

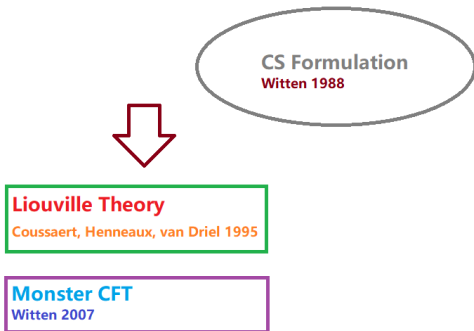
Modular Average and Weyl Anomaly in Two-Dimensional Schwarzian Theory

Chen-Te Ma (ISU)

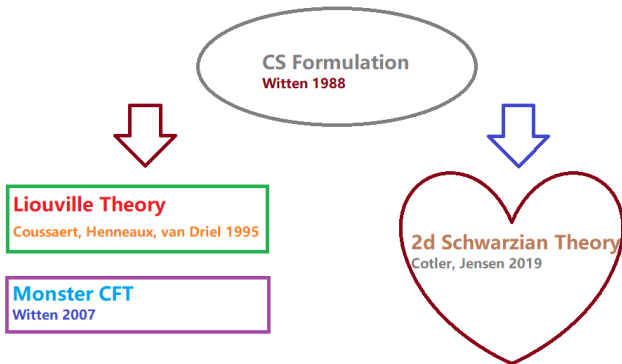
Xing Huang (Northwest U.)
arXiv:2303.17838 [hep-th]

February 27, 2024

History



History



Gauge Formulation of Einstein Gravity Theory

- SL(2)×SL(2) Chern-Simons formulation

$$\begin{aligned} & \frac{k}{2\pi} \int d^3x \operatorname{Tr} \left(A_t F_{r\theta} - \frac{1}{2} (A_r \partial_t A_\theta - A_\theta \partial_t A_r) \right) \\ & - \frac{k}{2\pi} \int d^3x \operatorname{Tr} \left(\bar{A}_t \bar{F}_{r\theta} - \frac{1}{2} (\bar{A}_r \partial_t \bar{A}_\theta - \bar{A}_\theta \partial_t \bar{A}_r) \right) \\ & - \frac{k}{4\pi} \int dt d\theta \operatorname{Tr} \left(\frac{E_t^+}{E_\theta^+} A_\theta^2 \right) + \frac{k}{4\pi} \int dt d\theta \operatorname{Tr} \left(\frac{E_t^-}{E_\theta^-} \bar{A}_\theta^2 \right) \end{aligned}$$

- $k \equiv 1/(4G_3)$, $1/l^2 \equiv -\Lambda$, where Λ is a cosmological constant
- bulk metric $g_{\mu\nu} = 2\operatorname{Tr}(e_\mu e_\nu)$, where $A - \bar{A} \sim e$, $A + \bar{A} \sim \omega$
- **boundary conditions** for the torus:

$$(E_\theta^+ A_t - E_t^+ A_\theta)|_{r \rightarrow \infty} = 0; (E_\theta^- \bar{A}_t - E_t^- \bar{A}_\theta)|_{r \rightarrow \infty} = 0 \quad (1)$$

Boundary Theory for Torus

- **two-dimensional Schwarzian theory** [Cotler, Jensen 2019]

$$\begin{aligned}
 & S_{\text{Gb}} \\
 = & \frac{k}{2\pi} \int dt d\theta \left(\frac{3}{2} \frac{(D_- \partial_\theta \mathcal{F})(\partial_\theta^2 \mathcal{F})}{(\partial_\theta \mathcal{F})^2} - \frac{D_- \partial_\theta^2 \mathcal{F}}{\partial_\theta \mathcal{F}} \right) \\
 & - \frac{k}{2\pi} \int dt d\theta \left(\frac{3}{2} \frac{(D_+ \partial_\theta \bar{\mathcal{F}})(\partial_\theta^2 \bar{\mathcal{F}})}{(\partial_\theta \bar{\mathcal{F}})^2} - \frac{D_+ \partial_\theta^2 \bar{\mathcal{F}}}{\partial_\theta \bar{\mathcal{F}}} \right), \quad (2)
 \end{aligned}$$

where

$$\mathcal{F} \equiv \frac{F}{E_\theta^+}, \quad \bar{\mathcal{F}} \equiv \frac{\bar{F}}{E_\theta^-}; \quad (3)$$

$$D_+ = \frac{1}{2} \partial_t - \frac{1}{2} \frac{E_t^-}{E_\theta^-} \partial_\theta, \quad D_- = \frac{1}{2} \partial_t - \frac{1}{2} \frac{E_t^+}{E_\theta^+} \partial_\theta \quad (4)$$

- measure $\int d\mathcal{F} d\bar{\mathcal{F}} (1/(\partial_\theta \mathcal{F} \partial_\theta \bar{\mathcal{F}}))$

Liouville Theory

- consider a general Weyl transformation $E^\pm \rightarrow \exp(\sigma)E^\pm$ on a **torus**
- boundary conditions become:

$$\begin{aligned}(E_\theta^+ A_t - E_t^+ A_\theta + E_t^+ A_\theta^2 J_2)|_{r \rightarrow \infty} &= 0; \\ (E_\theta^- \bar{A}_t - E_t^- \bar{A}_\theta + E_t^- \bar{A}_\theta^2 J_2)|_{r \rightarrow \infty} &= 0, \quad (5)\end{aligned}$$

where $A_\theta^2 = \bar{A}_\theta^2 = -\partial_t \sigma$

Liouville Theory

- consider a general Weyl transformation $E^\pm \rightarrow \exp(\sigma)E^\pm$ on a **torus**
- boundary conditions become:

$$\begin{aligned} (E_\theta^+ A_t - E_t^+ A_\theta + E_t^+ A_\theta^2 J_2)|_{r \rightarrow \infty} &= 0; \\ (E_\theta^- \bar{A}_t - E_t^- \bar{A}_\theta + E_t^- \bar{A}_\theta^2 J_2)|_{r \rightarrow \infty} &= 0, \end{aligned} \quad (5)$$

where $A_\theta^2 = \bar{A}_\theta^2 = -\partial_t \sigma$

- additional boundary term appears, and it is necessary for obtaining the Liouville theory **not** the same as in the claim of Ref. [Cotler, Jensen 2019]

$$-\frac{k}{4\pi} \int dt d\theta \operatorname{Tr}(A_\theta A_\theta + \bar{A}_\theta \bar{A}_\theta) + \frac{k}{8\pi} \int dt d\theta (A_\theta^2 A_\theta^2 + \bar{A}_\theta^2 \bar{A}_\theta^2) \quad (6)$$

Application

- implement the averaging of a modular group to compute the Rényi-2 mutual information for disjoint two-intervals
- non-perturbative effect kills the phase transition

Modular Averaging

- partition function of 2d Schwarzian theory on a **torus** is [Cotler, Jensen 2019]

$$Z_T(\tau) = |q|^{-\frac{c}{12}} \frac{1}{\prod_{n=2}^{\infty} |1 - q^n|^2}, \quad q \equiv e^{2\pi i \tau}, \quad c = \frac{3}{2G_3} + 13, \quad (7)$$

where τ is a complex structure

- summation is over the **SL(2, \mathbb{Z})** group (or modular group) in the boundary theory, which is equivalent to the path integration for all asymptotic **AdS₃** boundary [Maloney, Witten 2010]

$$Z_M(\tau) = \sum_{c_1, d_1; (c_1, d_1)=1} Z_T\left(\frac{a_1\tau + b_1}{c_1\tau + d_1}\right), \quad (8)$$

where $a_1 d_1 - b_1 c_1 = 1$, $a_1, b_1, c_1, d_1 \in \mathbb{Z}$

Rényi-2 Mutual Information

- Rényi entropy of order q is

$$R^{(q)} \equiv \frac{\ln \text{Tr} \rho^q}{1-q} = \frac{\ln Z^{(q)} - q \ln Z^{(1)}}{1-q}, \quad (9)$$

where ρ is a reduced density matrix, and $Z^{(q)}$ is a q -sheet partition function

- Rényi mutual information of order q is

$$I_{[u_1, v_1], [u_2, v_2]}^{(q)} = R_{[u_1, v_1]}^{(q)} + R_{[u_2, v_2]}^{(q)} - R_{[u_1, v_1] \cup [u_2, v_2]}^{(q)} \quad (10)$$

- Rényi-2 mutual information is [Headrick 2010]

$$I^{(2)} = \ln Z_T(\tau) + c \ln \left(\frac{\theta_2(\tau)}{2\eta(\tau)} \right) \quad (11)$$

Rényi-2 Mutual Information

- Rényi entropy of order q is

$$R^{(q)} \equiv \frac{\ln \text{Tr} \rho^q}{1-q} = \frac{\ln Z^{(q)} - q \ln Z^{(1)}}{1-q}, \quad (9)$$

where ρ is a reduced density matrix, and $Z^{(q)}$ is a q -sheet partition function

- Rényi mutual information of order q is

$$I_{[u_1, v_1], [u_2, v_2]}^{(q)} = R_{[u_1, v_1]}^{(q)} + R_{[u_2, v_2]}^{(q)} - R_{[u_1, v_1] \cup [u_2, v_2]}^{(q)} \quad (10)$$

- Rényi-2 mutual information is [Headrick 2010]

$$I^{(2)} = \ln Z_T(\tau) + c \ln \left(\frac{\theta_2(\tau)}{2\eta(\tau)} \right) \quad (11)$$

- modular averaging is equivalent to replacing Z_T with Z_M

Rényi-2 Mutual Information

$$c = 64$$

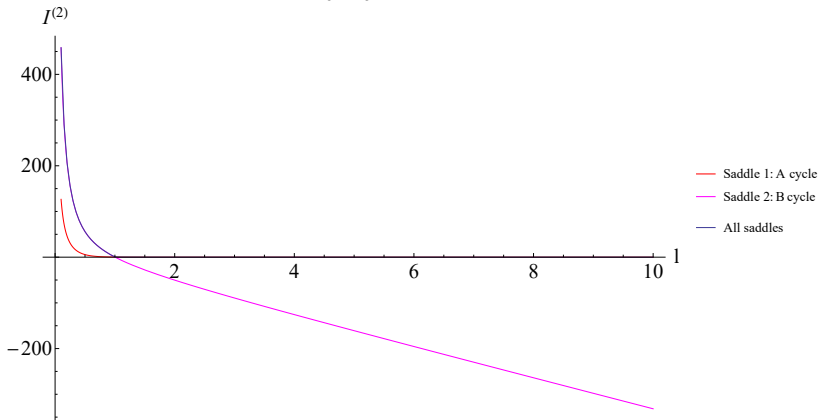


Figure: We show the Rényi-2 mutual information from summing all saddle points, A-cycle, and B-cycle with $c = 64$. Not all saddle points decay to zero when $l \equiv -i\tau \rightarrow \infty$ (like B-cycle), where τ is a complex structure.

Logarithmic Partition Function

$$c = 64$$

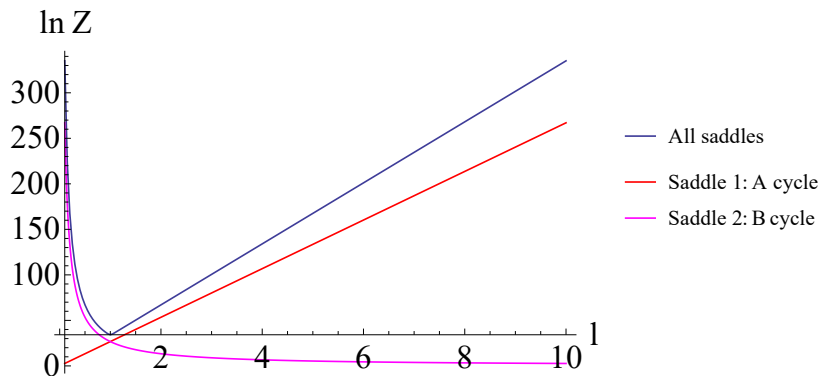


Figure: The logarithmic partition function is continuous for l .

First-Order Derivative

$$c = 64$$

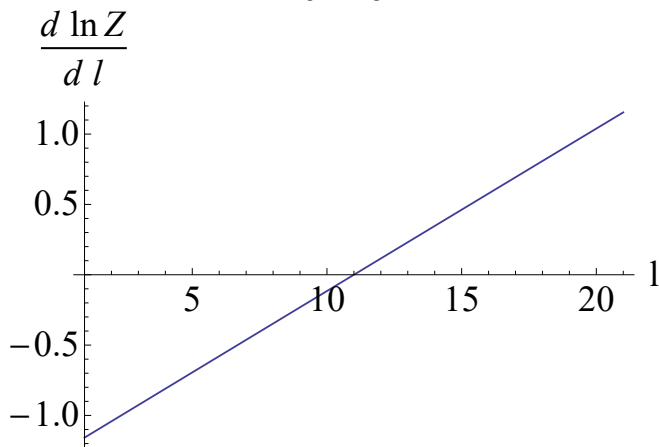


Figure: The first-order derivative of the logarithmic partition function is a continuous function for l .

Discussion and Conclusion

- agrees with the **Liouville Theory** for a torus manifold

Discussion and Conclusion

- agrees with the **Liouville Theory** for a torus manifold
- implement the averaging of a modular group to compute the **Rényi-2 mutual information** for disjoint two-intervals
- non-perturbative effect **kills** the phase transition

Discussion and Conclusion

- agrees with the **Liouville Theory** for a torus manifold
- implement the averaging of a modular group to compute the **Rényi-2 mutual information** for disjoint two-intervals
- non-perturbative effect **kills** the phase transition
- **boundary theory=resummation of perturbative gravity**

Thank you!

Cylinder \rightarrow Sphere

- respects the result from the **Cylinder** manifold to **Sphere** Manifold

$$e^a \rightarrow \exp(\sigma(t))e^a = \operatorname{sech}(\psi)e^a \quad (12)$$

but additional **finite** term appears

$$\frac{k}{2\pi} \int_{-\infty}^{\infty} d\psi \int_0^{2\pi} d\theta \partial_{\psi}^2 \sigma = -2k \quad (13)$$