

# Magnetically Confined Mountains on Neutron Stars in General Relativity

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## Outline

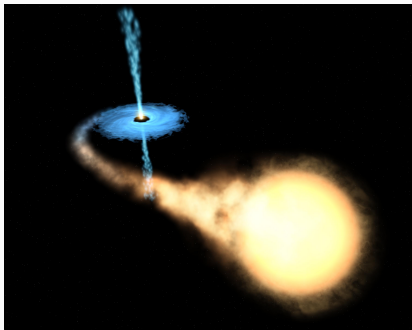
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1. Background and Motivation
2. Magnetically Confined Mountains on Neutron Stars
3. Hydromagnetic Structure
4. Continuous Gravitational Waves

# Background and Motivation

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# Spin-up of Neutron Stars



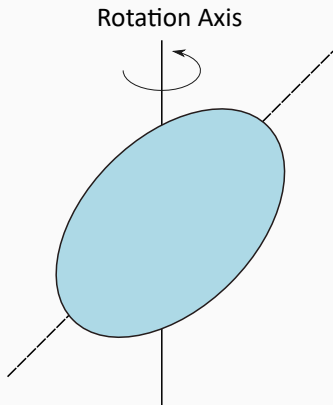
ESA, NASA, and Felix Mirabel

- Neutron stars (NSs) are sometimes found in binary systems such as Low Mass X-Ray Binaries (LMXBs);
- The infalling matter should spin-up NS close to their breaking frequency. But that is not experimentally verified<sup>1</sup>;
- NSs in accreting systems also show reduced magnetic fields.

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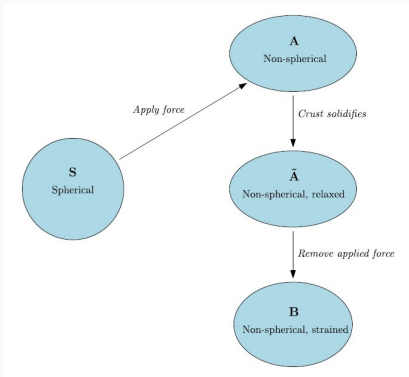
<sup>1</sup>Chakrabarty et al., 2003

# Spin-down Mechanisms



- We need mechanisms for the NS to lose angular momentum;
- One possibility is for the star to spin down due to GW emission;
- Therefore, we need a time-varying quadrupole;
- For a rotating star, this can be non-axisymmetric deformations, the so-called 'mountains'.

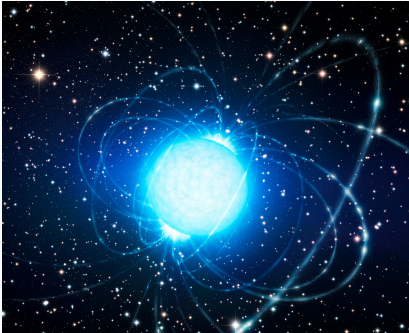
# Elastic of Mountains



Gittins, Andersson and Jones, 2021

- The two main categories of mountains are: elastic (see image) and magnetic;
- The elastic mountains come from elastic deformation of the crust;
- Mountain size is determined by the maximum stress the crust can sustain.

# Magnetic Mountains



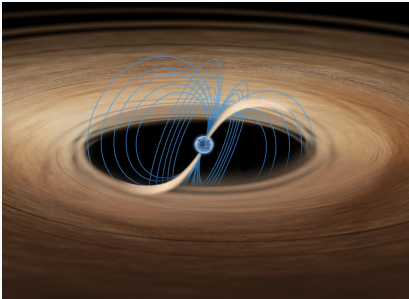
- Magnetic mountains are present in isolated and accreting systems;
- In isolated systems, the presence of magnetic field itself deforms the NS away from spherical symmetry;
- This is more relevant for magnetars and depends on the internal magnetic field.



# **Magnetically Confined Mountains on Neutron Stars**

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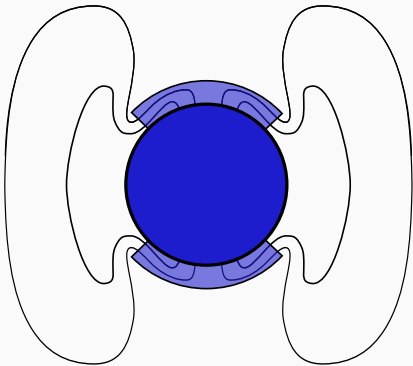
# Magnetically Driven Accretion



Source: NASA

- In the inner regions of the accretion disk, the fluid motion is dominated by the magnetic field (magnetosphere);
- Matter is accreted on the magnetic poles;
- The accreted plasma distorts the magnetic field of the star.

# Magnetically Confined Mountain



- Mountains start to form in the polar regions;
- Gravity tries to smooth out the mountain;
- The mountain deforms the magnetic field;
- Magnetic field resists and holds the mountains in place.

- The Newtonian analysis of this mechanism is already very explored in the literature <sup>2</sup>;
- The overall result is a decrease of the magnetic dipole moment of the star and an increase of the mass-quadrupole moment;
- These results partially help understand the observed properties of LMXBs.

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<sup>2</sup>See, for example, Uchida, 1981; Payne and Melatos, 2004; Vigelius and Melatos, 2008; Wette et al., 2010; Fujisawa et al., 2022.

# Formulation of the Problem in General Relativity

- The goal of my research was to approach this problem in General Relativity;
- The first principle equations are:

## Rest mass conservation

$$\nabla_a(\rho u^a) = 0$$

## Conservation of Energy Momentum

$$\nabla_a T^{ab} = \nabla_a(T_{fluid}^{ab} + T_{EM}^{ab}) = 0$$

## Einstein's Equations

$$G_{ab} = 8\pi T_{ab}$$

## Maxwell's Equations

$$d\mathbf{F} = 0$$

$$d\star\mathbf{F} = \star\mathbf{J}$$

## Relevant Approximations

- In the small accretion limit, we fix a background geometry of spacetime as Schwarzschild;

$$ds^2 = -e^{2\Phi} dt^2 + e^{-2\Phi} dr^2 + r^2 d\theta^2 + r^2 \sin^2 \theta d\phi^2$$

where

$$\Phi(r) = \frac{1}{2} \ln \left( 1 - \frac{2M_*}{r} \right)$$

- The fluid is a perfect conductor ( $\sigma \rightarrow \infty$ )

$$E_b = u^a F_{ab} = 0$$

## Relevant Approximations

- The system evolves through quasi-static steps of magnetostatic equilibrium;

$$u^a = e^{-\Phi} t^a$$

- The system is axisymmetric ( $\mathcal{L}_{\lambda} F = 0$ ), where  $\lambda_a = (d\phi)_a$  is the axisymmetric Killing vector;
- The magnetic field is poloidal;

$$F = \frac{2}{|\lambda|} d\phi \wedge d\psi$$



# General Relativistic Grad-Shafranov Equation

- Using these approximations in the first principle equations, we obtain:

## GR Grad-Shafranov Equation

$$\frac{1}{\lambda} \left( \nabla^a \nabla_a \psi - \frac{1}{\lambda} \nabla^a \lambda \nabla_a \psi \right) = -\rho h F'(\psi)$$

In spherical coordinates:

$$\begin{aligned} \frac{1}{r^2 \sin^2 \theta} \left[ \frac{\partial^2 \psi}{\partial r^2} - \frac{\partial}{\partial r} \left( \frac{2M}{r} \frac{\partial \psi}{\partial r} \right) + \frac{\sin \theta}{r^2} \frac{\partial}{\partial \theta} \left( \frac{1}{\sin \theta} \frac{\partial \psi}{\partial \theta} \right) \right] = \\ = -\rho h F'(\psi) \end{aligned}$$

# The Grad-Shafranov Equation

- The function  $F(\psi)$  is defined self-consistently by the magnetic flux freezing condition and imposing a rest mass-flux ratio.

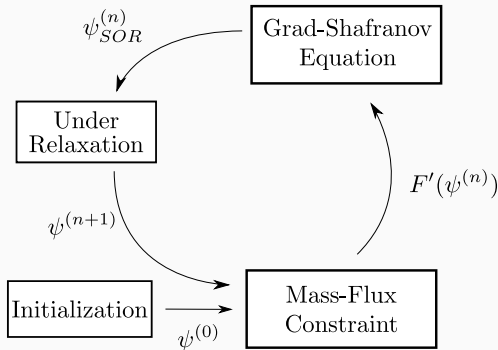
## GR self-consistent F

$$F(\psi) = \left( \frac{dM}{d\psi} \right)^{1+c_s^2} \left( \frac{2\pi}{c_s^2} \int_C r \sin \theta |\nabla \psi|^{-1} e^{-(\Phi-\Phi_0)/c_s^2} ds \right)^{-(1+c_s^2)}$$

- The general form of the problem is maintained in GR.

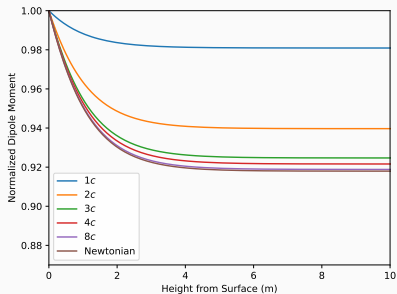
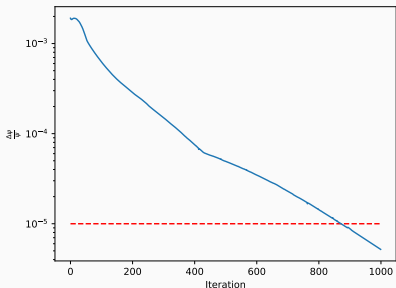
## Numerical Approach

- The GR Grad-Shafranov equation is solved with Successive Over Relaxation (SOR);
- The mass-flux constraint is maintained through an iterative process.



# Numerical Convergence

- The system does reach a final equilibrium.
- The relativistic results reproduce the Newtonian one in the appropriate limit ( $c \rightarrow \infty$ ).

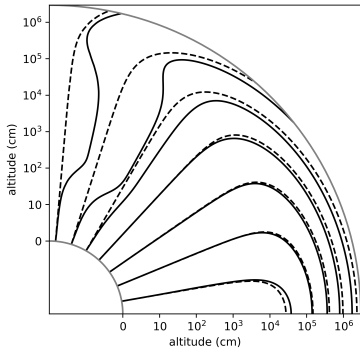


# Hydromagnetic Structure

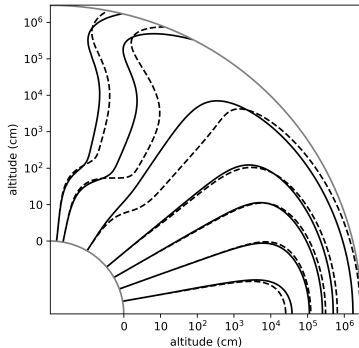
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# Numerical Results

- The distortion of the magnetic field lines is attenuated in GR.

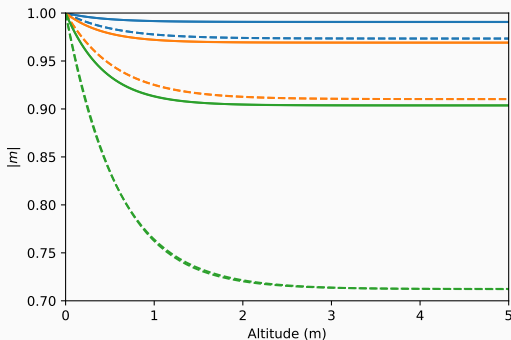


(a) GR vs. Dipole



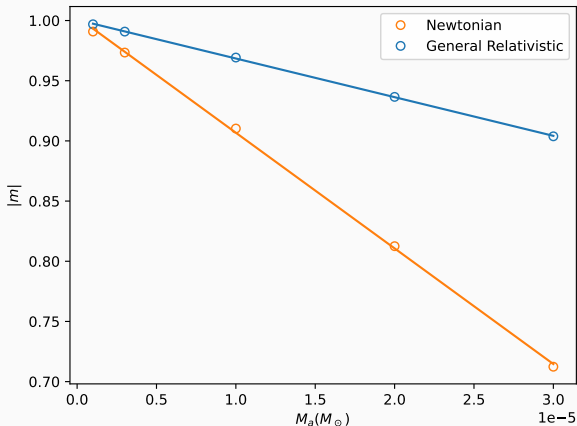
(b) GR vs. Newtonian

- The attenuation of the magnetic field is also found in the magnetic dipole moment:



Reduction of the magnetic dipole for different masses:  $10^{-6}M_{\odot}$  (blue),  $10^{-5}M_{\odot}$  (orange) and  $3 \times 10^{-5}M_{\odot}$  (orange). General relativistic results in solid lines, and Newtonian in dashed lines.

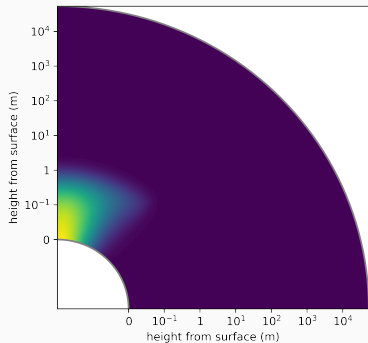
- And in the total magnetic dipole moment per accreted mass:



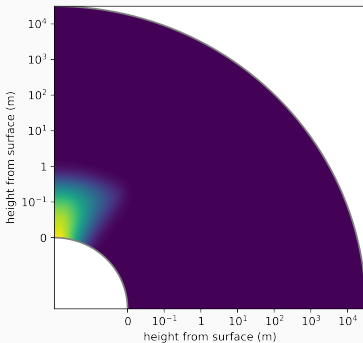
- In the plot above the GR curve is 3 times less steep than the Newtonian one!



- Because of the different deformation of the magnetic field, the density distribution also changes.



**(a)** Newtonian



**(b)** GR

- These results were summarized in a paper published at MNRAS.

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JOURNAL ARTICLE

### Magnetically confined mountains on accreting neutron stars in general relativity

Pedro H B Rossetto , Jörg Frauendiener, Ryan Brunet, Andrew Melatos

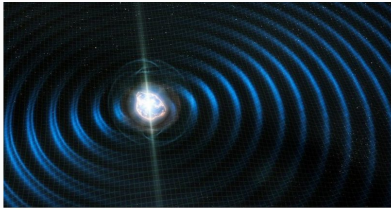
*Monthly Notices of the Royal Astronomical Society*, Volume 526, Issue 2, December 2023,  
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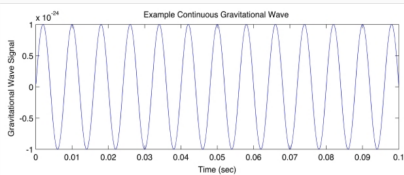
# Continuous Gravitational Waves

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# Continuous Gravitational Waves

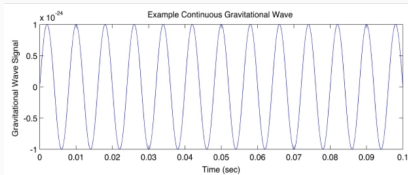
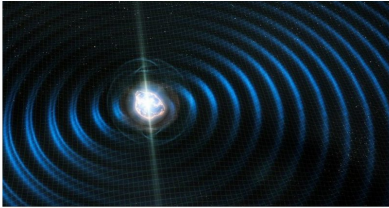


- Continuous gravitational waves (CGWs) are GW with constant frequency;
- The best source candidates for these waves are massive rotating non-axisymmetric objects;



[scienceblog.com/496929/](http://scienceblog.com/496929/) with alterations.

# Continuous Gravitational Waves



[scienceblog.com/496929/](http://scienceblog.com/496929/) with alterations.

- Continuous gravitational waves (CGWs) are GW with constant frequency;
- The best source candidates for these waves are massive rotating non-axisymmetric objects;
- That is: neutron stars with mountains!

- In accreting systems, the continuous aspect is due to the spin torque balance;

Source: OzGrav ARC Centre of Excellence

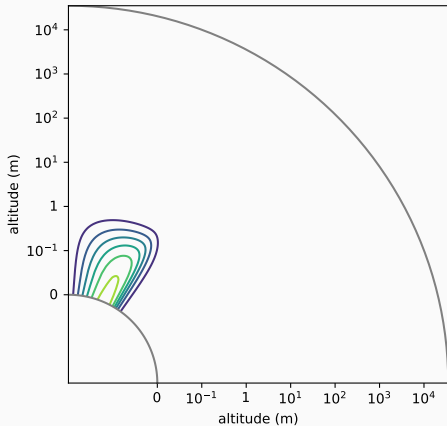
- Given the hydromagnetic structure of the solutions, we can calculate the mass ellipticity of the star:

$$\epsilon = \frac{\pi}{I_0} \int_{V'} \left( e + \frac{B^2}{2} \right) r'^4 \sin \theta (3 \cos^2 \theta - 1) dr' d\theta$$

- Which is related to the amplitude of continuous gravitational waves via:

$$h_0 = \frac{4\pi^2 I_0 f_{GW}^2}{r} \epsilon$$

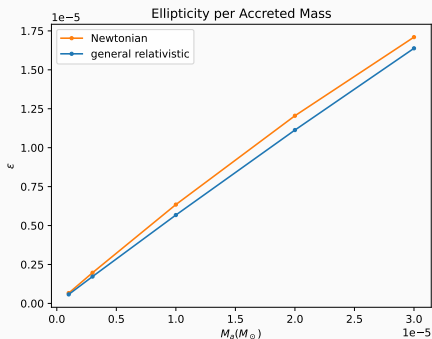
# Magnetic Contribution



- An interesting theoretical contribution is due to the magnetic field  $B^2/2$ ;
- In our case the ellipticities are small, in the order of  $\epsilon \sim 10^{11}$ ;
- This contribution is relevant for magnetars, but their rotation frequency is small which hinders gravitational wave emission.

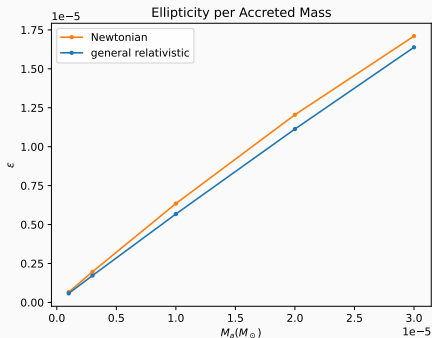


# Ellipticity



- We observe a reduction of the ellipticity of the star in GR;
- For some parameters of the model, the difference between relativistic and Newtonian can get to 10%;
- The ellipticity values are too high due to the simplifying assumptions of the model (isothermal equation of state and rigid stellar surface).

# Ellipticity



- We observe a reduction of the ellipticity of the star in GR;
- For some parameters of the model, the difference between relativistic and Newtonian can get to 10%;
- The ellipticity values are too high due to the simplifying assumptions of the model (isothermal equation of state and rigid stellar surface).

- GR has an important role to play in the model of magnetically confined mountains on neutron stars;
  - The relativistic effects on the magnetic properties of the star are large, yielding three times less magnetic screening;
  - Relativistic effects reduce the ellipticity of the star, reducing the amplitude of gravitational waves;
- The general relativistic model needs to be expanded.

- Several improvements can be done as future work in order to have more complete and astrophysically relevant model. They include:
  - Generalising the equation of state;
  - Consider time-dependent effects, mountain sinking and Ohmic diffusion;
  - Analyse the stability using GRMHD solvers, such as *GRHydro*;
  - Consider the effects of the neutron star rotation in the curvature of spacetime.

# Acknowledgements

I want to thank:

- The organizing committee for the invitation to give this presentation;
- My supervisors, Prof. Jörg Frauendiener and Prof. Andrew Melatos;
- And all of you have listened to the talk.

**Thank you!**

*Any questions?*