

Shimmering gravitons in the gamma-ray sky



Federico Urban

CEICO

Institute of Physics
Czech Academy of Sciences
Prague

Gravitation and Cosmology 2024
YITP Kyoto, Japan

Feb 21, 2024



Co-funded by
the European Union



MINISTRY OF EDUCATION,
YOUTH AND SPORTS

:: take home message ::

:: take home message ::



What is the most energetic gravitons we could detect?

:: Take home message ::



What is the most energetic gravitons we could detect?



Look at the inverse-Gertsenshtein effect

:: Take home message ::



What is the most energetic gravitons we could detect?



Look at the inverse-Gertsenshtein effect



We checked magnetars, the GMF, the IMF

:: Take home message ::



What is the most energetic gravitons we could detect?



Look at the inverse-Gertsenshtein effect



We checked magnetars, the GMF, the IMF



Our best bet is: just shy of a PeV, not beyond

:: Take home message ::



What is the **most energetic** gravitons we could detect?



Look at the **inverse-Gertsenshtein** effect



We checked magnetars, **the GMF**, the IMF



Our best bet is: **just shy of a PeV**, not beyond

JCAP06 (2023) 019

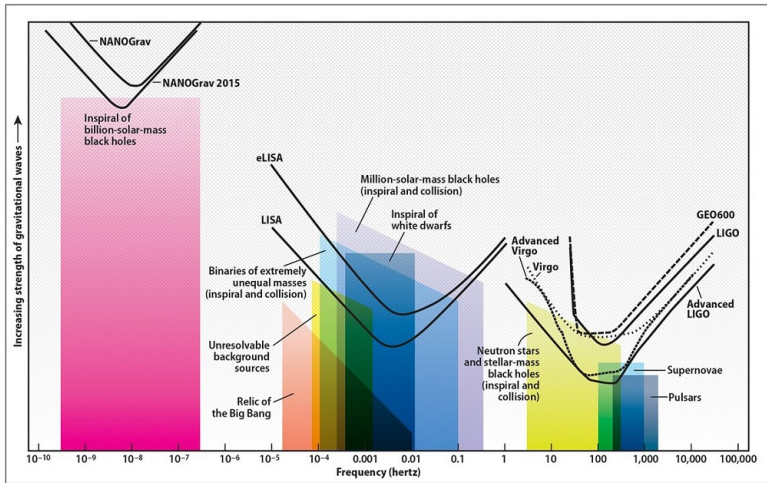
S. Ramazanov, R. Samanta, G. Trenkler, F. R. Urban



What we know



:: observations ::



:: expectations ::

:: expectations ::

- Experimental efforts to detect GWs at a **variety** of frequencies

:: expectations ::

- Experimental efforts to detect GWs at a **variety** of frequencies
- The **observed** nHz to kHz range of the SGWB and BBH mergers

:: expectations ::

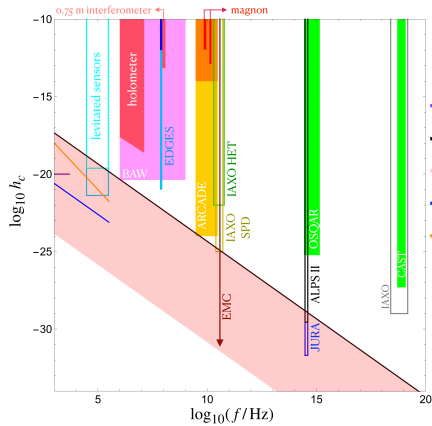
- Experimental efforts to detect GWs at a **variety** of frequencies
- The **observed** nHz to kHz range of the SGWB and BBH mergers
- The 10^{-16} Hz range of the **CMB** B-mode polarisation

:: expectations ::

- Experimental efforts to detect GWs at a **variety** of frequencies
- The **observed** nHz to kHz range of the SGWB and BBH mergers
- The 10^{-16} Hz range of the **CMB** B-mode polarisation
- The **MHz-GHz-THz** range of repurposed ALP detectors

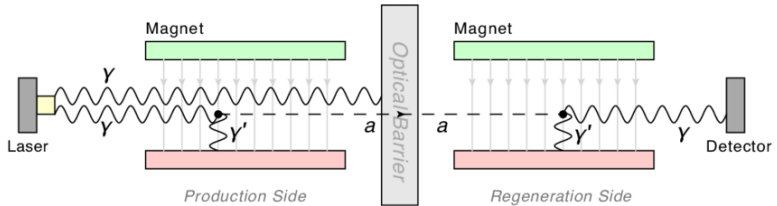
(Aggarwal + '20; Ita, Kohri, Nakayama '23)

:: very-ultra high freqs ::



10^{18} Hz is about 1 keV

:: light shining through the wall ::



© Ejlli et al., 2019

:: inverse Gertsenshtein ::

:: inverse Gertsenshtein ::

🌳 High-energy GWs can be converted to photons in an external magnetic field

:: inverse Gertsenshtein ::

- 🌳 High-energy GWs can be converted to photons in an external magnetic field
- 🌳 Good: MFs are everywhere from the Earth to galaxies to the IGM, and they can have huge correlation lengths

:: inverse Gertsenshtein ::

- 🌳 High-energy GWs can be converted to photons in an external magnetic field
- 🌳 Good: MFs are everywhere from the Earth to galaxies to the IGM, and they can have huge correlation lengths
- 🌳 Bad: MFs are weak and the coupling is Planck-suppressed

:: inverse Gertsenshtein ::

- 🌳 High-energy GWs can be converted to photons in an external magnetic field
- 🌳 Good: MFs are everywhere from the Earth to galaxies to the IGM, and they can have huge correlation lengths
- 🌳 Bad: MFs are weak and the coupling is Planck-suppressed

This is just good old SM

$$\mathcal{L} = \sqrt{|g|} g^{\mu\rho} g^{\nu\sigma} F_{\mu\nu} F_{\rho\sigma}$$

:: inverse Gertsenshtein ::

- 🌳 High-energy GWs can be converted to photons in an external magnetic field
- 🌳 Good: MFs are everywhere from the Earth to galaxies to the IGM, and they can have huge correlation lengths
- 🌳 Bad: MFs are weak and the coupling is Planck-suppressed

This is just good old SM

$$\mathcal{L} = \sqrt{|g|} g^{\mu\rho} g^{\nu\sigma} F_{\mu\nu} F_{\rho\sigma}$$

which generates the coupling

$$h \cdot A \cdot A$$

(Raffelt and Stodolsky 1987 and many more)

:: conversion probability ::

:: conversion probability ::

i. Neglect photon interactions

$$\mathcal{L}_{\text{int}} \propto q \bar{\Psi} \gamma^\mu A_\mu \Psi$$

:: conversion probability ::

i. Neglect photon interactions

$$\mathcal{L}_{\text{int}} \propto q \bar{\Psi} \gamma^\mu A_\mu \Psi$$

ii. The magnetic fields in galaxies and the IGM are

$$B \approx 1 - 100 \mu\text{G}, \quad 10^{-10} \mu\text{G} \lesssim B \lesssim 10^{-5} \mu\text{G}$$

:: conversion probability ::

i. Neglect photon interactions

$$\mathcal{L}_{\text{int}} \propto q \bar{\Psi} \gamma^\mu A_\mu \Psi$$

ii. The magnetic fields in galaxies and the IGM are

$$B \approx 1 - 100 \mu\text{G}, \quad 10^{-10} \mu\text{G} \lesssim B \lesssim 10^{-5} \mu\text{G}$$

iii. The correlation lengths for these fields are about

$$L \approx 0.1 - 10 \text{ kpc}, \quad L \approx \text{Mpc} - \text{Gpc}$$

:: conversion probability ::

i. Neglect photon interactions

$$\mathcal{L}_{\text{int}} \propto q \bar{\Psi} \gamma^\mu A_\mu \Psi$$

ii. The magnetic fields in galaxies and the IGM are

$$B \approx 1 - 100 \mu\text{G}, \quad 10^{-10} \mu\text{G} \lesssim B \lesssim 10^{-5} \mu\text{G}$$

iii. The correlation lengths for these fields are about

$$L \approx 0.1 - 10 \text{ kpc}, \quad L \approx \text{Mpc} - \text{Gpc}$$

iv. The conversion probability is similar to axion-photon

$$P_{h \rightarrow \gamma} = B^2 L^2 / 2M_{\text{P}}^2$$

∴ not so simple ∴

:: not so simple ::

 Photons interact via the Euler-Heisenberg Lagrangian

:: not so simple ::

☞ Photons interact via the Euler-Heisenberg Lagrangian

$$\text{☞} \frac{\alpha^2}{90m_e^4} \left[(F_{\mu\nu}F^{\mu\nu})^2 + \frac{7}{4} (\tilde{F}_{\mu\nu}F^{\mu\nu})^2 \right]$$

:: not so simple ::

☞ Photons interact via the Euler-Heisenberg Lagrangian

$$\text{☞} \frac{\alpha^2}{90m_e^4} \left[(F_{\mu\nu}F^{\mu\nu})^2 + \frac{7}{4} (\tilde{F}_{\mu\nu}F^{\mu\nu})^2 \right]$$

☞ Physically this gives three “mass” terms (\sim refraction index)

:: not so simple ::

☞ Photons interact via the Euler-Heisenberg Lagrangian

$$\text{☞} \frac{\alpha^2}{90m_e^4} \left[(F_{\mu\nu}F^{\mu\nu})^2 + \frac{7}{4} (\tilde{F}_{\mu\nu}F^{\mu\nu})^2 \right]$$

☞ Physically this gives three “mass” terms (\sim refraction index)

$$\text{☞} \Delta_{\text{pl}} = -\frac{\omega_{\text{pl}}^2}{2\omega} \approx -1.1 \cdot 10^{-13} \omega_{\text{peV}}^{-1} \text{ kpc}^{-1}$$

∴ not so simple ∴

☞ Photons interact via the Euler-Heisenberg Lagrangian

$$\text{☞} \frac{\alpha^2}{90m_e^4} \left[(F_{\mu\nu}F^{\mu\nu})^2 + \frac{7}{4} (\tilde{F}_{\mu\nu}F^{\mu\nu})^2 \right]$$

☞ Physically this gives three “mass” terms (\sim refraction index)

$$\text{☞} \Delta_{\text{pl}} = -\frac{\omega_{\text{pl}}^2}{2\omega} \approx -1.1 \cdot 10^{-13} \omega_{\text{PeV}}^{-1} \text{ kpc}^{-1}$$

$$\text{☞} \Delta_{\text{CMB}} = \frac{44\pi^2\alpha^2}{2025} \frac{T_{\text{CMB}}^4}{m_e^4} \omega \approx 8 \cdot 10^{-2} \omega_{\text{PeV}} \text{ kpc}^{-1} \quad \omega \lesssim 100 \text{ TeV}$$

∴ not so simple ∴

☞ Photons interact via the Euler-Heisenberg Lagrangian

$$\text{☞} \frac{\alpha^2}{90m_e^4} \left[(F_{\mu\nu}F^{\mu\nu})^2 + \frac{7}{4} (\tilde{F}_{\mu\nu}F^{\mu\nu})^2 \right]$$

☞ Physically this gives three “mass” terms (\sim refraction index)

$$\text{☞} \Delta_{\text{pl}} = -\frac{\omega_{\text{pl}}^2}{2\omega} \approx -1.1 \cdot 10^{-13} \omega_{\text{PeV}}^{-1} \text{ kpc}^{-1}$$

$$\text{☞} \Delta_{\text{CMB}} = \frac{44\pi^2\alpha^2}{2025} \frac{T_{\text{CMB}}^4}{m_e^4} \omega \approx 8 \cdot 10^{-2} \omega_{\text{PeV}} \text{ kpc}^{-1} \quad \omega \lesssim 100 \text{ TeV}$$

$$\text{☞} \Delta_{\text{QED}} = \frac{\alpha}{45\pi} \left(\frac{B}{B_{\text{cr}}} \right)^2 \omega \approx 0.15 \omega_{\text{PeV}} \left(\frac{B}{6\mu\text{G}} \right)^2 \text{ kpc}^{-1}$$

∴ not so simple ∴

☞ Photons interact via the Euler-Heisenberg Lagrangian

$$\text{☞} \frac{\alpha^2}{90m_e^4} \left[(F_{\mu\nu}F^{\mu\nu})^2 + \frac{7}{4} (\tilde{F}_{\mu\nu}F^{\mu\nu})^2 \right]$$

☞ Physically this gives three “mass” terms (\sim refraction index)

$$\text{☞} \Delta_{\text{pl}} = -\frac{\omega_{\text{pl}}^2}{2\omega} \approx -1.1 \cdot 10^{-13} \omega_{\text{PeV}}^{-1} \text{ kpc}^{-1}$$

$$\text{☞} \Delta_{\text{CMB}} = \frac{44\pi^2\alpha^2}{2025} \frac{T_{\text{CMB}}^4}{m_e^4} \omega \approx 8 \cdot 10^{-2} \omega_{\text{PeV}} \text{ kpc}^{-1} \quad \omega \lesssim 100 \text{ TeV}$$

$$\text{☞} \Delta_{\text{QED}} = \frac{\alpha}{45\pi} \left(\frac{B}{B_{\text{cr}}} \right)^2 \omega \approx 0.15 \omega_{\text{PeV}} \left(\frac{B}{6\mu\text{G}} \right)^2 \text{ kpc}^{-1}$$

☞ This suppresses the oscillations by $e^{i\Delta_{\gamma\gamma}d}$ ($\Delta_{\gamma\gamma} \sim \Delta_{\text{pl}} + \Delta_{\text{CMB}} + \Delta_{\text{QED}}$)

∴ not so simple ∴

☞ Photons interact via the Euler-Heisenberg Lagrangian

$$\text{☞} \frac{\alpha^2}{90m_e^4} \left[(F_{\mu\nu}F^{\mu\nu})^2 + \frac{7}{4} (\tilde{F}_{\mu\nu}F^{\mu\nu})^2 \right]$$

☞ Physically this gives three “mass” terms (\sim refraction index)

$$\text{☞} \Delta_{\text{pl}} = -\frac{\omega_{\text{pl}}^2}{2\omega} \approx -1.1 \cdot 10^{-13} \omega_{\text{PeV}}^{-1} \text{ kpc}^{-1}$$

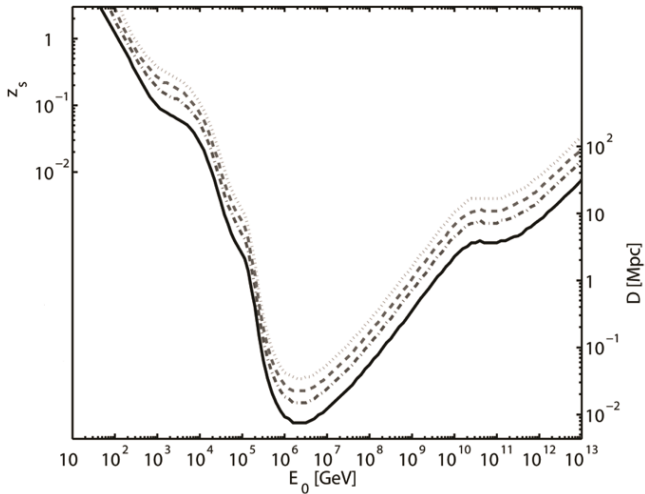
$$\text{☞} \Delta_{\text{CMB}} = \frac{44\pi^2\alpha^2}{2025} \frac{T_{\text{CMB}}^4}{m_e^4} \omega \approx 8 \cdot 10^{-2} \omega_{\text{PeV}} \text{ kpc}^{-1} \quad \omega \lesssim 100 \text{ TeV}$$

$$\text{☞} \Delta_{\text{QED}} = \frac{\alpha}{45\pi} \left(\frac{B}{B_{\text{cr}}} \right)^2 \omega \approx 0.15 \omega_{\text{PeV}} \left(\frac{B}{6\mu\text{G}} \right)^2 \text{ kpc}^{-1}$$

☞ This suppresses the oscillations by $e^{i\Delta_{\gamma\gamma}d}$ ($\Delta_{\gamma\gamma} \sim \Delta_{\text{pl}} + \Delta_{\text{CMB}} + \Delta_{\text{QED}}$)

$$\text{Beyond the PeV we find } P_{h \rightarrow \gamma} \simeq \frac{N_{\text{corr}} B^2}{M_{\text{P}}^2 \Delta_{\gamma\gamma}^2} \sim \frac{1}{\omega^2}$$

:: optical depth ::



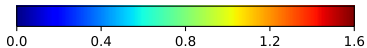
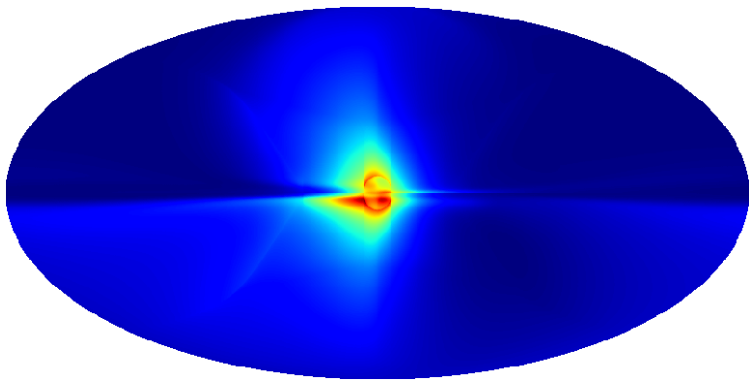


results



:: The flux ::

Photon flux [10^{-11} GeV/(cm² s sr)]

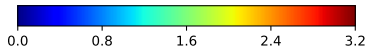
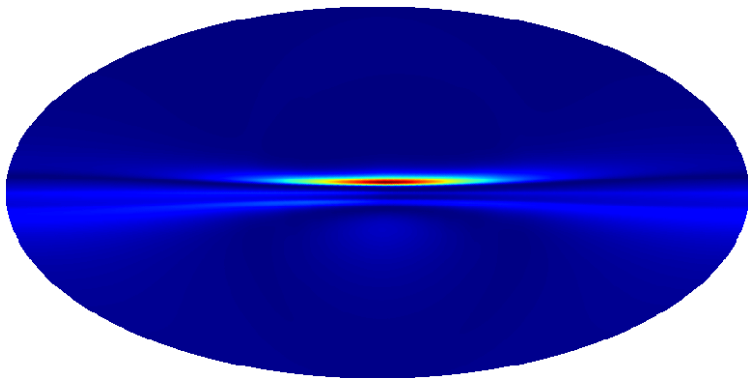


© Ramazanov, Samanta, Trenkler, FU, 2023

From the GMF of Jansson and Farrar, 2012

:: The flux ::

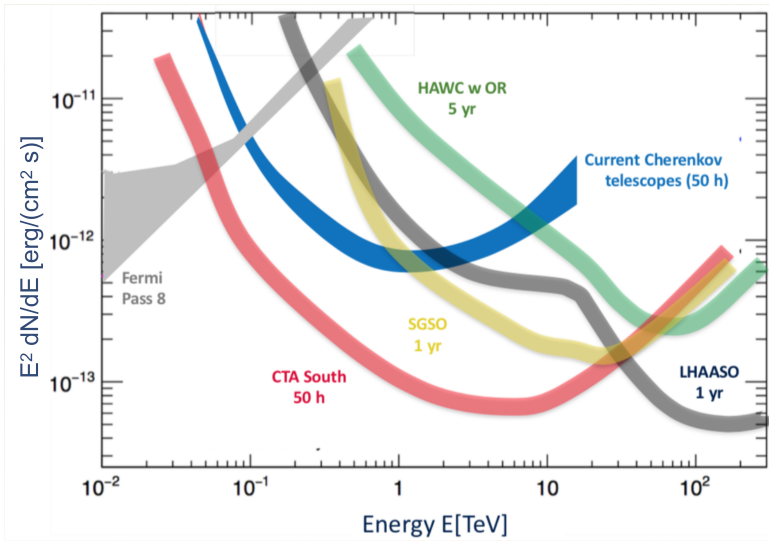
Photon flux [10^{-11} GeV/(cm² s sr)]



© Ramazanov, Samanta, Trenkler, FU, 2023

From the GMF of Pshirkov, Tinyakov, Kronberg, Newton-McGee, 2011

:: sensitivities ::



:: detect? ::

:: detect? ::



In either model the peak flux is around

$$\Phi_\gamma = \text{few} \times 10^{-11} \text{ GeV/cm}^2 \text{sec sr} \cdot \Omega_h h_0^2$$

:: detect? ::



In either model the peak flux is around

$$\Phi_\gamma = \text{few} \times 10^{-11} \text{ GeV/cm}^2\text{sec sr} \cdot \Omega_h h_0^2$$



At these energies, we know that LHAASO can do

$$\Phi_\gamma \approx 10^{-10} \text{ GeV/cm}^2\text{sec sr}$$

(Neronov, Semikoz '20)

:: detect? ::



In either model the peak flux is around

$$\Phi_\gamma = \text{few} \times 10^{-11} \text{ GeV/cm}^2\text{sec sr} \cdot \Omega_h h_0^2$$



At these energies, we know that LHAASO can do

$$\Phi_\gamma \approx 10^{-10} \text{ GeV/cm}^2\text{sec sr}$$

(Neronov, Semikoz '20)



So we would need $\Omega_h h_0^2 \approx 1$

We failed! But not miserably: $\Omega_h h_0^2 \approx 0.1 - 0.01$ next gen

:: detect? ::



In either model the peak flux is around

$$\Phi_\gamma = \text{few} \times 10^{-11} \text{ GeV/cm}^2\text{sec sr} \cdot \Omega_h h_0^2$$



At these energies, we know that LHAASO can do

$$\Phi_\gamma \approx 10^{-10} \text{ GeV/cm}^2\text{sec sr}$$

(Neronov, Semikoz '20)



So we would need $\Omega_h h_0^2 \approx 1$

We failed! But not miserably: $\Omega_h h_0^2 \approx 0.1 - 0.01$ next gen



The question now is: how to make these GWs?

We study the late decay of SHDM

:: produce ::

:: produce ::

Pre-recombination we can't have more than

$$\Omega_h h_0^2 \ll 10^{-6}$$

© Planck+BAO 2018

:: produce ::

Pre-recombination we can't have more than

$$\Omega_h h_0^2 \ll 10^{-6}$$

© Planck+BAO 2018

But post-recombination there can be much more

:: produce ::

Pre-recombination we can't have more than

$$\Omega_h h_0^2 \ll 10^{-6}$$

© Planck+BAO 2018

But post-recombination there can be much more

This is possible via late dark matter decay

$$S \rightarrow h + h$$

:: produce ::

Pre-recombination we can't have more than

$$\Omega_h h_0^2 \ll 10^{-6}$$

© Planck+BAO 2018

But post-recombination there can be much more

This is possible via late dark matter decay

$$S \rightarrow h + h$$

Two options:

:: produce ::

Pre-recombination we can't have more than

$$\Omega_h h_0^2 \ll 10^{-6}$$

© Planck+BAO 2018

But post-recombination there can be much more

This is possible via late dark matter decay

$$S \rightarrow h + h$$

Two options:

S decays completely by now (it is a small fraction of DM): $\Omega_h h_0^2 \approx f_{\text{DM}} \Omega_{\text{DM}} h_0^2$

:: produce ::

Pre-recombination we can't have more than

$$\Omega_h h_0^2 \ll 10^{-6}$$

© Planck+BAO 2018

But post-recombination there can be much more

This is possible via late dark matter decay

$$S \rightarrow h + h$$

Two options:

S decays completely by now (it is a small fraction of DM): $\Omega_h h_0^2 \approx f_{\text{DM}} \Omega_{\text{DM}} h_0^2$

S is still decaying (and is all of DM): $\Omega_h h_0^2 \approx 0.01$

:: Summary ::

:: Summary ::

☒ We look at the **graviton-photon** conversion in magnetic fields ☒

:: Summary ::

🔍 We look at the **graviton-photon** conversion in magnetic fields 🔍

🤖 We find a **sweet spot** at PeV energies 🤖

:: Summary ::

🔍 We look at the **graviton-photon** conversion in magnetic fields 🔍

🤖 We find a **sweet spot** at PeV energies 🤖

🤖 This is where the **conversion probability** is maximal 🤖

:: Summary ::

🔍 We look at the **graviton-photon** conversion in magnetic fields 🔍

🤖 We find a **sweet spot** at PeV energies 🤖

🤖 This is where the **conversion probability** is maximal 🤖

🔍 The Universe is **opaque** to PeV photons, so we look at the Galaxy only 🔍

:: Summary ::

🔗 We look at the **graviton-photon** conversion in magnetic fields 🔗

🤖 We find a **sweet spot** at PeV energies 🤖

🤖 This is where the **conversion probability** is maximal 🤖

🔗 The Universe is **opaque** to PeV photons, so we look at the Galaxy only 🔗

🤖 Realistic **relic abundances** of gravitons are out of reach, but not by much 🤖

:: Summary ::

🔗 We look at the **graviton-photon** conversion in magnetic fields 🔗

🤖 We find a **sweet spot** at PeV energies 🤖

🤖 This is where the **conversion probability** is maximal 🤖

🔗 The Universe is **opaque** to PeV photons, so we look at the Galaxy only 🔗

😊 Realistic **relic abundances** of gravitons are out of reach, but not by much 😊

😊 These gravitons come from **late** dark matter decay 😊

:: Summary ::

🔍 We look at the **graviton-photon** conversion in magnetic fields 🔍

🤖 We find a **sweet spot** at PeV energies 🤖

🤖 This is where the **conversion probability** is maximal 🤖

🔍 The Universe is **opaque** to PeV photons, so we look at the Galaxy only 🔍

😊 Realistic **relic abundances** of gravitons are out of reach, but not by much 😊

😊 These gravitons come from **late** dark matter decay 😊

👁️ In the future we could **detect** shimmering gravitons in the gamma-ray sky 👁️

:: Summary ::

🔗 We look at the **graviton-photon** conversion in magnetic fields 🔗

🤖 We find a **sweet spot** at PeV energies 🤖


🤖 This is where the **conversion probability** is maximal 🤖

🔗 The Universe is **opaque** to PeV photons, so we look at the Galaxy only 🔗

😊 Realistic **relic abundances** of gravitons are out of reach, but not by much 😊

😊 These gravitons come from **late** dark matter decay 😊

👻 In the future we could **detect** shimmering gravitons in the gamma-ray sky 👻



JCAP06 (2023) 019

Ideas welcome for how to make $h \rightarrow \gamma$ work for EeV+!