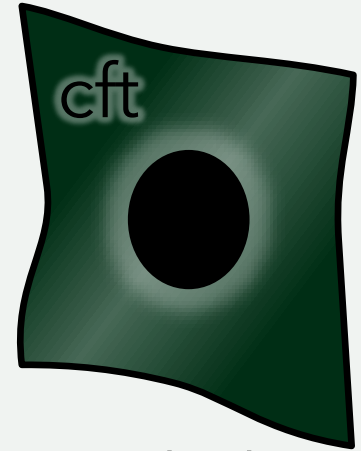
- Holographically: bulk gravity \equiv CFT
- Dynamical brane coupled to bulk gravity \equiv
Dynamical (boundary) geometry coupled to CFT – w/ backreaction

$$R_{ij} - \frac{1}{2}Rg_{ij} + \dots = 8\pi G_N \langle T_{ij} \rangle_{CFT}$$



CFT on *dynamical* black hole – holographically

$$R_{ij} - \frac{1}{2}Rg_{ij} + \dots = 8\pi G_N \langle T_{ij} \rangle_{CFT}$$



- Black hole on brane = Black hole + Hawking radiation, coupled

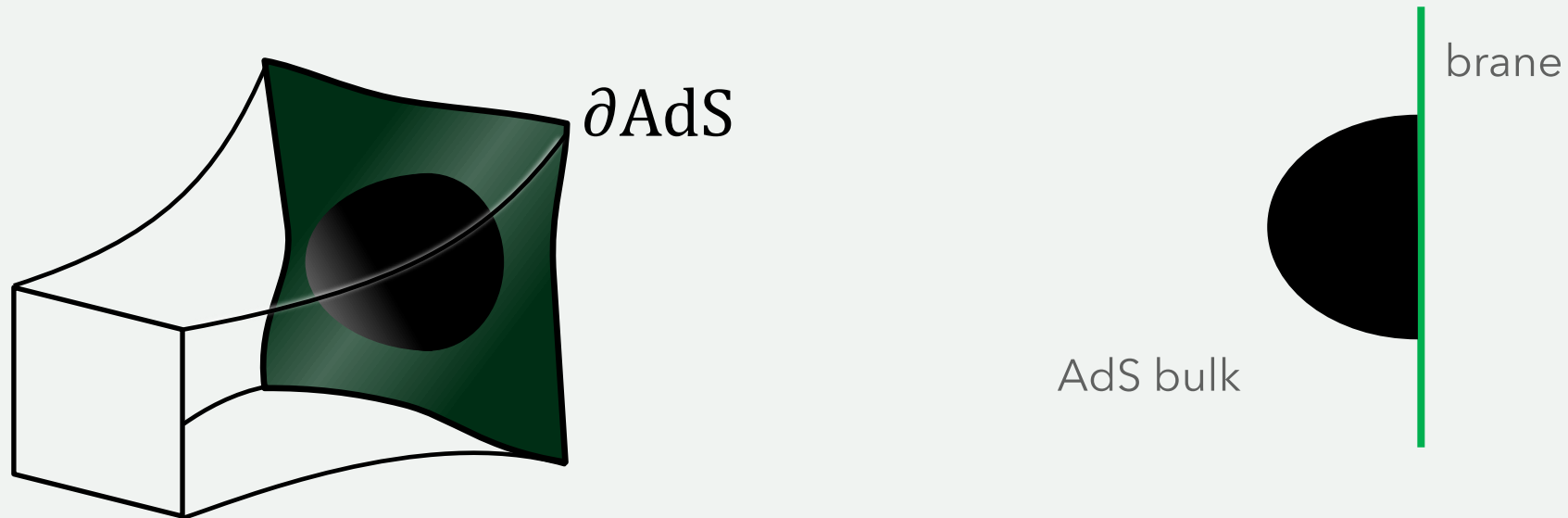
- Dynamics of black hole evaporation = Classical dynamics in bulk

Tanaka 2002

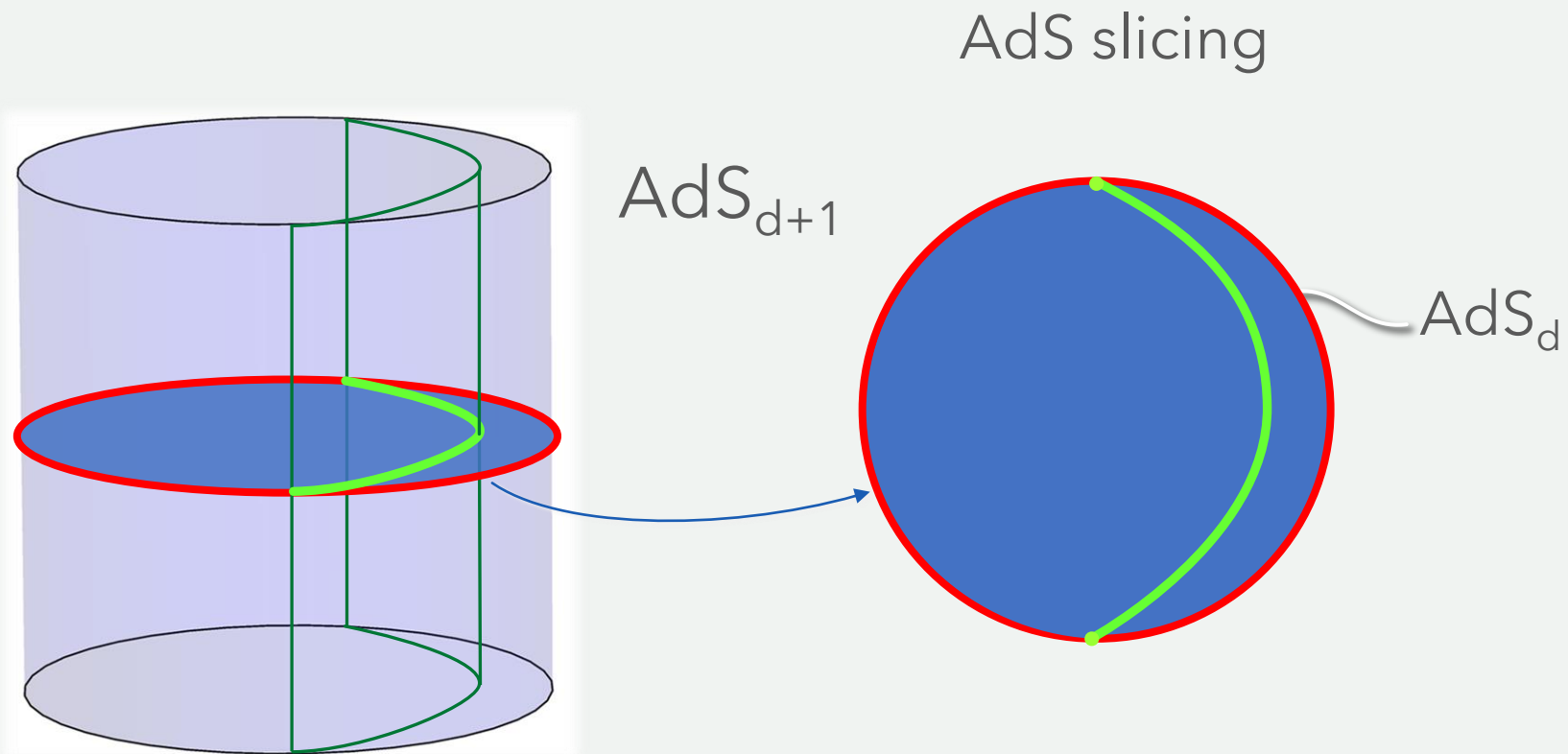
RE+Fabri+Kaloper 2002

Black droplets on a brane

- Black hole coupled to $\text{CFT}_d =$ black droplet on a brane in AdS_{d+1}

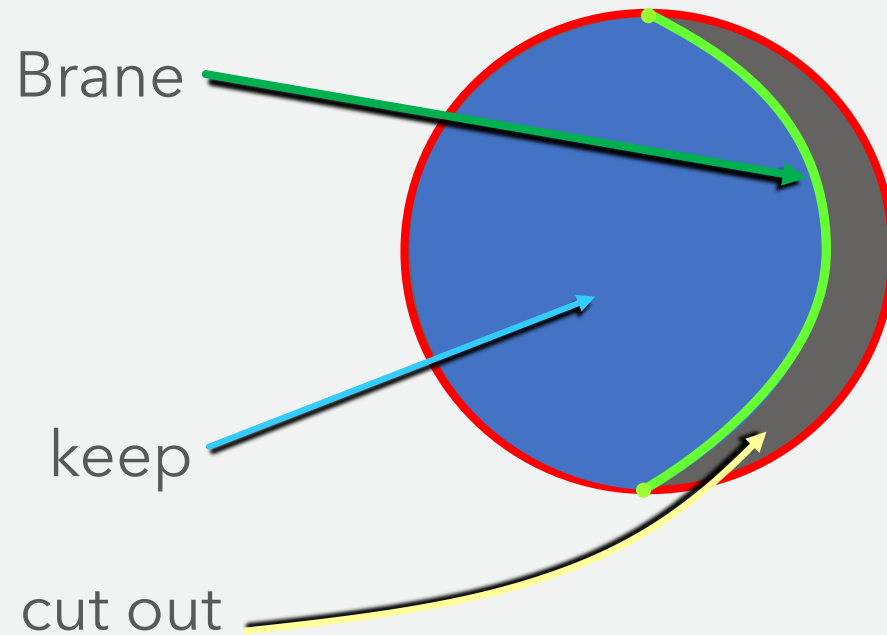


Slicing AdS



Braneworlds

Karch-Randall: AdS branes

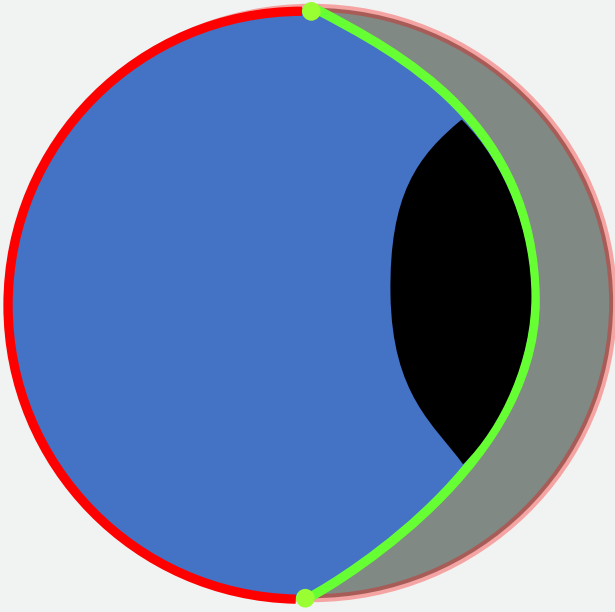


Randall+Sundrum 1999

Karch+Randall 2000

& paste to copy: 2-sided brane, \mathbb{Z}_2 orbifold

Black hole on a brane



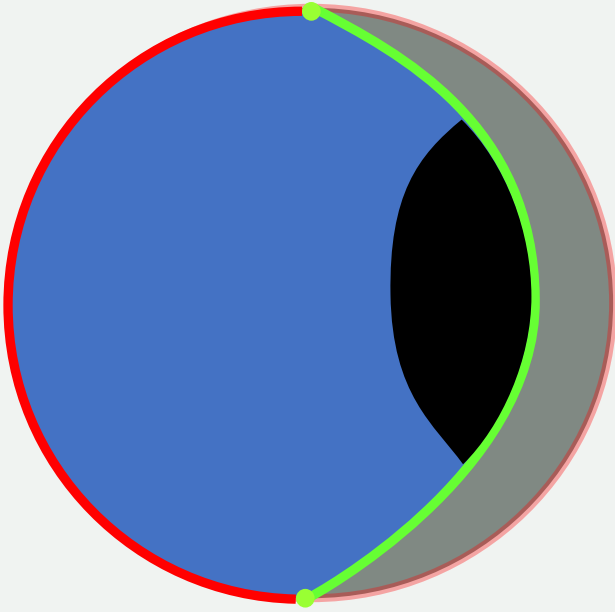
- Exact solutions in AdS_4
 - Numerical solutions in AdS_5 (hard)
- } time-independent

Figueras+Lucietti+Wiseman

- Analytical/easy numerics in large-D limit $\text{AdS}_{D \gg 1}$

& with time evolution

Black hole on a brane



- **Exact solutions in AdS_4**
 - Numerical solutions in AdS_5 (hard)
- Figueras+Lucietti+Wiseman
- Analytical/easy numerics **in large-D limit** $\text{AdS}_{D \gg 1}$
& with time evolution
- } time-independent

Exact Holographic
3D Quantum Black Holes

Black hole on the brane from AdS C-metric

Plebański + Demiański 1976

$$ds^2 = \frac{\ell^2}{(\ell + xr)^2} \left(-H(r)dt^2 + \frac{dr^2}{H(r)} + r^2 \left(\frac{dx^2}{G(x)} + G(x)d\phi^2 \right) \right)$$

$$H(r) = \frac{r^2}{\ell_3^2} + \kappa - \frac{\mu\ell}{r}$$

$$G(x) = 1 - \kappa x^2 - \mu x^3$$

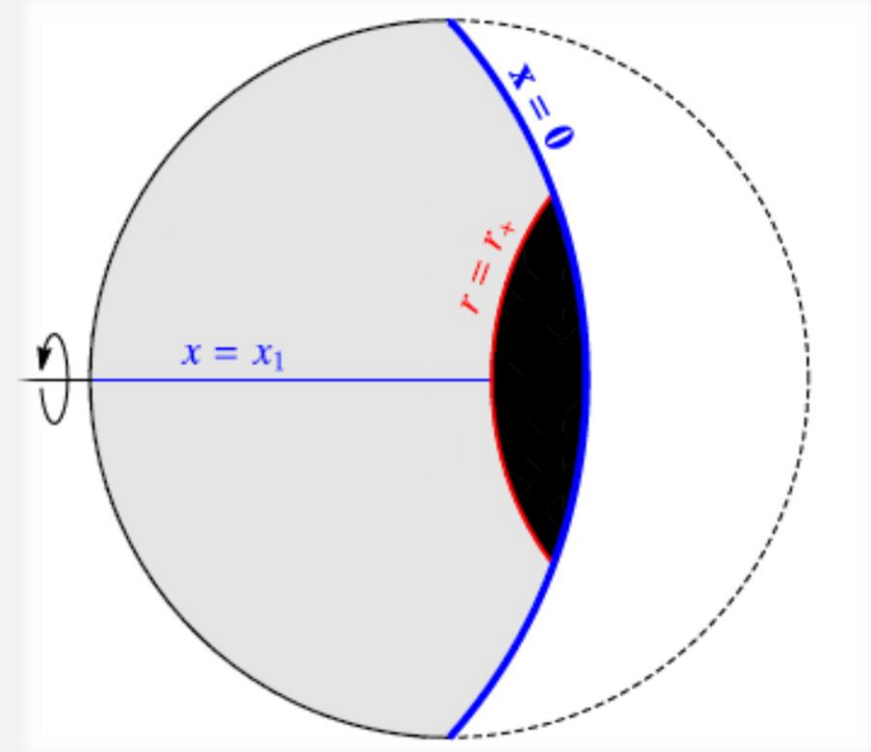
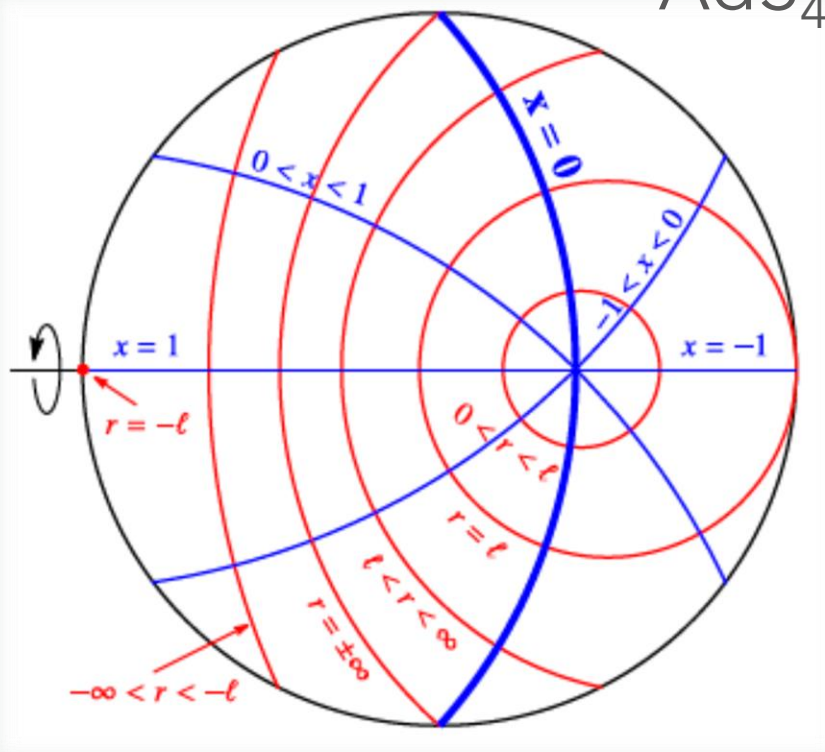
RE+Horowitz+Myers 1999

Adapted coordinates

$$ds^2 = \frac{\ell^2}{(\ell + xr)^2} \left(-H(r)dt^2 + \frac{dr^2}{H(r)} + r^2 \left(\frac{dx^2}{G(x)} + G(x)d\phi^2 \right) \right)$$

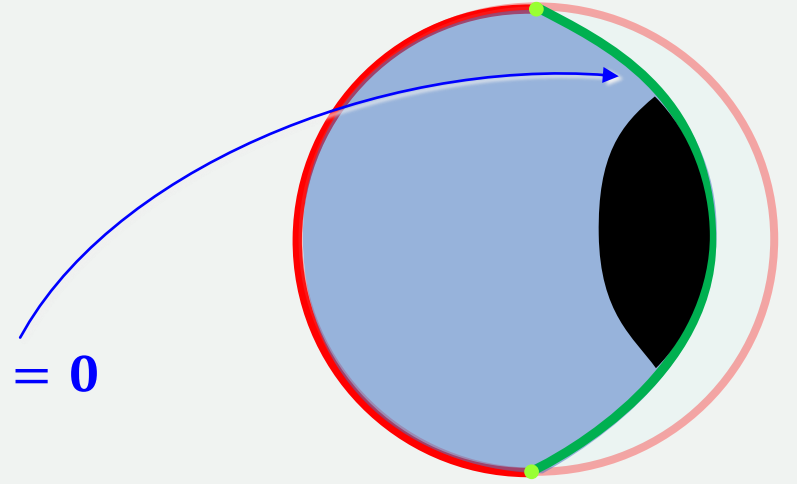
brane at $x = 0$

AdS₄



Quantum BTZ metric

3D metric induced on brane at $x = 0$



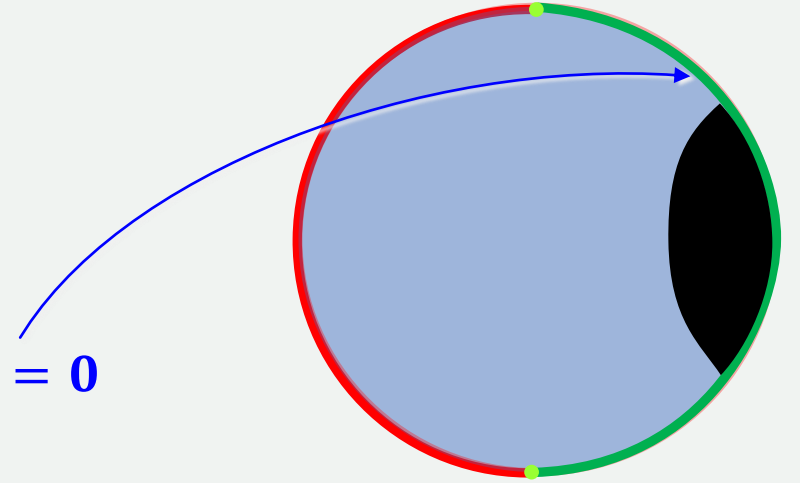
$$ds^2 \Big|_b = - \left(\frac{r^2}{\ell_3^2} - 8G_3 M - \ell \frac{F(M)}{r} \right) dt^2 + \frac{dr^2}{\frac{r^2}{\ell_3^2} - 8G_3 M - \ell \frac{F(M)}{r}} + r^2 d\phi^2$$

RE+Fabbri+Kaloper 2002

RE+Frassino+Way 2020

Quantum BTZ metric

3D metric induced on brane at $x = 0$

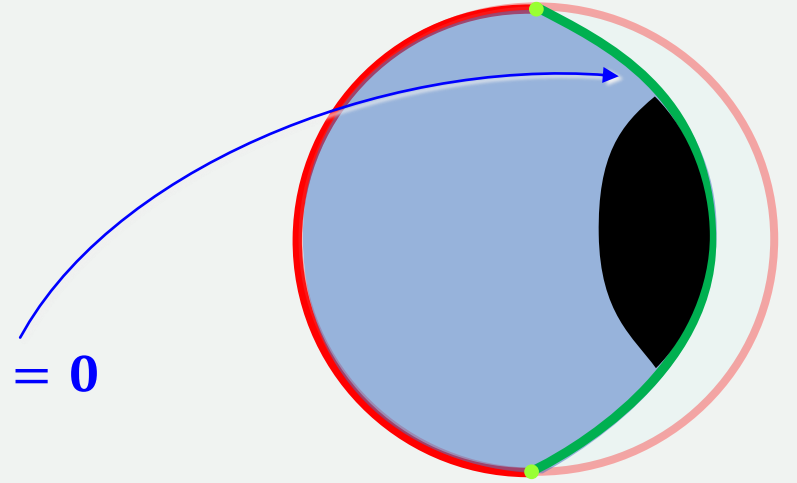


$$ds^2 \Big|_b = - \left(\frac{r^2}{\ell_3^2} - 8G_3 M - \ell \frac{F(M)}{r} \right) dt^2 + \frac{dr^2}{\frac{r^2}{\ell_3^2} - 8G_3 M - \ell \frac{F(M)}{r}} + r^2 d\phi^2$$

$\ell = 0$: BTZ black hole

Quantum BTZ metric

3D metric induced on brane at $\mathbf{x} = \mathbf{0}$



$$ds^2 \Big|_b = - \left(\frac{r^2}{\ell_3^2} - 8G_3 M - \ell \frac{F(M)}{r} \right) dt^2 + \frac{dr^2}{\frac{r^2}{\ell_3^2} - 8G_3 M - \ell \frac{F(M)}{r}} + r^2 d\phi^2$$

$\ell = 0$: BTZ black hole

$\ell > 0$: quantum-corrected BTZ

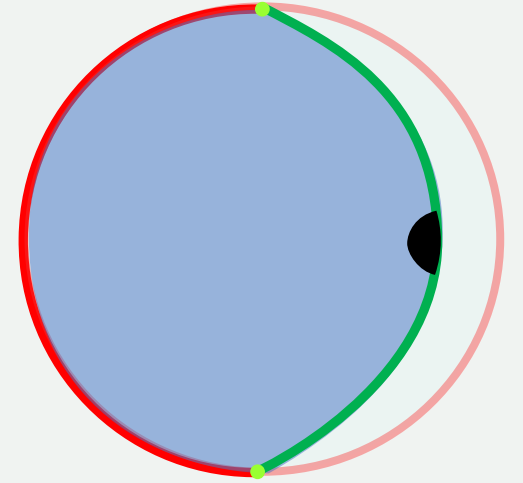
Small AdS₃ quantum black holes

$$ds^2 = - \left(\frac{r^2}{\ell_3^2} - 8G_3 M - \ell \frac{F(M)}{r} \right) dt^2 + \frac{dr^2}{\frac{r^2}{\ell_3^2} - 8G_3 M - \ell \frac{F(M)}{r}} + r^2 d\phi^2$$

$$-\frac{1}{8G_3} < M < 0 \quad \begin{array}{l} \ell = 0 : \text{conical defect} \\ \ell > 0 : \text{small black hole} = \text{dressed cone} \end{array}$$

Casimir energy on a cone dresses conical singularity

Quantum Cosmic Censorship



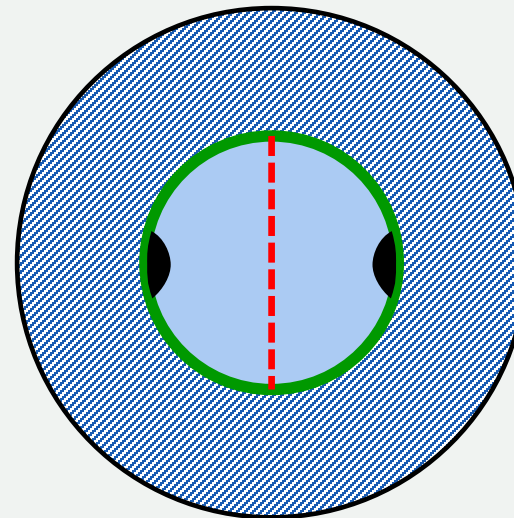
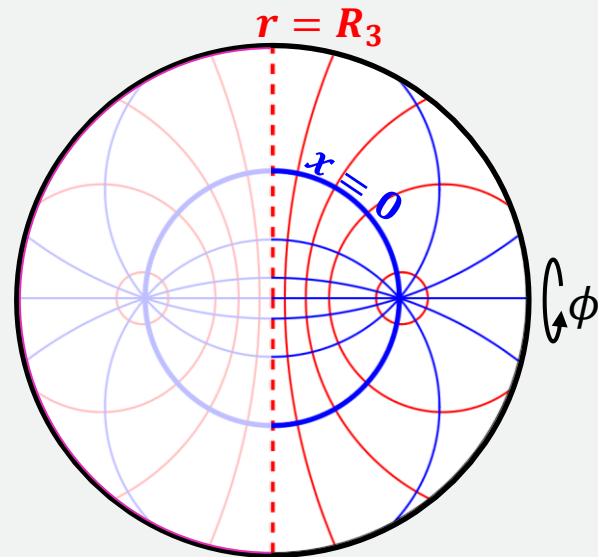
dS_3 quantum black holes

RE+Pedraza+Svesko+Tomašević+Visser 2022

$$ds^2 = - \left(1 - \frac{r^2}{\ell_3^2} - \ell \frac{\mu}{r} \right) dt^2 + \frac{dr^2}{1 - \frac{r^2}{\ell_3^2} - \ell \frac{\mu}{r}} + r^2 d\phi^2$$

$\ell = 0$: (conical defect in) dS_3

$\ell > 0$: black hole in dS_3



The take away (1): Holographic backreaction works

Quantum-backreacted black hole and CFT stress tensor
can be described *exactly* and *in detail*

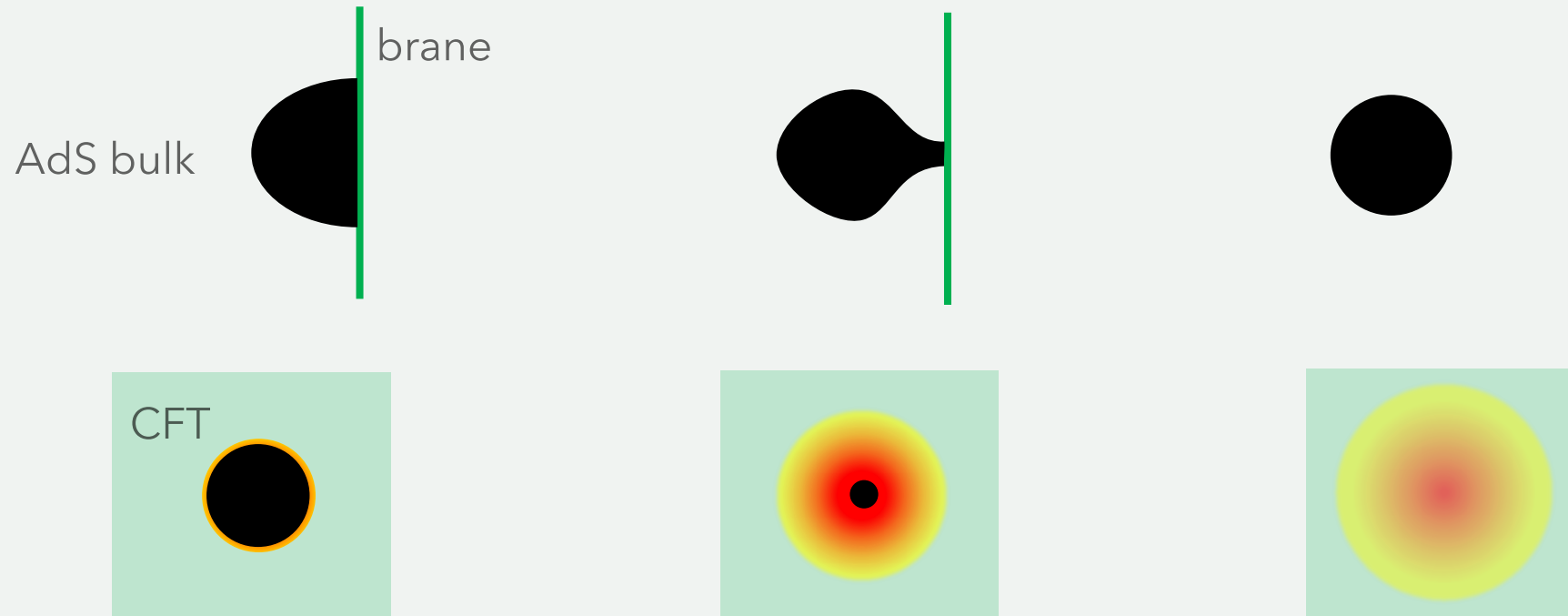
Efficient use of holography to solve a hard quantum problem

Holographic duals of evaporating black holes

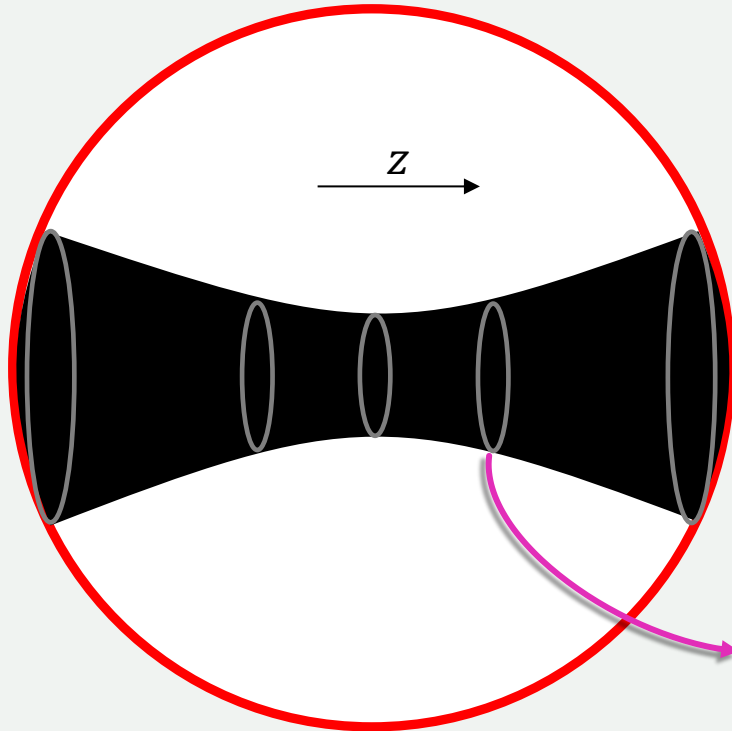
RE+Luna+Suzuki+Tomašević+Way 2023

Black hole coupled to CFT radiation - holographically

- Black hole = bulk horizon intersecting brane
- Thermal CFT radiation = horizon in bulk



Black string in AdS_D



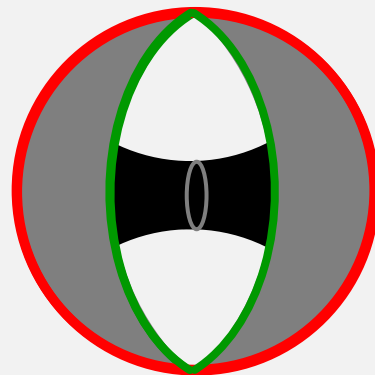
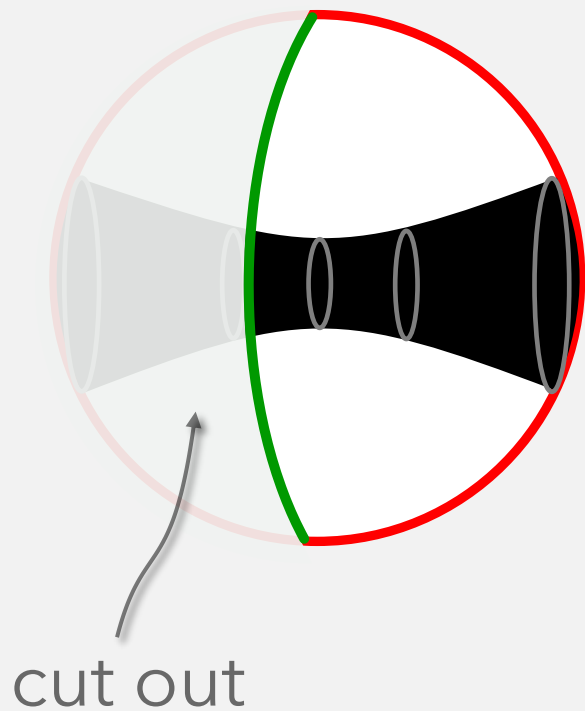
$$ds^2 = \frac{L^2}{\cos^2 z} \left(dz^2 + ds^2(\text{Schw-AdS}_{D-1}) \right)$$

Boundary geometry:

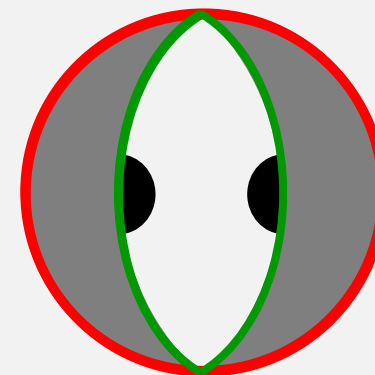
Sphere with **two black holes at antipodes**

Schw- AdS_{D-1}

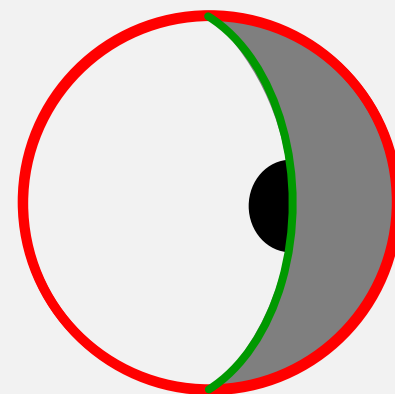
Braneworld



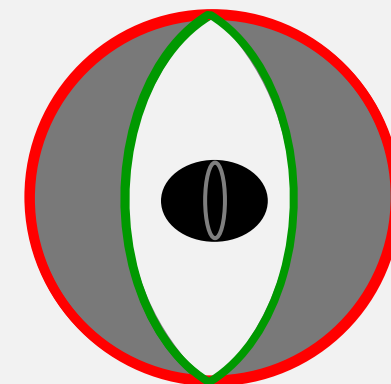
Black string/Funnel



Double droplet



Single droplet



Bulk black hole

Large D help

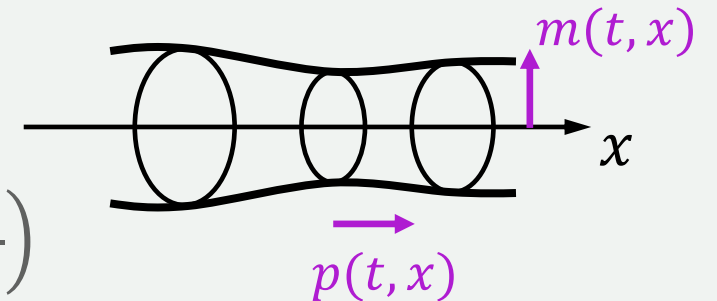
RE+Suzuki+Tanabe 2013, 2015

- $D \gg 1$: gravitational dynamics localized near horizon
- Radiation is suppressed $\sim e^{-D}$
- Effective theory for horizon dynamics

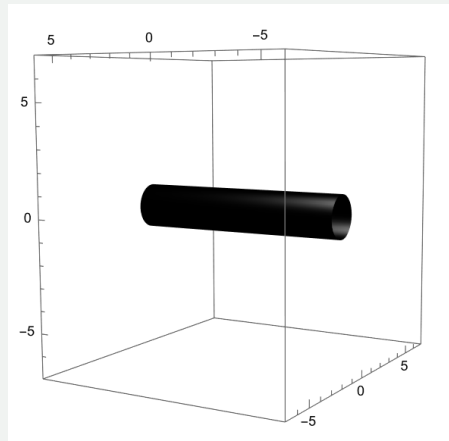
E.g., for AF black strings:

$$\partial_t m - \partial_x^2 m = -\partial_x p$$

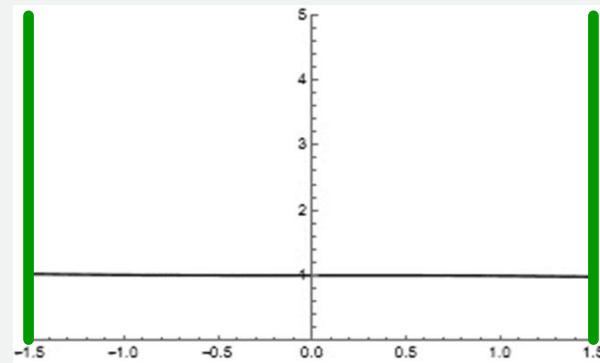
$$\partial_t p_i - \partial_x^2 p_i = \partial_x m - \partial_x \left(\frac{p^2}{m} \right)$$



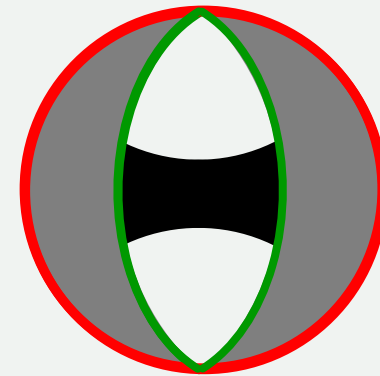
Some static phases



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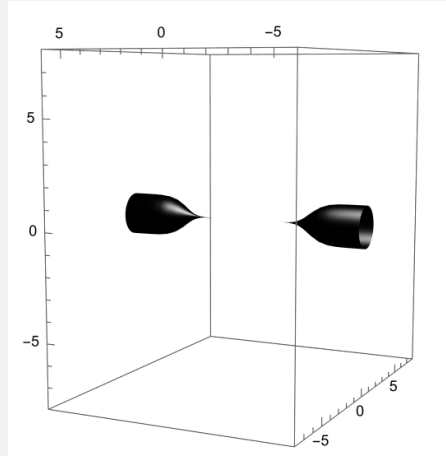


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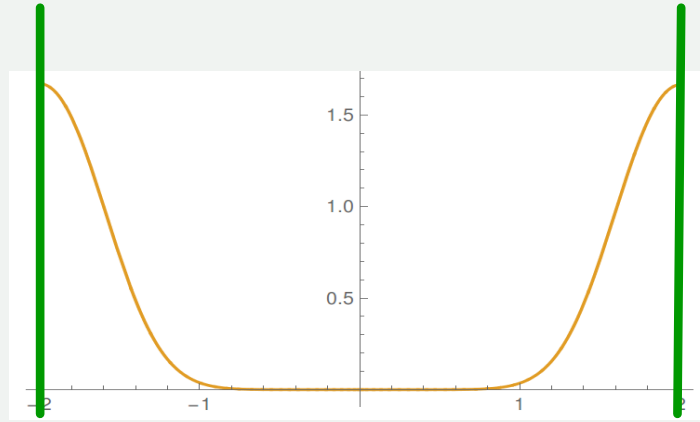


Black string/Uniform funnel

Black holes @ large D = gaussian blobs

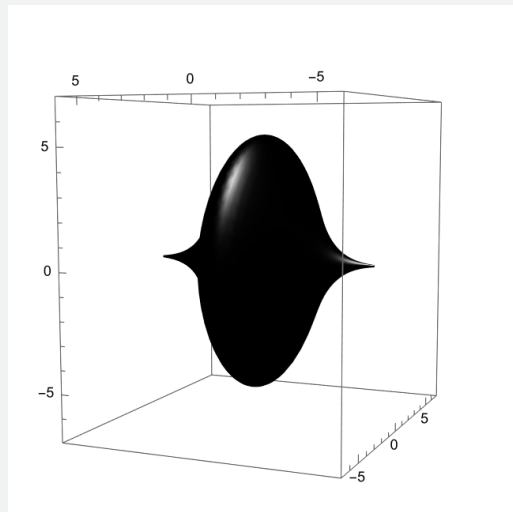
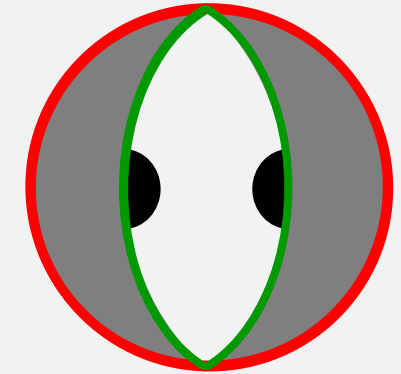


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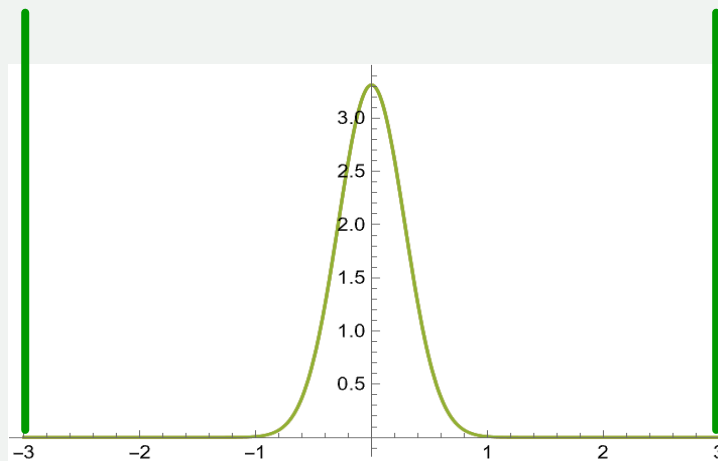


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Double droplet

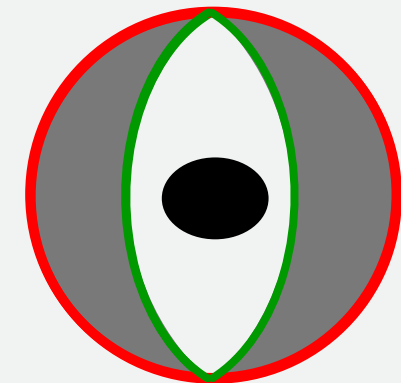


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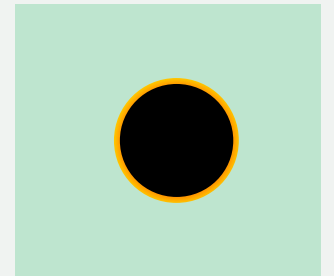
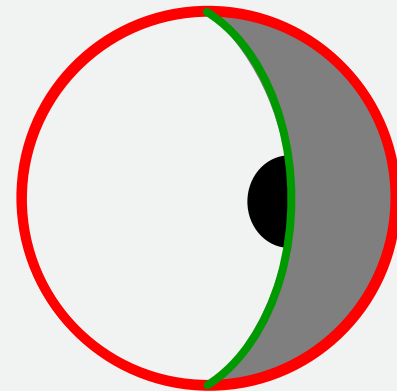
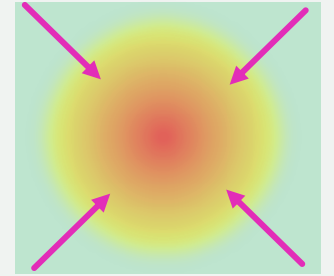
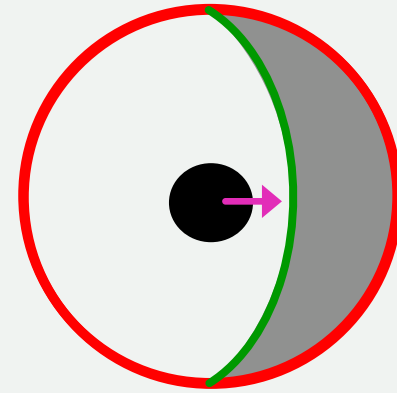
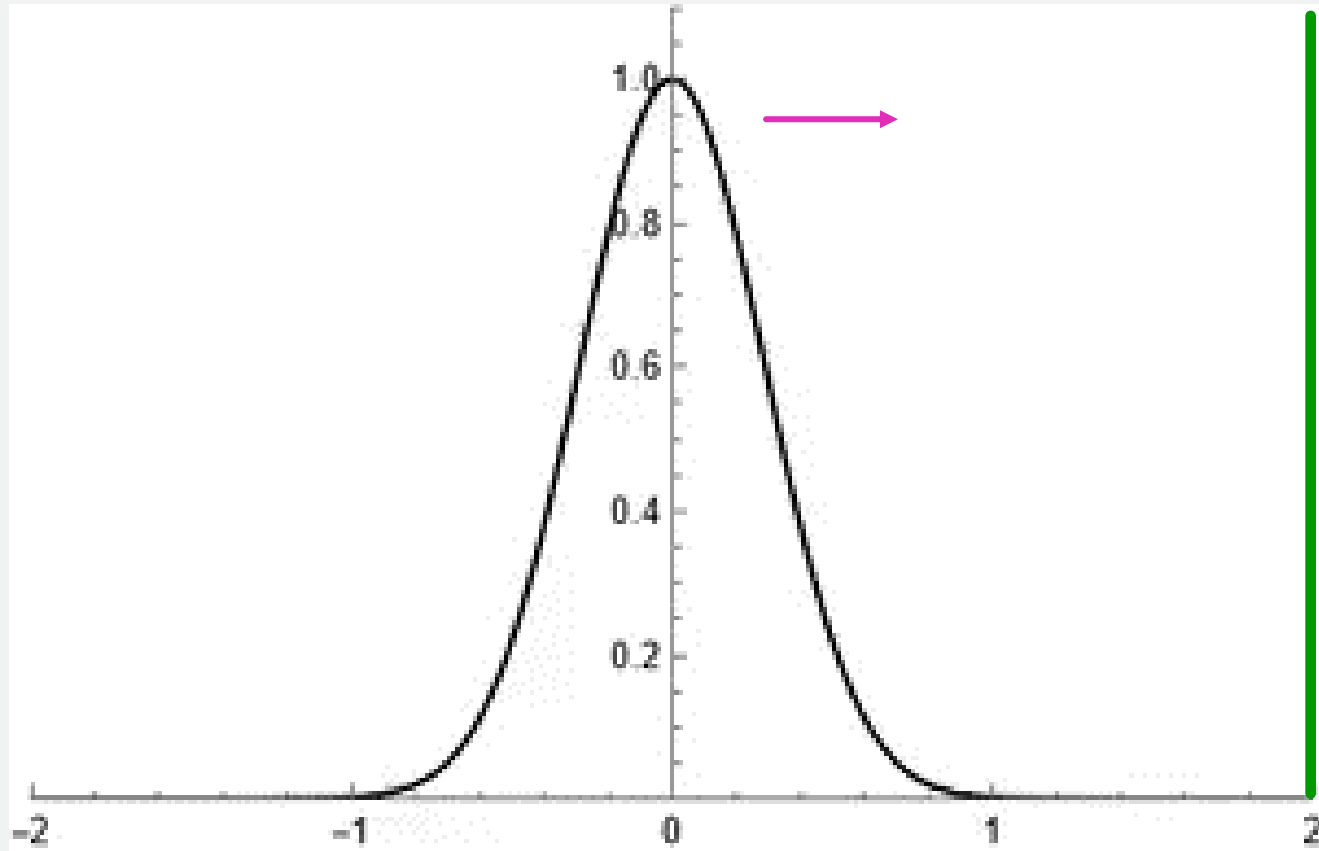
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Bulk black hole

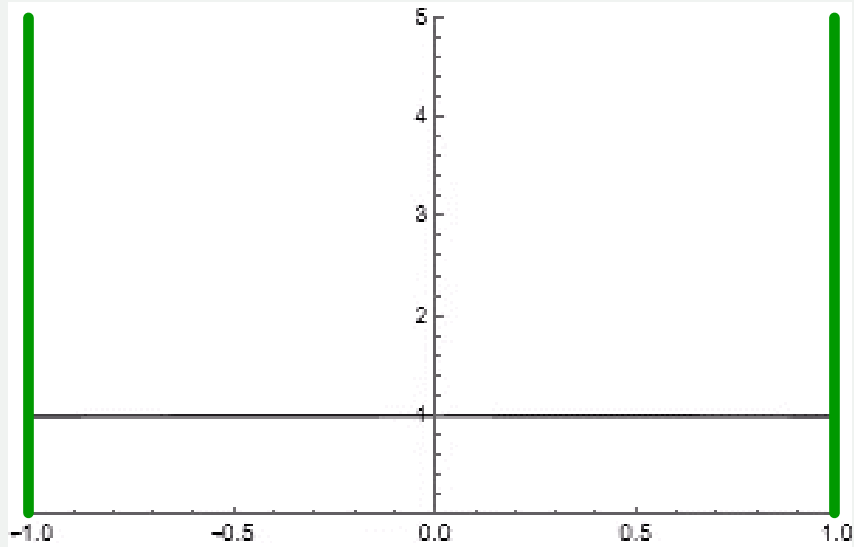


Dynamic play

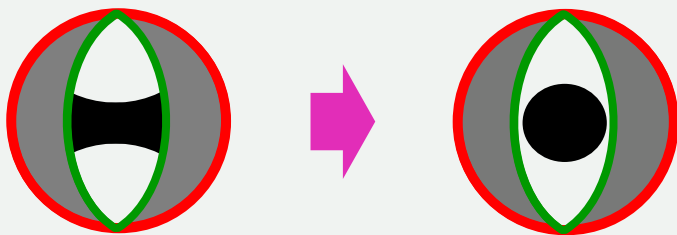
Collapse on the brane



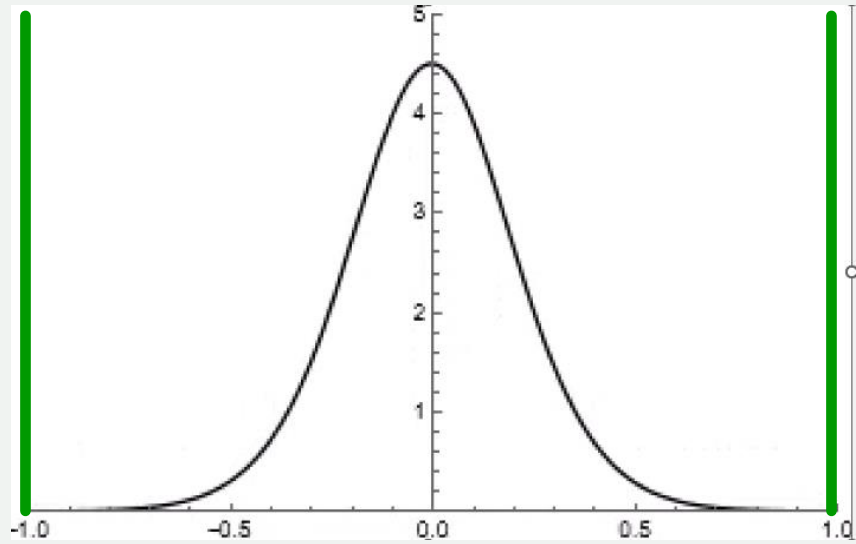
Evaporation



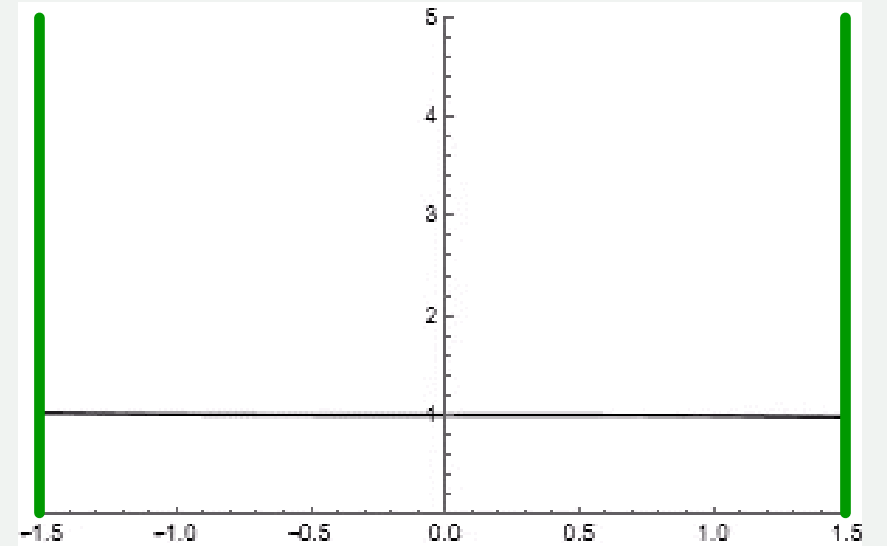
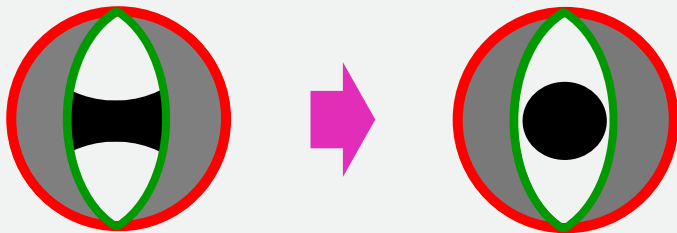
Evaporation as flow into bulk



Evaporation

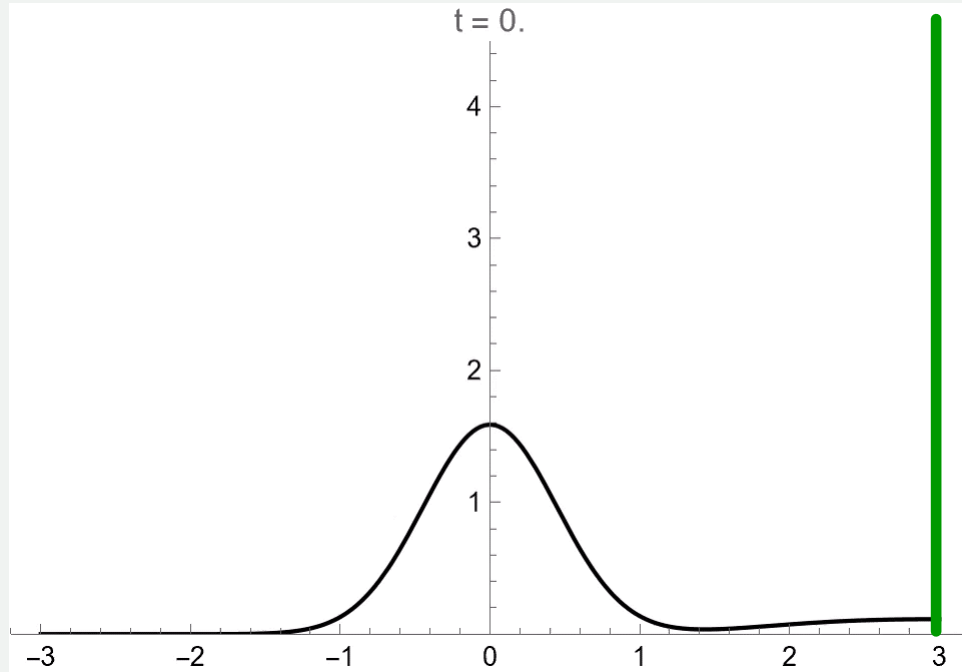


Evaporation as flow into bulk

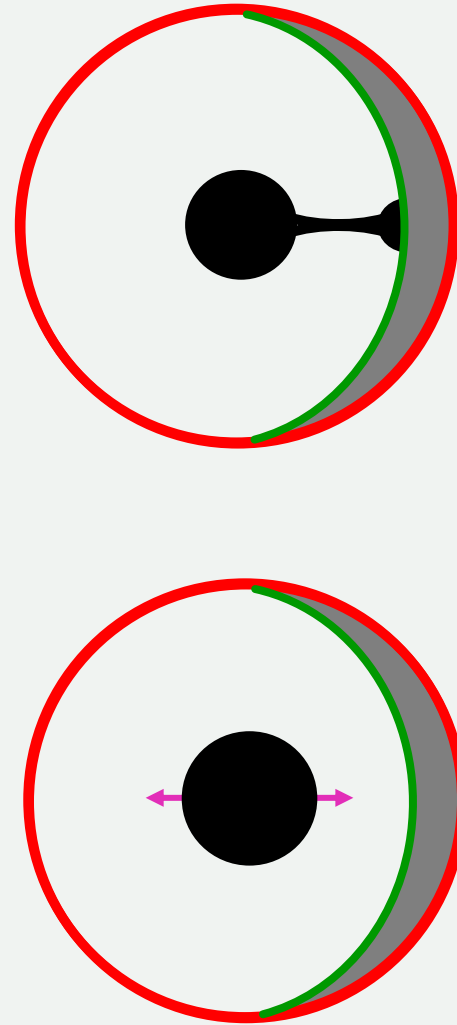


Evaporation as flow between black holes





Evaporation into a bulk black hole



Holographic black hole evaporation

How and Why

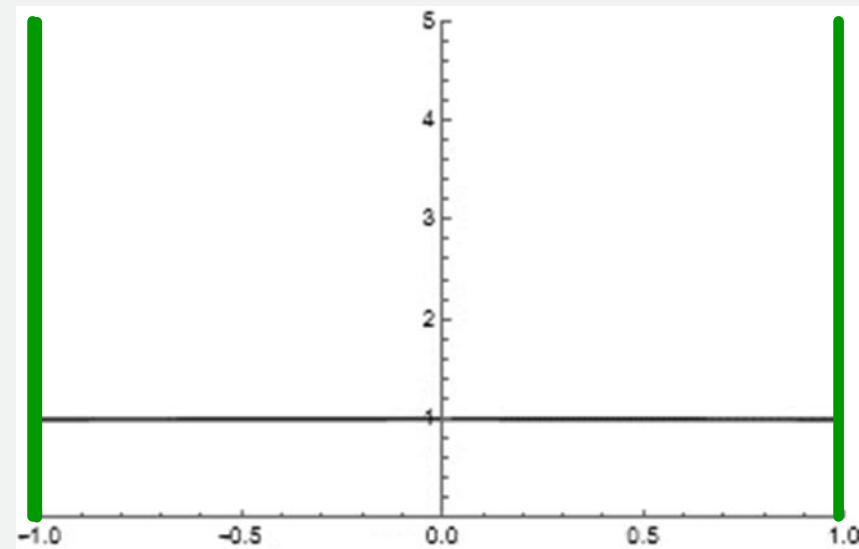
What makes holographic dual evaporation possible?

- Black hole on brane evaporates by funneling radiation into a colder horizon—in the bulk, on another brane, or at a boundary
"Heat (radiation) flows from the **hotter** system to the **colder** one"
- Requires access to thermalized degrees of freedom
- Requires entanglement channel in bulk: black funnel

$$r_0 < L_{\text{AdS}}$$

small AdS black hole
hotter when smaller

Twin black holes evaporating into colder bulk

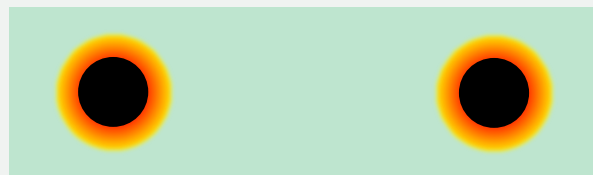
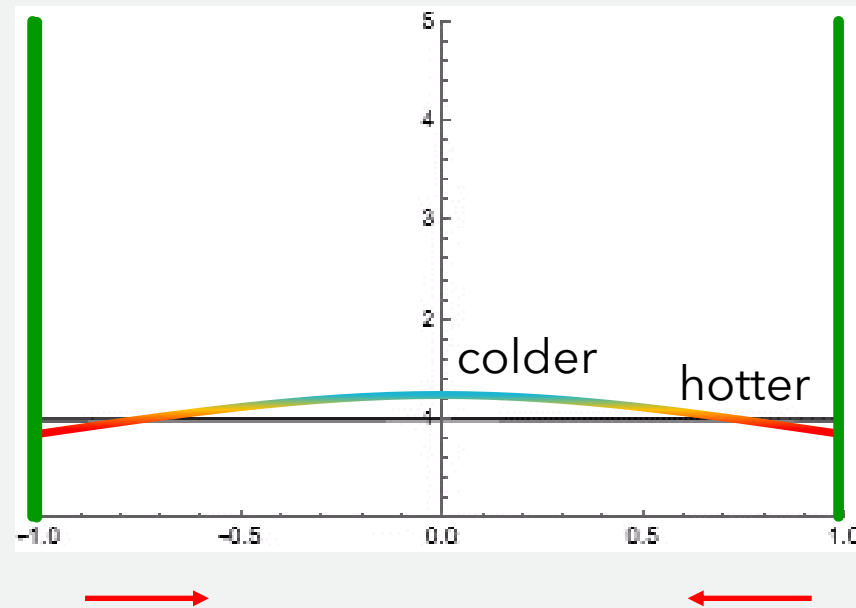


$$r_0 < L_{\text{AdS}}$$

small AdS black hole
hotter when smaller

Twin black holes evaporating into colder bulk

Radiation flows from
hotter to colder

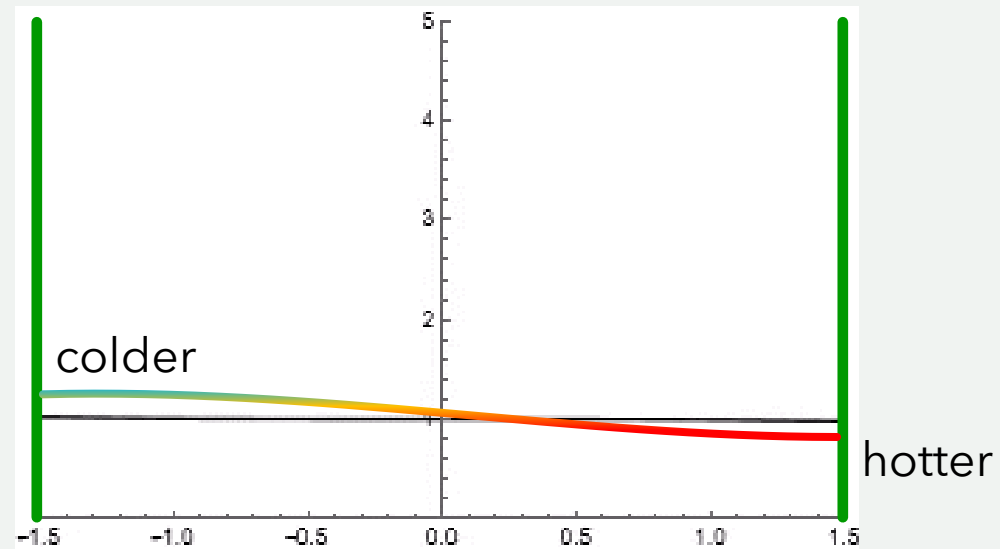


$$r_0 < L_{\text{AdS}}$$

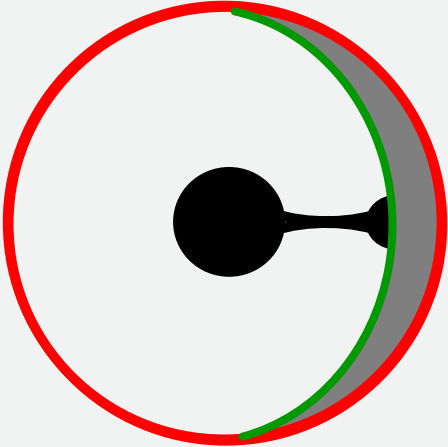
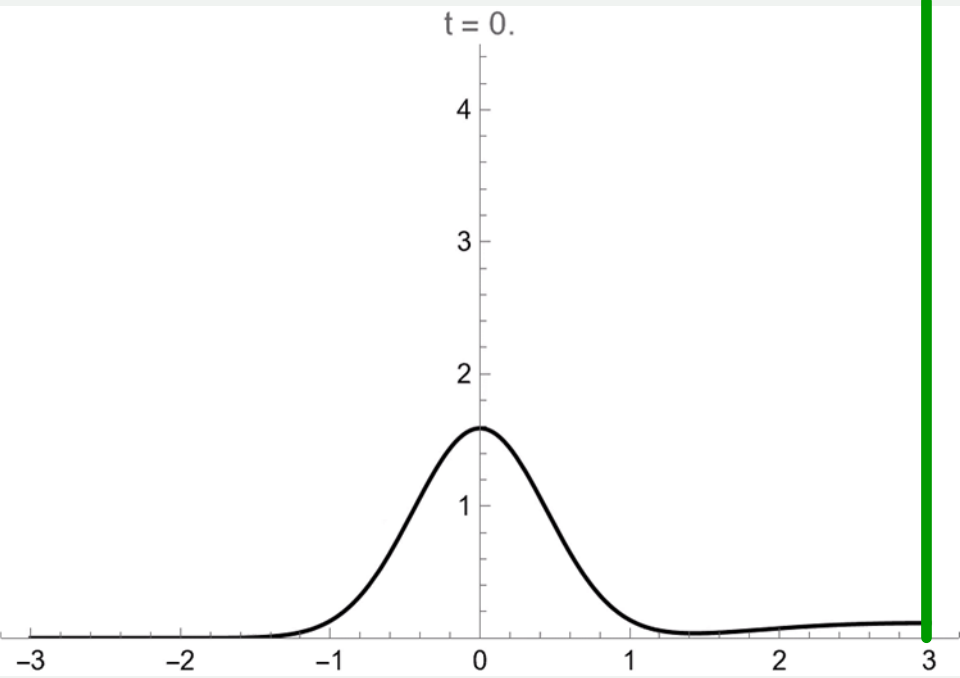
small AdS black hole
hotter when smaller

Hot black hole radiating to colder one

Radiation flows from
hotter to colder



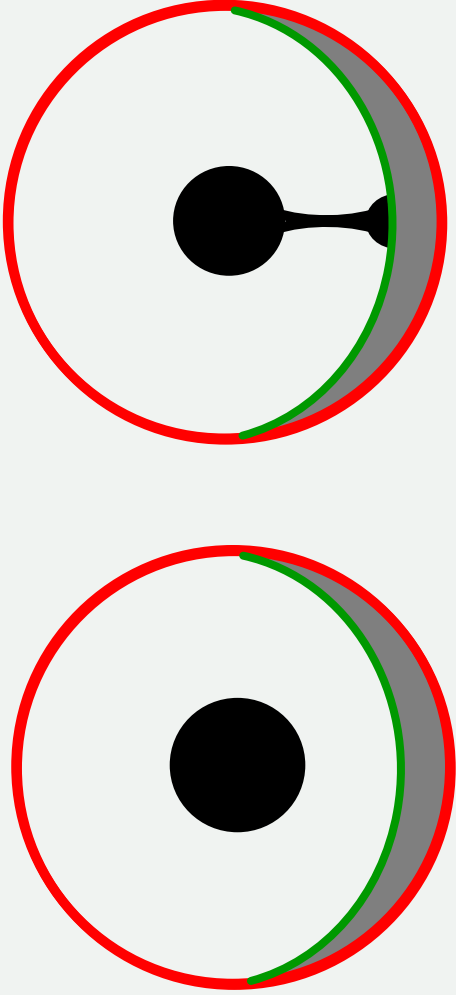
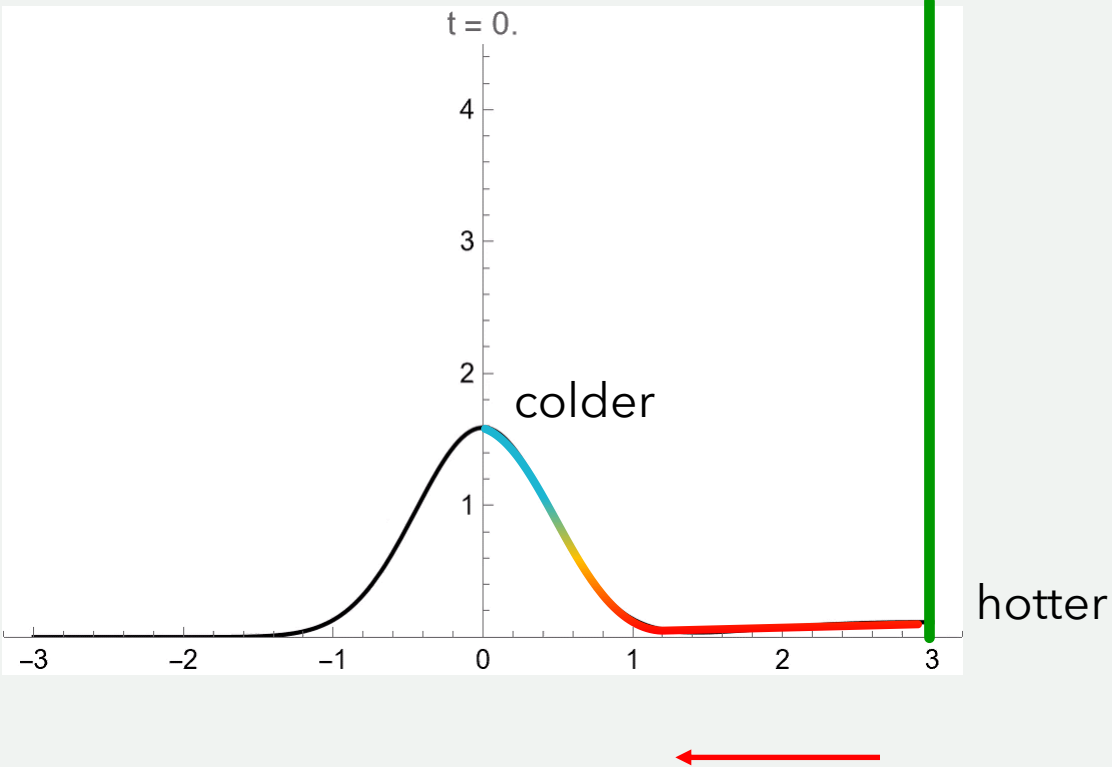
Hot black hole evaporates into colder bath



$r_0 < L_{\text{AdS}}$
small AdS black hole
hotter when smaller

Hot black hole evaporates into colder bath

$r_0 < L_{\text{AdS}}$
small AdS black hole
hotter when smaller



The take away (2): Classicalization of Hawking evaporation

- Dynamical holographic black hole evaporation through classical bulk evolution **is possible** – it is **dual to horizon instabilities**
- **Hot** black hole evaporates **through funnel** into **colder** thermal degrees of freedom

Conclusions

Quantum Black Holes as Braneworld Holograms

- Efficient solution of quantum backreaction = classical bulk problem in one more dimension - not easy but doable
 - New Quantum Black Hole solutions in 3D
 - Classicalization of Hawking evaporation
- Limited to infinite coupling+infinite-N QFT (with peculiarities of its own)
- Bottom-up holography: effective theory, not fully UV complete

Thank you