

COMPOSITE hybrid INFLATION

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COMPOSITE hybrid INFLATION

G. Cacciapaglia, D. Y. Cheong, A. Deandrea, W. Isnard and SCP
JCAP 10 (2023) 063 [2307.01852]

Dhong Yeon



Wanda



Lagrangian of Universe



Standard Model

new physics

$$\mathcal{L}_{\text{universe}} = \sqrt{g} [\kappa R + \Lambda_{CC} + \mathcal{L}_{SM} + \Delta \mathcal{L}]$$

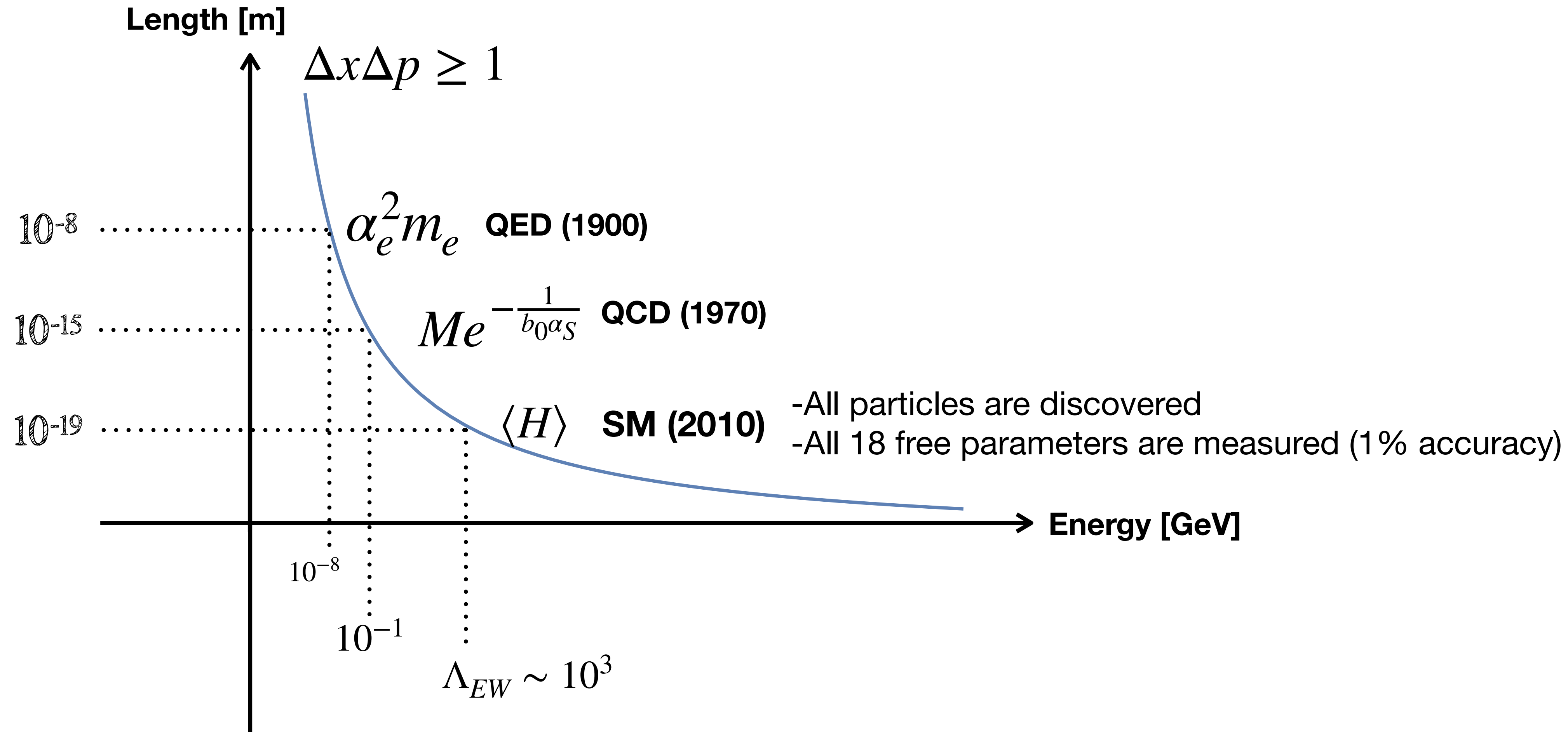
Dark energy

Gravity

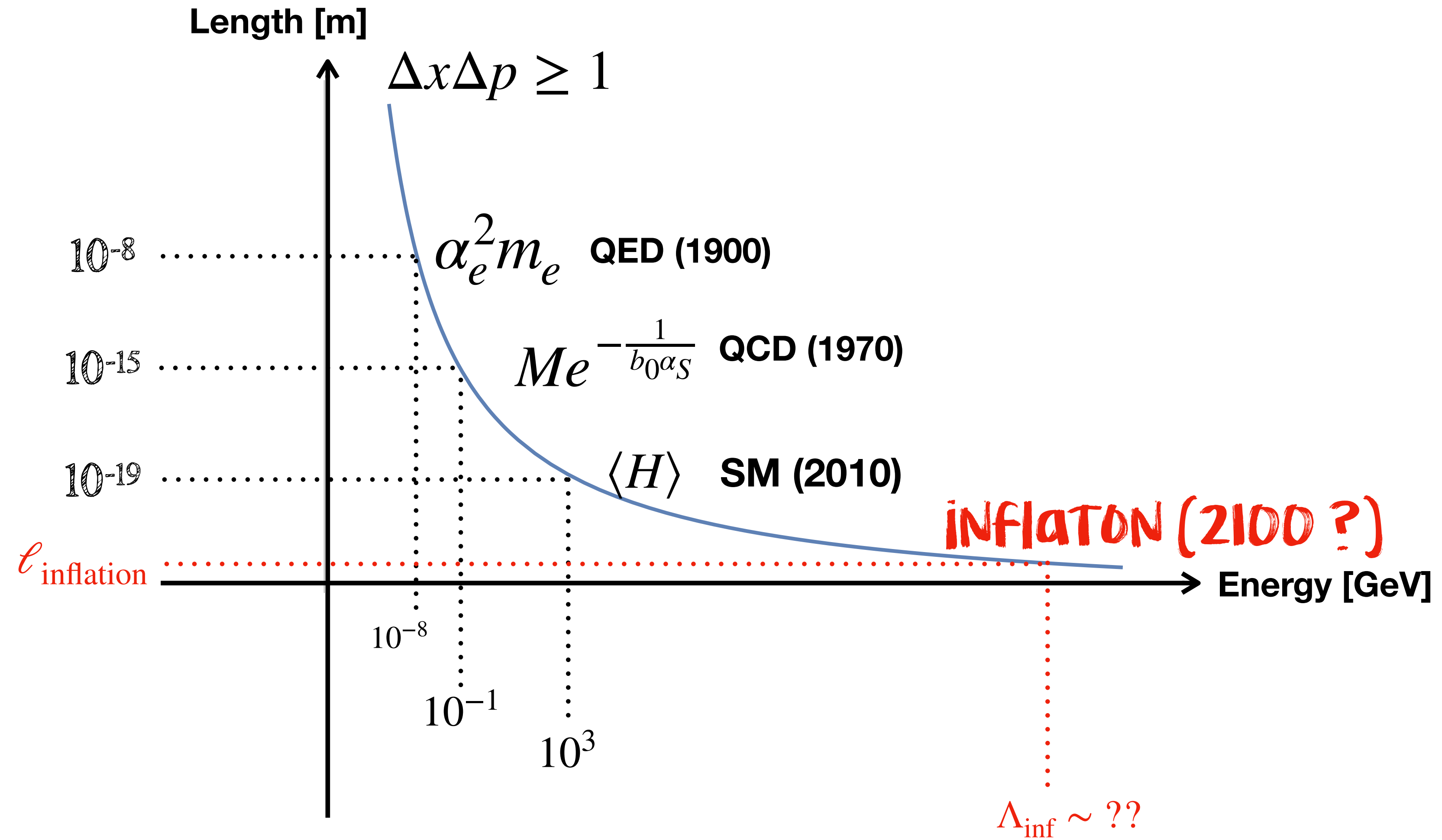
Team suit of Yonsei-HEP-COSMO group

<http://hepcosmo.yonsei.ac.kr>

Scales vs Physics



Scales vs Physics



Inflationary scale seems to be high

$$\Lambda_{inf} \gg \Lambda_{EW} \sim \text{TeV}$$

because **no inflaton** was found at the LHC

But, what's the inflaton?

How does she look like?

PDG BOOK V.2023

Citation: R.L. Workman *et al.* (Particle Data Group), Prog.Theor.Exp.Phys. **2022**, 083C01 (2022) and 2023 update

NON-COLORED

LEPTONS

LORENTZ SYMMETRY

e

$$J = \frac{1}{2}$$

ZERO if CHIRAL SYMMETRY

ZERO THANKS TO CPT SYMMETRY

$$\text{Mass } m = (548.579909065 \pm 0.000000016) \times 10^{-6} \text{ u}$$

$$\text{Mass } m = 0.51099895000 \pm 0.00000000015 \text{ MeV}$$

$$|m_{e^+} - m_{e^-}|/m < 8 \times 10^{-9}, \text{ CL} = 90\%$$

$$|q_{e^+} + q_{e^-}|/e < 4 \times 10^{-8}$$

Magnetic moment anomaly

$$(g-2)/2 = (1159.65218062 \pm 0.00000012) \times 10^{-6}$$

$$(g_{e^+} - g_{e^-}) / g_{\text{average}} = (-0.5 \pm 2.1) \times 10^{-12}$$

ZERO if CP SYMMETRY

$$\text{Electric dipole moment } d < 0.11 \times 10^{-28} \text{ e cm, CL} = 90\%$$

$$\text{Mean life } \tau > 6.6 \times 10^{28} \text{ yr, CL} = 90\% \text{ [a]}$$

ABSOLUTE STABLE BY GAUGE SYMMETRY

PDG BOOK V. 2100

Citation : SC Park et. al. Particle Data Group (2100 update)

INFLATON

ϕ_{inf}

Spin=0 ??

Mass =??

Charge= 0 ??

Lifetime =??

Elementary??

LORENTZ
SYMMETRY

NON-COLORED?

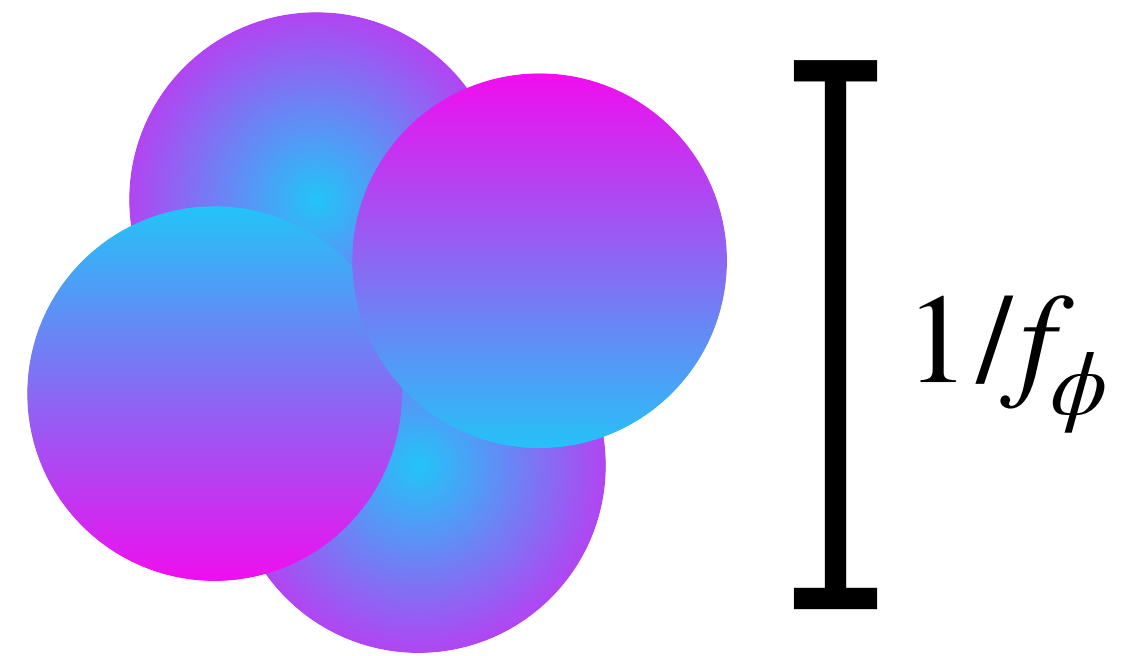
NOT STABLE
FOR REHEATING

Spin of inflaton

- Scalar ($s=0$) is favored due to isotropy
- $s=1/2$ condensate $\langle \psi \rangle$ generates Lorentz violation
- $s=1$ models tried [Golovnev, Mukhanov, Vanchurin (08)], [Natuko, Eiichiro, Yamaguchi (14)], but disfavored because of anisotropy. $\langle V^i \rangle$ breaks Lorentz.
- maybe in extra dimensions $A^M = (A^\mu, A^5)$ [Arkani-Hamed, Cheng, Creminelli, Randall (03)]

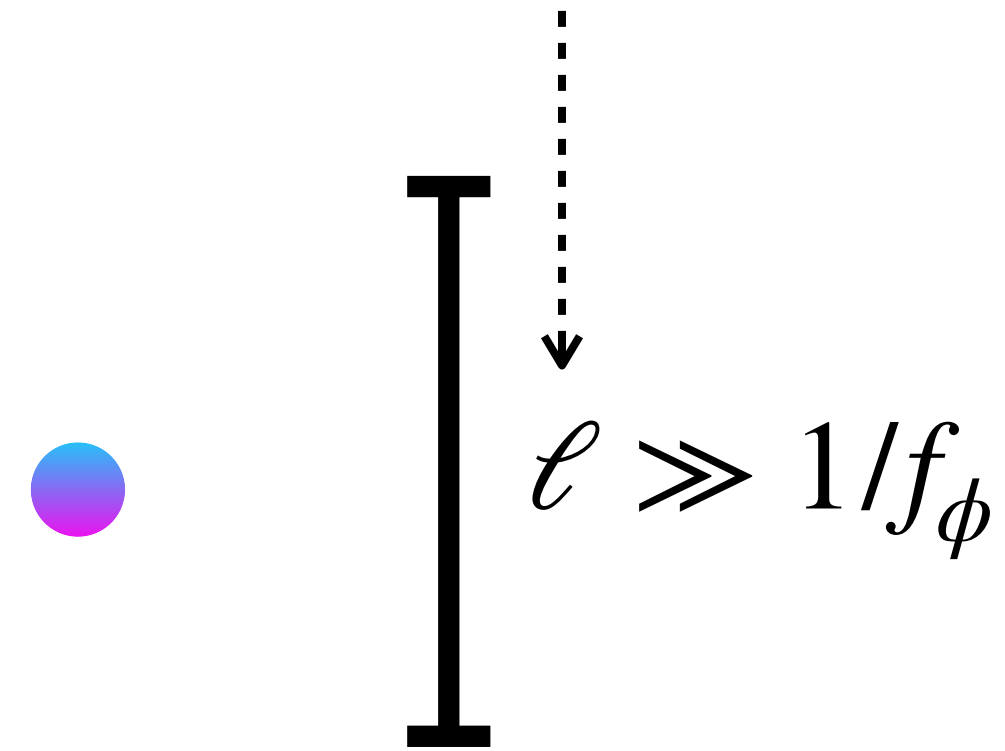
Composite or elementary?

- It is not certain whether inflaton is composite or elementary.
- Situation is similar to the SM Higgs: The Higgs could be a composite state. Technicolor in 4D (or RS in 5D) is regarded natural as
$$m_{scalar}^2 \sim \Lambda_{composite}^2 \ll \Lambda_{cut}^2$$
- Strong dynamics \implies organizing principle for multiple “composite states”
- (ex) pions, dilatons...+other mesons



\approx

Length of experiment



$$f_\phi = \frac{\Lambda_{composite}}{2\pi}$$

Looks elementary

This work

[G. Cacciapaglia, D. Y. Cheong, A. Deandrea, W. Isnard and SCP [2307.01852] *JCAP* 10 (2023) 063].

- We propose a composite inflaton model based on $SU(N_c)$ gauge theory coupled to N_f Dirac fields of fundamental representation.
- But, $SO(N)$, $SP(N)$, G_2 also work
- For a certain set of (N_c, N_f) , the theory will flow to an IR fixed point, hence generates a phase of nearly scale invariant dynamics (Walking regime) e.g.

[Dietrich and F. Sannino (06)], [F. Sannino (07)]

walking = slow-rolling

This work

[G. Cacciapaglia, D. Y. Cheong, A. Deandrea, W. Isnard and S. C. Park [2307.01852] *JCAP* 10 (2023) 063].

✓ Condensation $\langle \psi \bar{\psi} \rangle$ or $\langle GG \rangle$ generates spontaneous breaking for chiral and scale symmetry

✓ PNGBs (they are light!) are the pions $\phi(x) = \phi^a(x) T^a$ (chiral symmetry breaking) and dilaton $\chi(x)$ (scale symmetry breaking)

✓ Technically, it is convenient to use non-linear realization of pion by

$$U(x) = e^{i\phi^a T^a / f_\phi}.$$

The composite inflation model

$$\eta = (-1, +1, +1, +1), M_P = 1/\sqrt{8\pi G} = 1 \quad U(x) = e^{i\phi^a T^a / f_\phi}$$

$$(f_\phi, \lambda_\chi, f_\chi, \delta_1, \delta_2)$$

Order param. of Scale anomaly

dilaton kinetic term

pion kinetic term

$$A \sim 1/4 \text{ for } \langle \chi \rangle \doteq f_\chi$$

$$\begin{aligned} \frac{\mathcal{L}}{\sqrt{-g}} \supset & \frac{1}{2}R - \frac{1}{2}\partial_\mu\chi\partial^\mu\chi - \frac{f_\phi^2}{2}\left(\frac{\chi}{f_\chi}\right)^2 \text{Tr}[\partial_\mu U^\dagger\partial^\mu U] - \frac{\lambda_\chi}{4}\chi^4\left(\log\frac{\chi}{f_\chi} - A\right) \\ & + \frac{\lambda_\chi\delta_1 f_\chi^4}{2}\left(\frac{\chi}{f_\chi}\right)^{3-\gamma_m} \text{Tr}[U + U^\dagger] + \frac{\lambda_\chi\delta_2 f_\chi^4}{4}\left(\frac{\chi}{f_\chi}\right)^{2(3-\gamma_{4f})} \text{Tr}[(U - U^\dagger)^2] \\ & - V_0, \end{aligned}$$

C.C.

scaling dimensions

$$0 < \gamma_m, \gamma_{4f} \lesssim 1$$

Effective potential of pions
from chiral symmetry breaking

δ_1 : from bare mass term of confining fermion

δ_2 : from four-fermi interactions

Parameters

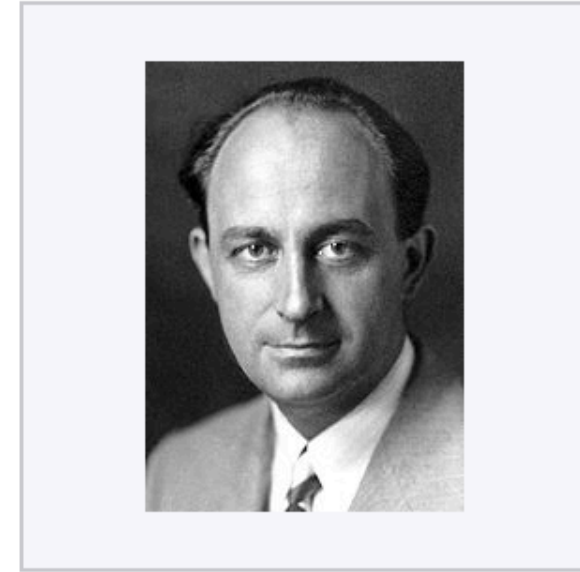
- Along $\phi/f_\phi = 0$, 4 Free parameters $(\lambda_\chi, f_\chi, \delta_1, \delta_2)$ describes the composite dynamics of dilaton
- $0 < \gamma_m, \gamma_{4f} \lesssim 1$: scaling dimensions of the induced operators suggested by lattice calculations [DeGrand (15)]
- $\text{Max}[\lambda_\chi \delta_1 f_\chi^4, \lambda_\chi \delta_2 f_\chi^4] \lesssim f_\phi^4$ such that the theory stays in a perturbative

regime ==> We request $\text{Max}(\delta_1, \delta_2) \lesssim \frac{f_\phi^4}{\lambda_\chi f_\chi^4}$

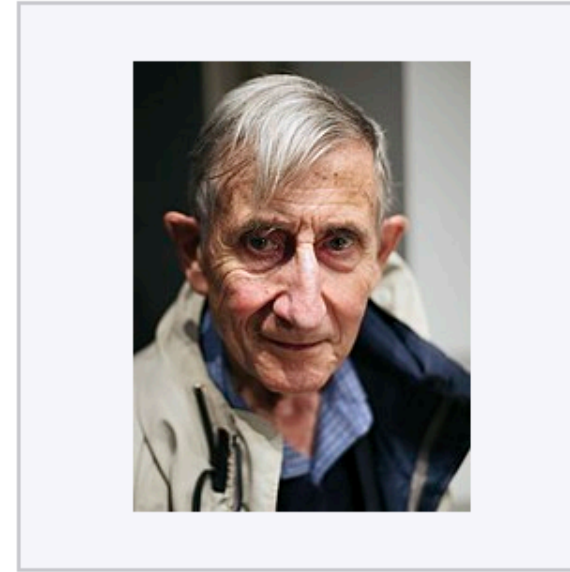
J. von Neumann's elephant



John von Neumann



Enrico Fermi



Freeman Dyson in 2005

I remember my friend Johnny von Neumann used to say,
"With four parameters I can fit an elephant,
and with five I can make him wiggle his trunk."

E. Fermi to F. Dyson in response to his calculation of meson-proton scattering (1951)

Four parameter fit to an elephant

Elephant Drawing

created by scp

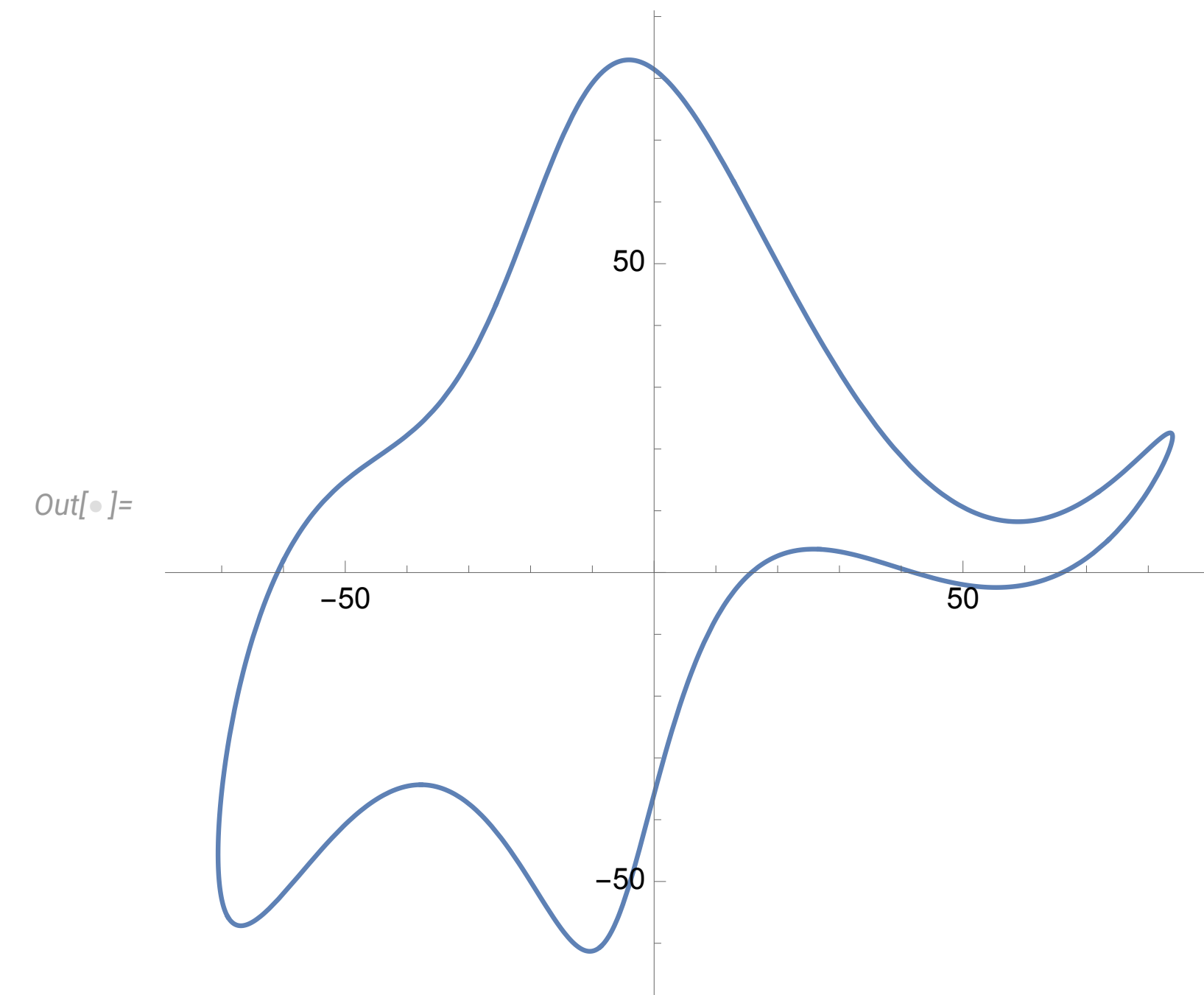
Ref) <http://aapt.scitation.org/doi/abs/10.1119/1.3254017?journalCode=ajp>

```
In[ ]:= p1 = 50. - 30. I; p2 = 18. + 8. I; p3 = 12. - 10. I; p4 = -14. - 60. I;
```

```
In[ ]:= cx[1] = Re[p1] * I;  
cx[2] = Re[p2] * I;  
cx[3] = Re[p3];  
cx[4] = 0;  
cx[5] = Re[p4];  
cy[1] = Im[p4] + I * Im[p1];  
cy[2] = Im[p2] * I;  
cy[3] = Im[p3] * I;  
cy[4] = 0;  
cy[5] = 0;
```

```
In[ ]:= Do[ax[k] = Re[cx[k]], {k, 1, 5}];  
Do[bx[k] = Im[cx[k]], {k, 1, 5}];  
Do[ay[k] = Re[cy[k]], {k, 1, 5}];  
Do[by[k] = Im[cy[k]], {k, 1, 5}];
```

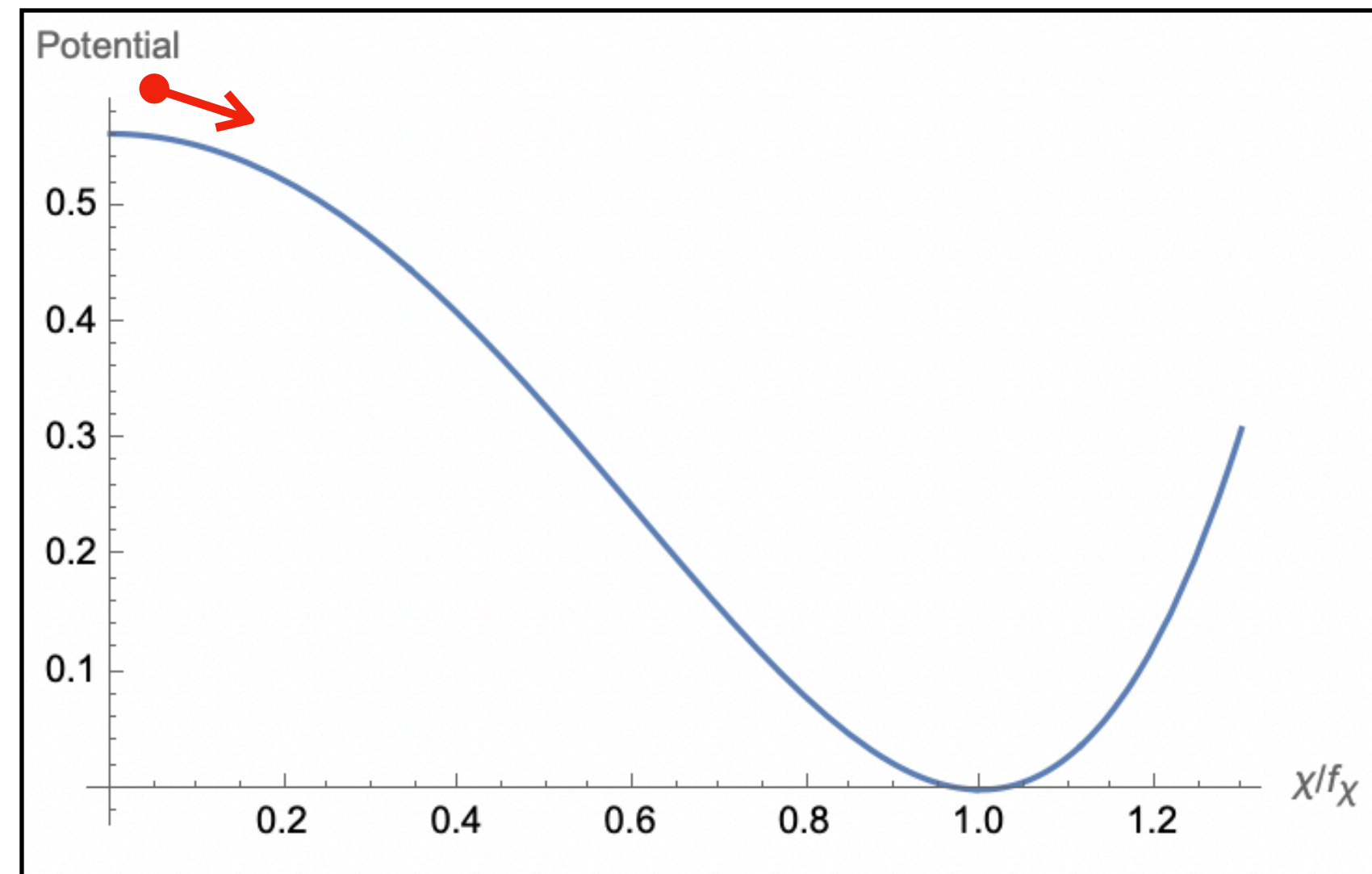
```
In[ ]:=  
myx = Sum[ax[k] * Cos[k * t] + bx[k] * Sin[k * t], {k, 1, 5}];  
myy = Sum[ay[k] * Cos[k * t] + by[k] * Sin[k * t], {k, 1, 5}];  
ParametricPlot[{myy, -myx}, {t, 0, 2 Pi}]
```



Dilaton as inflaton

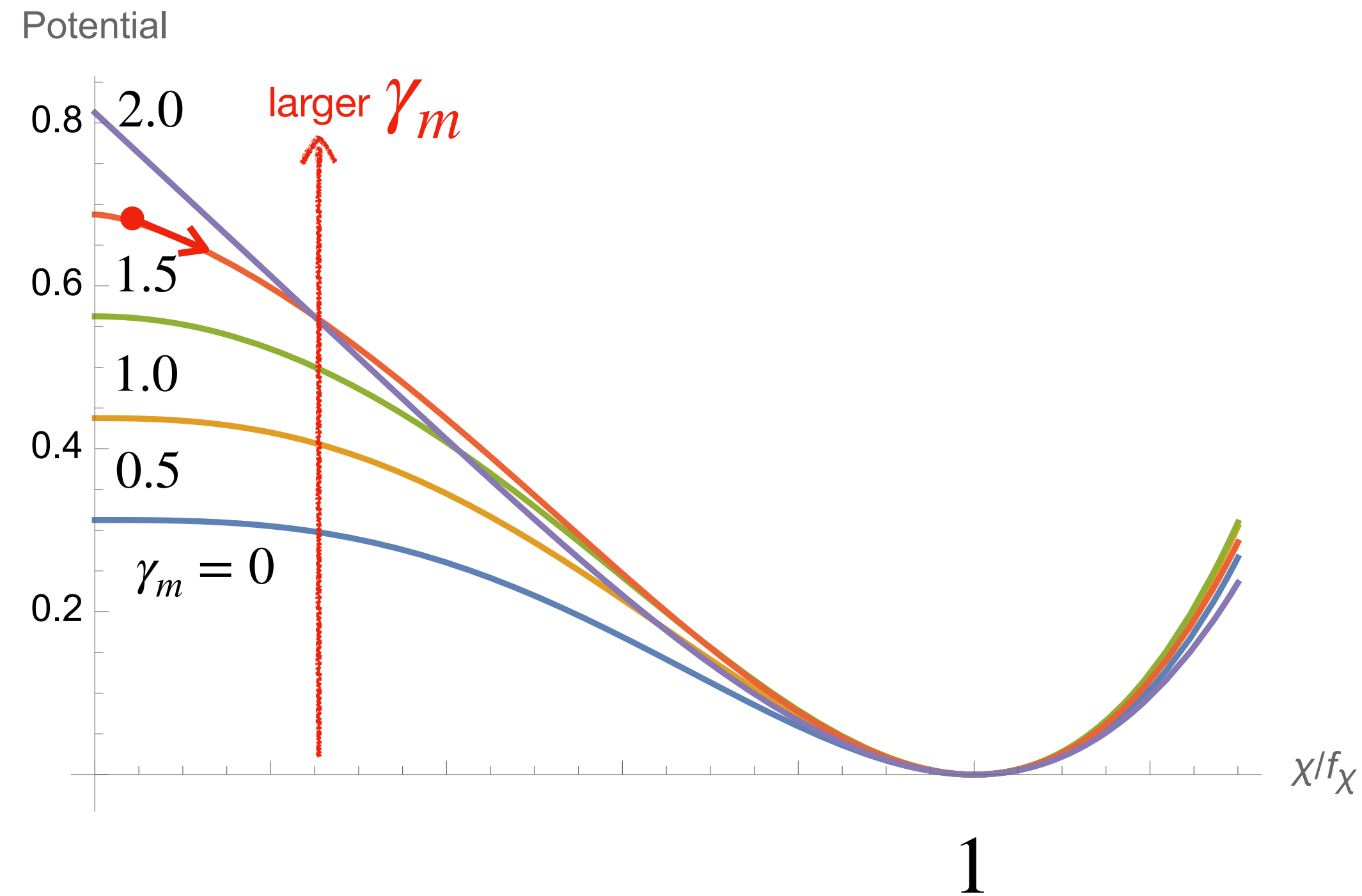
- The dilaton potential **along $\phi/f_\phi = 0$**

$$V(\phi = 0, \chi) = -\lambda_\chi \delta_1 f_\chi^4 \left(\frac{\chi}{f_\chi} \right)^{3-\gamma_m} + \frac{\lambda_\chi}{4} \chi^4 \left(\log \frac{\chi}{f_\chi} - A^{\text{single}} \right) + V_0^{\text{single}}$$



$$V(\phi = 0, \chi) = -\lambda_\chi \delta_1 f_\chi^4 \left(\frac{\chi}{f_\chi} \right)^{3-\gamma_m} + \frac{\lambda_\chi}{4} \chi^4 \left(\log \frac{\chi}{f_\chi} - A^{\text{single}} \right) + V_0^{\text{single}}$$

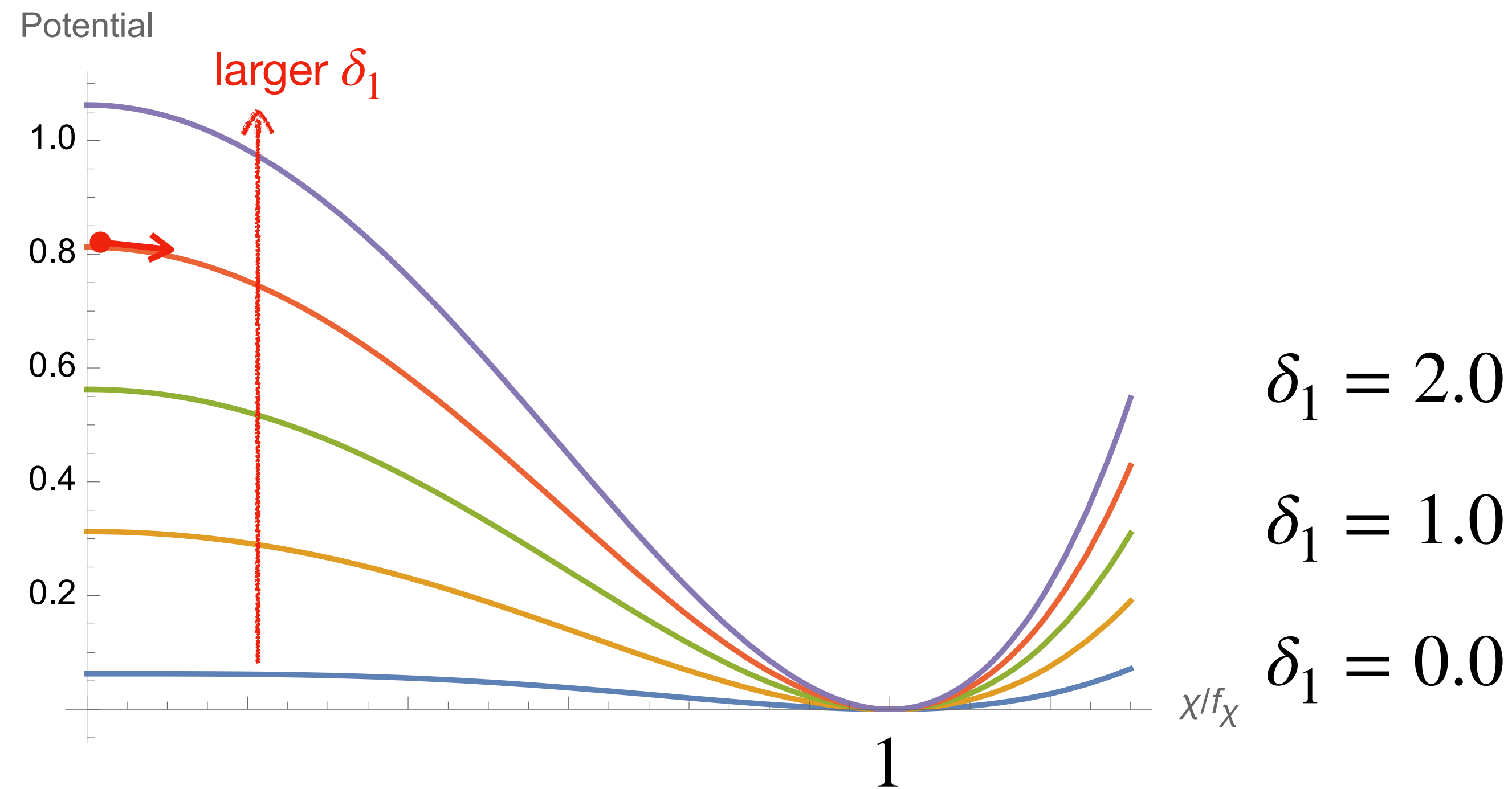
$\delta_1 = 1.0$ fixed



- $A^{\text{single}} = \frac{1}{4} + \delta_1(3 - \gamma_m)$ for $V_{,\chi} = 0$ at $\chi = f_\chi$
- $V_0 = \frac{\lambda_\chi f_\chi^4}{16} [1 + 4\delta_1(1 + \gamma_m)]$ for $V = 0$ at $\chi = f_\chi$

$$V(\phi = 0, \chi) = -\lambda_\chi \delta_1 f_\chi^4 \left(\frac{\chi}{f_\chi} \right)^{3-\gamma_m} + \frac{\lambda_\chi}{4} \chi^4 \left(\log \frac{\chi}{f_\chi} - A^{\text{single}} \right) + V_0^{\text{single}}$$

$\gamma_m = 1.0$ fixed



• $A^{\text{single}} = \frac{1}{4} + \delta_1(3 - \gamma_m)$ for $V_{,\chi} = 0$ at $\chi = f_\chi$

• $V_0 = \frac{\lambda_\chi f_\chi^4}{16} [1 + 4\delta_1(1 + \gamma_m)]$ for $V = 0$ at $\chi = f_\chi$

Precision test

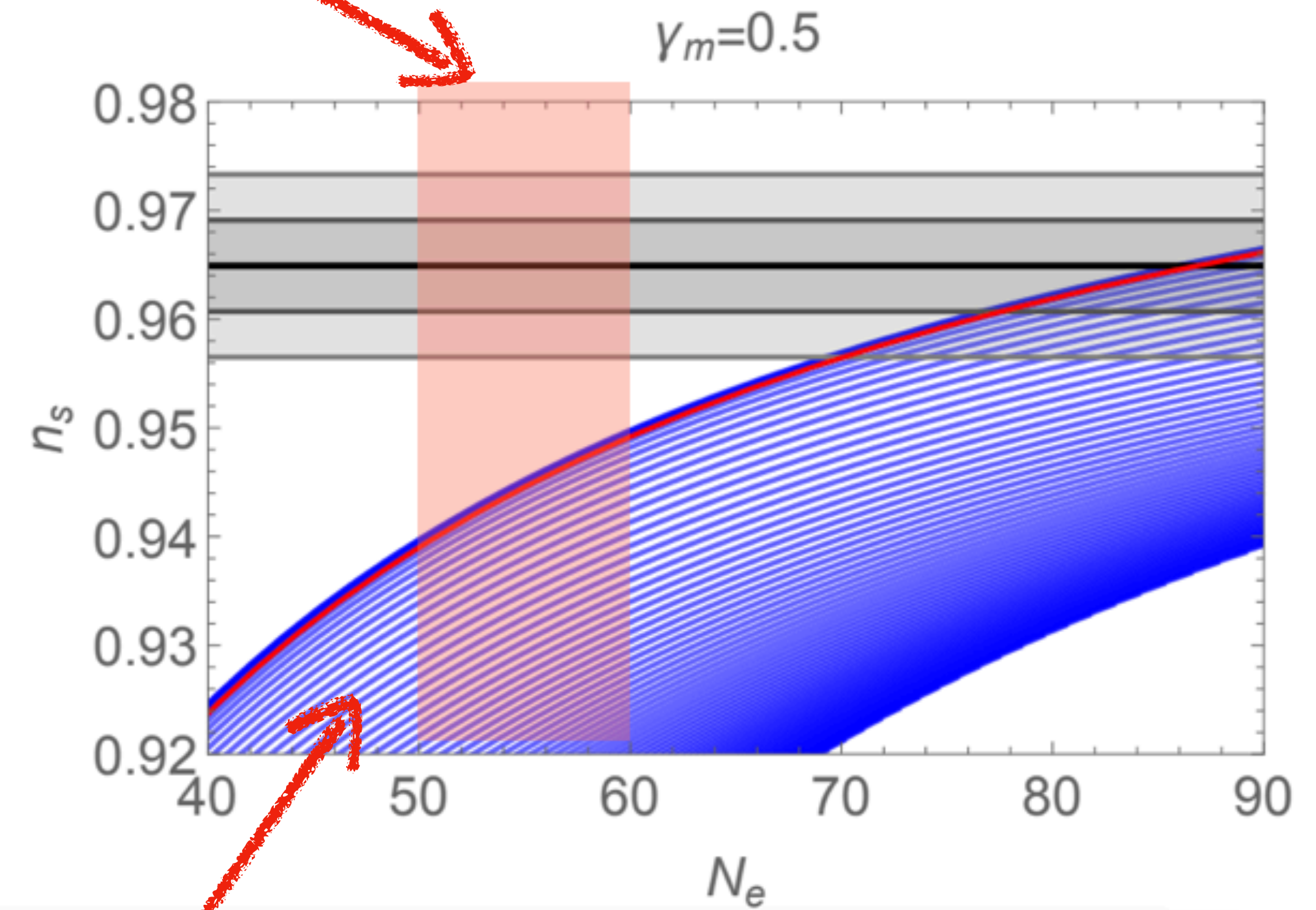
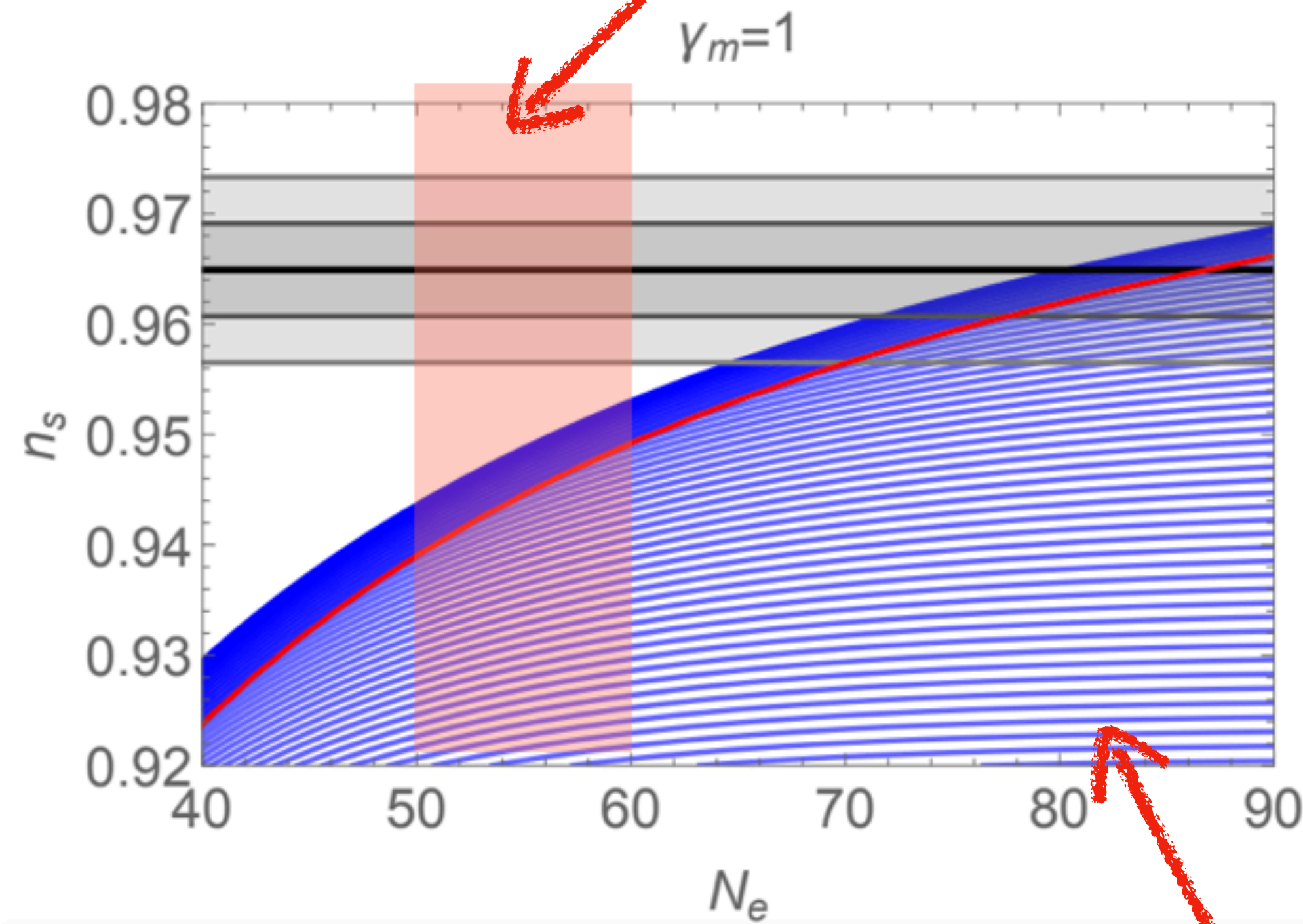
- PLANCK [[Planck 1807.06277](#)]

$$\begin{aligned}n_s &= 0.9649 \pm 0.0042, \\r &< 0.036, \\ \ln(10^{10} \mathcal{P}_\zeta(k_*)) &= 3.040 \pm 0.0016, \\ \frac{dn_s}{d \ln k} &= -0.0045 \pm 0.0067.\end{aligned}$$

- Many representative models were already ruled out
- We can test our composite model

Dilaton inflaton (n_s, N_e) $r \ll 1$ all the cases

Obs indicate here



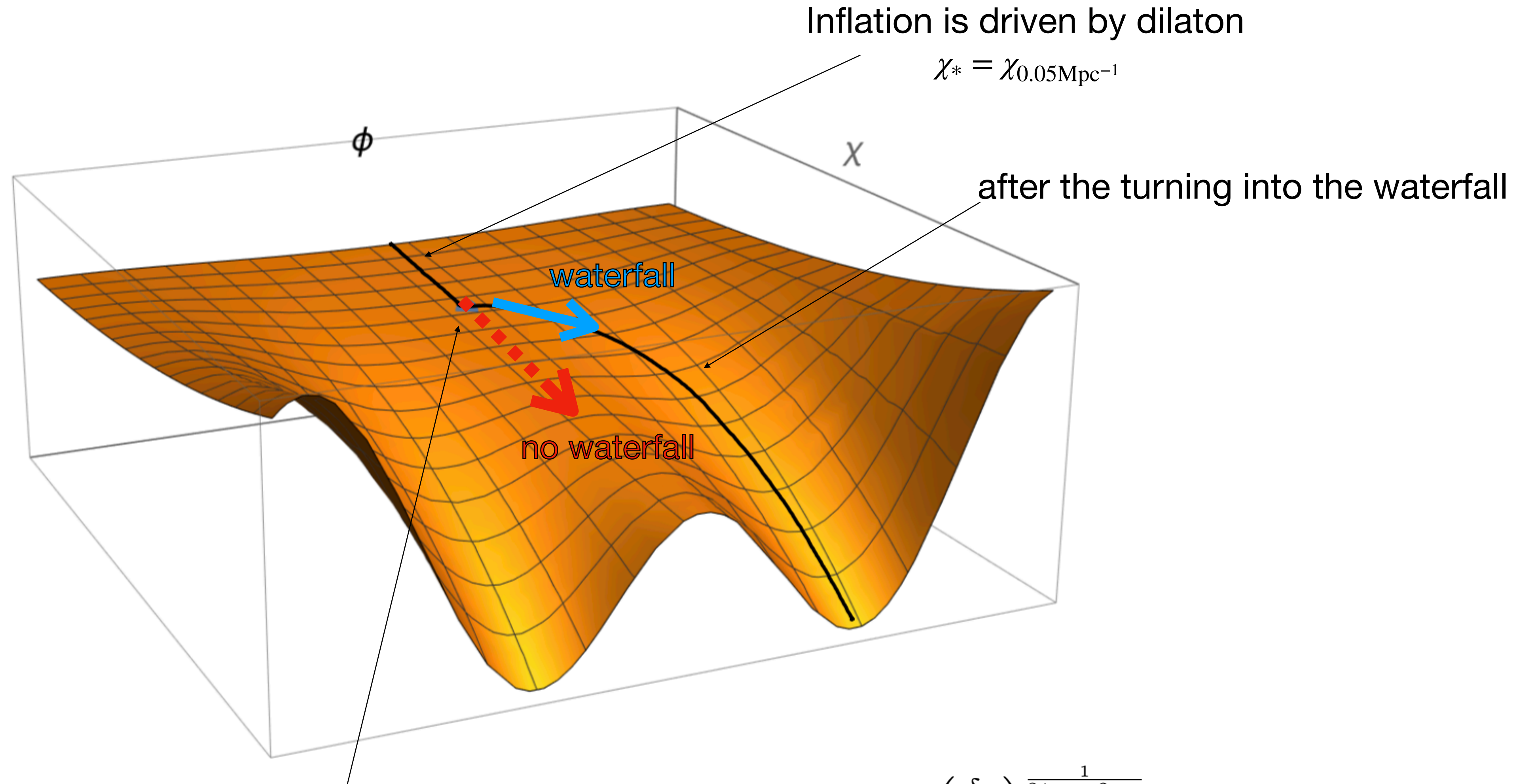
The dilaton model prediction here

Wait! Pions

$$m_{\phi}^2 = \lambda_{\chi} \frac{f_{\chi}^4}{f_{\phi}^2} \times \begin{cases} \delta_1 f_{\chi}^{\gamma_m - 3} \chi^{3 - \gamma_m} - 2\delta_2 f_{\chi}^{2(\gamma_{4f} - 3)} \chi^{2(3 - \gamma_{4f})} & \text{if } \phi = 0, \chi < \chi_c, \text{ along dilaton direction} \\ \left(2\delta_2 - \frac{\delta_1^2}{2\delta_2}\right) & \text{if } \phi = \phi_0, \chi = f_{\chi} \cdot \text{ at the minimum} \end{cases}$$

- At the turning point, $\chi = \chi_c$, m_{ϕ}^2 changes its sign then rapidly turns into the waterfall direction $\phi \rightarrow \phi_0$ then terminates the inflation
- Pions play the role of waterfall field!

Dilaton + pion potential



$V_{\phi\phi}$ changes the sign and 'Turn' made at $\chi = \chi_c \equiv \left(\frac{\delta_1}{2\delta_2}\right)^{\frac{1}{3+\gamma_m-2\gamma_4 f}} f_\chi$

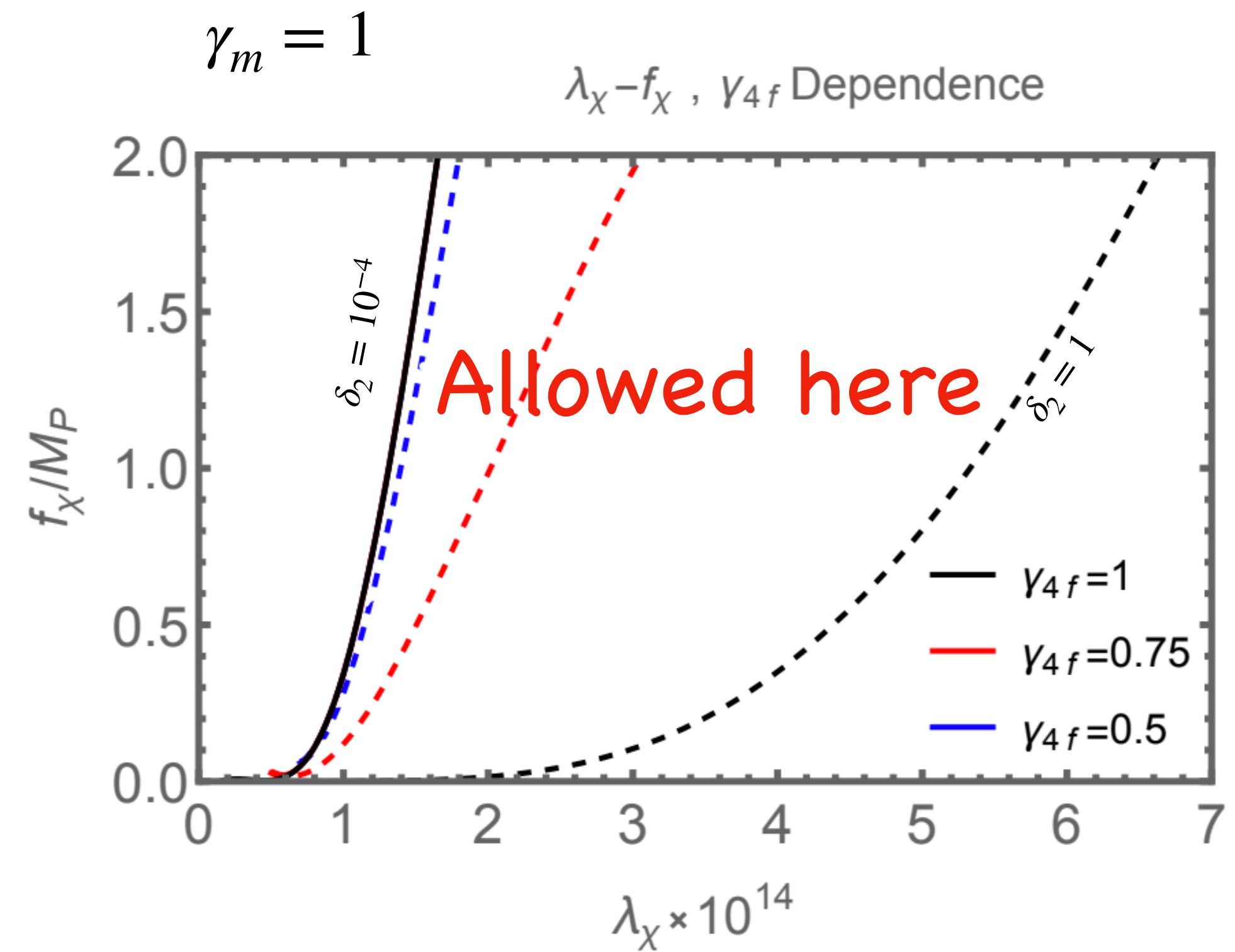
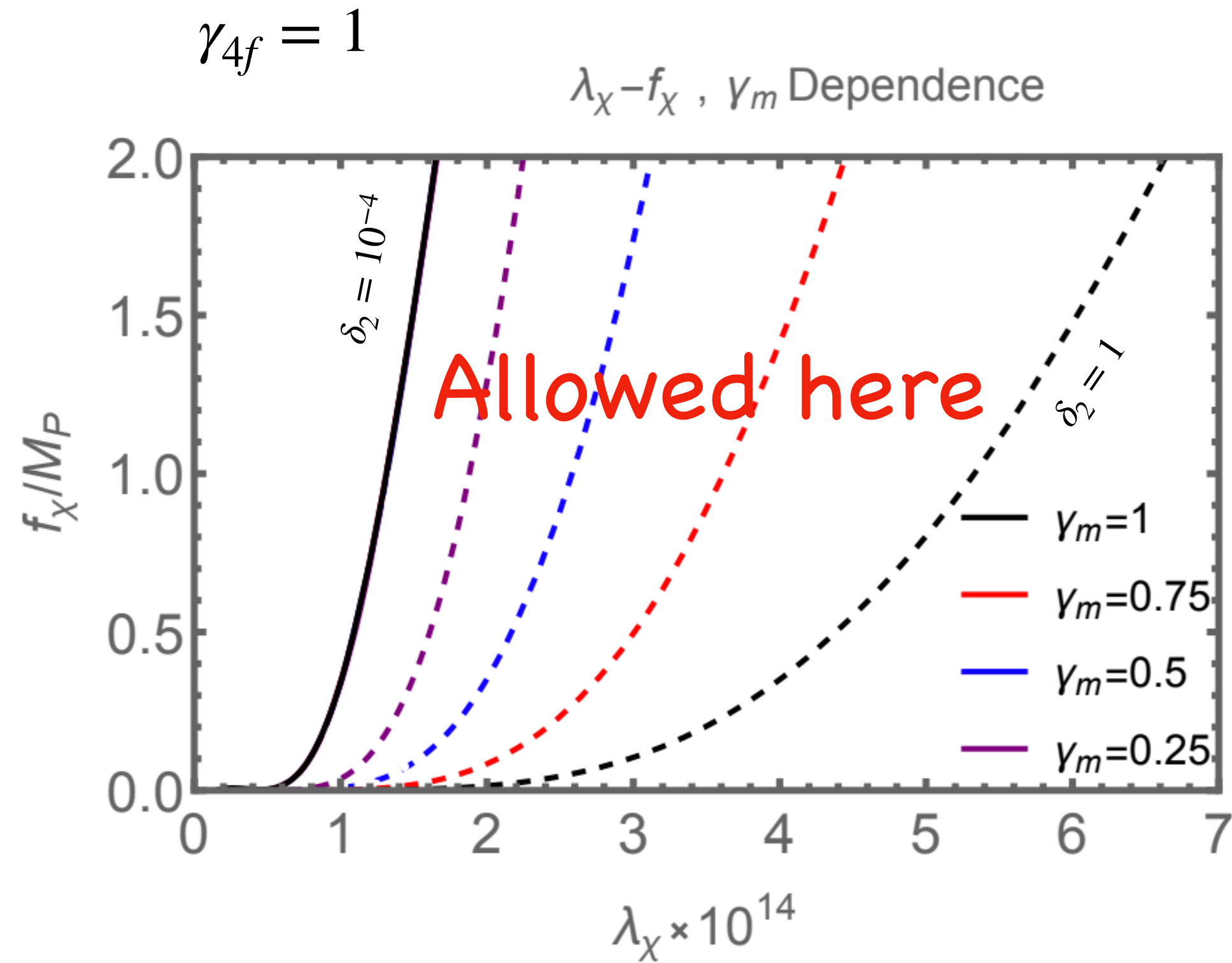
Allowed Parameter spaces

$$n_s = 0.9649 \pm 0.0042,$$

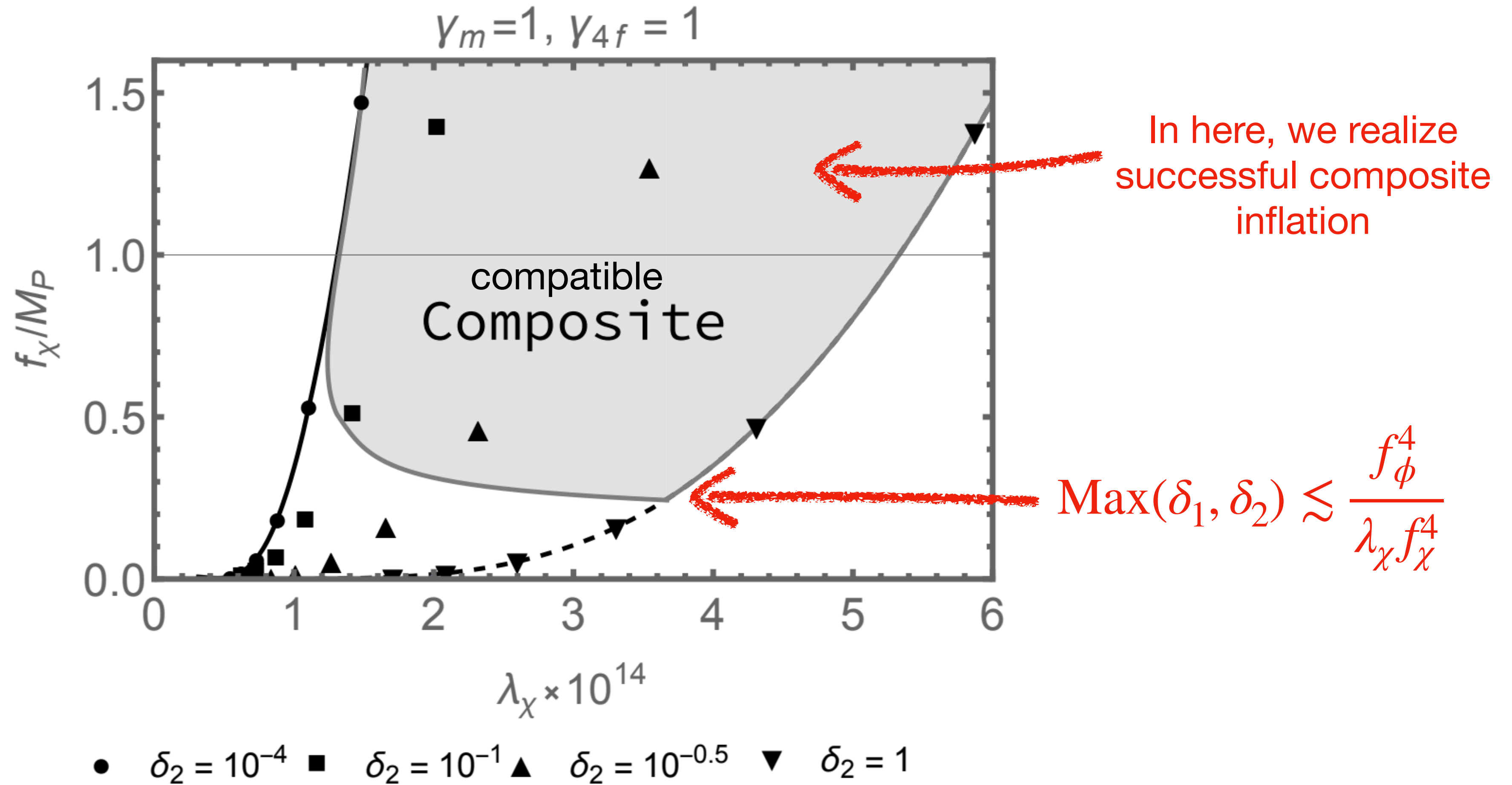
$$r < 0.036,$$

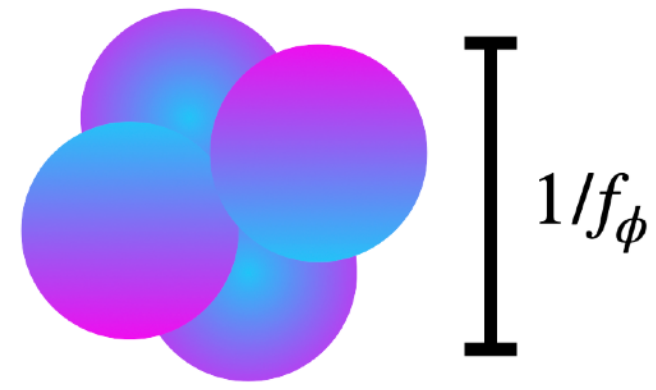
$$\ln(10^{10} \mathcal{P}_\zeta(k_*)) = 3.040 \pm 0.0016,$$

$$\frac{dn_s}{d \ln k} = -0.0045 \pm 0.0067.$$



CMB compatible space





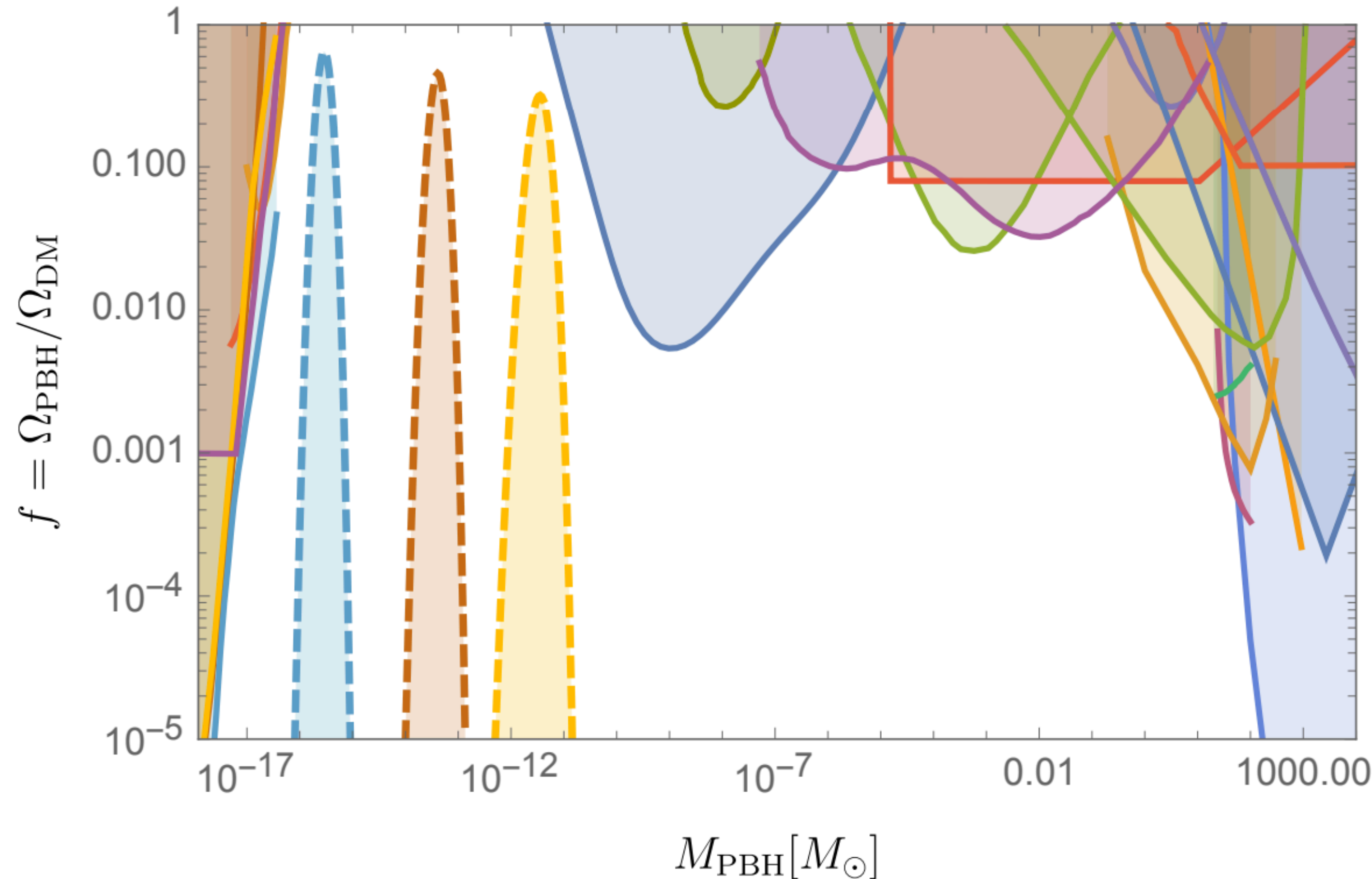
Conclusion & Outlook

- Composite inflation model is suggested based on a $SU(N_c)$ gauge symmetry
- The suggested Composite model is compatible with the Planck 18+ BK18 + BAO observation when dilation plays the role of inflaton and pion plays the role of waterfall.
- Inflation scale $H_{inf} \sim 10^{10}$ GeV, Composite (pion) scale $m_\phi \sim f_\phi \sim 10^{14}$ GeV : everything is below Planck scale \sim good EFT description
- More details of reheating is specified only after (pion-standard model) and (dilaton~standard model) interactions are given.
- Some of the pions could be stable and **WIMPZILLA** Dark matter (work in progress)
- Tachyonic instability generates **PBH & GWs** etc. (work in progress)

PBH from Higgs-R²

D.Y.Cheong, S.M.Lee, SCP JCAP01 (2021) 032 [1912.12032]

$$S_J = \int d^4x \sqrt{-g_J} \left[F(h, R_J) - \frac{1}{2} g^{\mu\nu} \nabla_\mu h \nabla_\nu h - \frac{\lambda}{4} h^4 \right], \quad F(h, R_J) = \frac{M_P^2}{2} \left(R_J + \frac{\xi h^2}{M_P^2} R_J + \frac{R_J^2}{6M^2} \right)$$



+RG running effect

$$\beta_\alpha = -\frac{1}{16\pi^2} \frac{(1+6\xi)^2}{18},$$

$$\beta_\xi = -\frac{1}{16\pi^2} \left(\xi + \frac{1}{6} \right) \left(12\lambda + 6y_t^2 - \frac{3}{2}g'^2 - \frac{9}{2}g^2 \right),$$

$$\beta_\lambda \beta_{\text{SM}} + \frac{1}{16\pi^2} \frac{2\xi^2 (1+6\xi)^2 M^4}{M_P^4},$$

+ R³

D.Y.Cheong, H.M.Lee, SCP
PLB 805 (2020) 135453 [2002.07981](#)

Tachyonic instability

D.Y.Cheong, K. Kohri, SCP
JCAP 10 (2022) 015 [2205.14813](#)