

Superradiant Instability of Magnetic Black Holes

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Introduction

- Gravitational wave detectors (LVK, LISA ...) probe strong gravity.
- Black holes become discovery engines: higher-derivative completions of GR, ultralight boson dark matter candidates, early universe...
- Crucial to understand BH physics: superradiance, GW spectrum ...
- *Today's talk*: magnetic BHs are multi-messenger sources via a new superradiant instability.

Magnetic Black Holes

- Magnetic monopoles are GUT predictions, and are ubiquitous in ST.

['t Hooft '74, Polyakov '74]

- Not observed nowadays: very hard to produce, inflation.

- May be found inside BHs, leading to *magnetic BHs*.

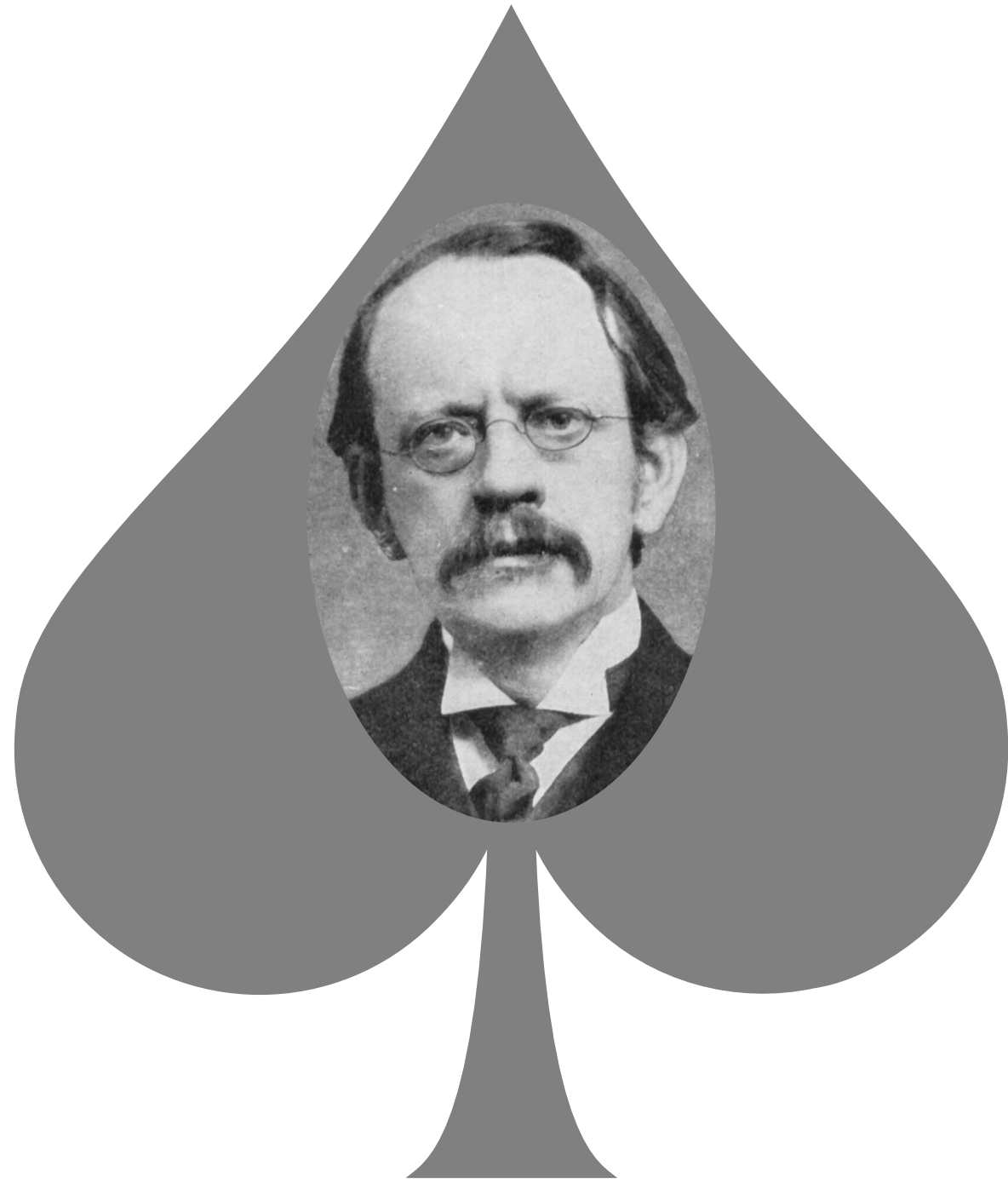
[Gibbons '77, Ortiz '92, Lee+ '92, Stojkovic+ '05, Maldacena '21...]

- What are their signatures?

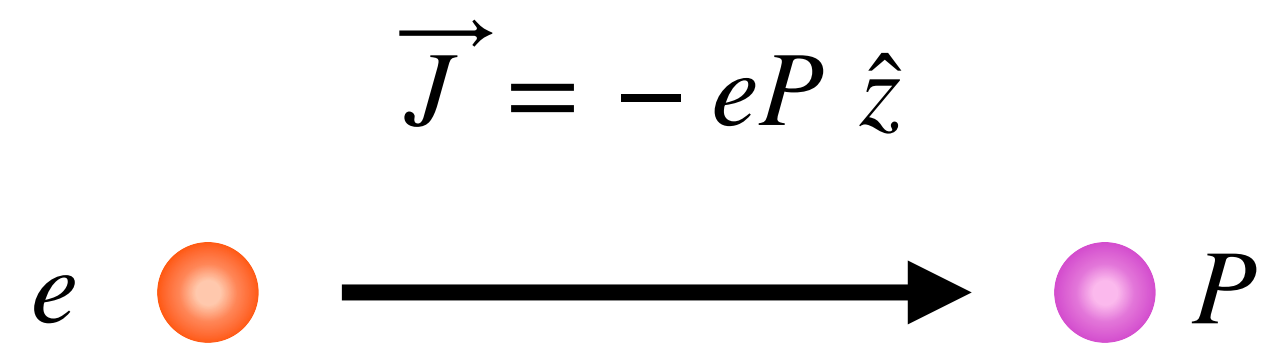
[Liu+ '20, Carullo+ '22, DP+ '23, de Felice+ '23, Kobayashi+ '23, Zhang+ '23, Wang+ '23...]

Magnetic charges are much more than a smooth deformation...

Lessons from Thomson, Dirac, Wu and Yang

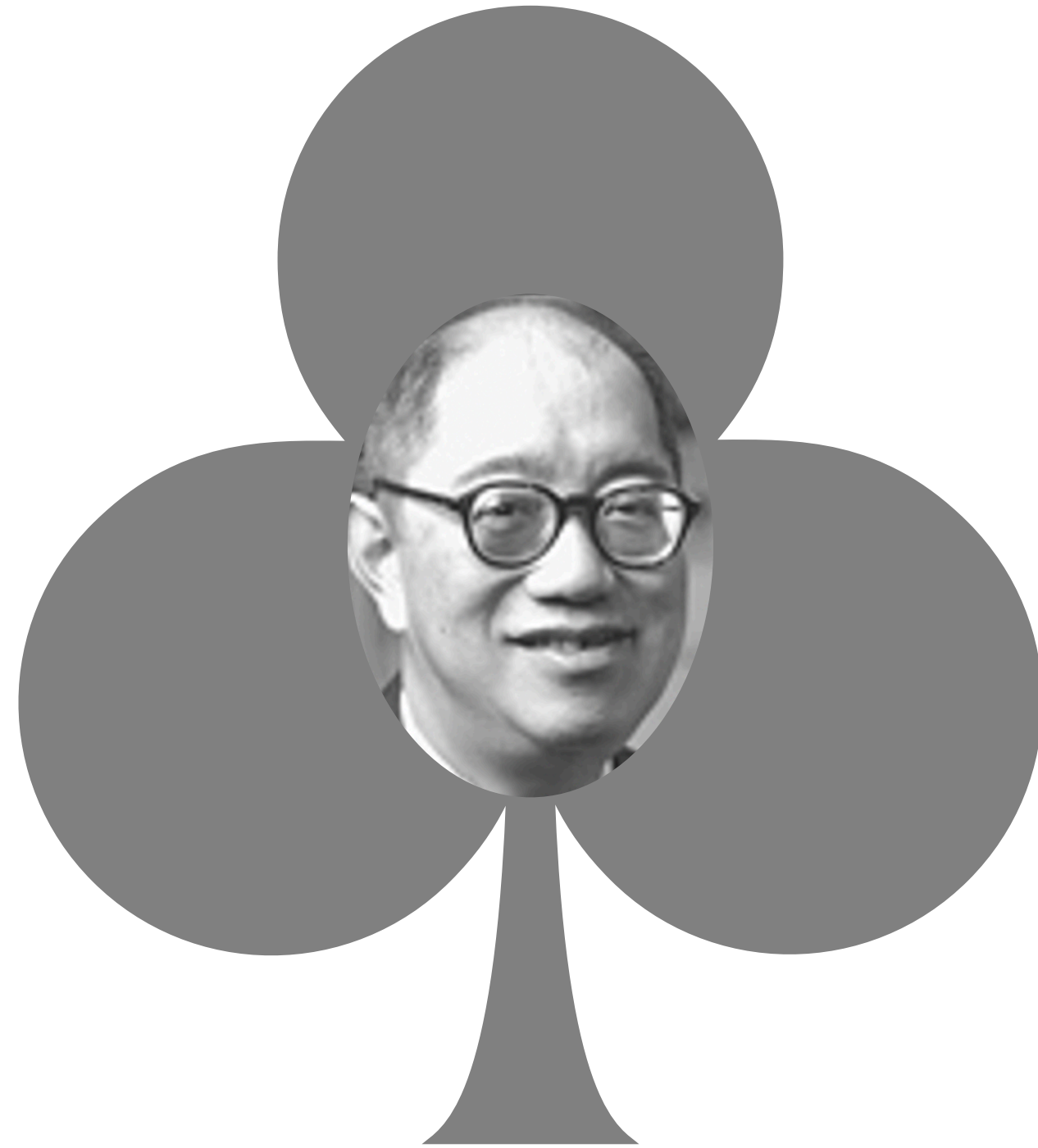


[Thomson 1904]

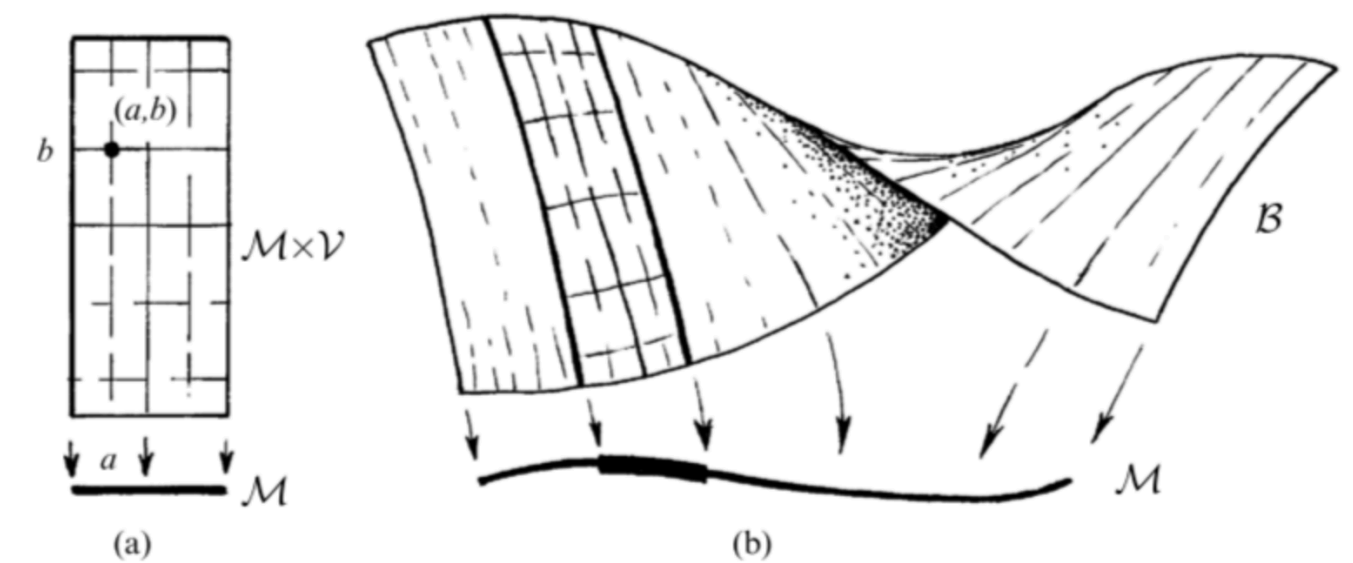


[Dirac '31]

$$2Pe = \pm N, \quad N = 0, 1, 2, 3, \dots$$



[Wu and Yang '76]

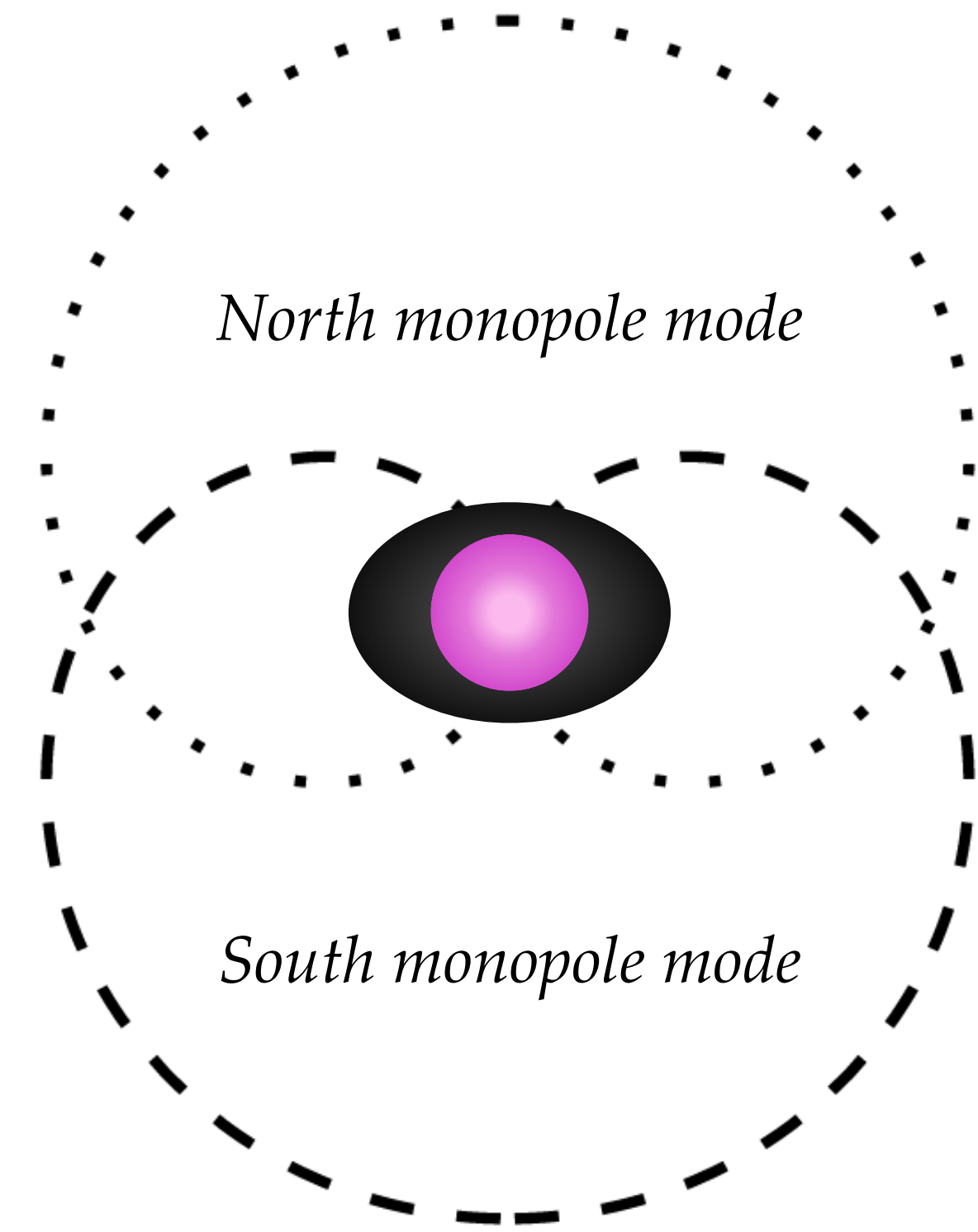


Monopole Spherical Harmonics

$$\mathbf{D} = \nabla + ieA \quad L_l = -i\epsilon_{ljk}x^j\mathbf{D}^k + eP\frac{x^l}{r}$$

$$L^2 Y_{q,\ell,m} = \ell(\ell + 1)Y_{q,\ell,m} \quad L_z Y_{q,\ell,m} = mY_{q,\ell,m} \quad (q \equiv -eP)$$

$$2|q| = 0, 1, 2, \dots \quad \ell = |q|, |q| + 1, \dots \quad m = -\ell, -\ell + 1, \dots, \ell$$

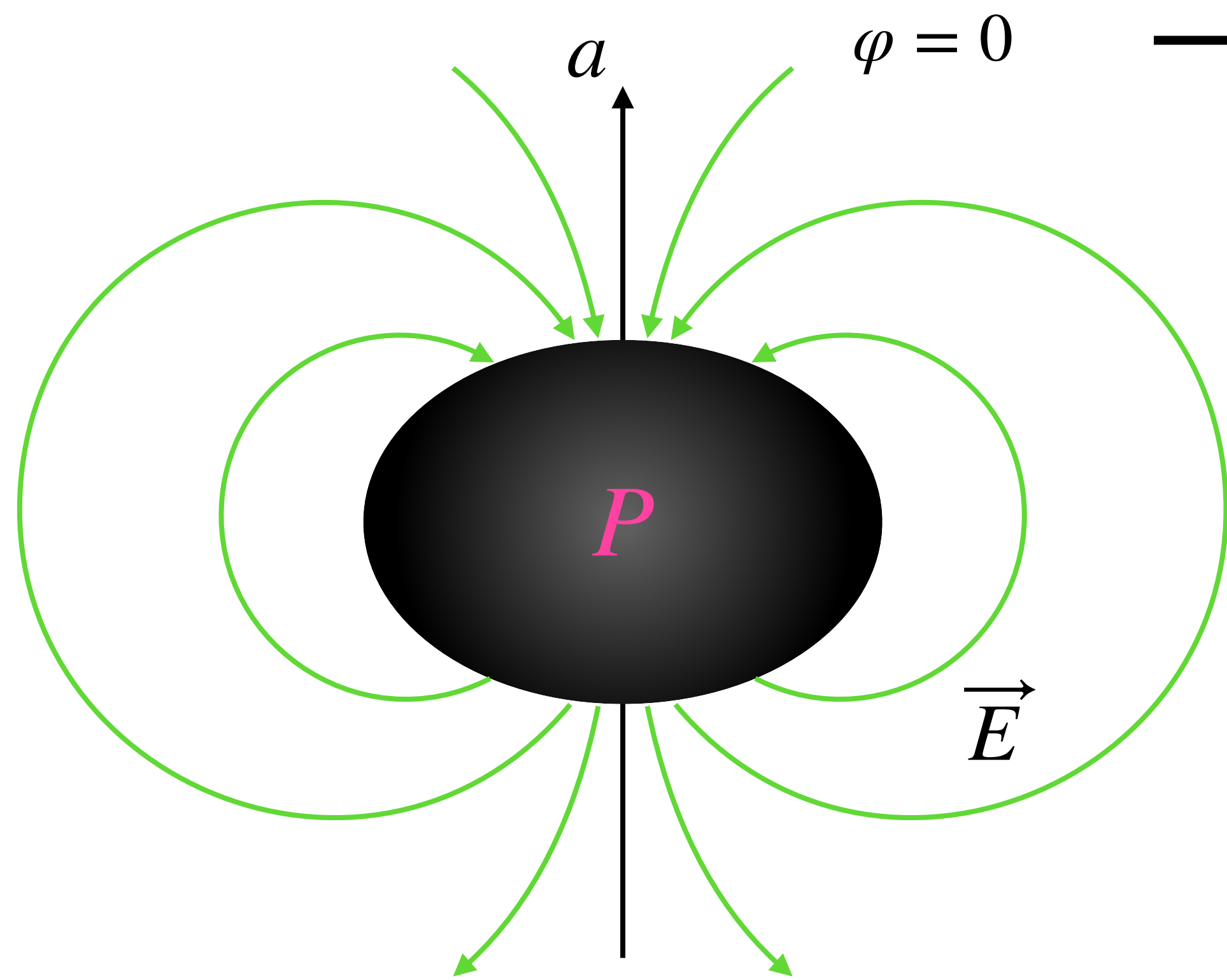


$$|q| = \ell = \pm m$$

$$Y_{q,|q|,\mp q}^{\pm} \sim (1 \pm \cos \theta)^{|q|}$$

Theory and Set Up

$$S[g, A, \varphi] = \frac{1}{16\pi} \int d^4x \sqrt{-g} (R - F^2) - \frac{1}{2} \int d^4x \sqrt{-g} (D_a \bar{\varphi} D^a \varphi + \mu^2 \bar{\varphi} \varphi) \quad (D = \nabla + ieA)$$



$\varphi = 0$



$$\varphi = e^{-i\omega t} e^{im\phi} \Psi(r, \theta)$$

(i) $\omega = \omega_R + i\omega_I$

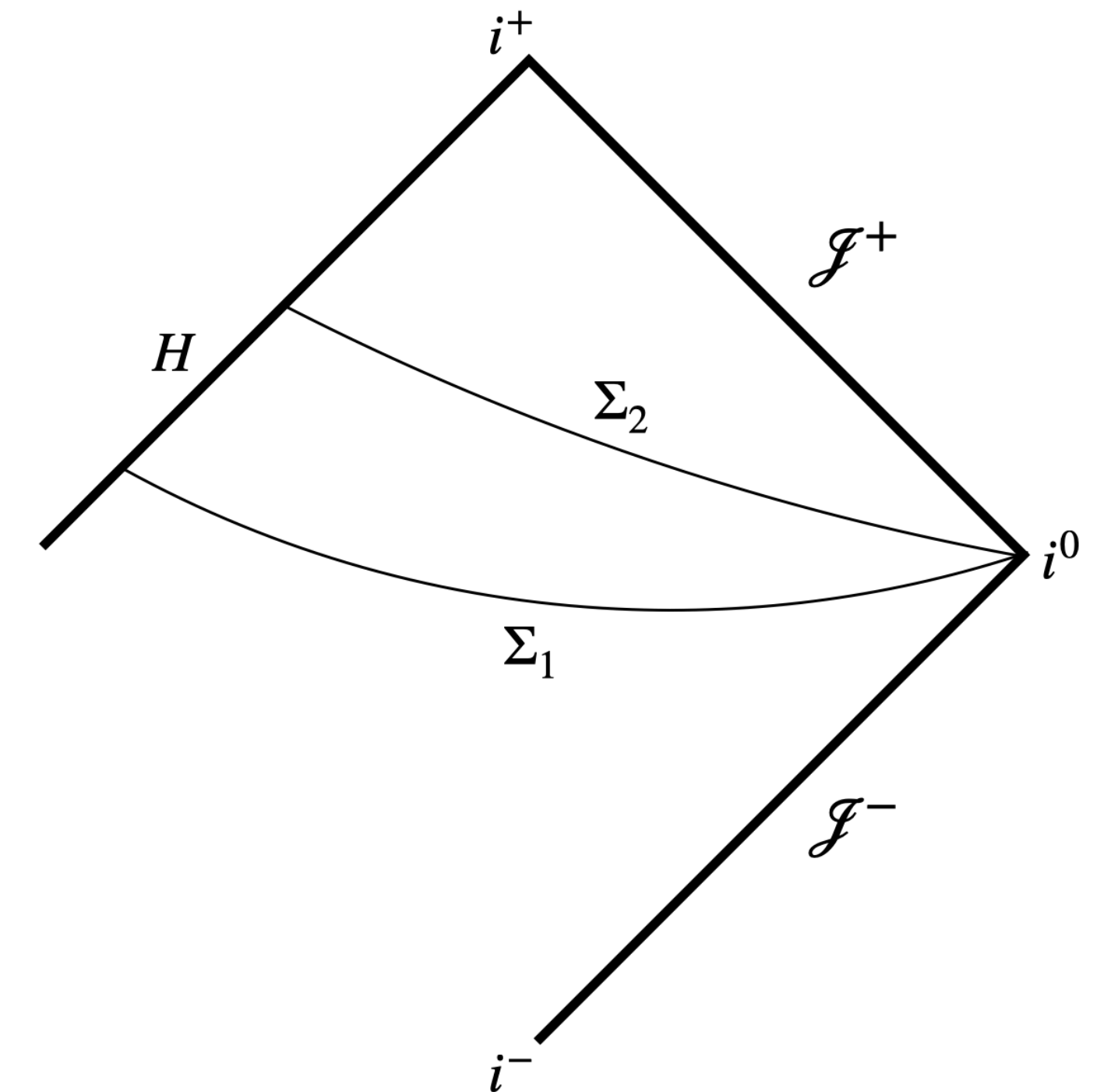
(ii) $\varphi \sim e^{-\sqrt{\mu^2 - \omega^2} r} \quad (r \rightarrow \infty)$

$$\Delta E = (\omega_I^2 + \omega_R \omega_*) \int_H e^{2\omega_I \nu} |\psi|^2 \tilde{\epsilon}$$

$$\Delta J = m\omega_* \int_H e^{2\omega_I \nu} |\psi|^2 \tilde{\epsilon}$$

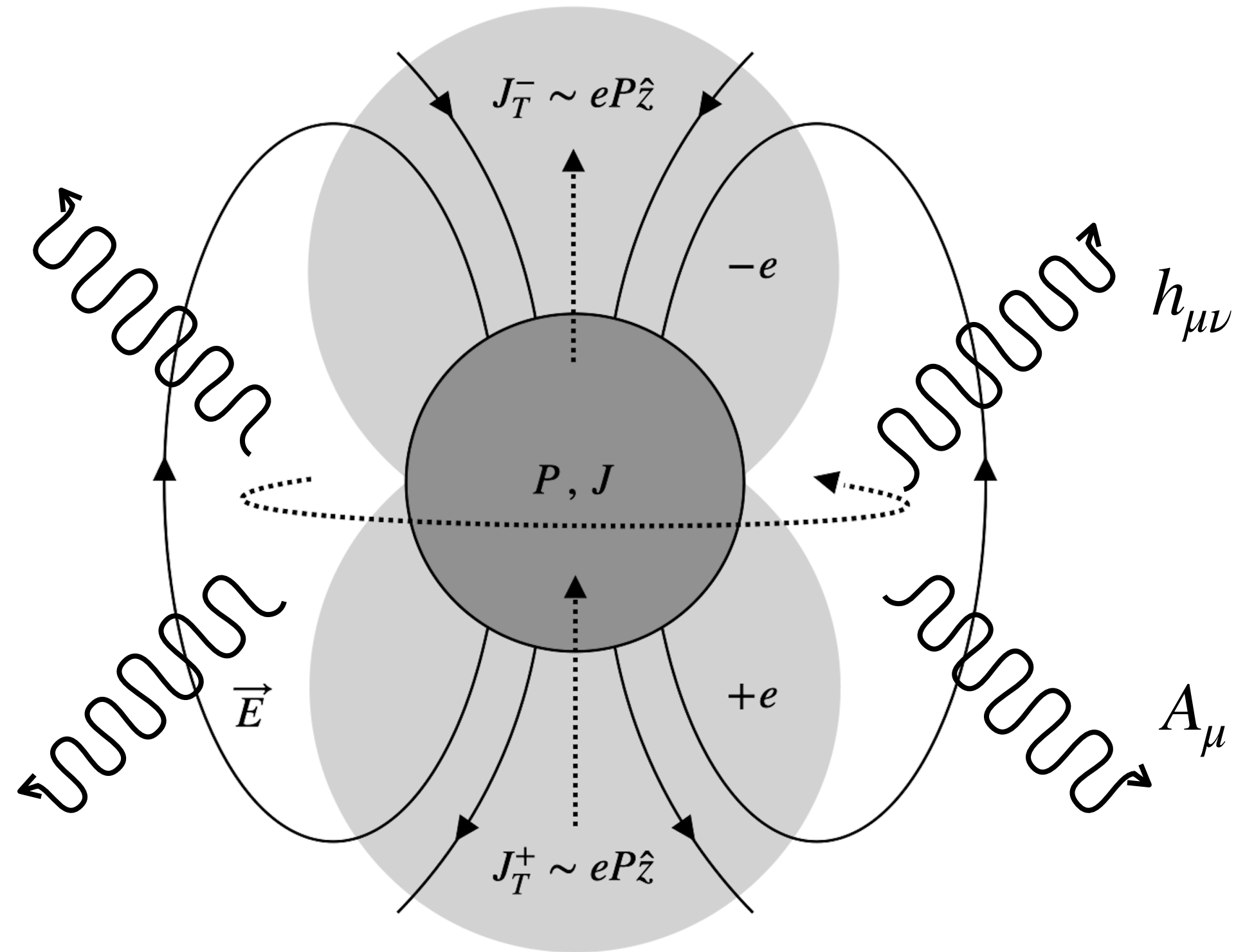
$$\Delta Q = -e\omega_* \int_H e^{2\omega_I \nu} |\psi|^2 \tilde{\epsilon}$$

$$\omega_* = \omega_R - m\Omega$$

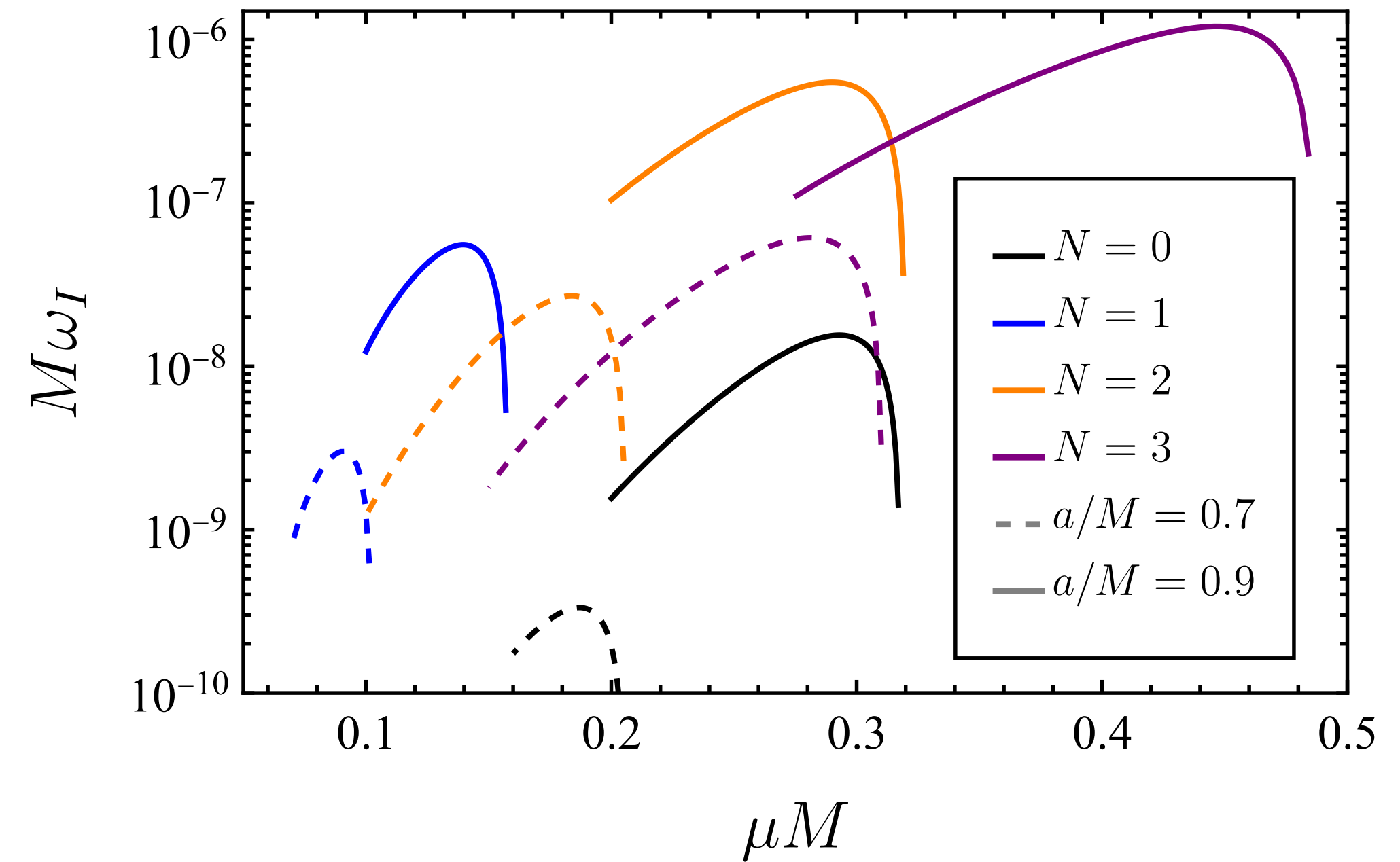


[Ortín and DP '22, Dyson and DP '23]

Superradiant Instability of Magnetic Black Holes



$$e/\mu = e/\mu(\pi^\pm) \quad N = 0, 1, 2, 3 \rightarrow (P/M \sim 10^{-19})$$



$$(a = 0.9, N = 0, \mu = \mu(\pi^0)) \longrightarrow M \sim 5 \times 10^{11} \text{kg}, \tau_{SI} \sim 10^{-16} \text{s} \longrightarrow \tau_{SI}/\tau(\pi^0) > 1$$



$$(a = 0.9, N = 3, \mu = \mu(\pi^\pm)) \longrightarrow M \sim 10^{12} \text{kg}, \tau_{SI} \sim 10^{-18} \text{s} \longrightarrow \tau_{SI}/\tau(\pi^\pm) = 10^{-10}$$



Conclusions

- Rotating magnetic black holes are superradiantly unstable.
- Novel mechanism for superradiant instability.
- With a *single* monopole the instability is relevant for SM fields.
- Rotating magnetic black holes are sources of (high-frequency) GWs and EM radiation.