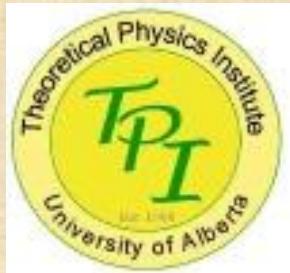


Principal Killing Strings

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GC2018, Yukawa Institute, Kyoto,
February 16, 2018



Based on:

"Principal Killing strings in higher-dimensional Kerr-NUT-(A)dS spacetimes", Jens Boos and V.F,
e-Print: arXiv:1801.00122 (2018);

"Stationary black holes with stringy hair",
Jens Boos and V.F., e-Print: arXiv:1711.06357
(2017), (to appear in PRD);

"Stationary strings and branes in the higher-dimensional Kerr-NUT-(A)dS spacetimes",
David Kubiznak and V.F., JHEP 0802 (2008) 007;
e-Print: arXiv:0711.2300.

"Black holes, hidden symmetries, and
complete integrability",

V.F., Pavel Krtous and David Kubiznak,

Living Rev.Rel. 20 (2017) no.1, 6;

e-Print: arXiv:1705.05482.

Solving stationary string equations in the Kerr-NUT-(A)dS background

Killing-Yano objects

Conformal KY tensor (CKY) of rank p in D dims:

$$\nabla_X \omega = \frac{1}{p+1} X \bullet (\nabla \wedge \omega) + \frac{1}{D-p+1} X \wedge (\nabla \bullet \omega) + 0,$$

Killing tensor (KYT): $\nabla \bullet \omega \sim \delta \omega = 0$;

Closed conformal KY tensor (CCKY): $\nabla \wedge \omega \sim d\omega = 0$.

$$\nabla_X (*\omega) = \frac{1}{p_* + 1} X \bullet (\nabla \wedge *\omega) + \frac{1}{D - p_* + 1} X \wedge (\nabla \bullet *\omega),$$

$$p_* = D - p$$

Properties:

- $*(KYT) = CCKY$
- $*(CCKY) = KYT$
- $CCKY \wedge CCKY = CCKY$

Principal tensor = non-degenerate rank 2 CCKY tensor

$$D = 2n + \varepsilon$$

$$\nabla_c h_{ab} = g_{ca} \xi_b - g_{cb} \xi_a, \quad \xi_a = \frac{1}{D-1} \nabla^b h_{ba}.$$

$$\mathbf{h} = r \mathbf{l}_+ \wedge \mathbf{l}_- + \sum_{\mu=1}^{n-1} x_\mu \mathbf{e}^\mu \wedge \hat{\mathbf{e}}^\mu,$$

$$\mathbf{g} = -\mathbf{l}_+ \mathbf{l}_- - \mathbf{l}_- \mathbf{l}_+ + \sum_{\mu=1}^{n-1} (\mathbf{e}^\mu \mathbf{e}^\mu + \hat{\mathbf{e}}^\mu \hat{\mathbf{e}}^\mu) + \varepsilon \hat{\mathbf{e}}^0 \hat{\mathbf{e}}^0.$$

$$(\mathbf{l}_+, \mathbf{l}_-) = -1, \quad (\mathbf{e}^\mu, \mathbf{e}^\mu) = (\hat{\mathbf{e}}^\mu, \hat{\mathbf{e}}^\mu) = (\hat{\mathbf{e}}^0, \hat{\mathbf{e}}^0) = 1.$$

Darboux basis: $(\mathbf{l}_+, \mathbf{l}_-, \mathbf{e}^\mu, \hat{\mathbf{e}}^\mu, \hat{\mathbf{e}}^0)$.

Non-degenerate: There are exactly n non-vanishing "eigenvalues" (r, x_μ) that are functionally independent in some domain. In this domain none of the gradients of them is a null vector.

A metric which admits a principal tensor
is off-shell Kerr-NUT-(A)dS metric.

It contains n arbitrary functions of 1 variable.

On-shell: Einstein equations are satisfied \Rightarrow
Kerr-NUT-(A)dS solution.

- ξ is a primary Killing vector: $L_\xi \mathbf{g} = L_\xi \mathbf{h} = 0$;
- $\mathbf{h}^{(j)} = \frac{1}{j!} \mathbf{h}^{\wedge j}$ is a CCKY $2j$ -form;
- $\mathbf{f}^{(j)} = * \mathbf{h}^{(j)}$ is a KY (D-2j) form
- $k_{(j)}^{ab} = \frac{1}{(D-2j-1)!} f_{c_1 \dots c_{D-2j-1}}^{(j)a} f^{(j)b c_1 \dots c_{D-2j-1}}$ is a rank 2 Killing tensor;
- $\zeta_{(j)} = \mathbf{k}_{(j)} \cdot \xi$, ($j = 0, \dots, n-1 + \varepsilon \equiv m$) are commuting (secondary) Killing vectors;
- $\mathbf{k}_{(0)} = g$;
- Frobenius theorem: $\xi = \partial_\tau$, $\zeta_{(j)} = \partial_{\psi_j}$;
- (r, x_μ, τ, ψ_j) are canonical coordinates;
- I_\pm are principal null directions; their integral lines are geodesics.

Geodesic equations are completely integrable:
There exist D integrals of motion for a free particle,
 $(n + \varepsilon)$ first order $\zeta_{(j)} \cdot p$ and n second order $p \cdot k_{(j)} \cdot p$.

Q: If instead of a particle one has a string:
Are Nambu-Goto string equations completely
integrable in the Kerr-NUT-(A)dS geometry?

A: In a general case - No.
If a string is stationary - Yes.

Stationary strings in Kerr-NUT-(A)dS

Killing vector $\xi = \partial_t$, coordinates - (t, y^i)

$$g_{ab} = p_{ab} + \frac{\xi_a \xi_b}{\xi^2}, \quad F = -\xi^2, \quad A_i = \frac{\xi_i}{\xi^2},$$

$$ds^2 = -F(dt + A_i dy^i)^2 + p_{ij} dy^i dy^j.$$

Nambu-Goto action for a string

$$I = -\Delta t E, \quad E = \mu \int \sqrt{F} dl = \mu \int d\sigma \sqrt{F p_{ij} \frac{dy^i}{d\sigma} \frac{dy^j}{d\sigma}} .$$

String configuration $y^i(\sigma)$ is a geodesic in $(D-1)$ -dimensional space with metric $\tilde{p}_{ij} = F p_{ij}$.

If metric g_{ab} admits the principal tensor it has $n + \varepsilon$ Killing vectors and n rank 2 Killing tensors. This gives $D = 2n + \varepsilon$ integrals of motion for a free particle.

The reduced ref-shifted metric \tilde{p}_{ij} does not admit a principal tensor. However, when ξ is a primary Killing vector, it has $n - 1 + \varepsilon$ Killing vectors and n rank 2 Killing tensors. This gives $D = 2n - 1 + \varepsilon$ integrals of motion. Thus, the stationary string equations are completely integrable.

[V.F. and David Kubiznak (2008)]

Principal Killing Strings

$\ell_{\pm} = \beta I_{\pm}$ is a tangent vector to a principal null geodesic in the affine parametrization:

$$\nabla_{\ell_{\pm}} \ell_{\pm} = 0, \quad [\ell_{\pm}, \xi] = 0.$$

Frobenius theorem implies that a ST is foliated by 2-D (Killing) surfaces Σ . Coordinates $Y^a = (z^A, y^i)$. Equations of a given Σ are $y^i = const.$

$z^A = (v, \lambda_{\pm})$ are coordinates on Σ , such that $\xi = \partial_v$, $\ell_{\pm} = \partial_{\lambda_{\pm}}$.

The Killing surface in the off - shell Kerr - NUT - (A)dS metric is minimal.

Induced 2-D metric: $d\gamma^2 = \gamma_{AB} dz^A dz^B = \xi^2 d\lambda_\pm^2 - 2d\nu d\lambda_\pm$.

$n_{(i)}^\mu$ are (D-2) mutually orthogonal unit vectors normal to Σ_\pm . The extrinsic curvature is:

$$\Omega_{(i)AB} = g_{ab} n_{(i)}^a Y^c,_A \nabla_c Y^b,_B, \quad \Omega_{(i)} = g_{ab} n_{(i)}^a Z^b = 0.$$

$$Z^b = \gamma^{AB} Y^c,_A \nabla_c Y^b,_B = - \left(\xi^a \nabla_a \ell_\pm^b + \ell_\pm^a \nabla_a \xi^b + \xi^2 \ell_\pm^a \nabla_a \ell_\pm^b \right).$$

$$Z^b = 2\ell_\pm^a \nabla_a \xi^b = -F_a^b \ell_\pm^a = -\kappa_\pm \ell_\pm^b.$$

$$F^a_b \ell_\pm^b = \kappa_\pm \ell_\pm^a, \quad \kappa_\pm = \frac{1}{2} \ell_\pm^a (\xi^2)_{;a}.$$

Incoming principal Killing string:

$$\lambda = -r, \quad \psi_j = -P_n^{(j)}(r), \quad x_\mu = \text{const},$$

$$P_n^{(j)} = \int \frac{r^{2(n-1-j)} dr}{X_n(r)}.$$

$$(\tau, r, x_\mu, \psi_k) \rightarrow (v, r, x_\mu, \hat{\phi}_k)$$

$$d\tau = dv - \sum_{j=1}^m a_j d\hat{\phi}_j - \frac{r^{2(n-1)}}{X_n} dr,$$

$$d\psi_j = a_j^{-1} d\hat{\phi}_j - \frac{1}{X_n} r^{2(n-1-j)} dr.$$

String equation in the null incoming coordinates

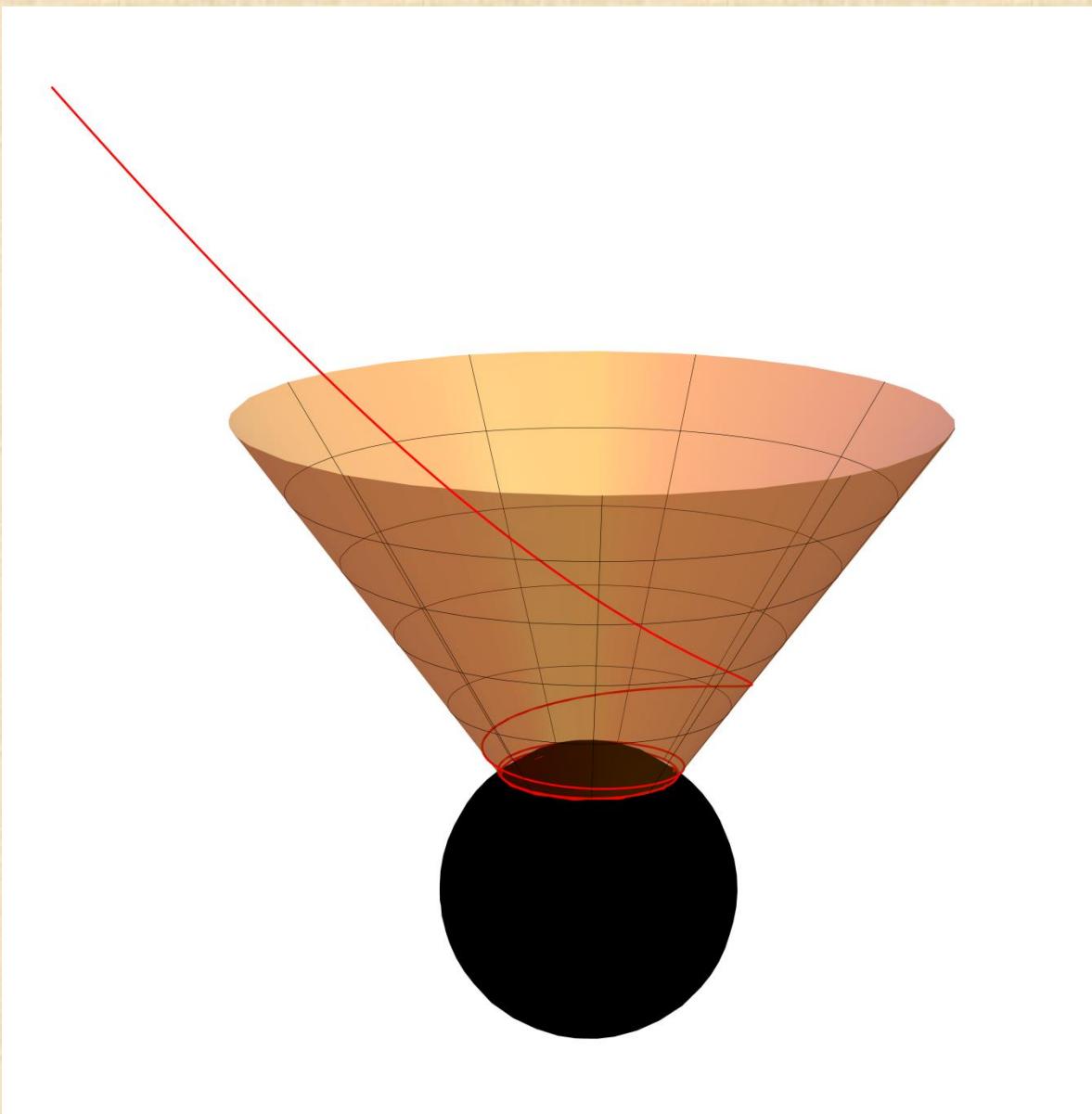
$$\hat{\phi}_j = \hat{\phi}_j^0 = \text{const}, \quad x_\mu = x_\mu^0 = \text{const}.$$

String's stress-energy tensor in in-coming coordinates

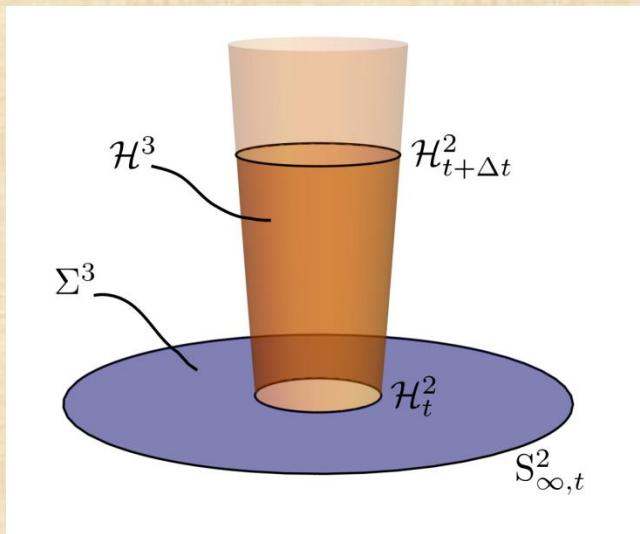
$$T^{ab} = -\frac{\mu_s}{\sqrt{-g}} \left(2\xi^{(a} \ell^{b)} + \xi^2 \ell^a \ell^b \right) q,$$

$$q = q(x_\mu, \hat{\phi}_j \mid x_\mu^0, \hat{\phi}_j^0) \equiv$$

$$\prod_{\mu=1}^{n-1} \delta(x_\mu - x_\mu^0) \prod_{j=1}^m \delta(\hat{\phi}_j - \hat{\phi}_j^0).$$



Applications to Myers-Perry ST



$$-16\pi \frac{D-3}{D-2} M = \oint_B \nabla^a \xi^b dS_{ab},$$

$$-16\pi J_{(i)} = \oint_B \nabla^a \zeta_{(i)}^b dS_{ab}.$$

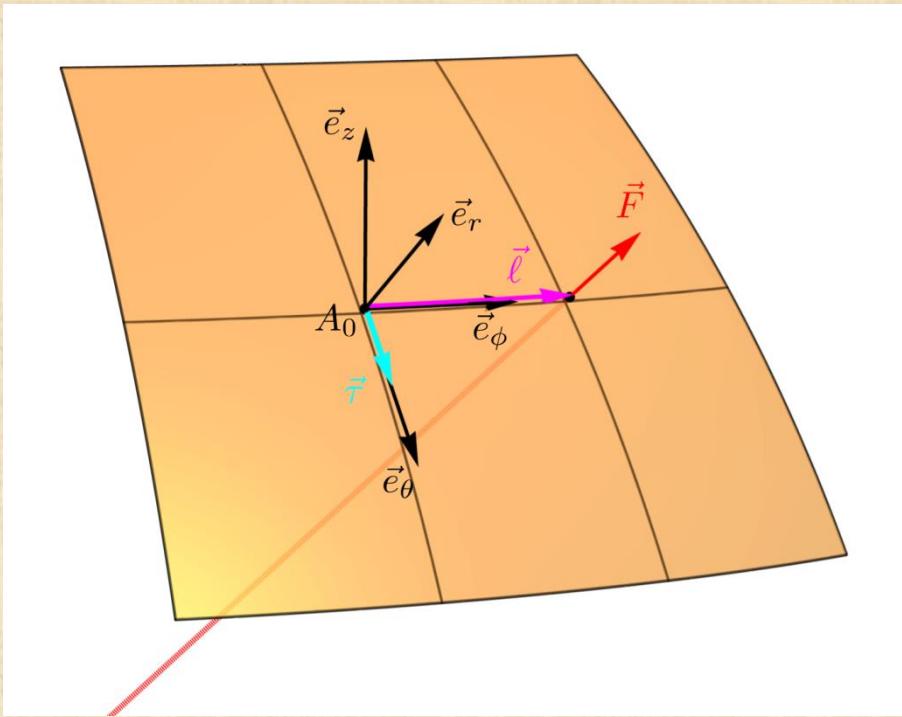
$$J_{(i)} = \frac{2}{D-2} Ma_i.$$

$$\Delta E = \int_{H^{D-1}} T^a{}_b \xi^b d\Sigma_a, \quad \Delta J_i = \int_{H^{D-1}} T^a{}_b \zeta_i^b d\Sigma_a$$

$$\dot{M} = 0, \quad j_{(i)} = -\mu_s a_i \left(\mu_i^0 \right)^2.$$

$$J_{(i)} = J_{(i)}^0 \exp(-v/v_i), \quad v_i = \frac{2M}{D-2} \frac{1}{\mu_s (\mu_i^0)^2}.$$

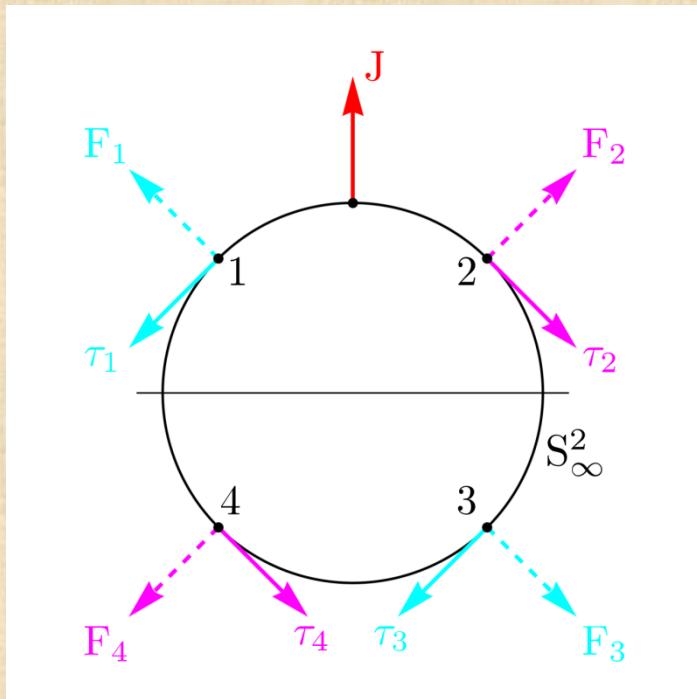
$\mathbf{j} = \text{Projection}(\boldsymbol{\tau})$



e_i and \hat{e}_i are unit vectors along x_i and y_i ; ω^i and $\hat{\omega}^i$ are dual unit forms. $\mathbf{J} = \sum_{i=1}^m \frac{2Ma_i}{D-2} \omega^i \wedge \hat{\omega}^i; x_i^0 = 0 \Rightarrow \boldsymbol{\delta} = \sum_{i=1}^m a_i \mu_i^0 e_i$,

$$\mathbf{F} = \frac{\mu_s}{r} \left[\sum_{i=1}^m \left(x_i^0 e_i + y_i^0 \hat{e}_i \right) + (1-\varepsilon) z^0 \hat{e}_z \right].$$

$$\boldsymbol{\tau} \equiv \boldsymbol{\delta} \wedge \mathbf{F} \equiv \mu_s \left[\sum_{i,j=1}^m a_i \mu_i^0 \mu_j^0 \omega^i \wedge \hat{\omega}^j + \sum_{i=1}^m a_i \mu_i^0 \mu_0^0 (1-\varepsilon) z^0 \omega^i \wedge \omega^0 \right].$$



2^n string segments = 2^{n-1} "infinite captured strings"

Instead of Summary...

Happy Birthday

