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Darkness without dark matter and energy – generalized unimodular gravity

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Introduction: darkness in cosmology

DE and DM – rich playground for modifications of GR

DE and DM from pure gravity sector

- not local DoF but global (mechanical) DoF

Reduction of local gauge (diffeomorphism) invariance (Horava-Lifshitz gravity – UV consistent renormalizable QG?)

Unimodular gravity (UMG): # local DoF = 2 + one global DoF (x), $p=w^2$, $w = -1 \neq w(z)$

Lorentz violation + UMG = Generalized UMG: $w=w(z) \neq const$

Plan

Intro: darkness in cosmology

Unimodular gravity

Generalized unimodular gravity (GUMG) -- source of dark fluid

Reduced diffeomorphism invariance and fluid dynamics in a comoving frame

Canonical formalism of GUMG

Discussion: GUMG, cosmological initial conditions and curvature density parameter Ω_{K}

Unimodular gravity

$$\det g_{\mu\nu} = -1$$

$$R_{\mu\nu} - \frac{1}{2}g_{\mu\nu}R = -\Lambda g_{\mu\nu}$$
$$\partial_{\mu}\Lambda = 0$$

x = *const* -- spacetime constant of integration of EoM

$$p = -\varepsilon, \quad \varepsilon = M_P^2 \Lambda$$
 $M_P^2 = (8\pi G)^{-1}$

Local invariance – volume preserving diffeomorphisms:

$$x^{\mu} \to x^{\mu} + \xi^{\mu}(x), \quad \partial_{\mu}\xi^{\mu} = 0$$

Violation of Lorentz invariance – Horava-Lifshitz gravity theory:

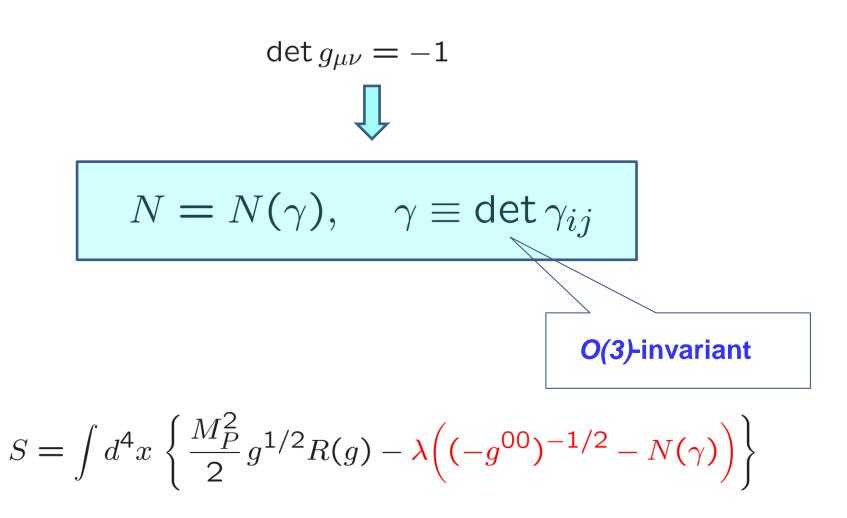
Search for UV consistent quantum gravity – unitary and perturbatively renormalizable

ADM (3+1)-decomposition of 4-metric:

$$g_{\mu\nu}dx^{\mu}dx^{\nu} = (N_i N^i - N^2) dt^2 + 2N_i dt dx^i + \gamma_{ij}dx^i dx^j$$

$$N = \frac{1}{\sqrt{-g^{00}}}, \quad N_i = g_{0i}, \quad N^i = \gamma^{ij}N_j$$
 Lapse and shift functions

Generalized unimodular gravity (GUMG) -- source of dark fluid



$$R^{\mu\nu} - \frac{1}{2}g^{\mu\nu}R = \frac{1}{M_P^2}T^{\mu\nu},$$

$$T^{\mu\nu} \equiv -\frac{2}{g^{1/2}}\frac{\delta}{\delta g_{\mu\nu}}\int d^4x \,\lambda \left((-g^{00})^{-1/2} - N(\gamma)\right)$$

$$= \varepsilon \,u^{\mu}u^{\nu} + p\left(g^{\mu\nu} + u^{\mu}u^{\nu}\right)$$
Perfect fluid
$$u^{\mu} = -g^{0\mu}N$$

$$\varepsilon = \frac{\lambda}{2\sqrt{\gamma}}, \quad p = \frac{\lambda}{\sqrt{\gamma}}\left(\frac{\gamma}{N}\frac{dN}{d\gamma}\right)$$

$$p = w\varepsilon, \quad w = 2\frac{\gamma}{N}\frac{dN}{d\gamma} = 2\frac{d\ln N}{d\ln \gamma}$$
Equation of state

Reduced diffeomorphism invariance

$$\delta_{\xi}g^{\mu\nu} = -\nabla^{\mu}\xi^{\nu} - \nabla^{\nu}\xi^{\mu}, \quad \xi^{\mu} = \begin{bmatrix} \xi^{0} \\ \xi^{i} \end{bmatrix}$$

$$\delta_{\xi}(N - N(\gamma)) \Big|_{N = N(\gamma)} = N \Big[\partial_{t}\xi^{0} - (1 + w) N^{i} \partial_{i}\xi^{0} - w \partial_{i}\xi^{i} \Big] = 0$$

$$\xi^{\mu}_{\perp} = \begin{bmatrix} 0 \\ \xi^{i}_{\perp} \end{bmatrix} \quad \text{two purely spatial diffeos with} \text{ a transverse 3-vector parameter} \quad \partial_{i}\xi^{i}_{\perp}(t, \mathbf{x}) = 0$$

$$\xi^{\mu}_{\parallel} = \begin{bmatrix} \xi^{0} \\ \xi^{i}_{\parallel} [\xi^{0}] \end{bmatrix} \quad \text{diffeo with a timelike parameter} \quad \xi^{0}(t, \mathbf{x})$$

$$\xi^{i}_{\parallel} [\xi^{0}] = \partial^{i}\frac{1}{\Delta}\frac{\partial_{t}\xi^{0} - (1 + w)N^{l}\partial_{l}\xi^{0}}{w}$$

spatially nonlocal dependence of ""II on "

4

Dynamics of dark fluid in a comoving frame

Special case of dark dust – "projectable" GR

$$N = \text{const}, w = 0$$



initial conditions function *S(x)* survives:

$$\varepsilon = \frac{S(\mathbf{x})}{N\sqrt{\gamma}} \equiv \frac{\tilde{S}(\mathbf{x})}{\sqrt{\gamma}}$$

DM? Analogue of mimetic model. Non-clustering DE = global (mechanical) mode Initial conditions for it and its quantization?

Need a physical sector of GUMG – canonical formalism

Lagrangian quantization? Problem of open algebra of gauge generators and BV formalism (Faddeev-Popov technique does not apply beyond one loop):

$$\delta_{\boldsymbol{\xi}} g^{a} = R^{a}_{\mu} \xi^{\mu}, \quad R^{b}_{\mu} \frac{\delta R^{a}_{\nu}}{\delta g^{b}} - R^{b}_{\nu} \frac{\delta R^{a}_{\mu}}{\delta g^{b}} = C^{\lambda}_{\mu\nu} R^{a}_{\lambda},$$
$$C^{\lambda}_{\mu\nu} = C^{\lambda}_{\mu\nu} [g] \neq \text{const}$$

field-dependent and nonlocal

Canonical formalism of GUMG

$$S[\gamma_{ij}, \pi^{ij}; N^i] = \int d^4x \left(\pi^{ij} \dot{\gamma}_{ij} - N(\gamma)H - N^i H_i\right)$$

What are the canonical generators of diffeos?

$$\xi_{||}^{i} = \partial^{i} \frac{1}{\Delta} \frac{\partial_{t} \xi^{0}}{w} + \dots \sim \partial_{t} \xi^{0} + \dots$$

S

Hamiltonian and momentum constraint functions

primary constraints with Lagrange multiplies

Hamiltonian
$$\mathcal{H} = \int d^3x \left(N(\gamma)H + N^iH_i + v^ip_i \right)$$

$$\begin{aligned} \frac{\delta S}{\delta v^{i}} &= -p_{i} = 0\\ \dot{p}_{i} &= \{p_{i}, \mathcal{H}\} = -H_{i} = 0\\ \dot{H}_{i} &= \{H_{i}, \mathcal{H}\} = T_{i} \equiv \partial_{i}(NwH) = 0\\ \dot{T}_{i} &= S_{i} \equiv \partial_{i}(\Omega \partial_{k}N^{k} + X) = 0 \end{aligned} \qquad \text{all constraints}\\ \dot{S}_{i} &= \partial_{i}(\Omega \partial_{k}v^{k} + ...) = 0 \qquad \text{equations on Lagrange multipliers}\\ \Omega &\equiv 1 + w + 2\frac{d\ln w}{d\ln \gamma} \end{aligned}$$

 $N_{\perp}^{\kappa}, v_{\perp}^{\kappa}$ 4 transverse components are not constrained – gauge fixing

GUMG:
$$T_i \equiv \partial_i (NwH) = 0 \Rightarrow NwH = E(t)$$

global (mechanical) mode

Compare:

GR: $H \equiv \frac{\pi_{ij}^2 - \frac{1}{2}\pi^2}{\sqrt{\gamma}} - \sqrt{\gamma}^3 R = 0$ UMG: $N = \frac{1}{\sqrt{\gamma}}, w = -1 \Rightarrow \frac{H}{\sqrt{\gamma}} = \Lambda = \text{const}$

Algebra of constraints and # of DoF

eight
constraints
$$p_i = 0, H_i = 0, T_i = 0, S_i = 0$$

four 1st class constraint functions

$$p_i^{\perp}, H_i^{\perp}$$

four 2nd class constraint functions $p_i^{||}, H_i^{||}, T_i^{||}, S_i^{||}$

2f # of DoF =

- **#** of phase space variables
- 2£ # of 1st class constraints
- # of 2^{nd} class constraints = $18 8 4 = 2 \pm 3$

Discussion

 Coupled field theory and mechanical (finite-dimensional) system – the problem of disentangling the global mode in the Hamiltonian and symplectic form

• Breakdown of the 3D diffeo invariance -- local invariance is too ugly. Recovery of 3D invariance (like in Horava gravity)

$$N - N(\gamma) = 0 \Rightarrow \Delta N - \kappa R = 0$$

Anisotropic stress tensor:

dimensionless parameter -- to avoid extra scale

$$\varepsilon = -\Delta\lambda$$

$$T_{ij} = -\frac{2}{N} \left(\kappa G_{ij} - \kappa (\nabla_i \nabla_j - \gamma_{ij} \Delta) + \nabla_{(i} N \nabla_{j)} - \frac{1}{2} \gamma_{ij} \nabla^n N \nabla_n \right) \lambda$$

Perfect fluid on homogeneous FRW:

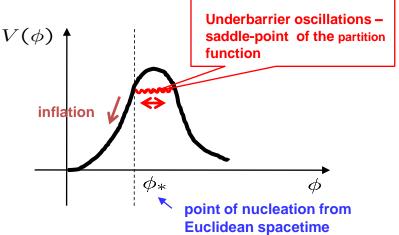
$$p = \frac{\kappa \lambda}{3N}, \quad \varepsilon = 0$$

violation of energy conditions!

Motivation – model of microcanonical initial conditions for CFT driven cosmology

a) Density matrix generalization of the Hartle-Hawking state

- b) Initial conditions for inflation at the top of inflaton potential
- c) Quasi-thermal primordial power spectrum -- revival of a mild "big bang" preceding inflation
- New type of hill-top inflation as a source of Higgs or *R*²-inflation model



 e) Origin of the Universe is a semiclassical subplanckian phenomenon – intrinsic mechanism of suppression for high energy scales due to hidden sector of conformal fields

BUT ONLY FOR K=+1!

$$\Omega_K = -\frac{K}{H^2 a^2} < 0, \quad \Omega_K + \Omega_\Lambda + \Omega_m = 1$$

Planck, lensing and BAO data:

$$\Omega_K = 0.000 \pm 0.005$$

 $\Omega_K \rightarrow 0$

Observational tendency:

Imitation of $K \neq 0$ by GUMG in
flat space Friedmann equation $H^2 = \frac{\Omega_m + \Omega_A + \Omega_K}{3M_P^2}$

$$w = -\frac{1}{3}, \ \varepsilon_K = -\frac{3M_P^2 K}{a^2} \sim \frac{1}{\gamma^{(1+w)/2}}, \ \Omega_K = \frac{\varepsilon_K}{3M_P^2 H^2}$$

THANK YOU!