

# Hyperuniformity in Chiral Active Fluids

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Yuta Kuroda  
Nagoya University, Japan

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Advances in Fluctuating Hydrodynamics

# Outline

- ▶ **Introduction**
  - Active Matter
  - Fluctuations in Active Matter
  - Hyperuniformity in Chiral Active Fluids
- ▶ **Theory for HU in 2D Chiral Active Fluid**
- ▶ **Density Fluctuations in 3D Chiral Active Fluid**
  - Microscopic Model and Hydrodynamic Equations
  - Numerical Simulation
- ▶ **Summary**

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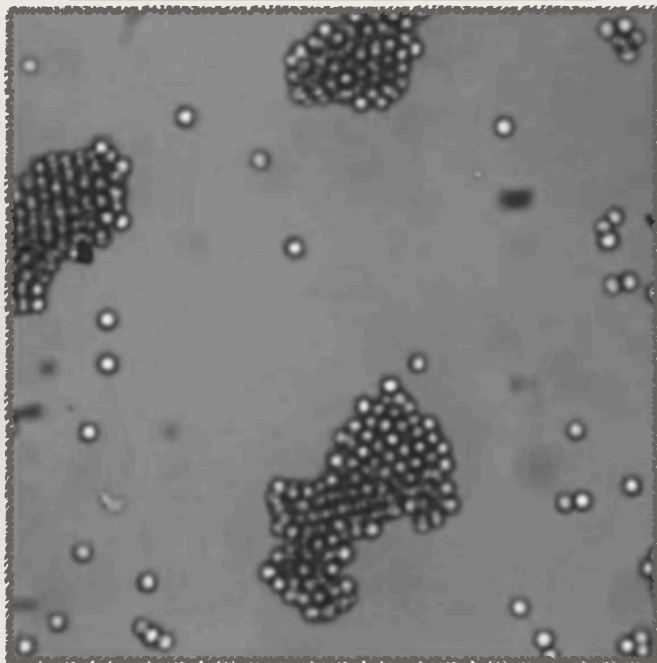
- ▶ **Introduction**
  - **Active Matter**
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- ▶ **Theory for HU in 2D Chiral Active Fluid**
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# Introduction

## ○ What is Active Matter ?

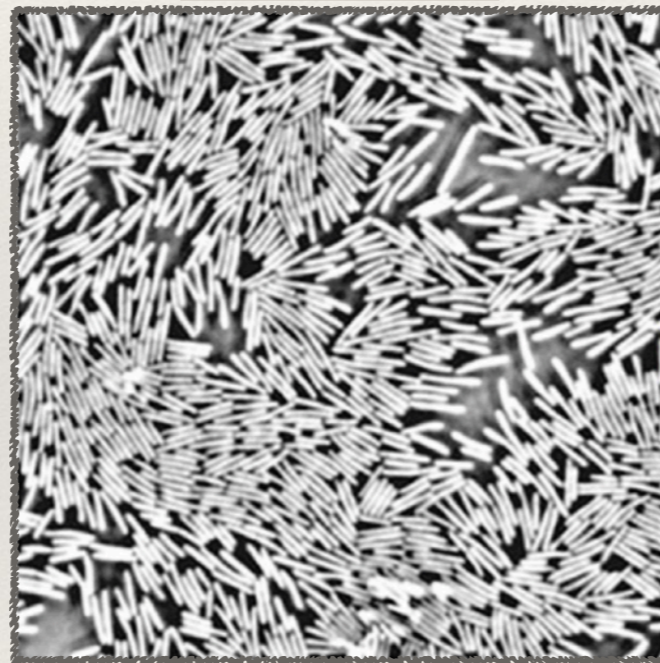
**Active matter  $\equiv$  any collection of self-propelled objects**

**Janus particle**



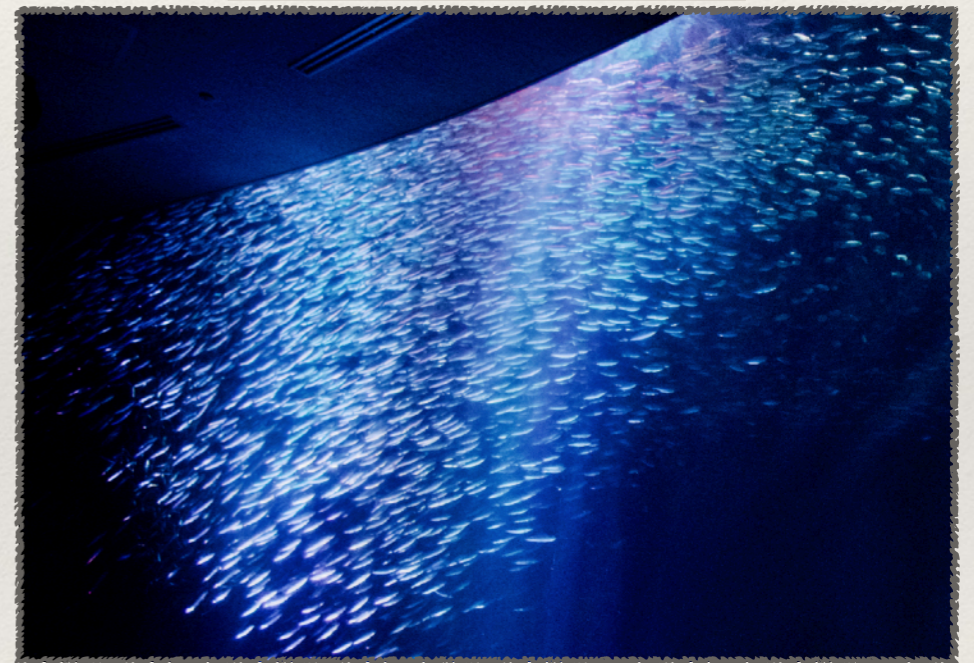
I. Buttinoni, *et al.*, (2015)

**Bacteria**



H. P. Zhang, *et al.*, (2010)

**School of fish**

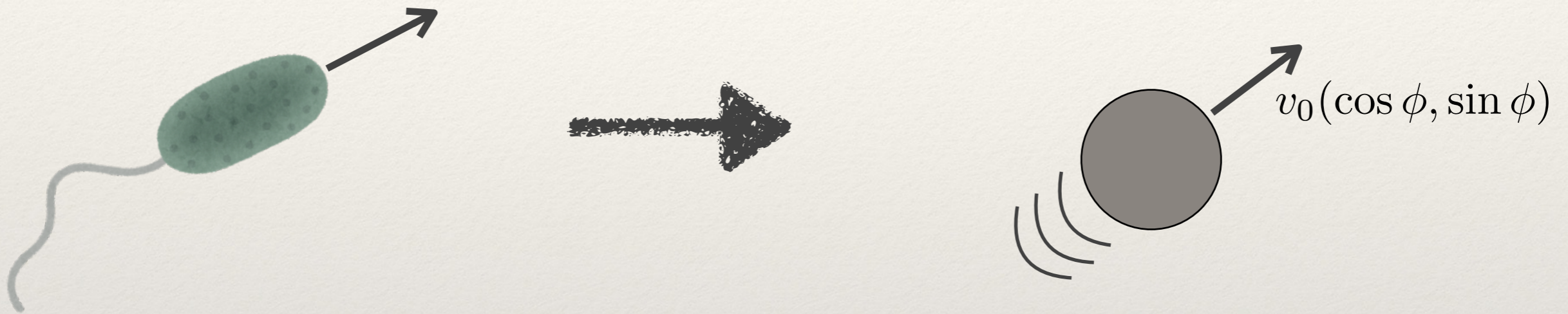


Taken at Nagoya public aquarium

# Introduction

## Modeling of Active Matter

- ☑ Mostly active particles are modeled as straight swimmers



### Active Brownian particle

$$\dot{\mathbf{r}}_j(t) = \underbrace{\mu \mathbf{F}_j^{\text{int}}}_{\text{interaction}} + \underbrace{v_0 \mathbf{e}(\phi_j(t))}_{\text{self-propelling}}$$

$$\dot{\phi}_j(t) = \sqrt{2D} \underbrace{\eta_j(t)}_{\text{noise}}$$

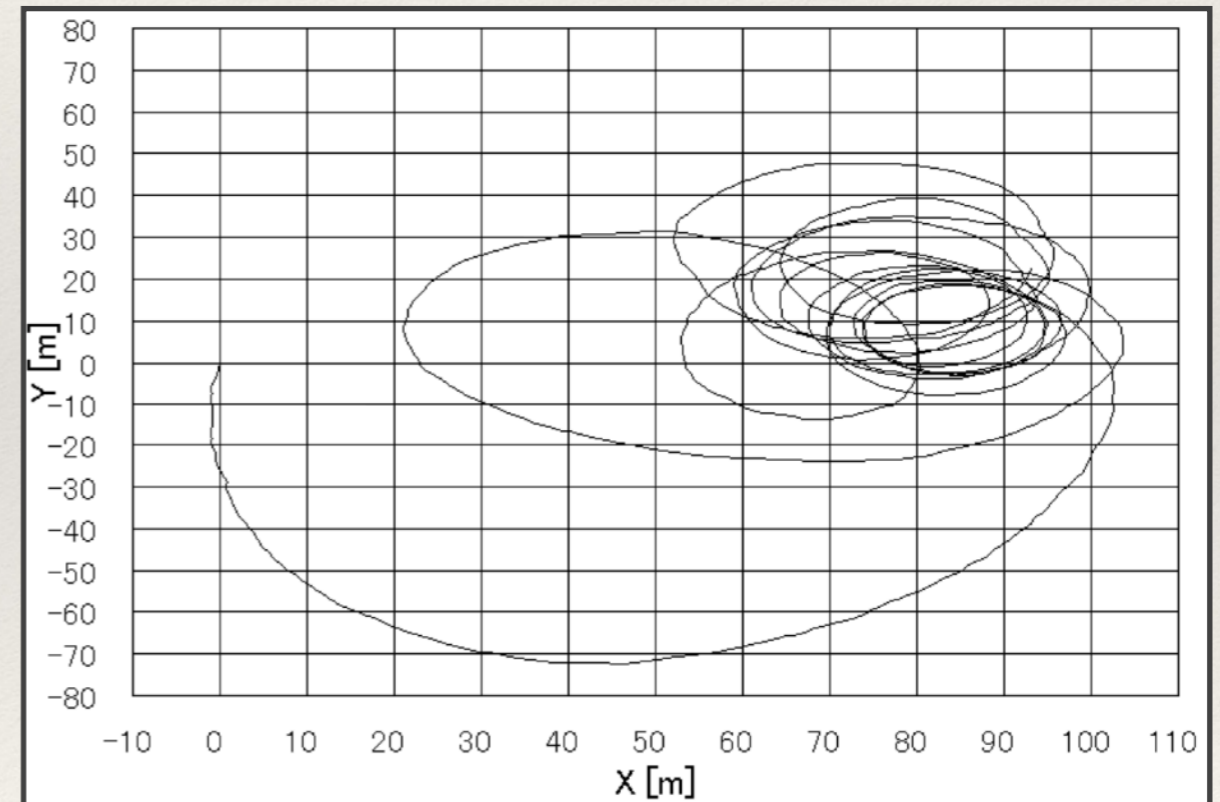
- ☑ ABP performs almost straight motion for the time scale  $1/D$

# Introduction

## Chiral Active Matter

- ☑ Most active objects *can't* go straight!

A walking guy who can't see anything draws a circular trajectory



T. Obata *et al.*, (2005)

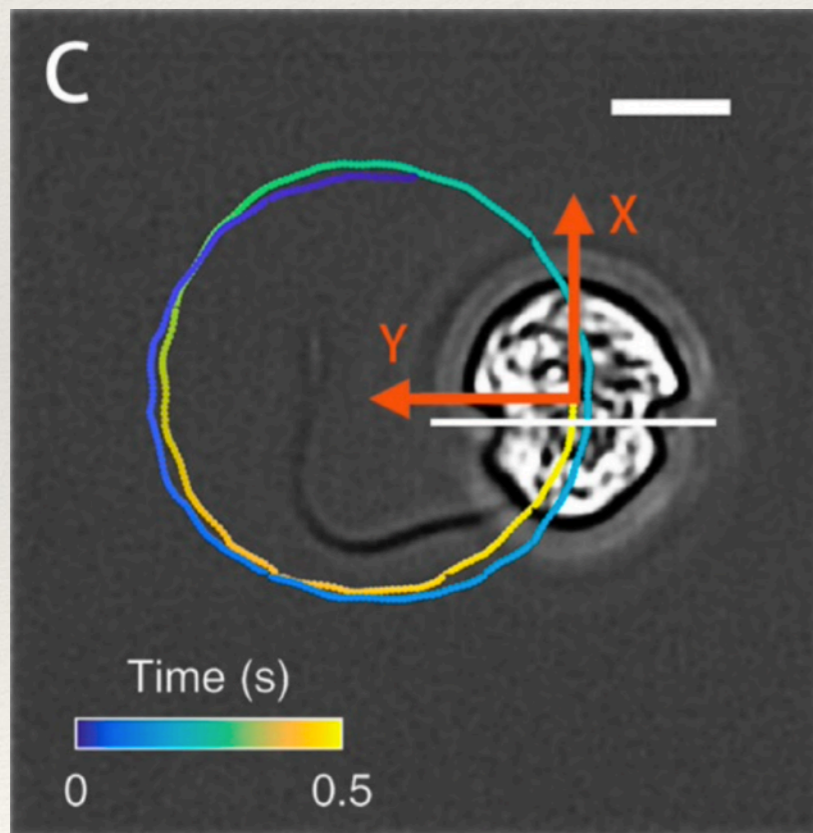
# Introduction

## Chiral Active Matter

### Chiral active matter:

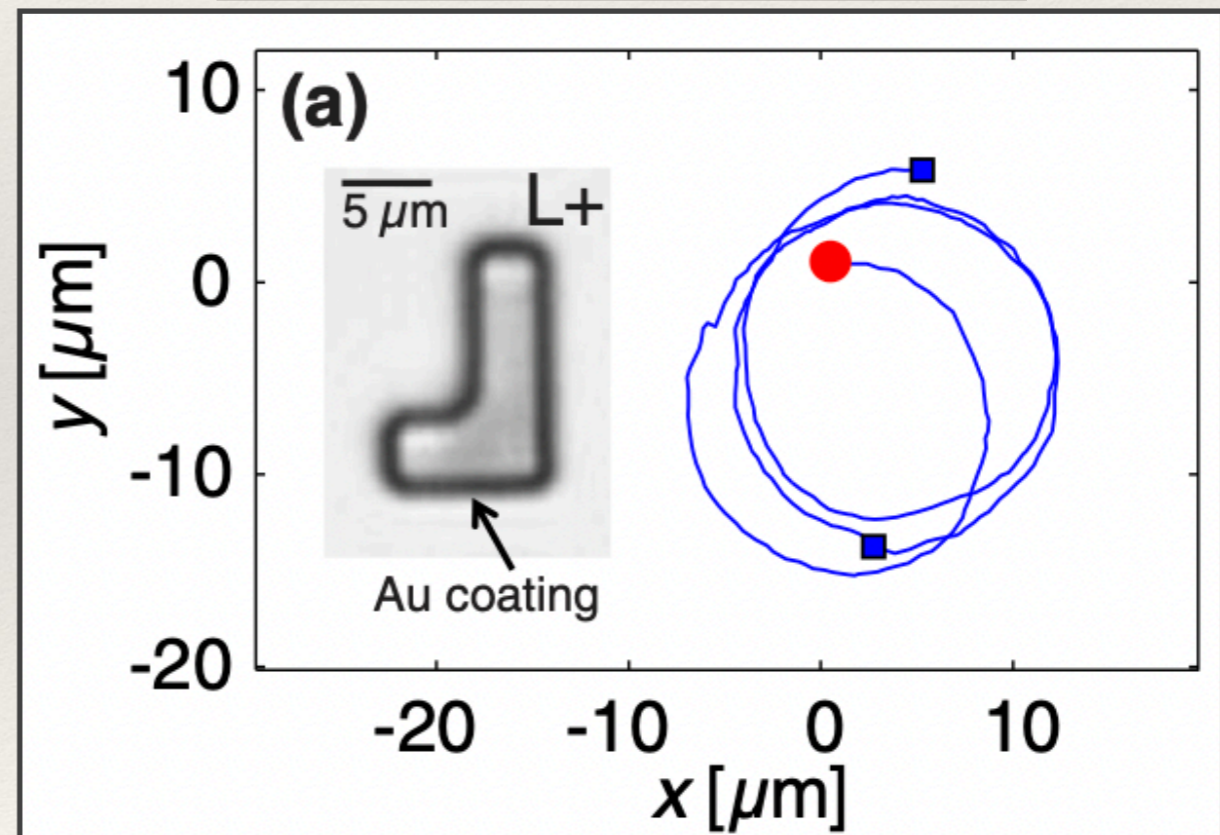
A collection of swimmers whose motion don't have mirror symmetry

Marine algae



M. Huang *et al.*, (2021)

L-shaped particle



F. Kümmel *et al.*, (2013)

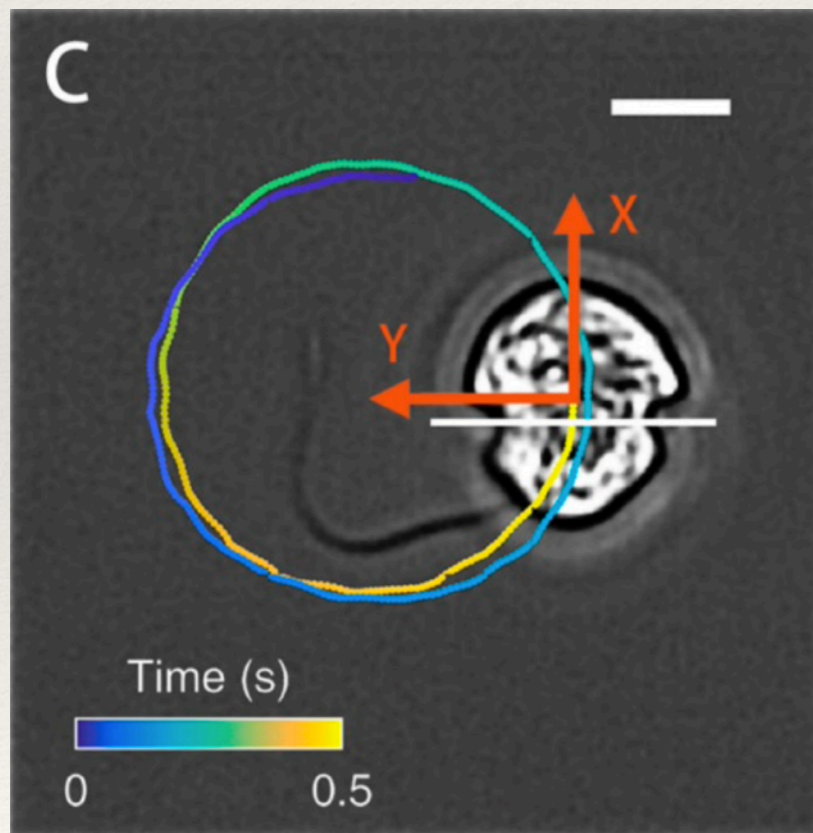
# Introduction

## Chiral Active Matter

- Collective phenomenon we're interested in here:

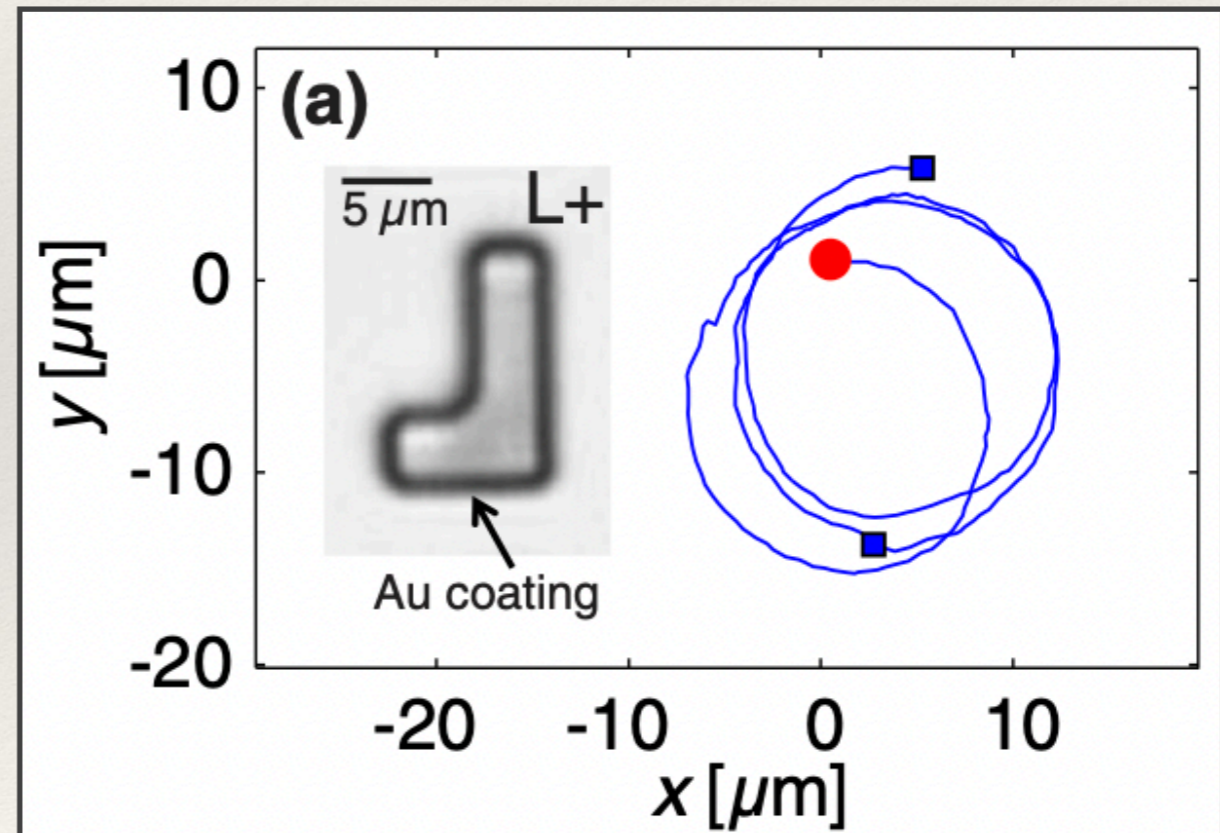
### *Density fluctuations in chiral active fluids*

Marine algae



M. Huang *et al.*, (2021)

L-shaped particle



F. Kümmel *et al.*, (2013)



# Introduction

## Fluctuations in Active Matter: Giant Number Fluctuations (GNF)

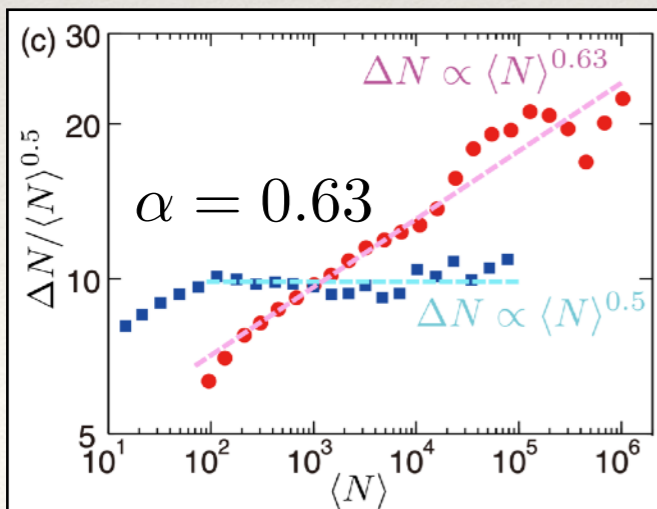
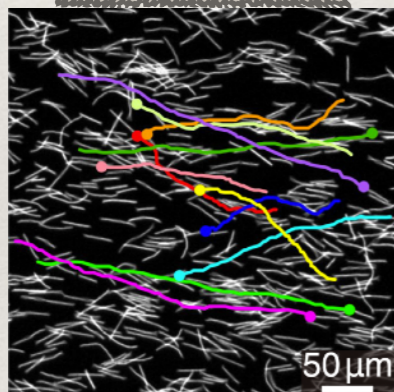
● Number fluctuation  $\Delta N = \sqrt{\langle (N - \langle N \rangle)^2 \rangle}$

☑ **GNF:**  $\Delta N \propto \langle N \rangle^\alpha$  with  $\alpha > 0.5$  (in equilibrium systems,  $\alpha = 0.5$ )

→  $S(q) = \frac{1}{N} \langle \delta\rho(\mathbf{q})\delta\rho^*(\mathbf{q}) \rangle \sim \frac{1}{q^\beta}$  with  $\beta = 4\alpha - 2$

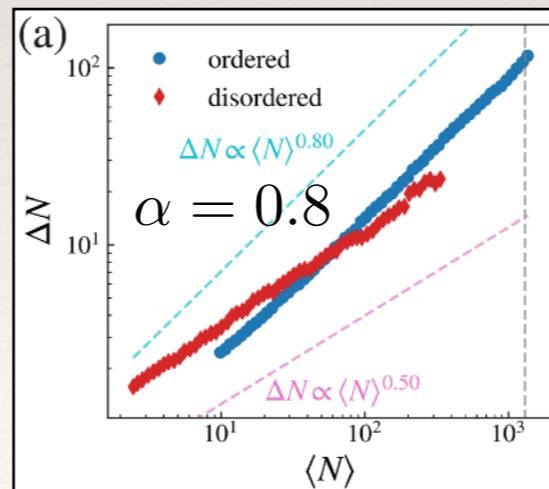
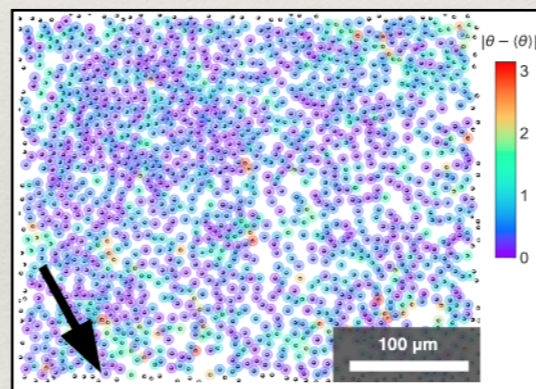
**Increases in density fluctuations!**

Bacteria



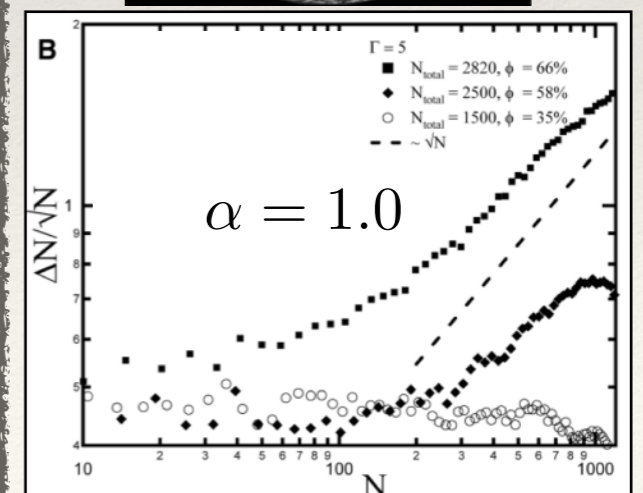
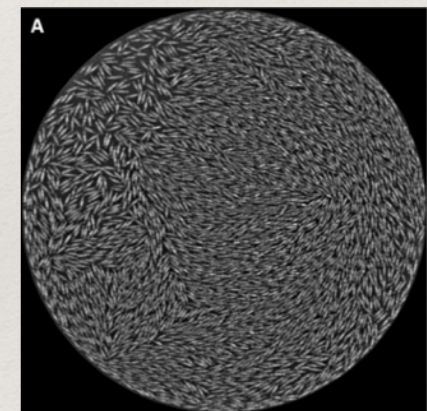
D. Nishiguchi *et al.*, (2017)

Janus particles



J. Iwasawa *et al.*, (2021)

Shaken rods



V. Narayan *et al.*, (2007)

# Introduction

## Fluctuations in Active Matter: Giant Number Fluctuations (GNF)

● Number fluctuation  $\Delta N = \sqrt{\langle (N - \langle N \rangle)^2 \rangle}$

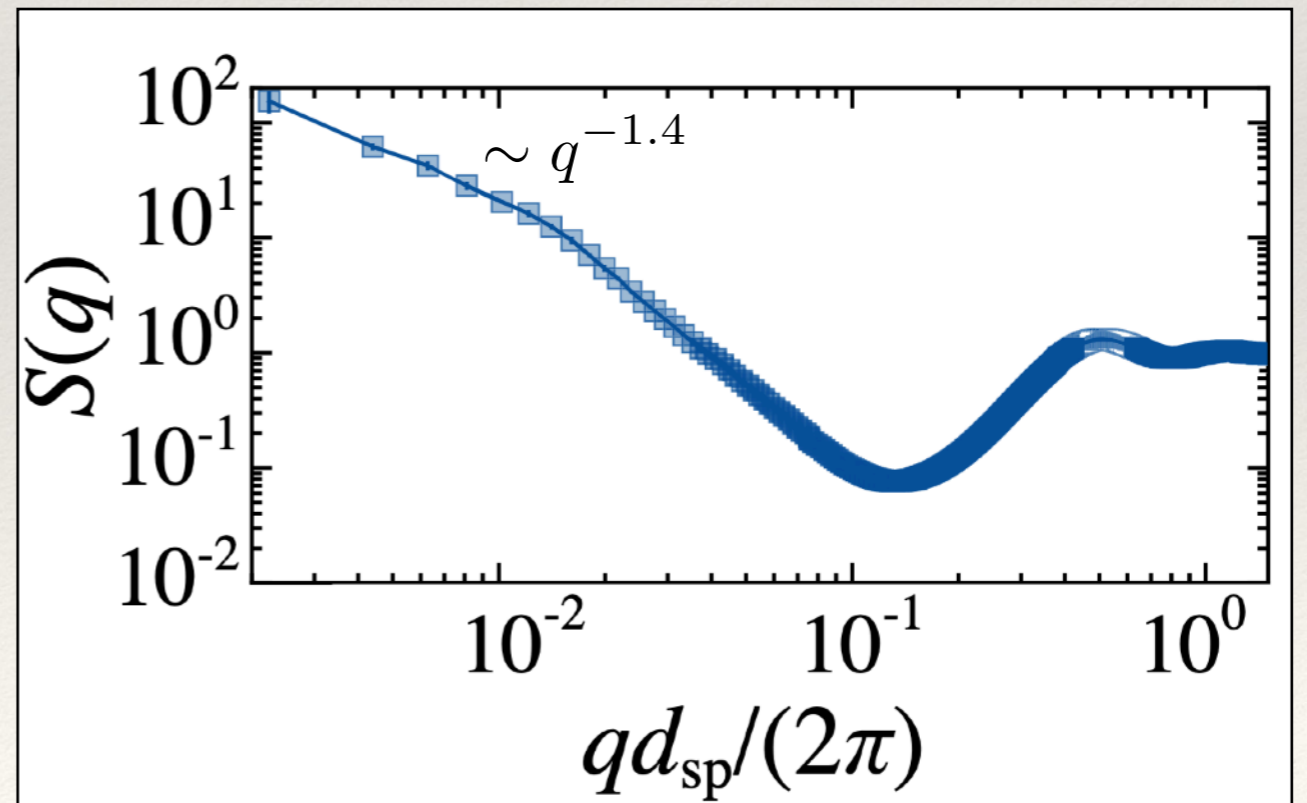
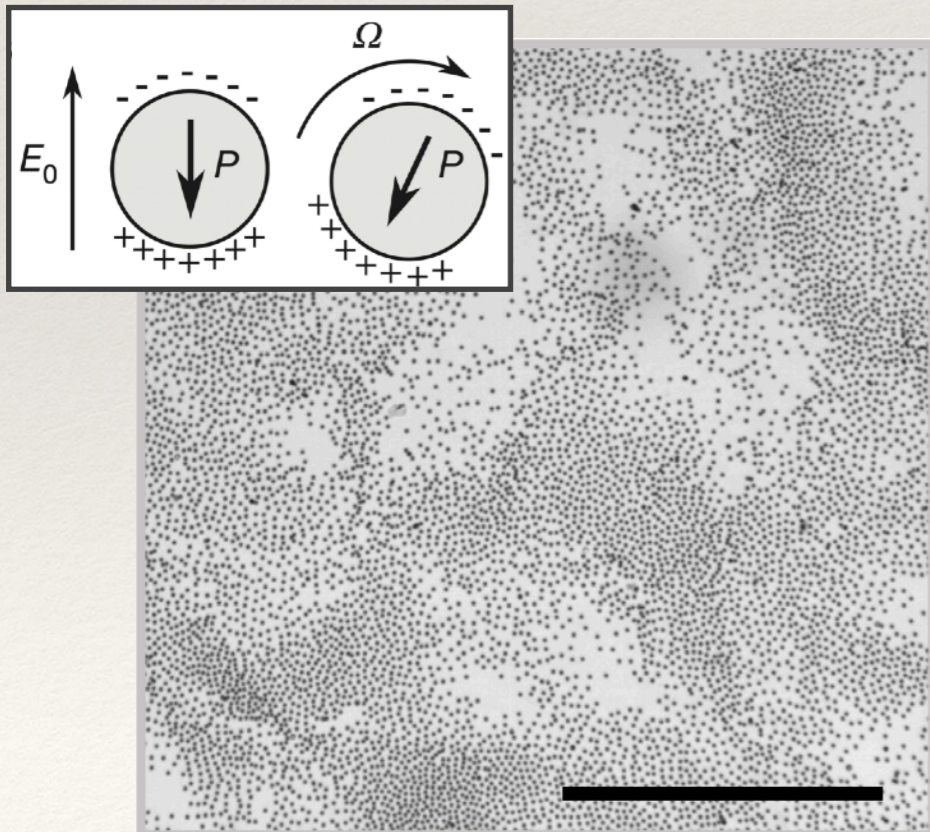
✓ **GNF:**  $\Delta N \propto \langle N \rangle^\alpha$  with  $\alpha > 0.5$  (in equilibrium systems,  $\alpha = 0.5$ )

→  $S(q) = \frac{1}{N} \langle \delta\rho(\mathbf{q})\delta\rho^*(\mathbf{q}) \rangle \sim \frac{1}{q^\beta}$  with  $\beta = 4\alpha - 2$

**Increases in density fluctuations!**

### Flocking state of the Quincke rollers

B. Zhang *et al.*, (2022)

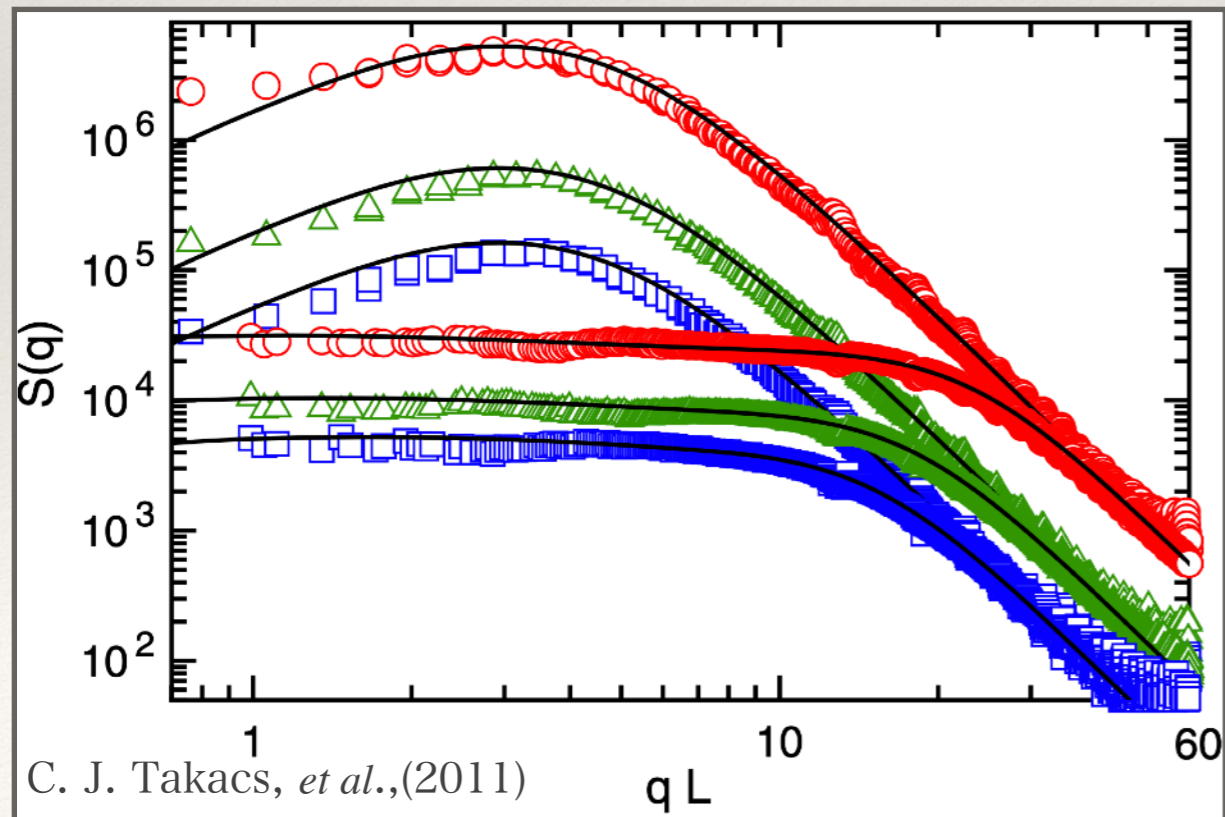
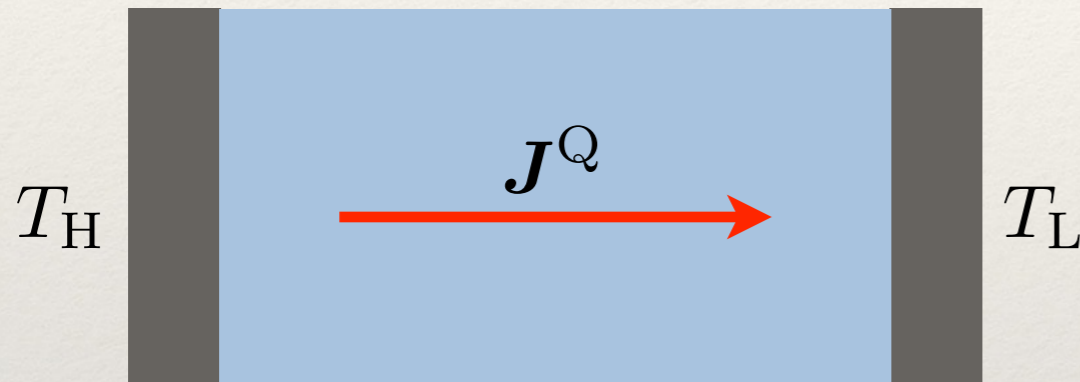


# Introduction

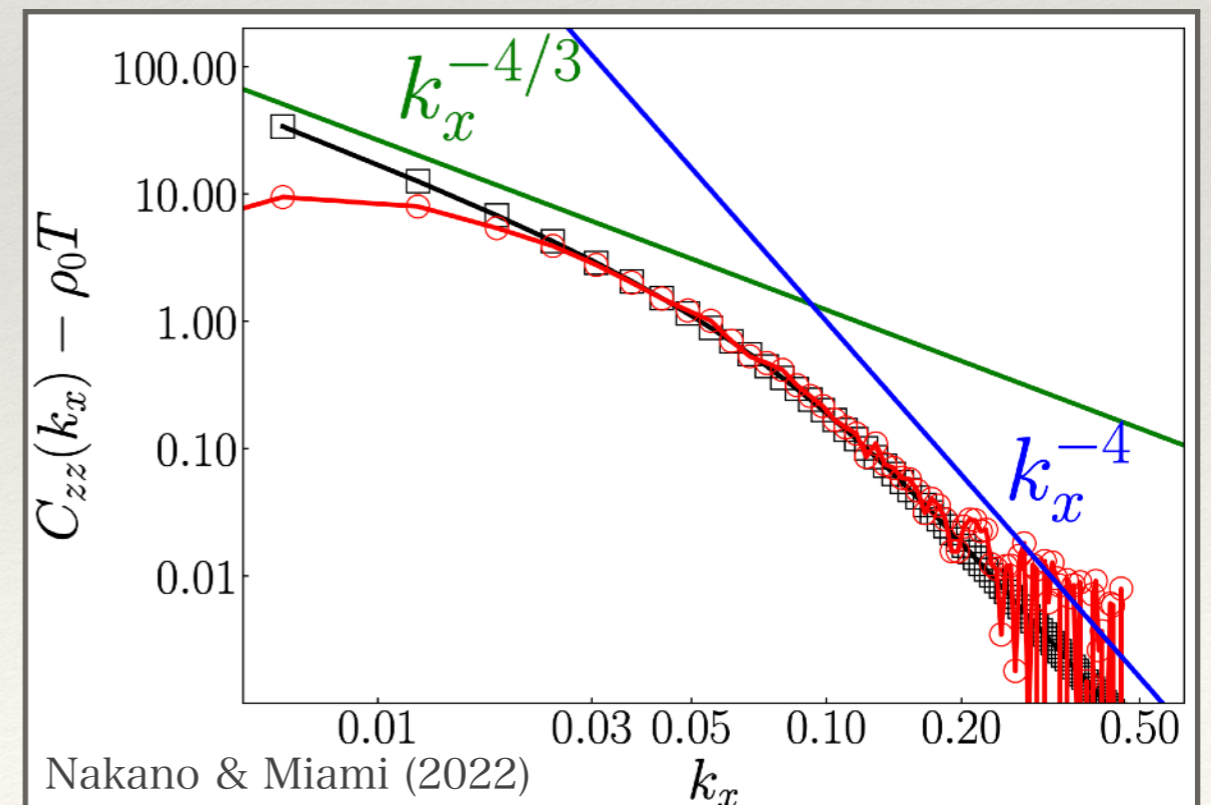
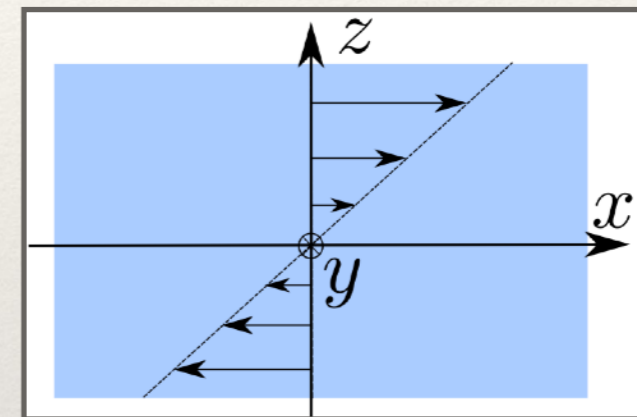
## Large Fluctuations in Non-equilibrium Systems

- Large fluctuations are abundant in systems out of equilibrium

### Temperature gradient



### Steady shear flow



# Introduction

## Hyperuniformity in 2D Chiral Active Fluid

**Q.** What about 2D chiral active fluids?

**A.** They show *hyperuniformity*, not GNF

**Hyperuniformity (HU):**  $S(q) \propto q^\gamma$  with  $\gamma > 0$

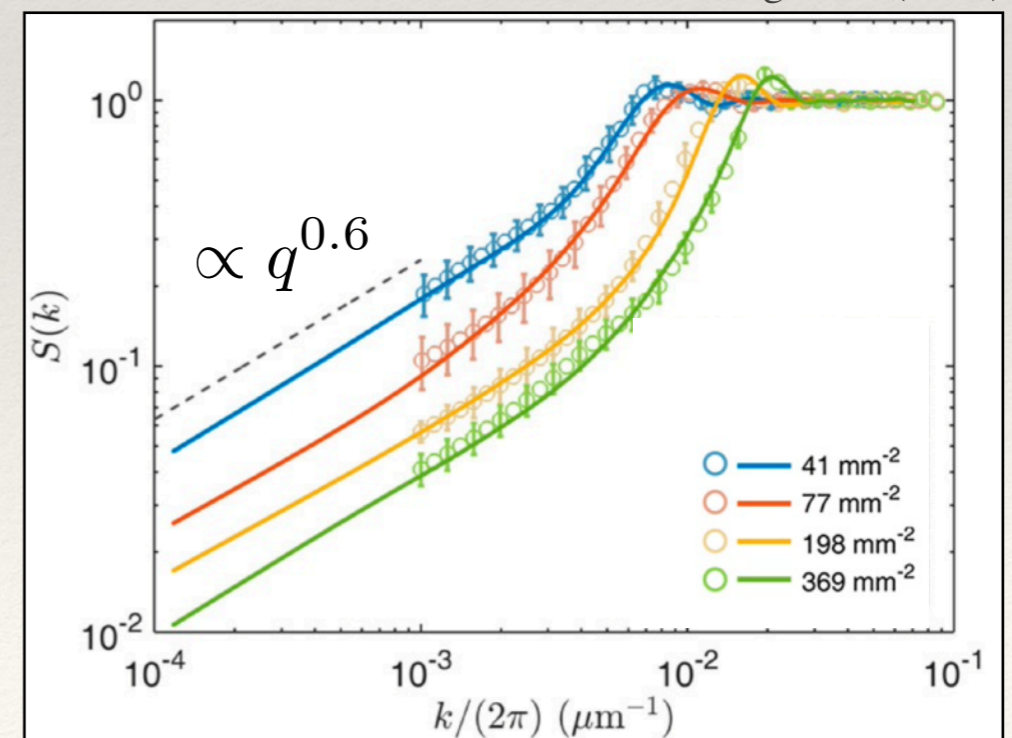
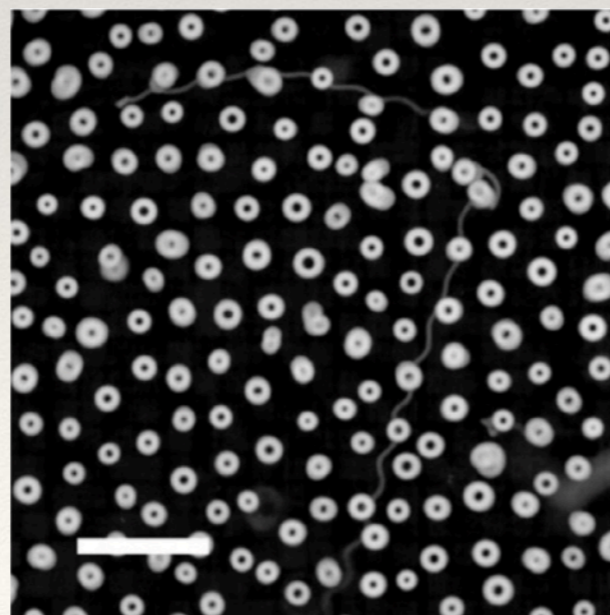
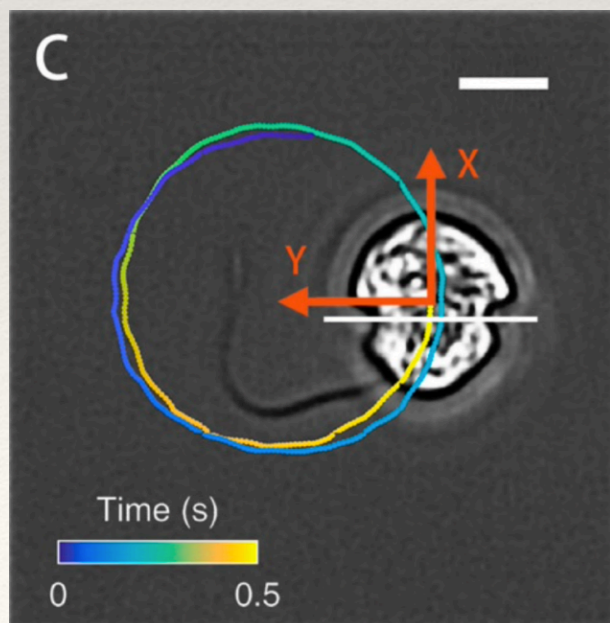
cf.) GNF means

$$S(q) \propto \frac{1}{q^\beta}$$

➔ Anomalous **suppression** of density fluctuations (opposite to GNF!)

### Experiment: Algae system

M. Huang *et al.*, (2021)



# Introduction

## A Numerical Study on Hyperuniformity in 2D Chiral Active Fluid

✓ HU in chiral active fluids has been observed numerically

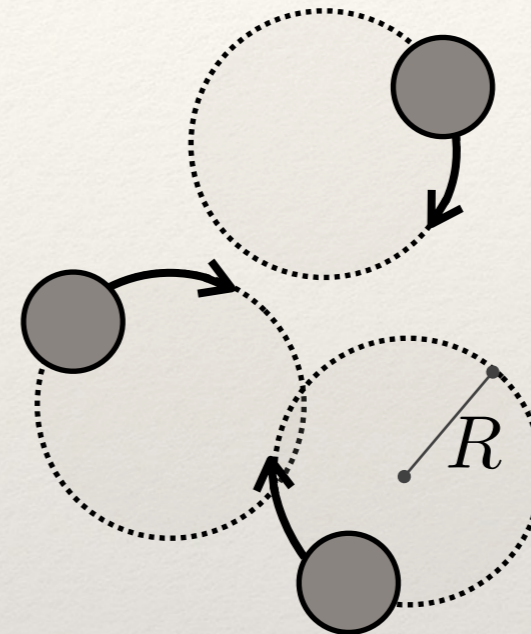
### Chiral active Brownian particles

$$\dot{\mathbf{r}}_j(t) = \mu \mathbf{F}_j^{\text{int}} + v_0 \mathbf{e}(\phi_j(t))$$

interaction self-propelling

$$\dot{\phi}_j(t) = \Omega + \sqrt{2D} \eta_j(t)$$

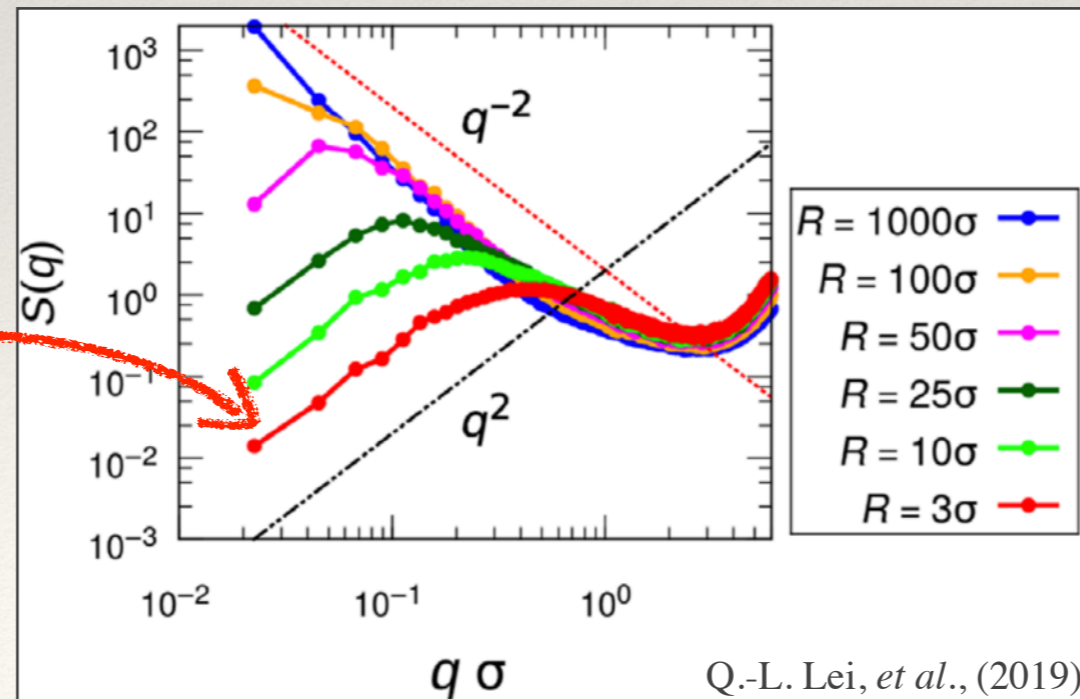
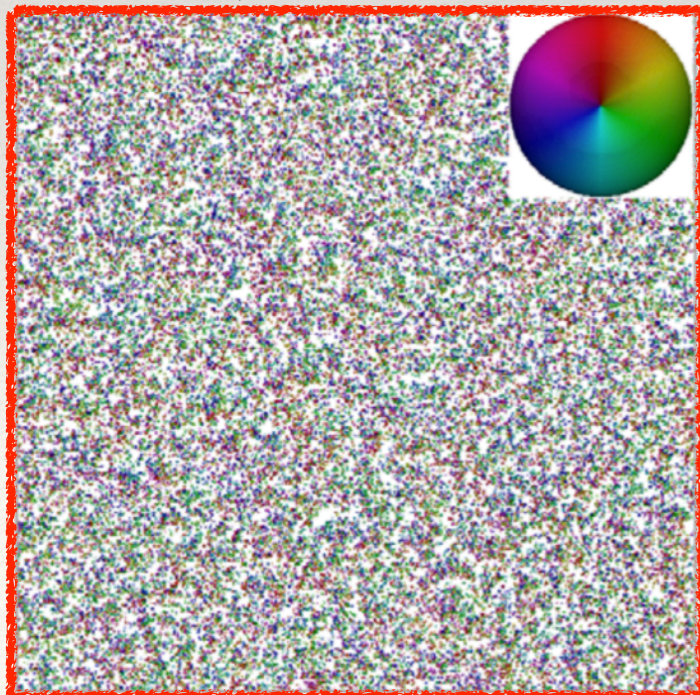
torque noise



Orbital radius

$$R = \frac{v_0}{\Omega}$$

$$D = 0, \varphi = 0.4, R = 3\sigma$$



Hyperuniformity  
characterized by

$$S(q) \propto q^2$$

# Introduction

## Summary of Backgrounds

- Most active matter systems (w/o chirality) show increases in density fluctuations at large scales
- 2D chiral active fluids exhibit the suppression of density fluctuations called hyperuniformity (observed numerically & experimentally)

## Questions

Q. Can we theoretically understand hyperuniformity in 2D chiral active fluids from a microscopic point of view?

→ **Yes!**

Q. What about density fluctuations in 3D chiral active fluid?

→ **Emergence of a singular correlation**

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## ▶ Summary

# Theory for HU in 2D Chiral Active Fluid

## Fluctuating Hydrodynamic Equations

$$\dot{\mathbf{r}}_j(t) = -\mu \sum_{k=1}^N \nabla_j U(\mathbf{r}_j - \mathbf{r}_k) + v_0 \mathbf{e}(\phi_j)$$

$$\dot{\phi}_j(t) = \Omega + \sqrt{2D} \eta_j(t)$$

$$\rho(\mathbf{r}, t) = \sum_{j=1}^N \delta(\mathbf{r} - \mathbf{r}_j(t))$$

**:Density**

$$\mathbf{p}(\mathbf{r}, t) = \sum_{j=1}^N \mathbf{e}(\phi_j(t)) \delta(\mathbf{r} - \mathbf{r}_j(t))$$

**:Polarization**

$$\partial_t \rho(\mathbf{r}, t) = -\nabla \cdot \mathbf{J}(\mathbf{r}, t)$$

$$\partial_t \mathbf{p}(\mathbf{r}, t) = -\nabla \cdot \left( \frac{\mathbf{J}(\mathbf{r}, t) \mathbf{p}(\mathbf{r}, t)}{\rho(\mathbf{r}, t)} \right) - D \mathbf{p}(\mathbf{r}, t) + \boldsymbol{\Omega} \times \mathbf{p}(\mathbf{r}, t) + \sqrt{D \rho(\mathbf{r}, t)} \boldsymbol{\Upsilon}(\mathbf{r}, t)$$

$$\mathbf{J}(\mathbf{r}, t) = -\mu \nabla P(\mathbf{r}, t) + v_0 \mathbf{p}(\mathbf{r}, t)$$

“Pressure”:  $\nabla P(\mathbf{r}, t) := \rho(\mathbf{r}, t) \nabla \frac{\delta \mathcal{F}[\rho(\cdot, t)]}{\delta \rho(\mathbf{r}, t)}$   $\mathcal{F}[\rho(\cdot, t)] = \frac{1}{2} \int_V d^2 \mathbf{r} \int_V d^2 \mathbf{r}' \rho(\mathbf{r}, t) \rho(\mathbf{r}', t) U(\mathbf{r} - \mathbf{r}')$



# Theory for HU in 2D Chiral Active Fluid

## Fluctuating Hydrodynamic Equations

$$\partial_t \rho(\mathbf{r}, t) = -\nabla \cdot \mathbf{J}(\mathbf{r}, t)$$

$$\partial_t \mathbf{p}(\mathbf{r}, t) = -\nabla \cdot \left( \frac{\mathbf{J}(\mathbf{r}, t) \mathbf{p}(\mathbf{r}, t)}{\rho(\mathbf{r}, t)} \right) - D \mathbf{p}(\mathbf{r}, t) + \boldsymbol{\Omega} \times \mathbf{p}(\mathbf{r}, t) + \sqrt{D \rho(\mathbf{r}, t)} \boldsymbol{\Upsilon}(\mathbf{r}, t)$$

$$\mathbf{J}(\mathbf{r}, t) = -\mu \nabla P(\mathbf{r}, t) + v_0 \mathbf{p}(\mathbf{r}, t)$$

**Linearization:**

$$\rho(\mathbf{r}, t) = \rho + \delta \rho(\mathbf{r}, t)$$

$$\mathbf{p}(\mathbf{r}, t) = \mathbf{0} + \delta \mathbf{p}(\mathbf{r}, t)$$

**Assumption:**

$$\nabla P(\mathbf{r}, t) \simeq \frac{1}{\rho \chi} \nabla \delta \rho(\mathbf{r}, t)$$

$$\chi^{-1} := \rho \left. \frac{\partial P}{\partial \rho} \right|_{\rho(\mathbf{r})=\rho}$$

$$\partial_t \delta \rho(\mathbf{r}, t) = b \nabla^2 \delta \rho(\mathbf{r}, t) - v_0 \nabla \cdot \delta \mathbf{p}(\mathbf{r}, t) \quad * b := \mu / (\rho \chi)$$

$$\partial_t \delta \mathbf{p}(\mathbf{r}, t) = -D \delta \mathbf{p}(\mathbf{r}, t) + \boldsymbol{\Omega} \times \delta \mathbf{p}(\mathbf{r}, t) + \sqrt{D \rho} \boldsymbol{\Upsilon}(\mathbf{r}, t)$$

# Theory for HU in 2D Chiral Active Fluid

## Static Structure Factor

$$\partial_t \delta \rho(\mathbf{r}, t) = b \nabla^2 \delta \rho(\mathbf{r}, t) - v_0 \nabla \cdot \delta \mathbf{p}(\mathbf{r}, t)$$

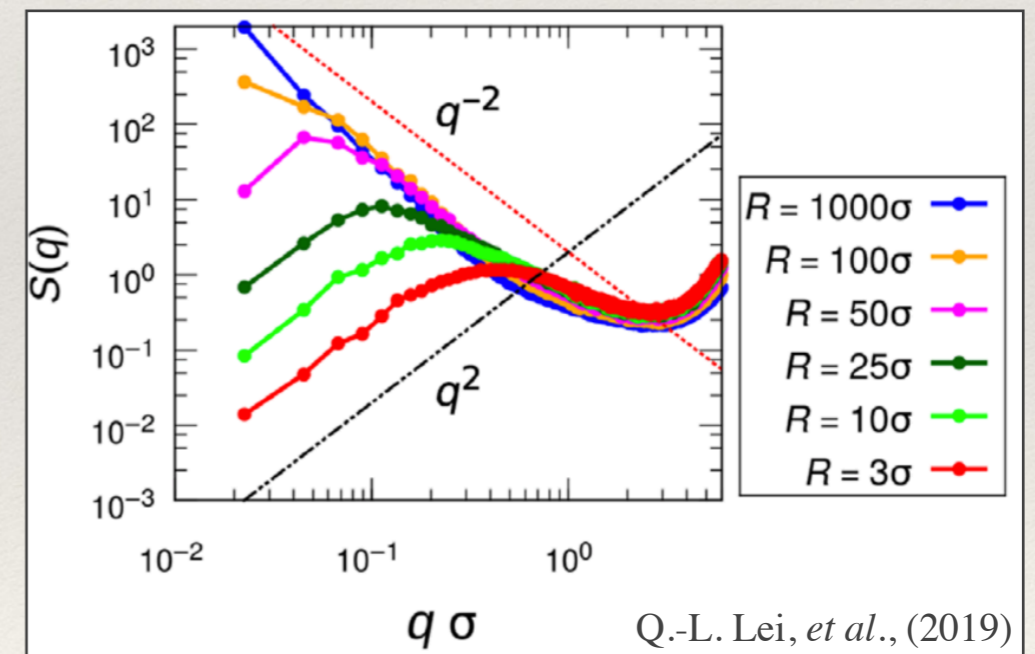
$$\partial_t \delta \mathbf{p}(\mathbf{r}, t) = -D \delta \mathbf{p}(\mathbf{r}, t) + \boldsymbol{\Omega} \times \delta \mathbf{p}(\mathbf{r}, t) + \sqrt{D \rho} \boldsymbol{\Upsilon}(\mathbf{r}, t)$$

● **Static structure factor:**  $S(q) = \frac{1}{N} \langle |\delta \rho(\mathbf{q}, t = 0)|^2 \rangle = \frac{v_0^2}{2b} \cdot \frac{D + bq^2}{\Omega^2 + (D + bq^2)^2}$

☑ In the limit  $D \rightarrow 0$ ,

$$S(q) = \frac{v_0^2}{2\Omega^2} \cdot \frac{q^2}{1 + b^2 q^4 / \Omega^2}$$
$$= \frac{1}{2} (Rq)^2 + O(q^6) \quad R = v_0 / \Omega$$

**Hyperuniformity!**



☑ We succeeded in deriving HU in 2D chiral active fluids!

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