Decoherence of Bubble Universes

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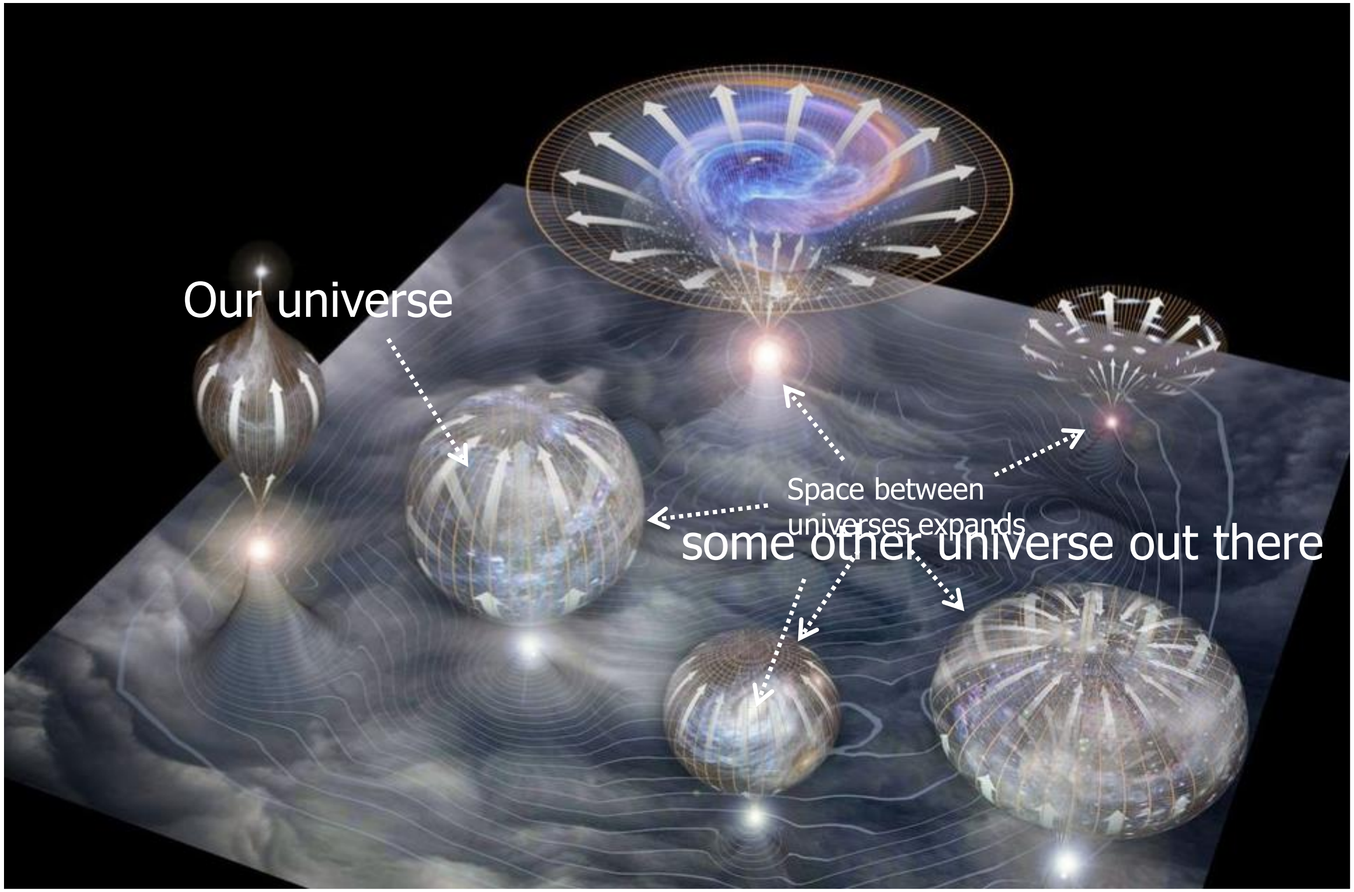
Inflationary cosmology/String landscape

Inflationary cosmology/String landscape suggest that our universe may not be the only universe but is part of a vast complex of universes that we call the multiverse.



These universes may be highly entangled initially.

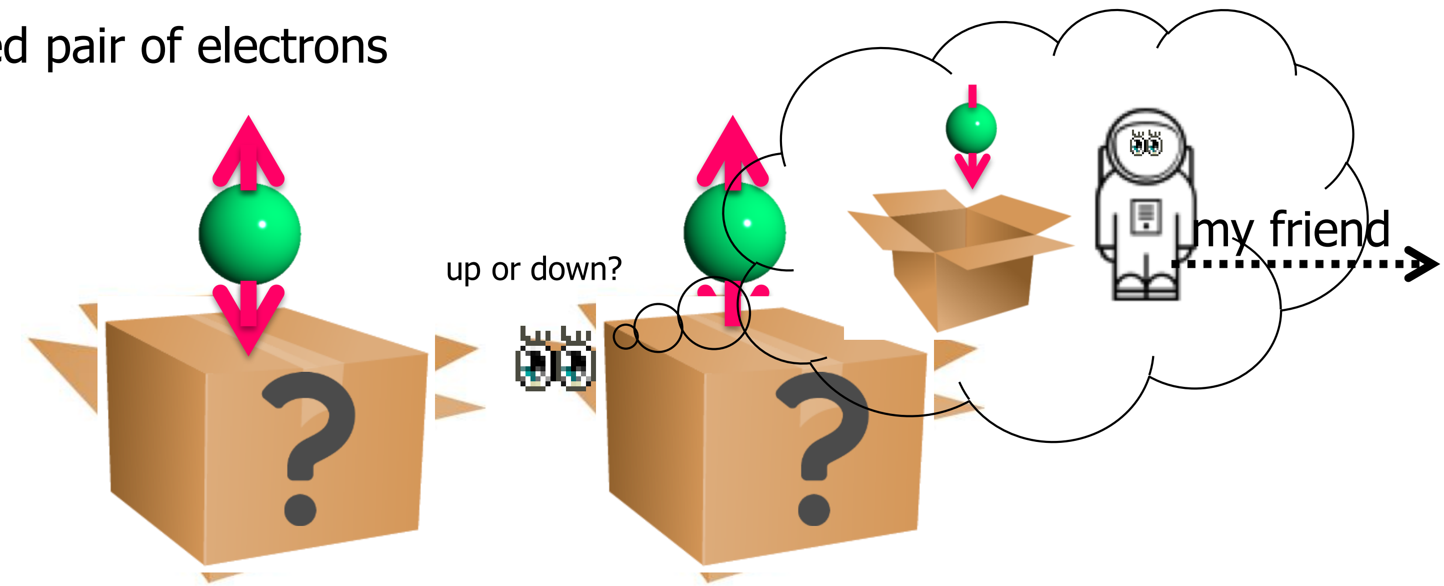
Quantum entanglement?

The most fascinating aspect: *Einstein-Podolsky-Rosen paradox*

To affect the outcome of local measurements instantaneously once a local measurement is performed.

The information travels faster than the speed of light?

Eg) An entangled pair of electrons



The instant measurement determines exactly what the other will find.

Quantum entanglement is purely an aspect of quantum physics that is so different from classical physics.

But my friend doesn't know it until he measures it by himself.

Causality remains intact.

Entanglement exists in arbitrary large distances

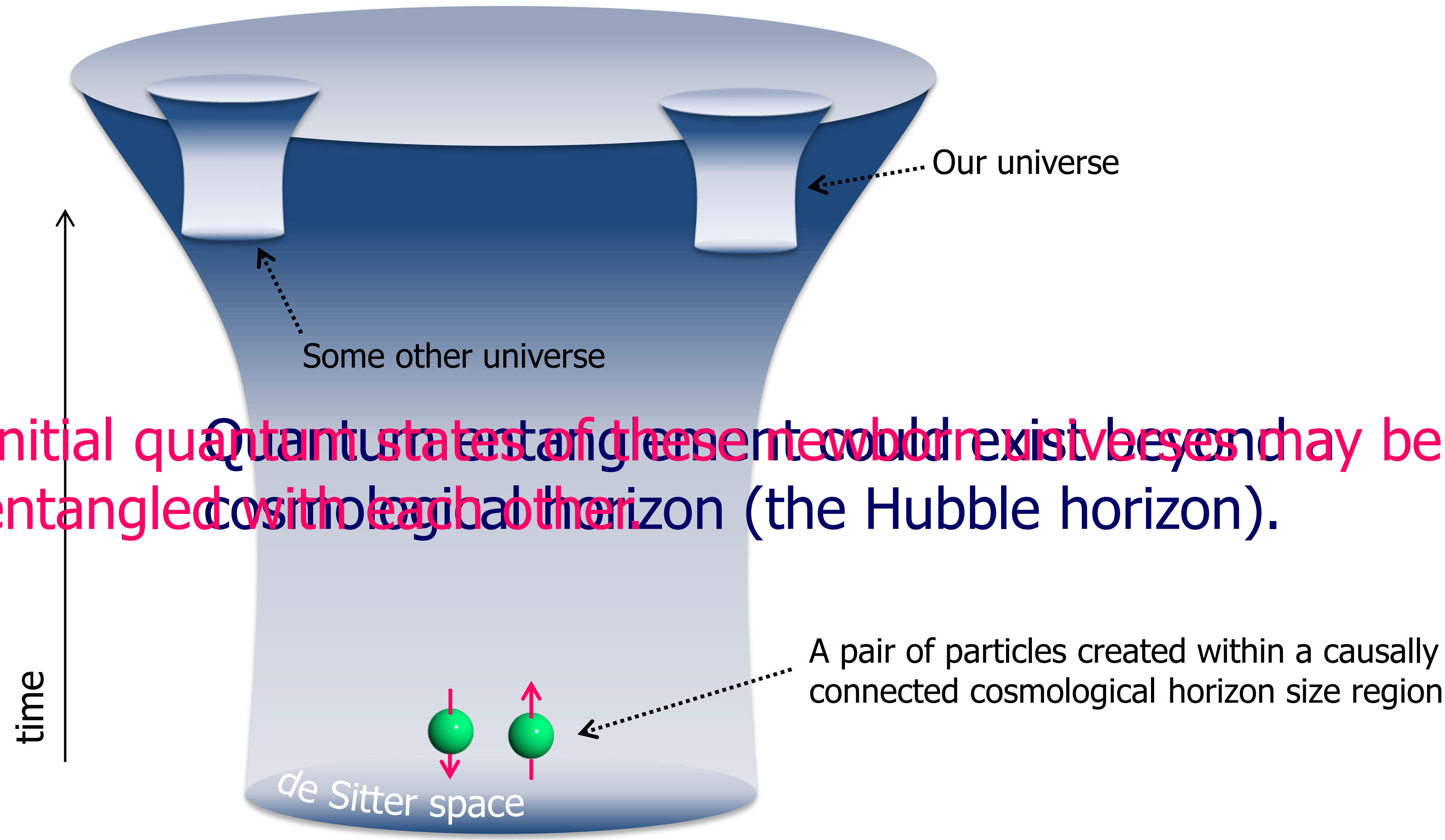
In principle, you gain information about the partner particle by measuring your own particle wherever the partner particle goes, if they are entangled.

a causally disconnected different universe

Naive expectation

Inflationary universe is approximated by a de Sitter space.

In quantum mechanics, vacuum state is full of virtual particles in entangled pairs.

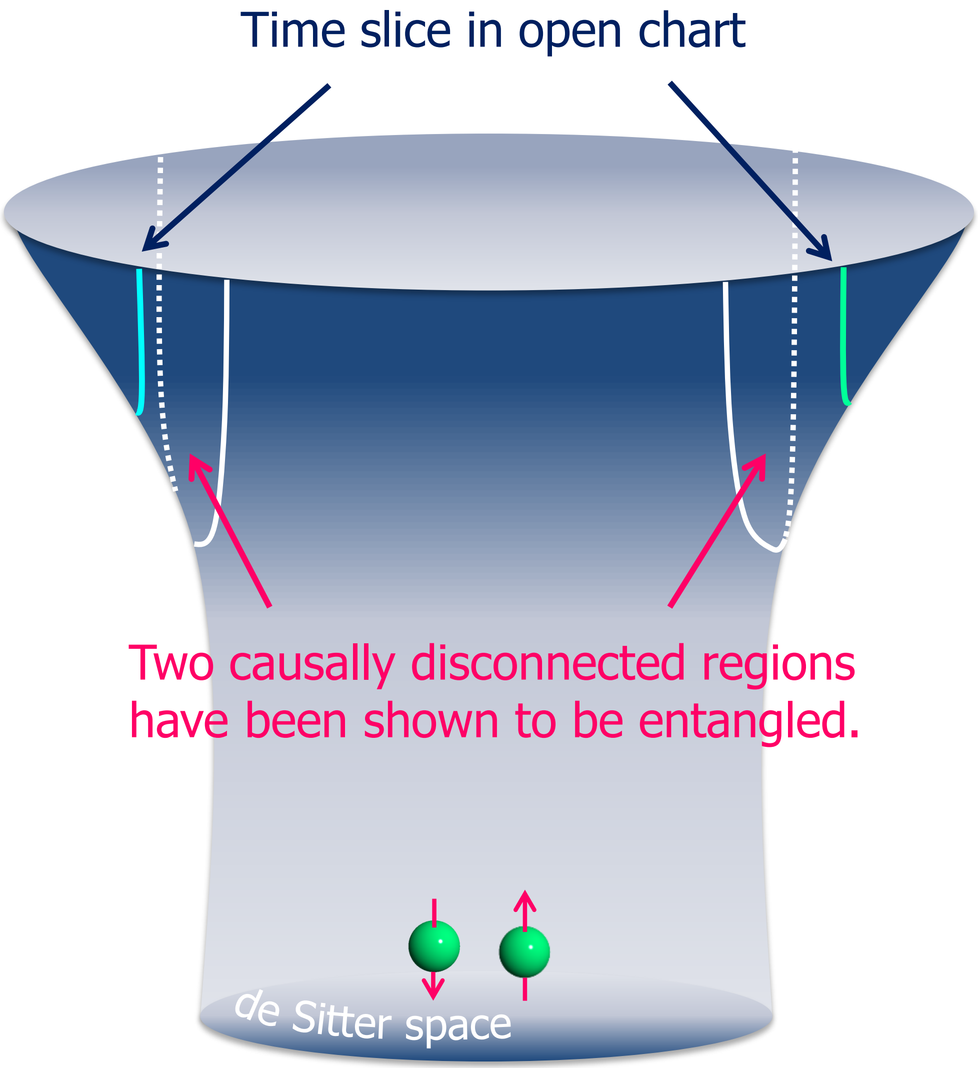


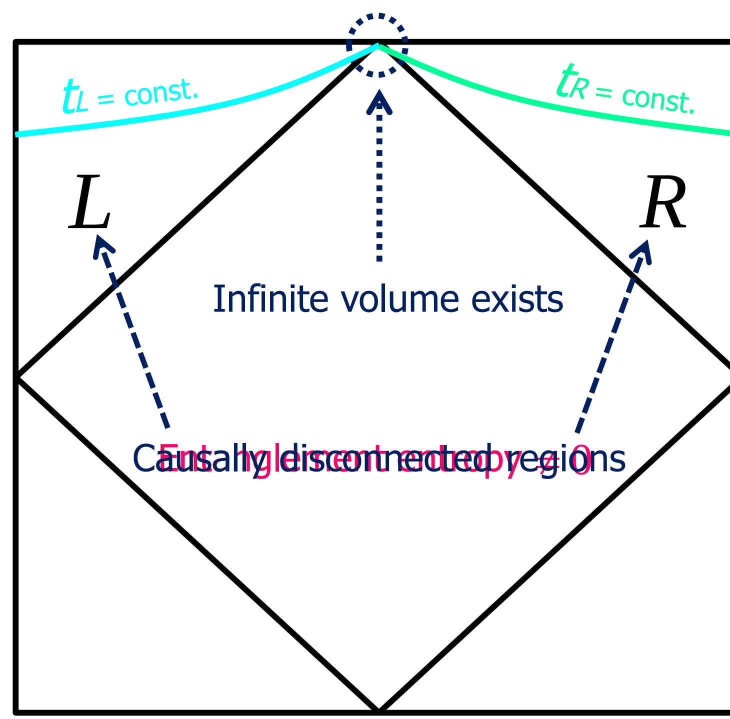
Initial quantum states of these newborn universes may be entangled with each other.

Quantum entanglement could exist beyond a cosmological horizon (the Hubble horizon).

Good news

Maldacena & Pimentel (2013)





Open chart

String/cosmic landscape

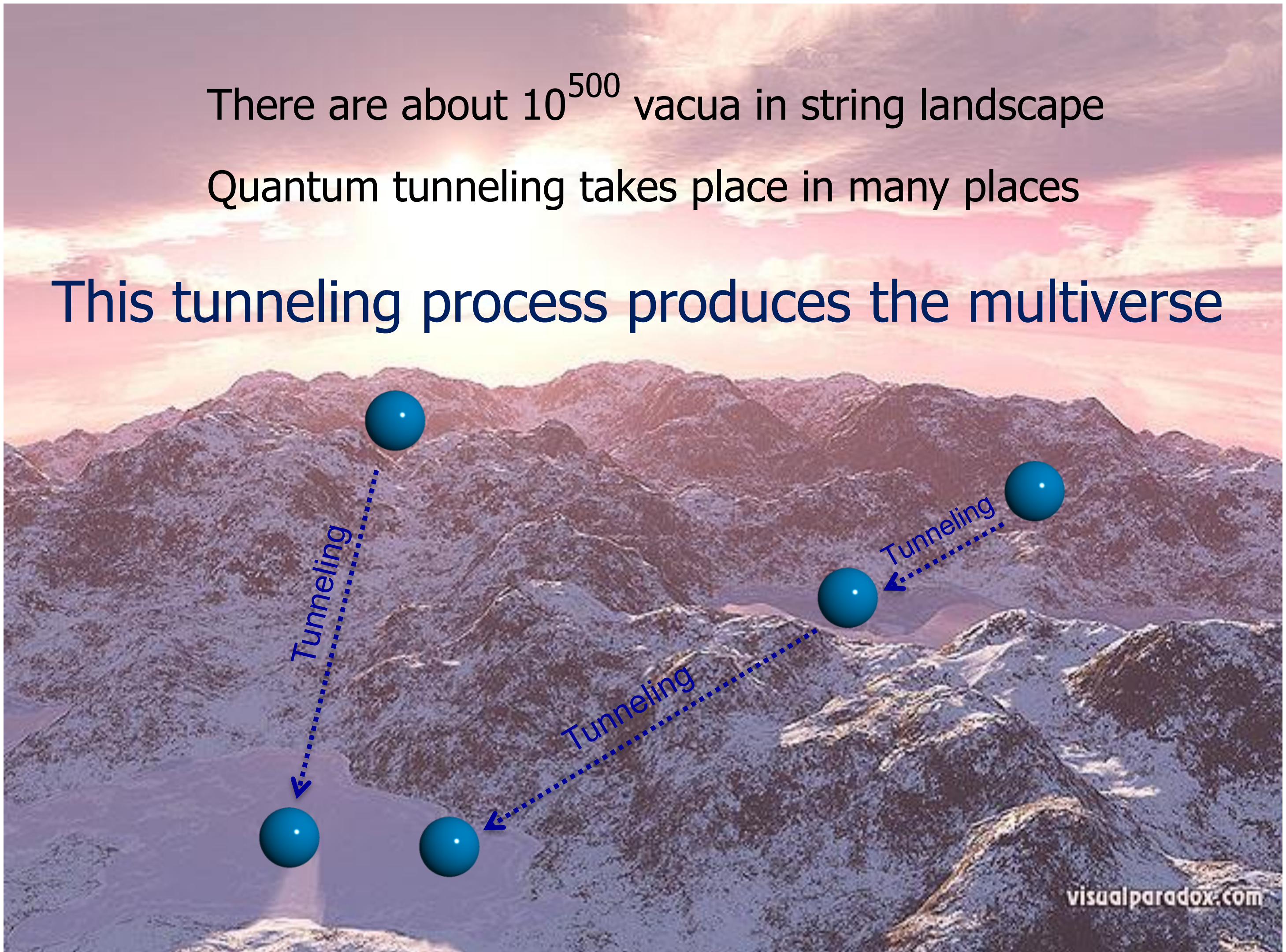
Sato et al. (1981), Vilenkin (1983), Linde (1986), Bousso & Polchinski (2000), Susskind (2003)

The configuration space of all possible values of scalar fields with all possible potentials.

There are about 10^{500} vacua in string landscape

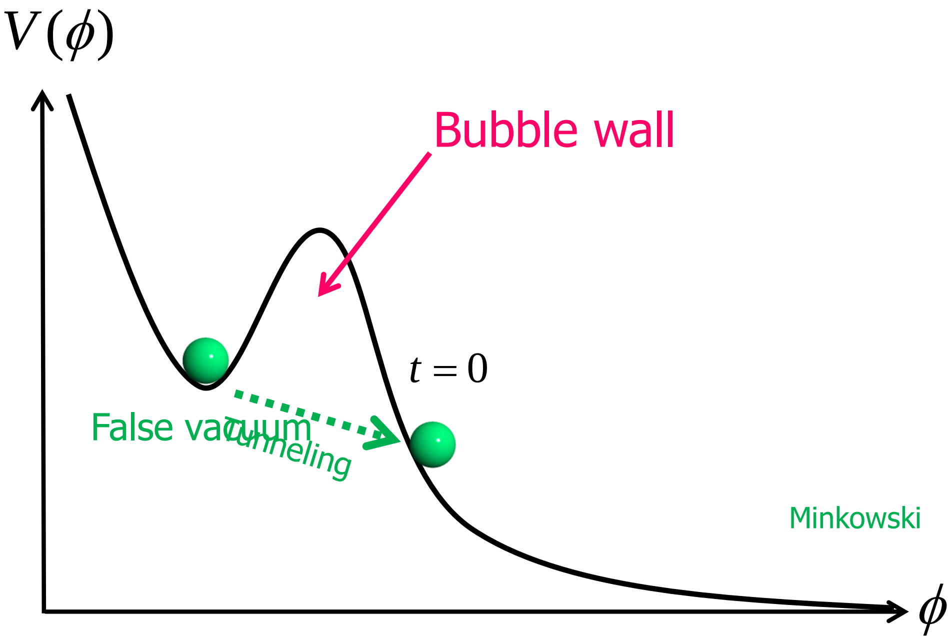
Quantum tunneling takes place in many places

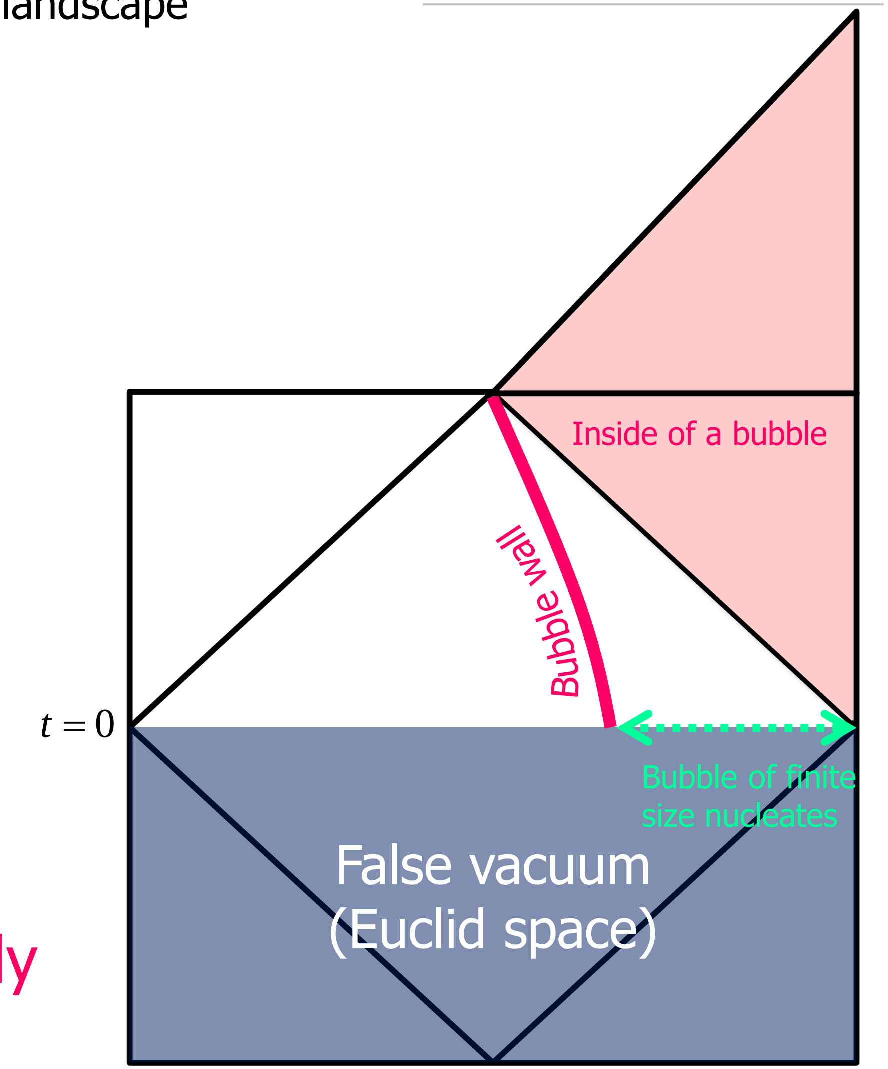
This tunneling process produces the multiverse



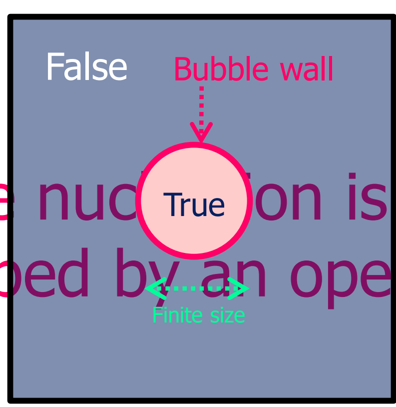
Open chart describes bubble nucleation

Eg) One type of potential of a scalar field in the landscape





Open chart



Bubble nucleation is conveniently described by an open chart.

Review of Maldacena & Pimentel's computation

Action (No bubble wall)

$$S=\int d^4x\sqrt{-g}\left[-\frac{1}{2}g^{\mu\nu}\partial_\mu\phi\partial_\nu\phi-\frac{m^2}{2}\phi^2\right]$$

Metric in each *R* and *L* region

$$ds^2=H^{-2}\left[-dt_{R/L}^2+\sinh^2 t_{R/L}\left(dr_{R/L}^2+\sinh^2 r_{R/L}\,d\Omega^2\right)\right]$$

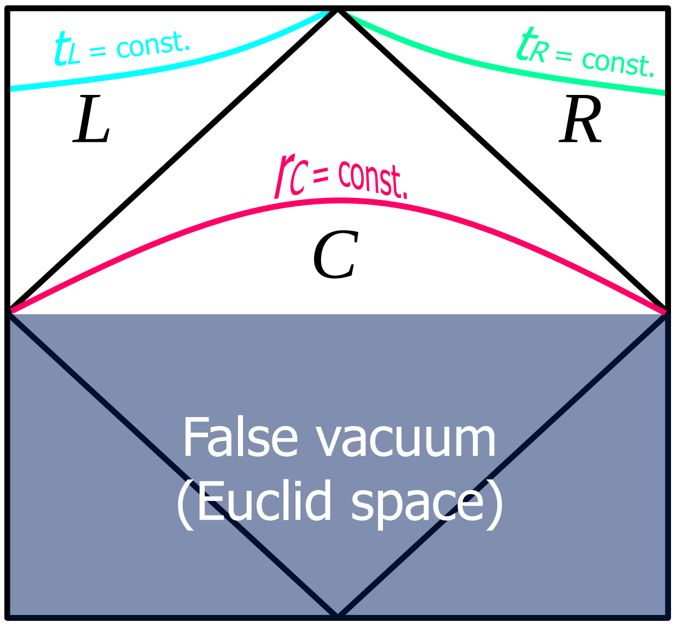
The Hubble radius^2 of de Sitter space

and *C* region

$$ds^2=H^{-2}\left[dt_C^2+\cos^2 t_C\left(-dr_C^2+\cosh^2 r_C\,d\Omega^2\right)\right]$$

can be obtained by analytic continuation from the Euclidean metric

$$ds^2=H^{-2}\left[-d\tau^2+\cos^2\tau\,\left(d\rho^2+\sin^2\rho\,d\Omega^2\right)\right]$$



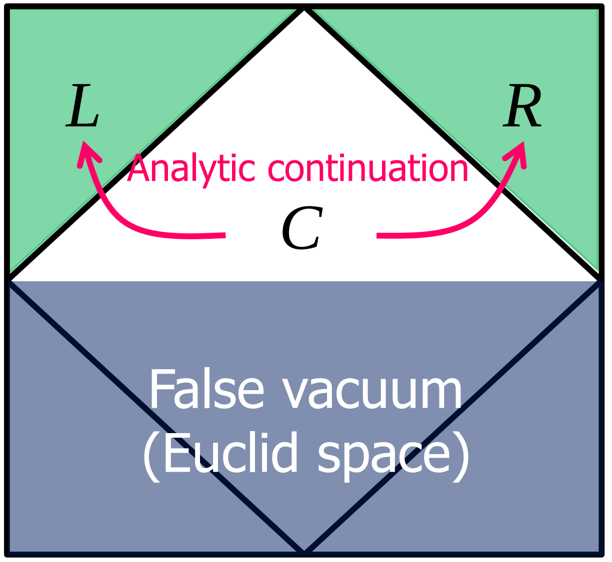
Open chart

The positive freq. mode in the Euclidean vacuum

Separation of variables

$$f_{p\ell m}(t_C, r_C, \Omega) \propto \frac{H}{\cos t_C}\chi_p(t_C)\, Y_{p\ell m}(r_C, \Omega)$$

Harmonic functions on the 3-dim hyperbolic space



The solutions of the mode function in the *C* region

$$\chi_p(t_C) = P^{ip}_{\nu-1/2}(\sin t_C) \qquad \nu = \sqrt{\frac{9}{4} - \frac{m^2}{H^2}}$$

The associated Legendre function

Mass parameter

We want the positive frequency mode functions supported both on *R* and *L* regions which are relevant for the bubble universes through false vacuum decay.

We require regularity in the lower hemisphere of the Euclidean de Sitter space when it is analytically continued to those regions.

Euclidean vacuum (Bunch-Davies vacuum) solutions

Sasaki, Tanaka & Yamamoto (1995)

Solutions supported both on the *R* and *L* regions

$$\chi_p^R(t) = \begin{cases} P_{\nu-1/2}^{ip}(\cosh t_R) \\ \dfrac{\cos\pi\nu}{i\sinh\pi p} P_{\nu-1/2}^{ip}(\cosh t_L) + e^{-\pi p} \dfrac{\cos\pi(ip+\nu)}{i\sinh\pi p} \dfrac{\Gamma(\nu+1/2+ip)}{\Gamma(\nu+1/2-ip)} P_{\nu-1/2}^{-ip}(\cosh t_L) \end{cases}$$

$$\chi_p^L(t) = \begin{cases} \dfrac{\cos\pi\nu}{i\sinh\pi p} P_{\nu-1/2}^{ip}(\cosh t_R) + e^{-\pi p} \dfrac{\cos\pi(ip+\nu)}{i\sinh\pi p} \dfrac{\Gamma(\nu+1/2+ip)}{\Gamma(\nu+1/2-ip)} P_{\nu-1/2}^{-ip}(\cosh t_R) \\ P_{\nu-1/2}^{ip}(\cosh t_L) \end{cases}$$

These factors come from the requirement of analyticity of Euclidean hemisphere

The Euclidean vacuum (Bunch-Davies vacuum) is selected as the initial state.

Bogoliubov transformation and entangled state

The Fourier mode field operator is

Positive freq. mode in the past in the Euclidean vacuum

$$\phi(t) = a_\sigma \chi^\sigma + a_\sigma^\dagger \chi^{\sigma *} \qquad a_s |0\rangle_{\text{ED}} = 0 \qquad \left[a_\sigma, a_{\sigma'}^\dagger\right] = \delta_{\sigma\sigma'}$$

$\sigma = R, L$

$$j^R \sim P_{n-1/2}^{ip}(\cosh t_R)$$
$$j^L \sim P_{n-1/2}^{ip}(\cosh t_L)$$

Positive freq. mode in the past in each *R* or *L* vacuum

$$= b_q \varphi^q + b_q^\dagger \varphi^{q*} \qquad b_R |0\rangle_R = 0, \qquad b_L |0\rangle_L = 0, \qquad \left[b_q, b_{q'}^\dagger\right] = \delta_{qq'}$$

$q = R, L$

$q = R, L$

Then the operators $\left(a_\sigma, a_\sigma^\dagger\right)$ and $\left(b_q, b_q^\dagger\right)$ are related by a Bogoliubov transformation.

The Euclidean vacuum can be constructed from the *R*, *L* vacua as

$$|0\rangle_{\text{ED}} \propto \exp\left(\frac{1}{2} \sum_{i,j=R,L} m_{ij} b_i^\dagger b_j^\dagger\right) |0\rangle_R |0\rangle_L$$

: Entangled state of the $\mathcal{H}_R \otimes \mathcal{H}_L$ Hilbert space

Symmetric matrix $\begin{pmatrix} m_{RR} & m_{RL} \\ m_{LR} & m_{LL} \end{pmatrix}$ which should consist of the Bogoliubov coefficients

Bogoliubov coefficients

The condition $a_s |0\rangle_{\mathrm{ED}} = 0$ determines m_{ij}

Conformal invariance ($\nu = 1/2$)
Masslessness ($\nu = 3/2$)

Unimporatant phase factor

$$m_{ij} = e^{i\theta} \frac{\sqrt{2} e^{-p\pi}}{\sqrt{\cosh 2\pi p + \cos 2\pi\nu}} \begin{pmatrix} \cos \pi\nu & i \sinh p\pi \\ i \sinh p\pi & \cos \pi\nu \end{pmatrix}$$

The density matrix $\rho = |0\rangle_{\mathrm{ED}\;\mathrm{ED}}\langle 0|$ is not diagonal in the $|0\rangle_R |0\rangle_L$ basis.

It is difficult to trace out the degree of freedom in, say, the *L* space later in order to calculate the entanglement entropy.

We perform a further Bogoliubov transformation to get a diagonalized form.

Bogoliubov transformation 2

We perform a further Bogoliubov transformation in each *R* and *L* region

R region: $c_R = u\, b_R + v\, b_R^\dagger$

L region: $c_L = u^* b_L + v^* b_L^\dagger$

$$|u|^2 - |v|^2 = 1 \qquad \left[c_i, c_j^\dagger \right] = \delta_{ij}$$

This transformation does not mix the operators in $\mathcal{H}_R$ space and those in $\mathcal{H}_L$ space and thus does not affect the entangled state between $\mathcal{H}_R$ and $\mathcal{H}_L$.

And try to obtain the relation

$$\left|0\right\rangle_{\text{ED}} = N_{g_p}^{-1} \exp\left(g_p\, c_R^\dagger c_L^\dagger \right) \left|0\right\rangle_{R\text{¢}} \left|0\right\rangle_{L\text{¢}} \qquad N_{\gamma_p}^2 = \left| \exp\left(\gamma_p\, c_R^\dagger c_L^\dagger \right) \right|^2 = \left(1 - |\gamma_p|^2 \right)^{-1}$$

we want to know

Reduced density matrix

The consistency condition $c_R|0\rangle_{\rm ED} = \gamma_p c_L^\dagger |0\rangle_{\rm ED}$ and $c_L|0\rangle_{\rm ED} = \gamma_p c_R^\dagger |0\rangle_{\rm ED}$ determines

$$\gamma_p = i\frac{\sqrt{2}}{\sqrt{\cosh 2\pi p + \cos 2\pi\nu} + \sqrt{\cosh 2\pi p + \cos 2\pi\nu + 2}}$$

$$\xrightarrow{\nu = 1/2,\, 3/2} \propto e^{-\pi p}$$

Conformal invariance ($\nu = 1/2$)
Masslessness ($\nu = 3/2$)

Finally, the reduced density matrix after tracing out *L* region is found to be diagonalized as

$$\rho_R = \mathrm{Tr}_L |0\rangle_{\rm ED\ ED}\langle 0| = \left(1 - |\gamma_p|^2\right) \sum_{n=0}^{\infty} |\gamma_p|^{2n} |n; p\ell m\rangle\langle n; p\ell m|$$

$|\gamma_p|^2 \propto e^{-2\pi p}$, $|\gamma_p|^{2n} \propto e^{-2\pi p n}$

: Thermal state $T = \dfrac{H}{2\pi}$

$$|n; p\ell m\rangle = \frac{1}{\sqrt{n!}}\left(c_R^\dagger\right)^n |0\rangle_{R'}$$

: *n* particle excitation states

Thermal state: $\dfrac{1}{e^{\varepsilon/T} - 1}$

The de Sitter space has some peculiar property for the conformal and massless cases.

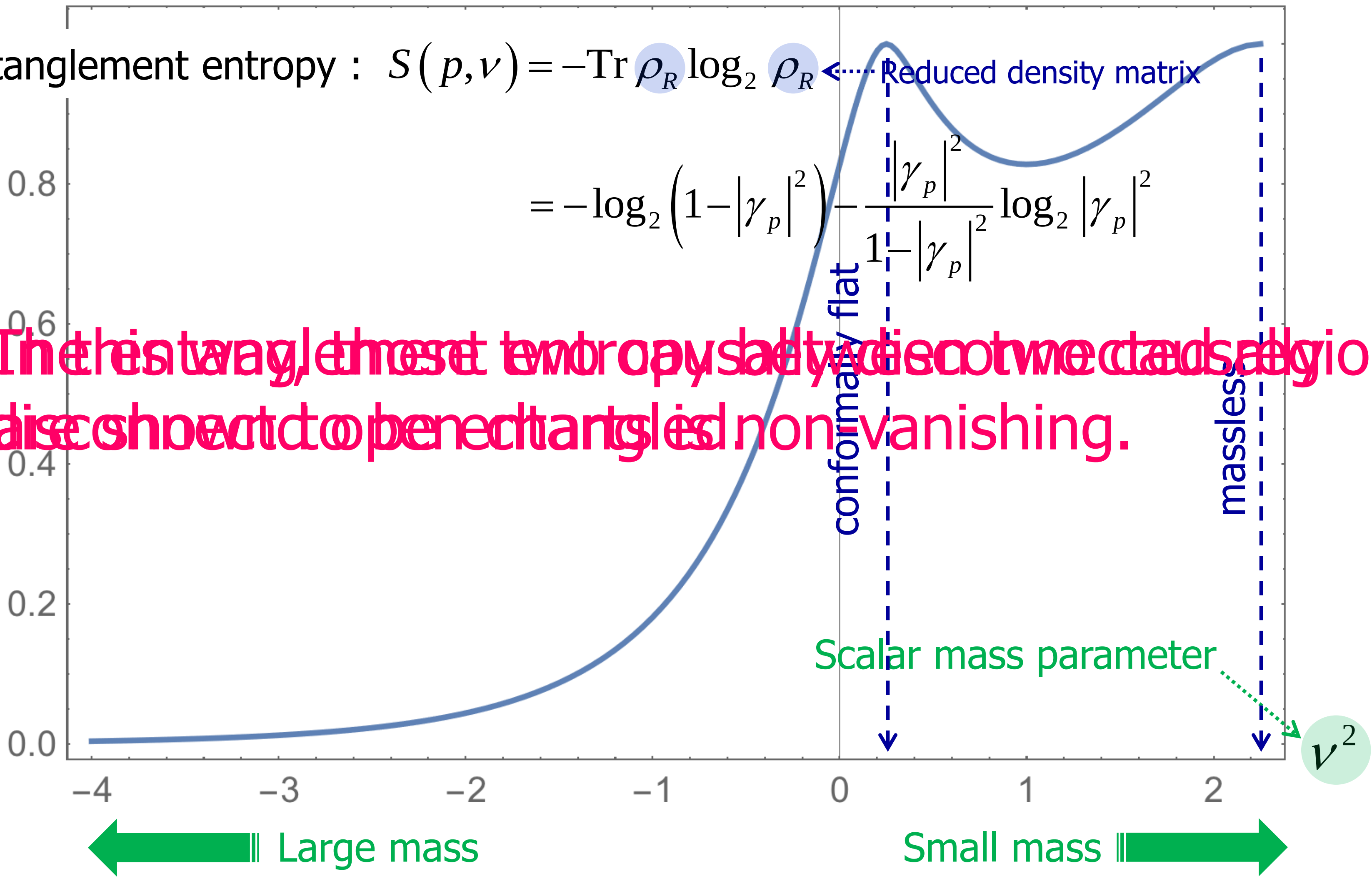
Entanglement entropy between *R* and *L* regions

Maldacena & Pimentel (2013)

They calculated the entanglement entropy between those *R* and *L* regions.

Entanglement entropy : $S(p,\nu) = -\mathrm{Tr}\,\rho_R \log_2 \rho_R$

$$= -\log_2\left(1-|\gamma_p|^2\right) - \frac{|\gamma_p|^2}{1-|\gamma_p|^2}\log_2|\gamma_p|^2$$



The entanglement entropy between those two causally disconnected regions is non-vanishing.

In this way, these two causally disconnected regions are shown to be entangled.

Now going back to the original question

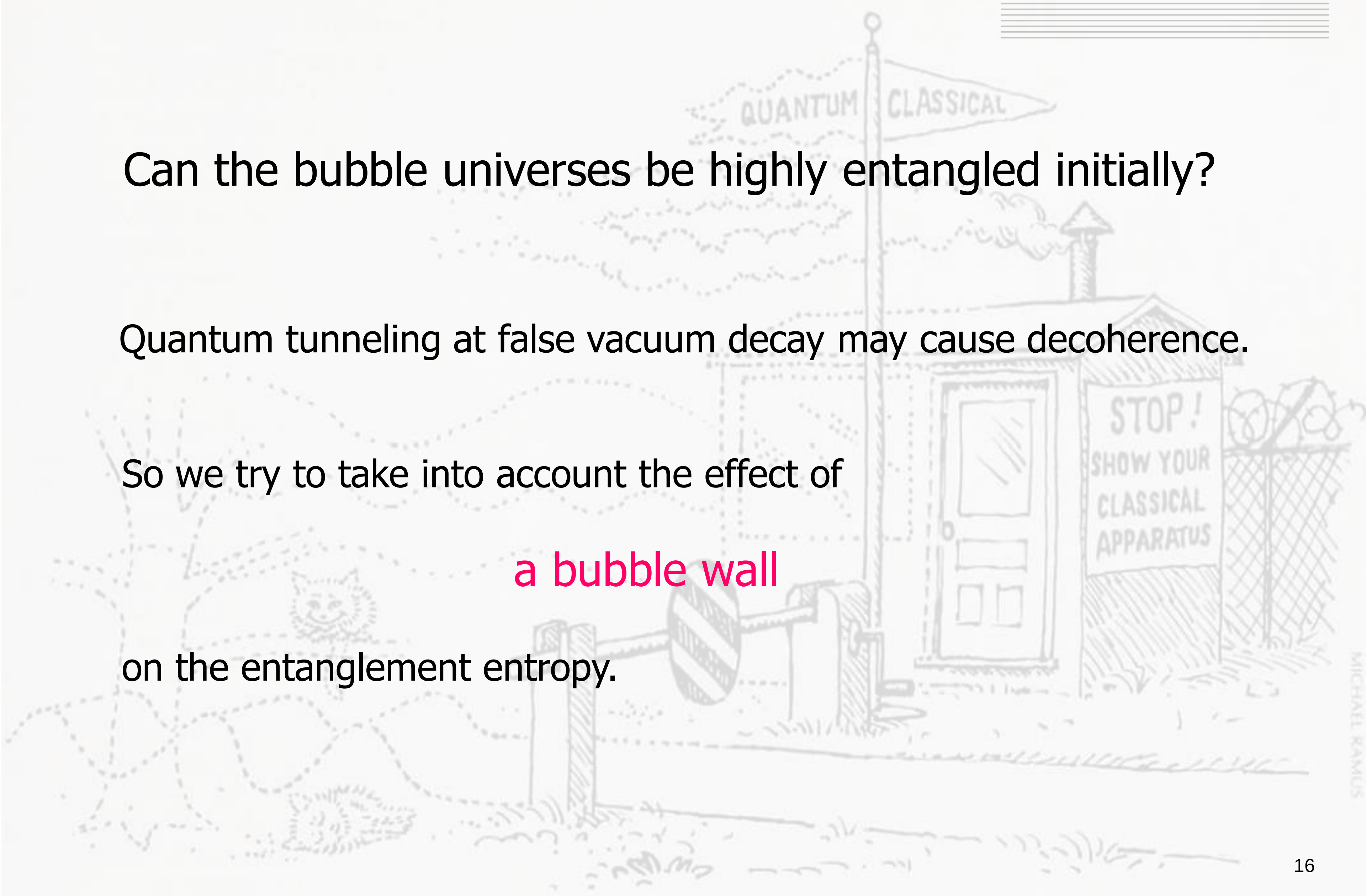
Can the bubble universes be highly entangled initially?

Quantum tunneling at false vacuum decay may cause decoherence.

So we try to take into account the effect of

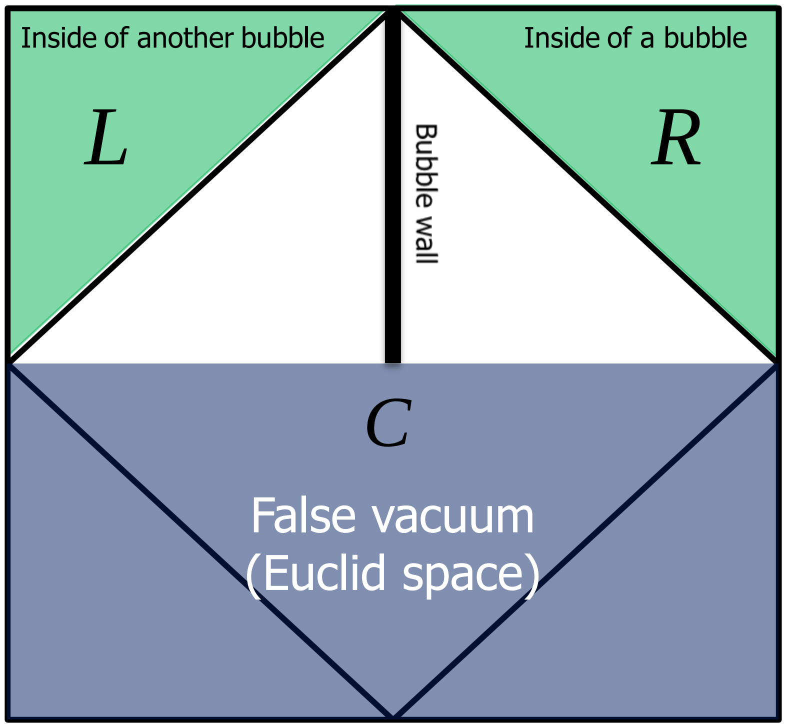
a bubble wall

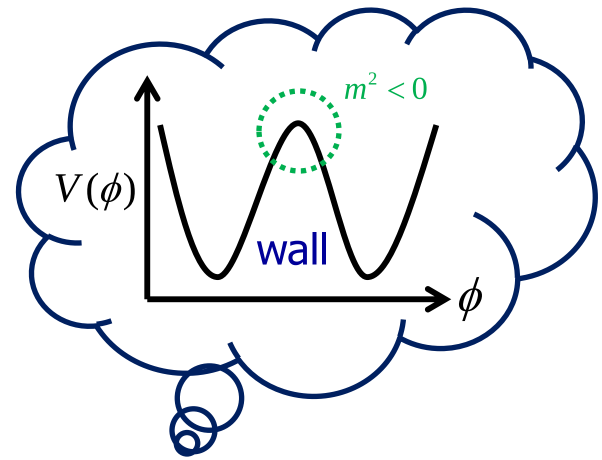
on the entanglement entropy.



Our setup

We assume there is a delta-functional wall between two open charts *R* and *L*.





A delta-functional wall of hight (depth) Λ

Action

This can be thought of as a model of pair creation of identical bubble universes separated by a bubble wall.

$$S = \int d^4x \sqrt{-g} \left[-\frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - \frac{m^2 - \Lambda \delta(t_C)}{2} \phi^2 \right]$$

The ED vacuum solutions in the presence of a wall

The positive frequency mode functions for the Euclidean vacuum in the presence of the bubble wall

These factors come from the requirement of analyticity of Euclidean hemisphere

$$\chi_p^R(t) = \begin{cases} P_{\nu-1/2}^{ip}(\cosh t_R) \\ \left(A_p C_p + B_p D_{-p}\right) P_{\nu-1/2}^{ip}(\cosh t_L) + e^{\pi p}\left(A_p D_p + B_p C_{-p}\right) P_{\nu-1/2}^{-ip}(\cosh t_L) \end{cases}$$

$$\chi_p^L(t) = \begin{cases} \left(A_p C_p + B_p D_{-p}\right) P_{\nu-1/2}^{ip}(\cosh t_R) + e^{\pi p}\left(A_p D_p + B_p C_{-p}\right) P_{\nu-1/2}^{-ip}(\cosh t_R) \\ P_{\nu-1/2}^{ip}(\cosh t_L) \end{cases}$$

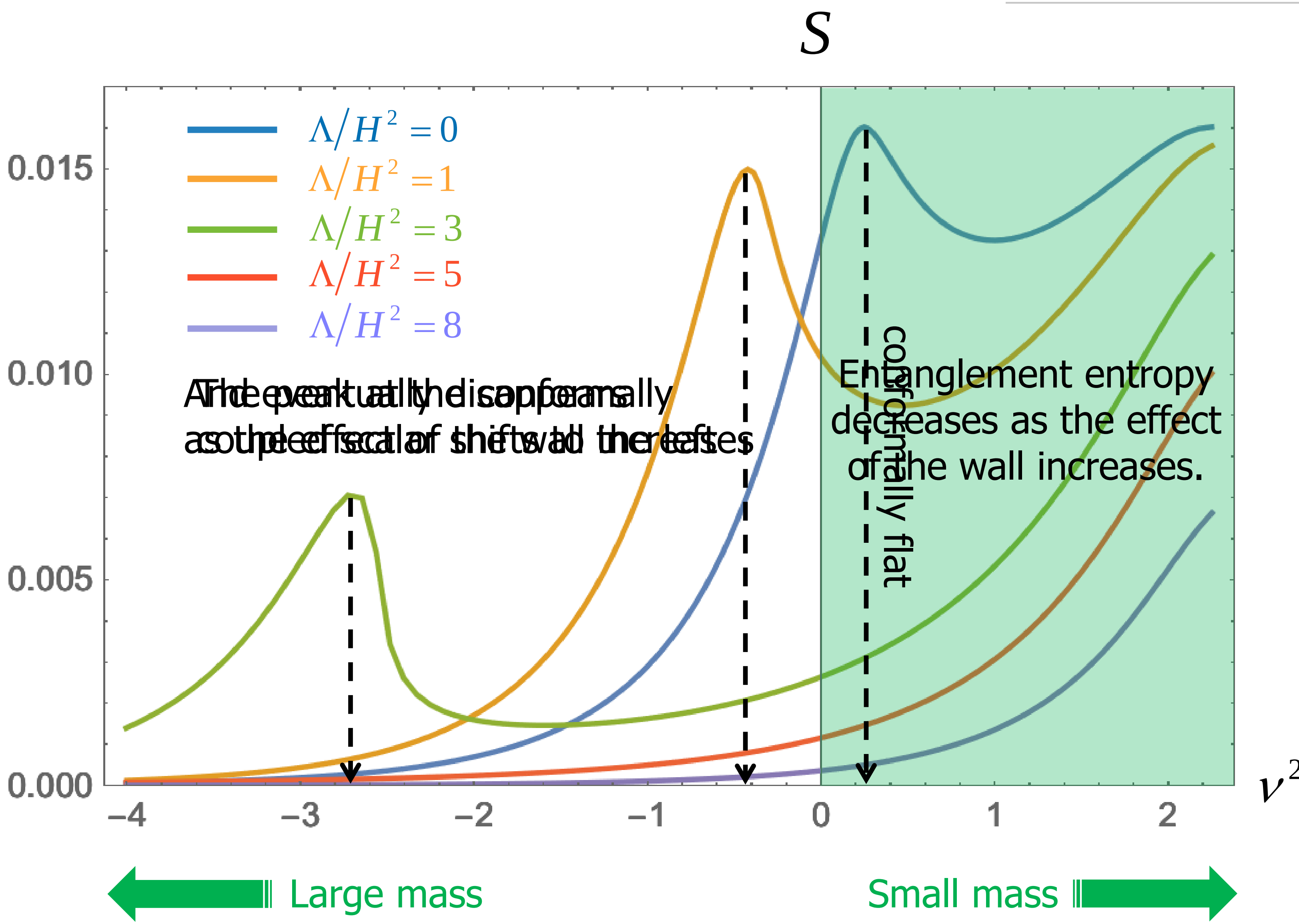
Wall effect

$$A_p = 1 + \frac{\pi}{2i\sinh \pi p}\frac{\Lambda}{H^2} P_{\nu-1/2}^{ip}(0)\, P_{\nu-1/2}^{-ip}(0) \qquad C_p = \frac{\cos \pi\nu}{i\sinh \pi p}$$

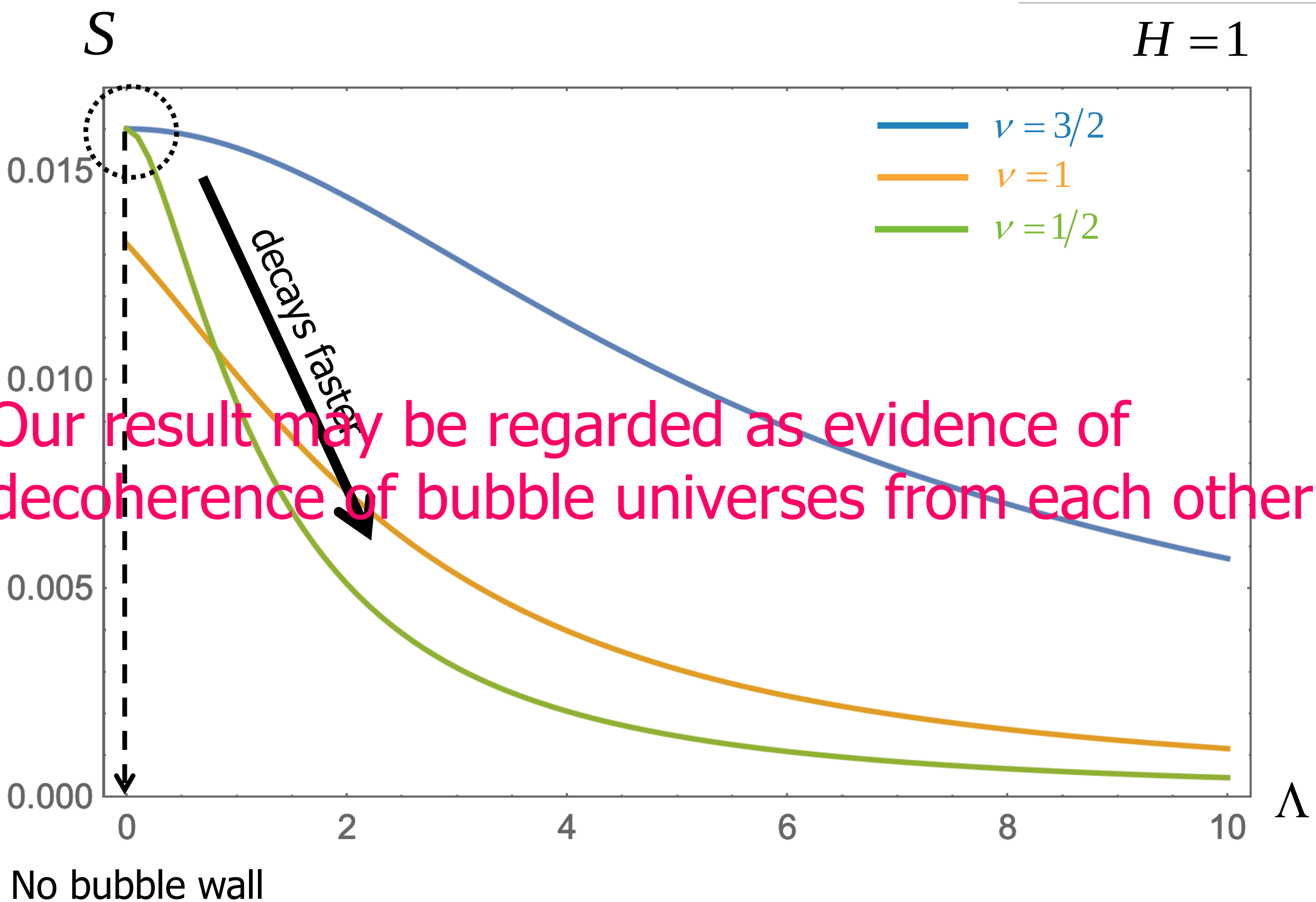
$$B_p = -\frac{\pi}{2i\sinh \pi p}\frac{\Lambda}{H^2}\left(P_{\nu-1/2}^{ip}(0)\right)^2 \qquad D_p = -e^{-2\pi p}\frac{\cos\left(\nu + ip\right)\pi}{i\sinh \pi p}\frac{\Gamma\left(1/2+\nu+ip\right)}{\Gamma\left(1/2+\nu-ip\right)}$$

We can expect the effect of the wall would appear in the entanglement entropy.

Entanglement entropy between R and L regions



Wall dependence of the entanglement entropy



Our result may be regarded as evidence of decoherence of bubble universes from each other !?

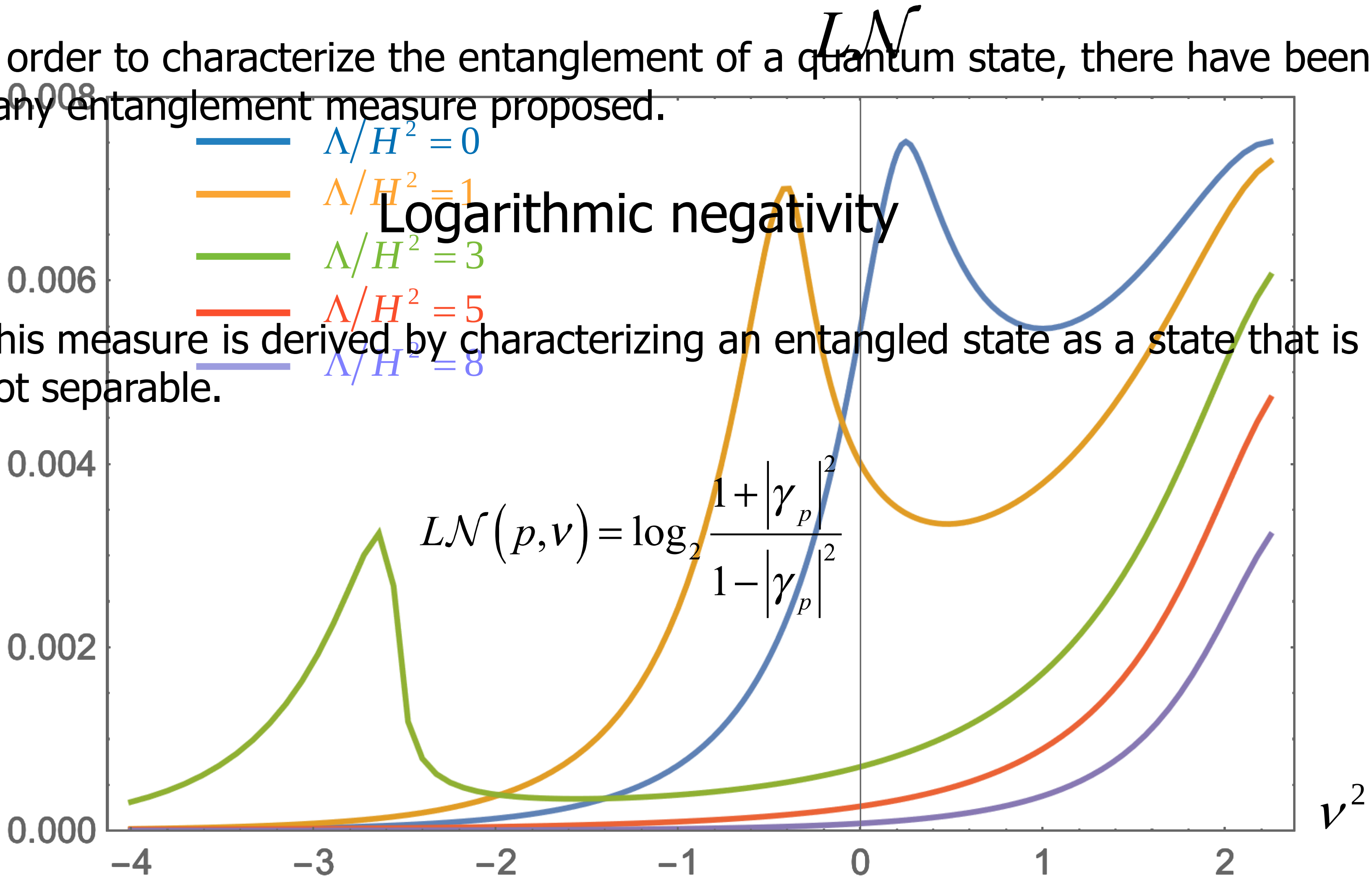
Logarithmic negativity between *R* and *L* regions

In order to characterize the entanglement of a quantum state, there have been many entanglement measure proposed.

Logarithmic negativity

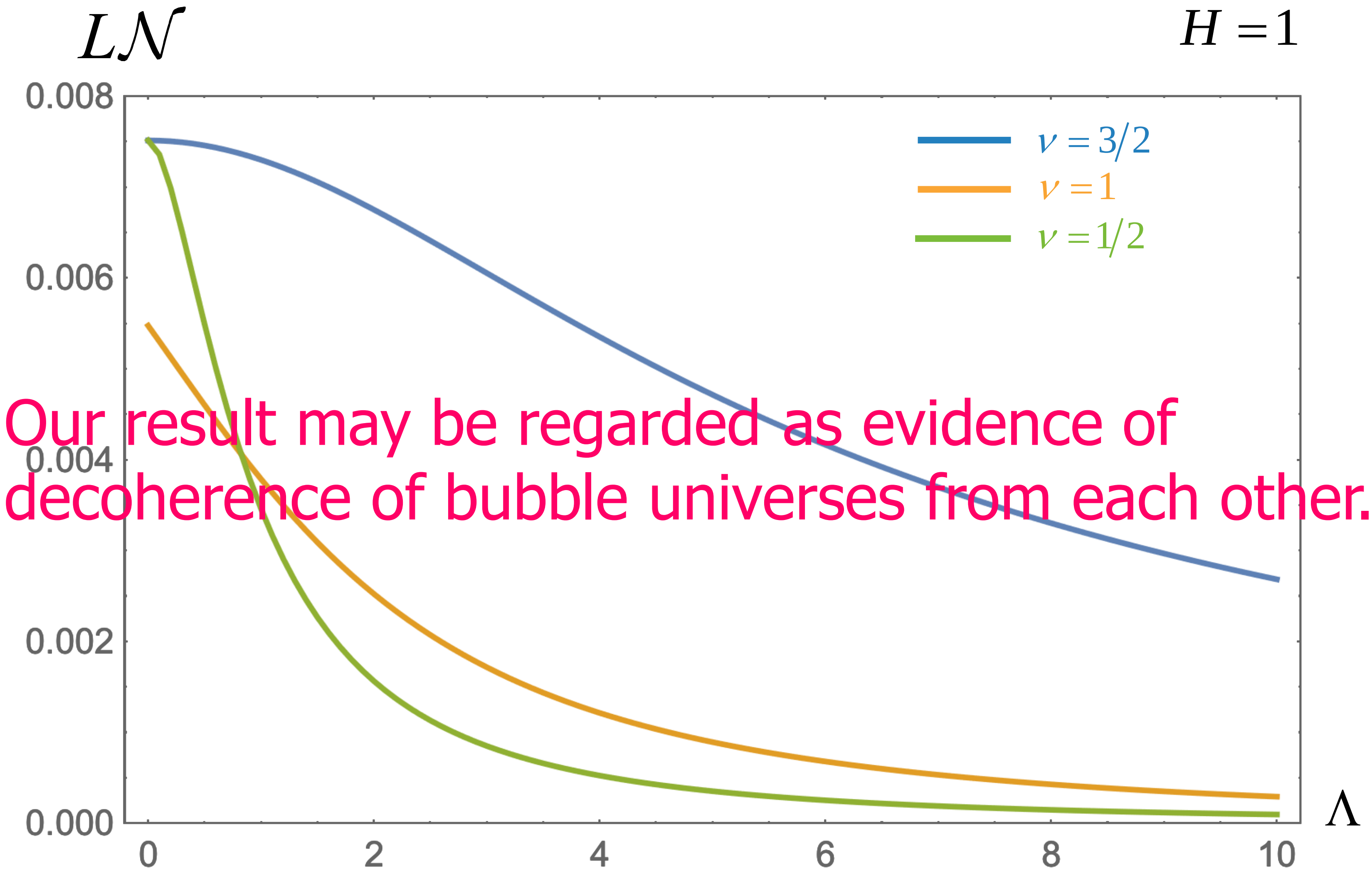
This measure is derived by characterizing an entangled state as a state that is not separable.

$$\mathcal{LN}(p,\nu)=\log_2\frac{1+|\gamma_p|^2}{1-|\gamma_p|^2}$$



Qualitative features are the same as the result of entanglement entropy.

Wall dependence of the logarithmic negativity



Qualitative features are the same as the result of entanglement entropy.

Summary

We studied the effect of a bubble wall on the quantum entanglement of a free massive scalar field between two causally disconnected open charts in de Sitter space.

We assumed there is a delta-functional wall between them.

Our model may be regarded as a model describing the pair creation of identical bubble universes separated by a bubble wall.

We computed the entanglement entropy and logarithmic negativity of the scalar field and compared the result with the case of no bubble wall.

We found that larger the wall leads to less entanglement.

Our result may be regarded as evidence of decoherence of bubble universes from each other.