

CONFORMAL SYMMETRY IN STANDARD MODEL AND GRAVITY

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Stefano Lucat and T. Prokopec, arXiv:1705.00889 [gr-qc]; 1709.00330 [gr-qc];1606.02677 [hep-th]
T. Prokopec, Leonardo da Rocha, Michael Schmidt, Bogumila Swiezewskae-Print: arXiv:1801.05258 [hep-ph] +

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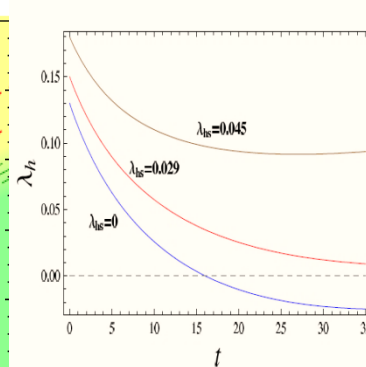
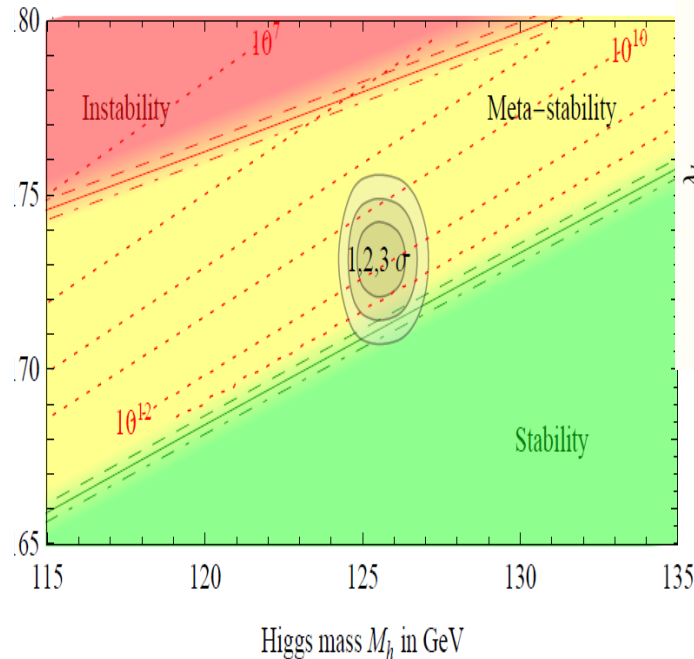
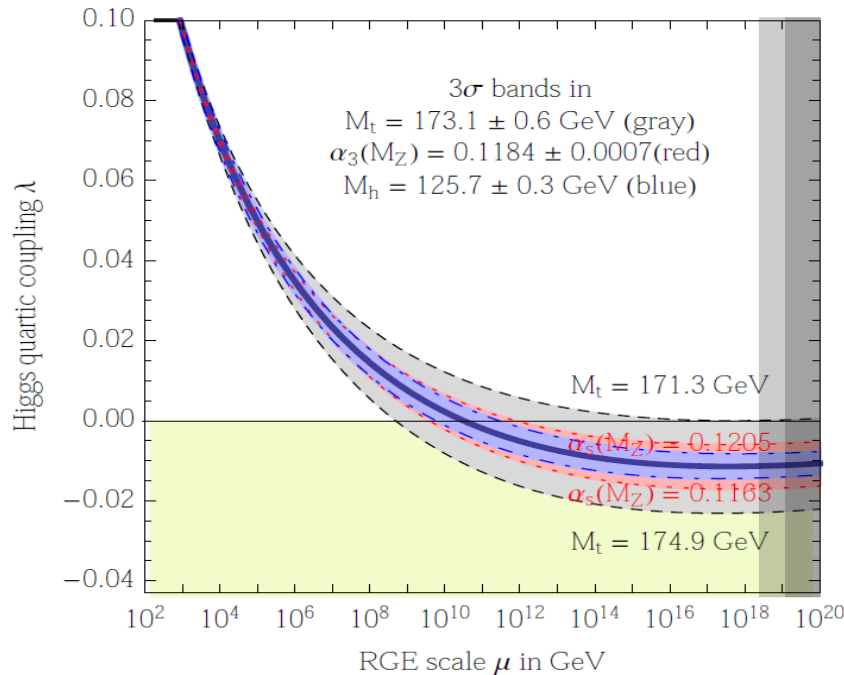
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PHYSICAL MOTIVATION

PHYSICAL MOTIVATION

- AT LARGE ENERGIES THE STANDARD MODEL IS ALMOST CONFORMALLY INVARIANT.
- HIGGS MASS AND KINETIC TERMS BREAK THE SYMMETRY
- OBSERVED HIGGS MASS: $m_H = 125.3\text{GeV}$ is close to the stability bound
- STABILITY BOUND: $m_H \approx 130\text{GeV}$: CAN BE ATTAINED BY ADDING SCALAR

Oleg Lebedev, e-Print: arXiv:1203.0156 [hep-ph]



THEORETICAL MOTIVATION

THEORETICAL MOTIVATION IN SM⁶

- **HIGGS MASS TERM** RESPONSIBLE FOR GAUGE HIERARCHY PROBLEM
- IF WE COULD FORBID IT BY SYMMETRY, THE GHP WOULD BE SOLVED
- THIS SYMMETRY COULD BE **WEYL SYMMETRY** IMPOSED CLASSICALLY
- HIGGS MASS COULD BE GENERATED DYNAMICALLY BY THE COLEMAN-WEINBERG (CW) MECHANISM

THEORETICAL MOTIVATION IN GRAVITY

- THE SYMMETRY IS BROKEN BY THE NEWTON CONSTANT AND COSMOLOGICAL TERM, G & Λ .
- G & Λ ARE RESPONSIBLE FOR GRAVITATIONAL HIERARCHY PROBLEM.
- SCALAR DILATON & CARTAN TORSION CAN RESTORE WEYL SYMMETRY IN CLASSICAL GRAVITY.
- G & Λ CAN BE GENERATED BY DILATON CONDENSATION INDUCED BY QUANTUM EFFECTS akin to THE COLEMAN-WEINBERG MECHANISM.
- IF GRAVITY IS CONFORMAL IN UV, IT MAY BE FREE OF SINGULARITIES (BOTH COSMOLOGICAL AND BLACK HOLE).

WEYL SYMMETRY IN CLASSICAL GRAVITY

CARTAN EINSTEIN THEORY

- POSITS THAT FERMIONS (& SCALARS) SOURCE SPACETIME TORSION.
- TORSION IS CLASSICALLY A CONSTRAINT FIELD (NOT DYNAMICAL, DOES NOT PROPAGATE)
⇒ CARTAN EQUATION CAN BE INTEGRATED OUT, RESULTING IN THE KIBBLE-SCIAMA THEORY

Lucat, Prokopec, e-Print: arXiv:1512.06074 [gr-qc]

⇒ THIS THEORY PROVIDES ADDITIONAL SOURCE TO STRESS-ENERGY, WHICH CAN CHANGE BIG-BANG SINGULARITY TO A BOUNCE.

- CARTAN-EINSTEIN THEORY CAN BE MADE CLASSICALLY CONFORMAL!

Lucat & Prokopec, arxiv:1606.02677 [hep-th]

CLASSICAL WEYL SYMMETRY

°10°

- WEYL TRANSFORMATION ON THE METRIC TENSOR

$$g_{\mu\nu} \rightarrow \tilde{g}_{\mu\nu} = e^{2\theta(x)} g_{\mu\nu} \quad d\tau \rightarrow d\tilde{\tau} = e^{\theta(x)} d\tau$$

- GENERAL CONNECTION Γ , TORSION TENSOR T , CHRISTOFFEL CON $\overset{\circ}{\Gamma}$

$$\Gamma^\lambda_{\mu\nu} = T^\lambda_{\mu\nu} + T_{\mu\nu}{}^\lambda + T_{\nu\mu}{}^\lambda + \overset{\circ}{\Gamma}^\lambda_{\mu\nu}$$

$$\delta\Gamma^\mu_{\alpha\beta} = \delta^\mu_{(\alpha} \partial_{\beta)} \theta, \text{ ASSUME: } \delta\Gamma^\mu_{\alpha\beta} = \delta^\mu_{\alpha} \partial_{\beta} \theta \Rightarrow \delta T^\mu_{\alpha\beta} = \delta^\mu_{[\alpha} \partial_{\beta]} \theta$$

- RIEMANN TENSOR IS INVARIANT: $\delta R^\alpha_{\beta\gamma\delta} = 0$

- THIS IMPLIES THAT THE VACUUM EINSTEIN EQUATION IS WEYL INV:

$$G_{\mu\nu} = 0, \quad \delta G_{\mu\nu} = 0$$

GEOMETRIC VIEW OF TORSION

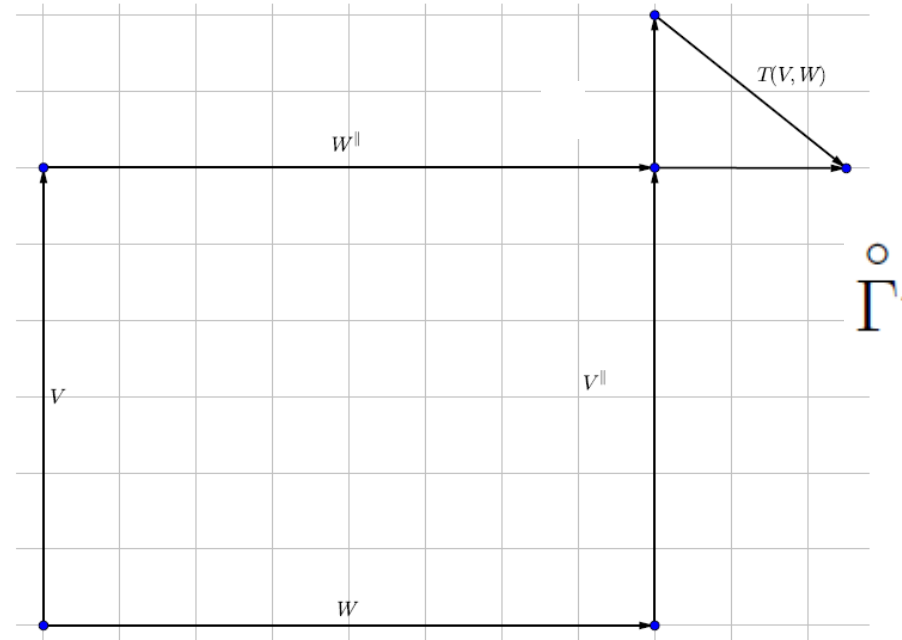
11°

- (VECTORIAL) TORSION TRACE 1-FORM:

$$\mathcal{T} \equiv \mathcal{T}_\mu dx^\mu = \frac{2}{D-1} T^\lambda{}_{\lambda\mu} dx^\mu$$

- TRANSFORMS AS A VECTOR FIELD:

$$\mathcal{T} \rightarrow \mathcal{T} + d\theta$$



- WHEN A VECTOR IS PARALLEL-TRANSPORTED, TORSION TRACE INDUCES A LENGTH CHANGE: **CRUCIAL** IN WHAT FOLLOWS

PARALLEL TRANSPORT AND JACOBI EQUATION

- GEODESIC EQUATION:

$$\nabla_{\dot{\gamma}} \frac{dx^\mu}{d\tau} \equiv \frac{dx^\lambda}{d\tau} \nabla_\lambda \frac{dx^\mu}{d\tau} = 0$$

$$\Gamma^\lambda_{\mu\nu} = T^\lambda_{\mu\nu} + T_{\mu\nu}^\lambda + T_{\nu\mu}^\lambda + \overset{\circ}{\Gamma}^\lambda_{\mu\nu}$$

$$\overset{\circ}{\Gamma} = \text{LEVI-CIVITA}$$

→ TRANSFORMS MULTIPLICATIVELY (as $1/d\tau^2$)

$$T[X, Y] = -\frac{1}{2}(\nabla_X Y - \nabla_Y X - [X, Y])$$

$$T^\lambda_{\mu\nu} = \Gamma^\lambda_{[\mu\nu]} = \frac{1}{2}(\Gamma^\lambda_{\mu\nu} - \Gamma^\lambda_{\nu\mu})$$

$$\nabla_{\dot{\gamma}} \frac{dx^\mu}{d\tau} = 0 \Rightarrow e^{-2\theta(x)} \nabla_{\dot{\gamma}} \frac{dx^\mu}{d\tau} = 0$$

NB: TRANSFORMATION OF $d\tau$ COMPENSATED BY TRANSFORMATION OF Γ !

- JACOBI EQUATION (JACOBI FIELDS $J \perp \dot{\gamma}$) AND RAYCHAUDHURI EQ:

$$\nabla_{\dot{\gamma}} \nabla_{\dot{\gamma}} J + 2\nabla_{\dot{\gamma}} T[\dot{\gamma}, J] = R[\dot{\gamma}, J]\dot{\gamma}$$

→ ALSO TRANSFORMS MULTIPLICATIVELY (as $1/d\tau^2$) UNDER WEYL TR

- SUGGESTS TO DEFINE A GAUGE INVARIANT PROPER TIME:

$$(d\tau)_{g.i.} = \exp\left(-\int_{x_0}^x T_\mu dx^\mu\right) d\tau := \text{PHYSICAL TIME OF COMOVING OBSERVERS!}$$

CONFORMAL SYMMETRY AND OBSERVATIONS

CONFRONTING OBSERVATIONS

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$$\Gamma_{EFF} \supset - \int d^4x \sqrt{-g} \{ [\alpha(\mu) \bar{R}^2 + \beta(\mu) \phi^2 R] + \gamma(\mu) T_{\alpha\beta} T^{\alpha\beta} \}$$

► $T_{\alpha\beta}$ = TORSION (TRACE) FIELD STRENGTH

EARLY COSMOLOGY

- INFLATIONARY MODELS GENERATED BY CONDENSATION OF SCALARON, DILATON OR TORSION TRACE MAY HAVE SPECIFIC FEATURES.
- PRELIMINARY RESULTS: CAN GET (quasi)de SITTER UNIVERSE AND NEARLY SCALE INVARIANT SCALAR SPECTRUM.
- STRONG 1st ORDER EW PT \Rightarrow GW PRODUCTION & BARYOGENESIS

J. REZACEK, B. SWIEZEWSKA and T. PROKOPEC, in progress

LATE COSMOLOGY

- CAN BE TESTED BY STUDYING e.g. DARK ENERGY AND STRUCTURE FORMATION, POSSIBLY DARK MATTER CANDIDATE.
- TORSION TRACE (AND MIXED TORSION) CAN BE DETECTED BY CONVENTIONAL GRAVITATIONAL WAVE DETECTORS

GRAVITATIONAL DETECTORS

GRAVITATIONAL WAVES

- GRAVITATIONAL WAVES

$$\frac{d^2 J^i}{dt^2} = \frac{1}{2} \ddot{h}_{ij}(t, \vec{x}) J^j$$

Plus polarization: $h_{xx} = -h_{yy} = h_+ \cos(\omega t - kz)$

$$J^x(t, z) = J_{(0)}^x \left[1 + (h_+/2) \cos(\omega t - kz) \right]$$

Cross polarization: $h_{xy} = h_{yx} = h_\times \cos(\omega t - kz)$

$$J^x(t, z) = J_{(0)}^x + (h_\times/2) J_{(0)}^y \cos(\omega t - kz) \Big]$$

DETECTORS FOR TORSION WAVES

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► **GW INTERFEROMETERS such as aLIGO/VIRGO**

● TORSION TRACE

$$\ddot{j}^i = J^0 \dot{\mathcal{T}}^i + J^j \partial_j \mathcal{T}^i \quad \mathcal{T}^i = \mathcal{T}_{(0)}^i \cos(\omega t - kz)$$

► LONGITUDINAL $\mathcal{T}_{(0),L}^i = \delta_z^i \frac{\omega}{m}$, $\mathcal{T}_{(0),L}^0 = -\frac{\|\vec{k}\|}{m}$

○ DETECTOR RESPONSE

$$\Delta J_{(0)}^z = 0, \quad \Delta J_{(0)}^{x,y} = -\frac{c^2 k}{\omega^2} \mathcal{T}_{(0),T}^{x,y} J_{(0)}^z \approx -\frac{c}{m} \mathcal{T}_{(0),T}^{x,y} J_{(0)}^z$$

► TRANSVERSE $\mathcal{T}_{(0),T}^i = \frac{1}{\sqrt{2}} (\delta_x^i \pm \delta_y^i)$, $\mathcal{T}_{(0),T}^0 = 0$

○ DETECTOR RESPONSE

$$\Delta J_{(0)}^z = -\frac{c^2 k}{\omega^2} \mathcal{T}_{(0),L}^z J_{(0)}^z \approx -\frac{c}{\omega} \mathcal{T}_{(0),L}^z J_{(0)}^z, \quad \Delta J_{(0)}^{x,y} = 0.$$

● GRAVITATIONAL WAVES vs TORSION WAVES: a comparison

- PHASE SHIFT $\frac{1}{4}$ PERIOD
- FREQUENCY DEPENDENCE
- TORSION TRACE (L) COUPLES TO

TORSION SOURCES

- E.G.: TORSION TRACE: LONGITUDINAL MODE $\mathcal{T}_\mu = \partial_\mu \theta$
 - ▶ ITS MASS IS PROTECTED BY THE CONFORMAL WARD=TAKAHASHI, (see talk of Stefano Lucat)

$$\square \theta = \frac{8\pi G_N}{c^4} \frac{T_\mu^\mu}{6}, \quad \square h_{ij} = \frac{8\pi G_N}{c^4} T_{ij}$$

- ▶ THIS IMPLIES ABOUT 1 order of magnitude suppression when compared with the amplitude of gravitational waves, i.e.

$$\frac{\theta}{h_{ij}} \sim \frac{e^2}{2}$$

- e=sources excentricity (can be 0.5)

- ▶ DETECTABLE BY THE NEXT GENERATION OF OBSERVATORIES such as **EINSTEIN TELESCOPE**.

CONCLUSIONS AND OUTLOOK

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- CHALLENGE 1: USE FRG METHODS TO STUDY HOW THIS THEORY DIFFERS FROM THE USUAL GRAVITY, i.e. WHETHER IT IS ASYMPTOTICALLY SAFE/ADMITS UV COMPLETION.
- CHALLENGE 2: CONFRONT THIS THEORY AS MUCH AS POSSIBLE WITH OBSERVATIONS
- CHALLENGE 3: CAN WE GET RID OF (COSMOLOGICAL AND BLACK HOLE) SINGULARITIES?

HINT: RECALL: $(d\tau)_{g.i.} = \exp\left(-\int_{x_0}^x T_\mu dx^\mu\right) d\tau :=$ PHYSICAL TIME OF COMOVING OBSERVERS