

*Tightening geometric and dynamical
constraints on dark energy and gravity
with galaxy intrinsic alignment*

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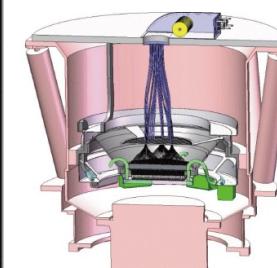
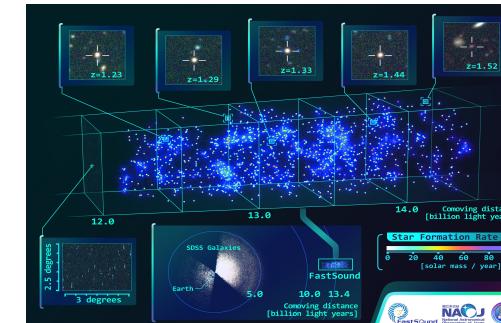
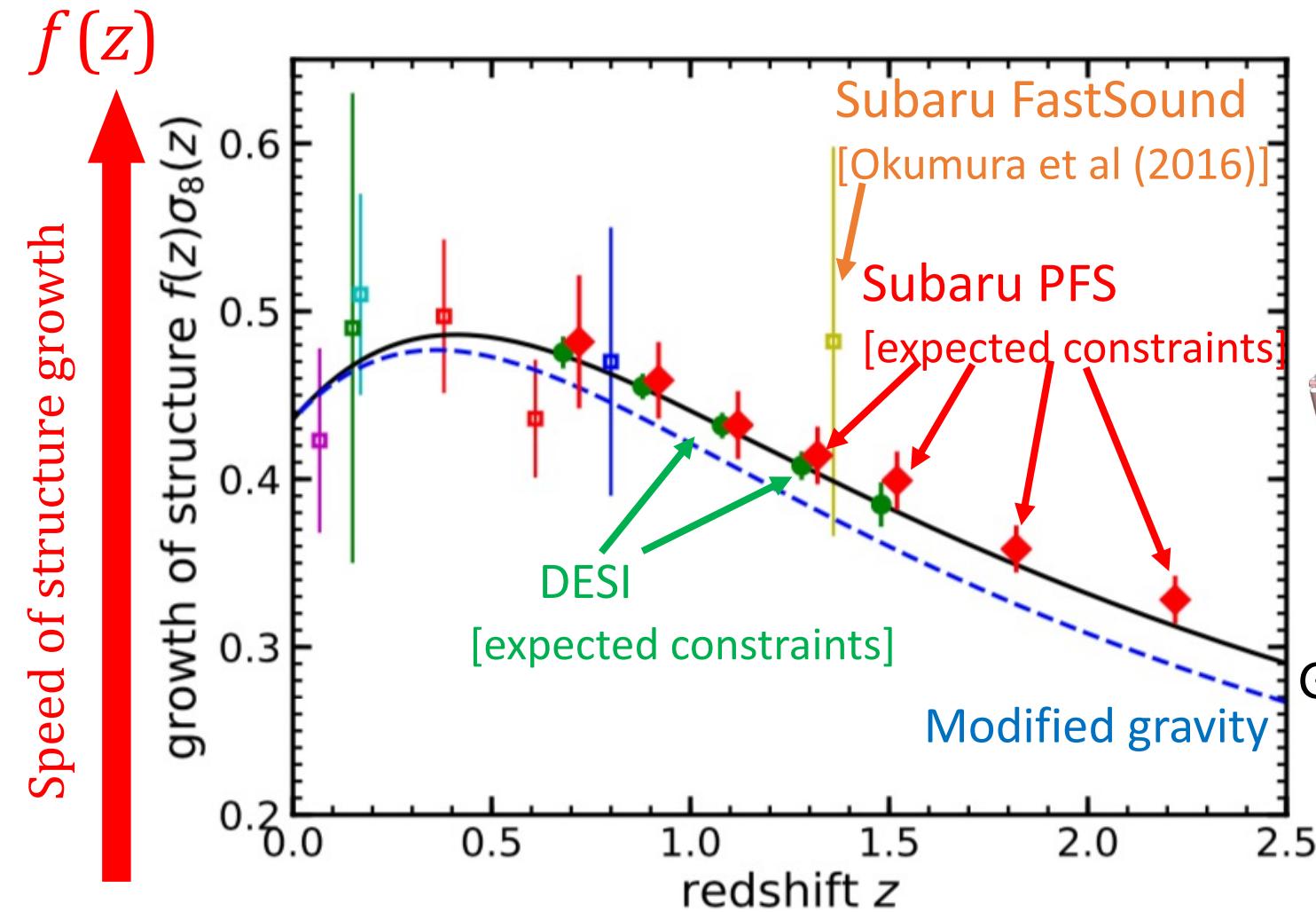
IA cosmology workshop, YITP, Nov. 29 – Dec. 3, 2021

References:

Okumura & Taruya, arXiv: 2110.11127

Tonegawa & **Okumura**, arXiv: 2109.14297

Observational constraints on the growth and expansion history of the universe



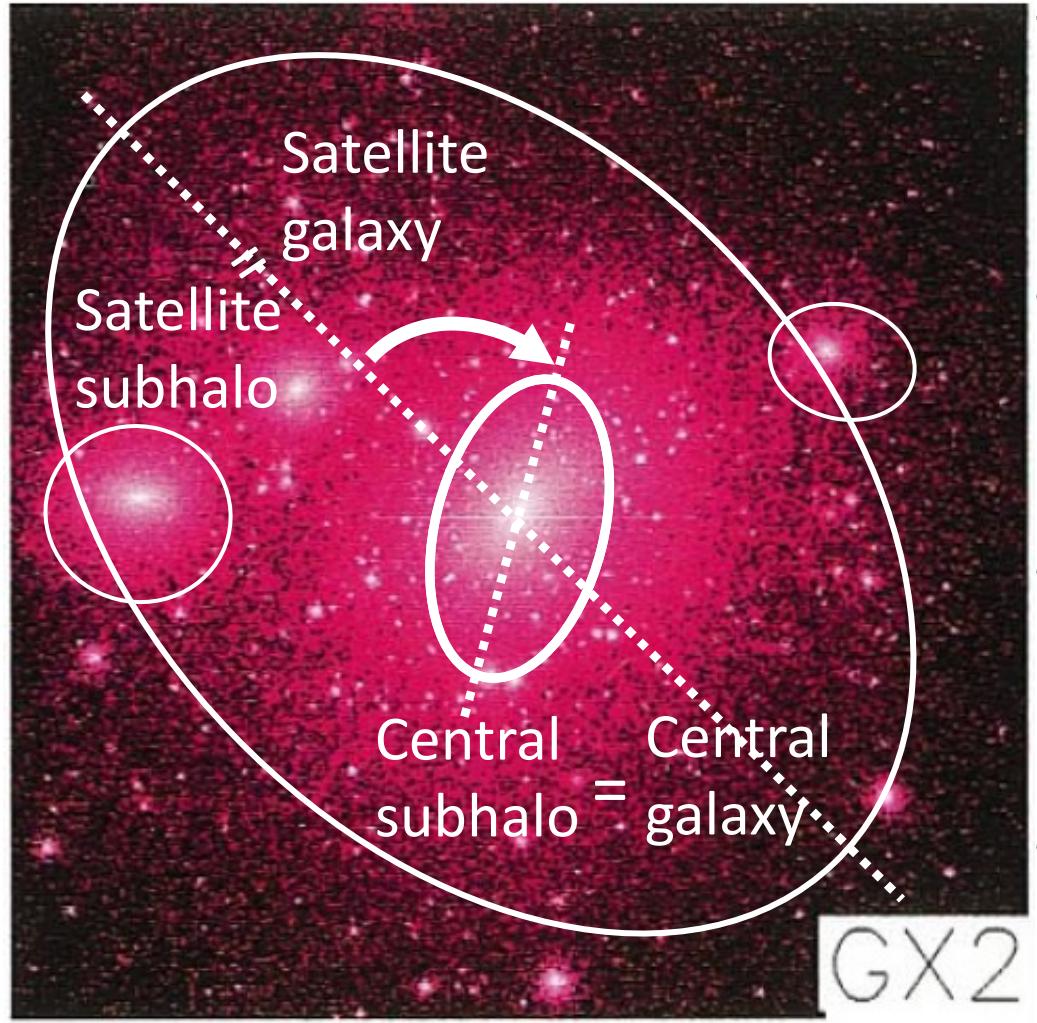
- $f(z)$ constraints
- $H(z)$ marginalized or fixed

Figure made by Ryu Makiya at ASIAA
(PFS Cosmology WG chair)

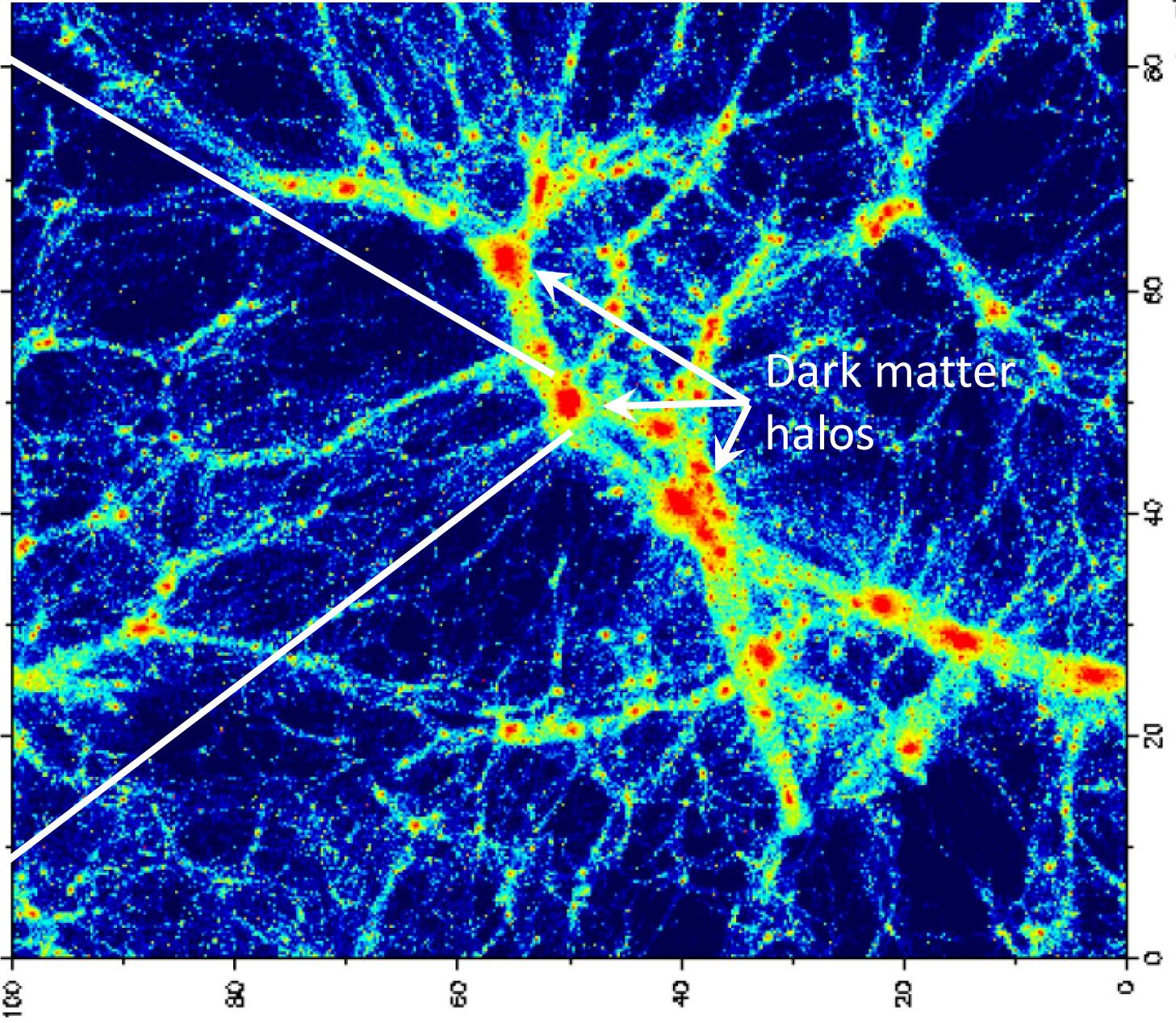
Outline

- Galaxy intrinsic alignment (IA) as a dynamical and geometric probe
- Kinetic Sunyaev-Zel'dovich (kSZ) effect as a dynamical and geometric probe
- Fisher matrix forecast with galaxy clustering + IA + kSZ
 - Geometric and dynamical constraints
 - Cosmological parameter constraints
 - Deep vs wide galaxy surveys
- Measurement of IA of red galaxies at $z > 1$

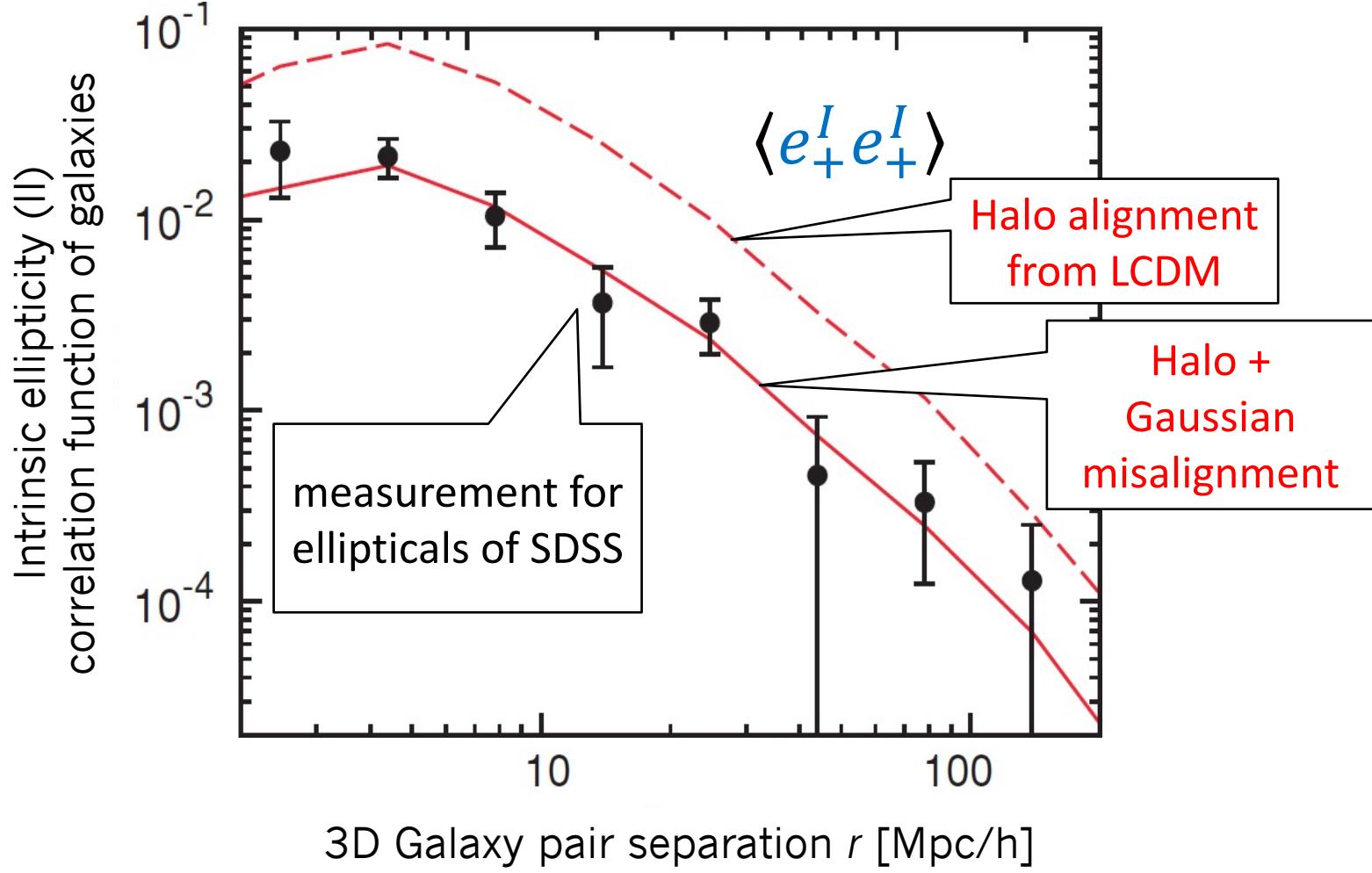
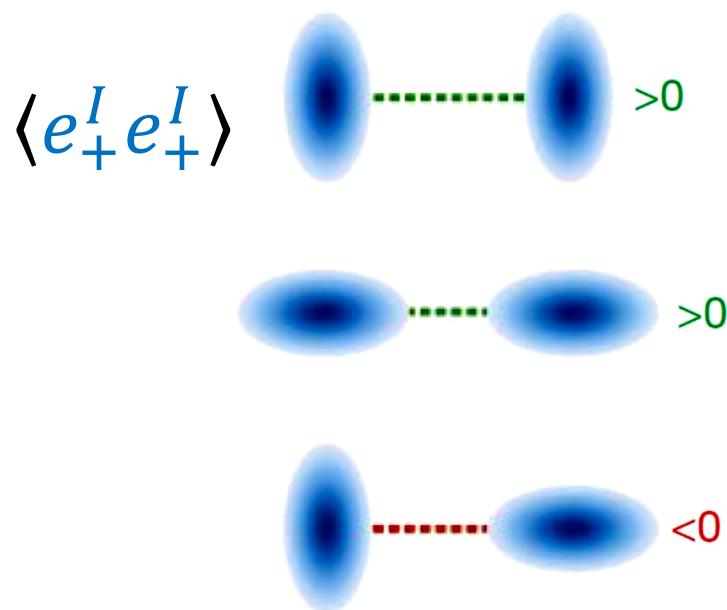
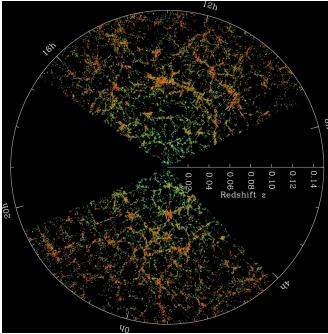
Intrinsic alignment (IA) of galaxy/halo shapes



GX2

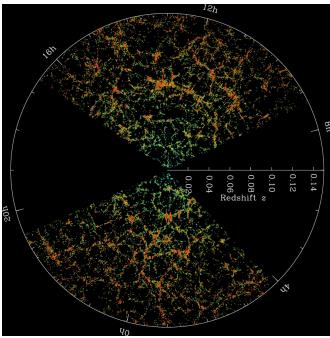


Intrinsic ellipticity auto correlation (II) of elliptical galaxies and the host halos

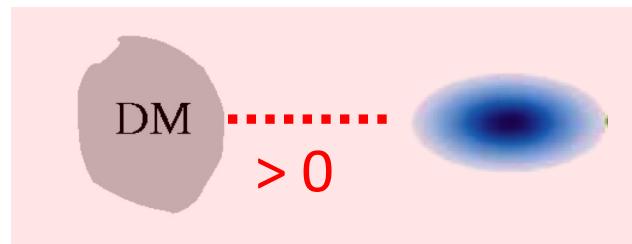


Okumura, Jing & Li (2009)

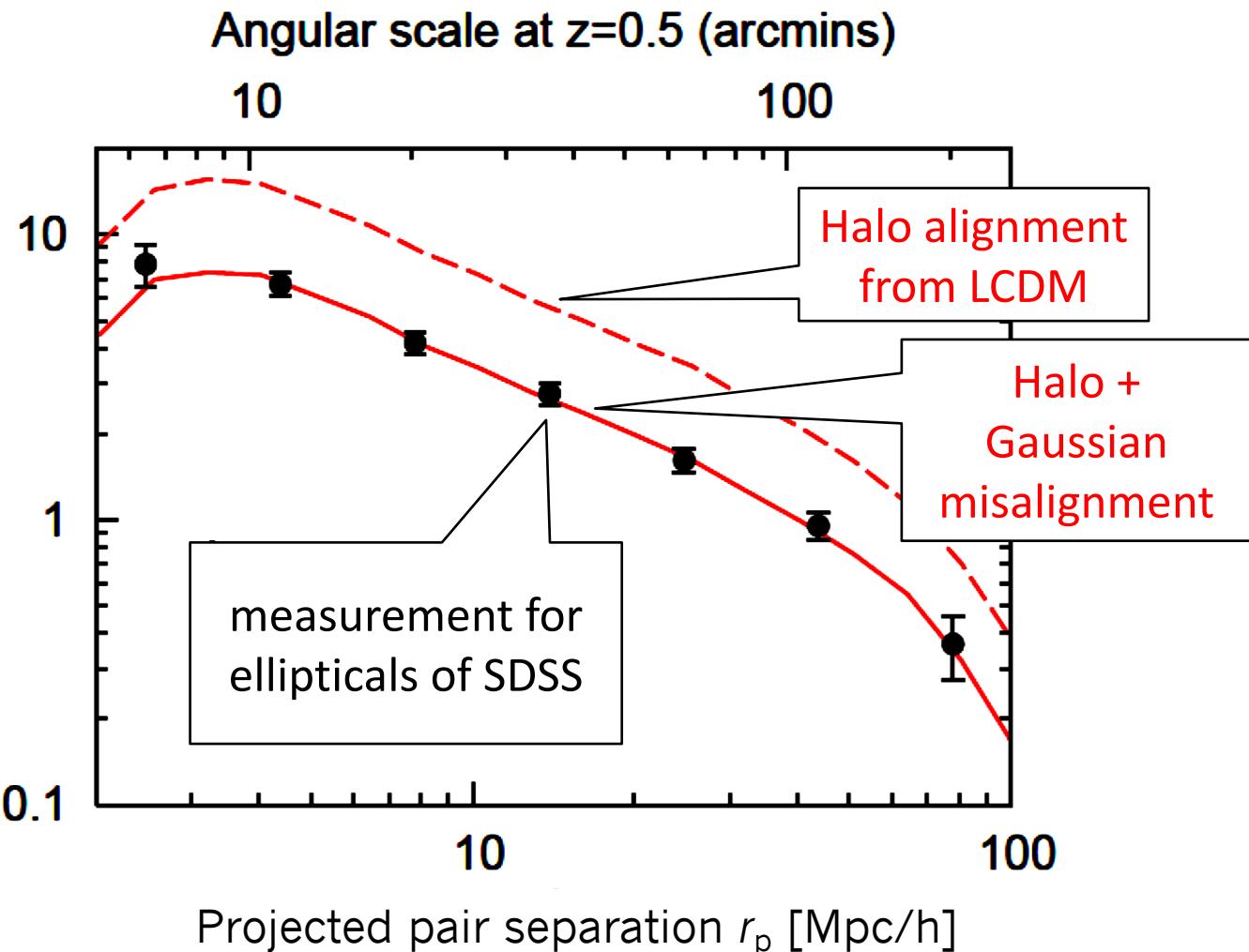
Cross-correlation function between ellipticity and density (GI)



$$\langle \delta e_+^I \rangle$$



Gravitational shear-intrinsic
Ellipticity (GI) correlation function



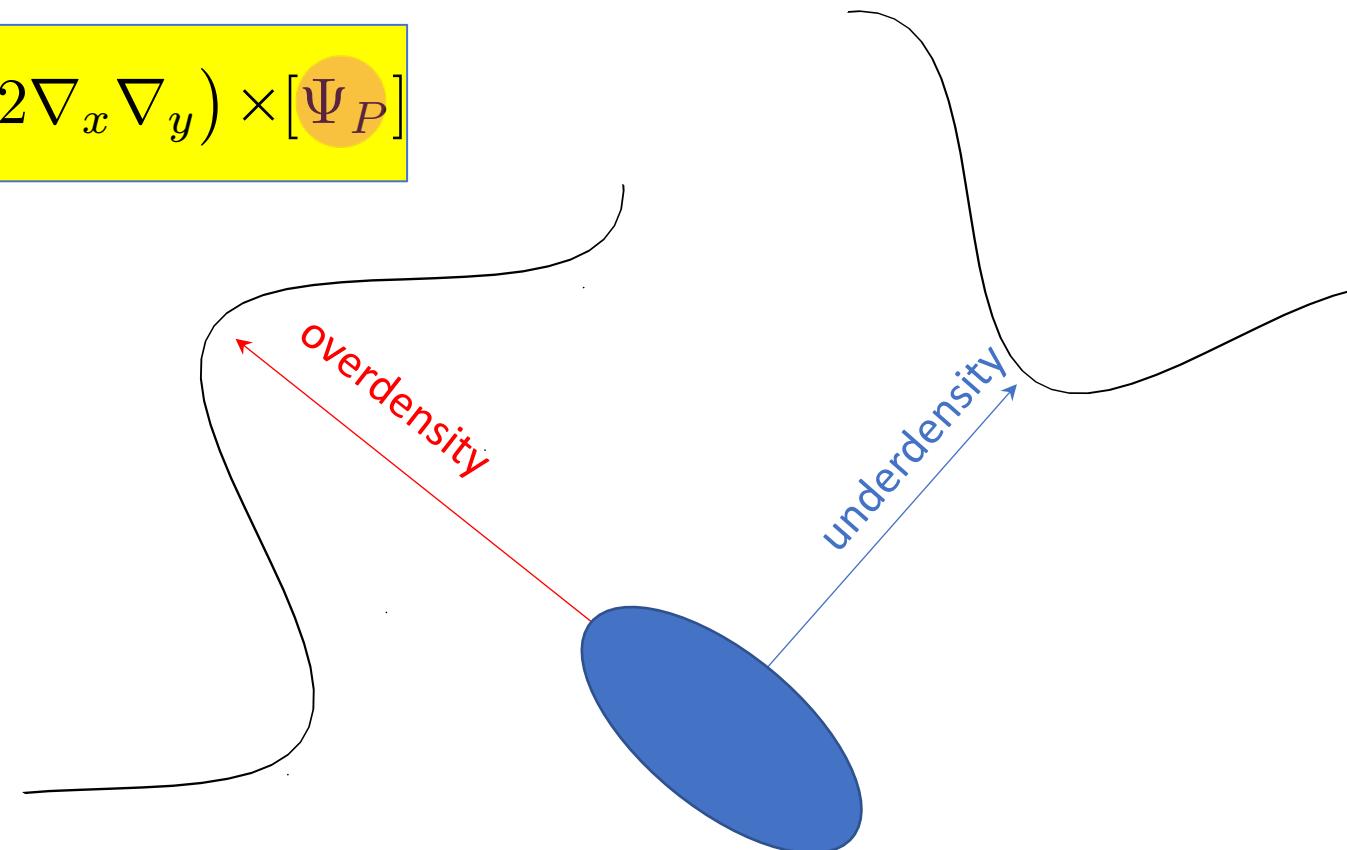
Linear alignment (LA) model

Catelan, Kamionkowski & Blandford (2001)
Hirata & Seljak (2004)

- Relates linear tidal field with galaxy/halo shape

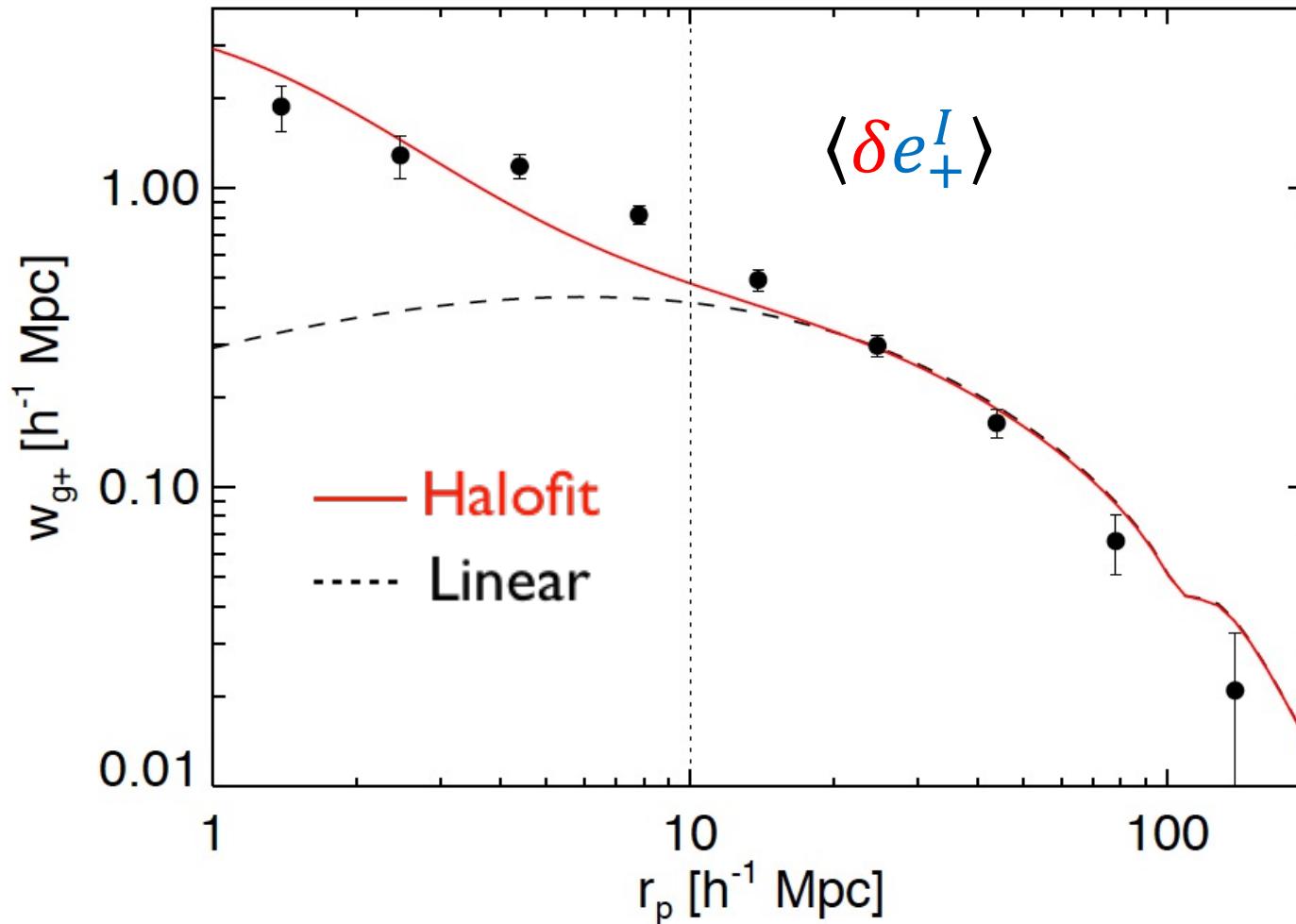
$$\gamma_{(+,\times)}(\mathbf{x}) = -\frac{C_1}{4\pi G} (\nabla_x^2 - \nabla_y^2, 2\nabla_x \nabla_y) \times [\Psi_P]$$

- Ψ_P : (Linear) Newton potential
- C_1 has to be determined by observation/simulation
(this parameter absorbs misalignment and other uncertainties)



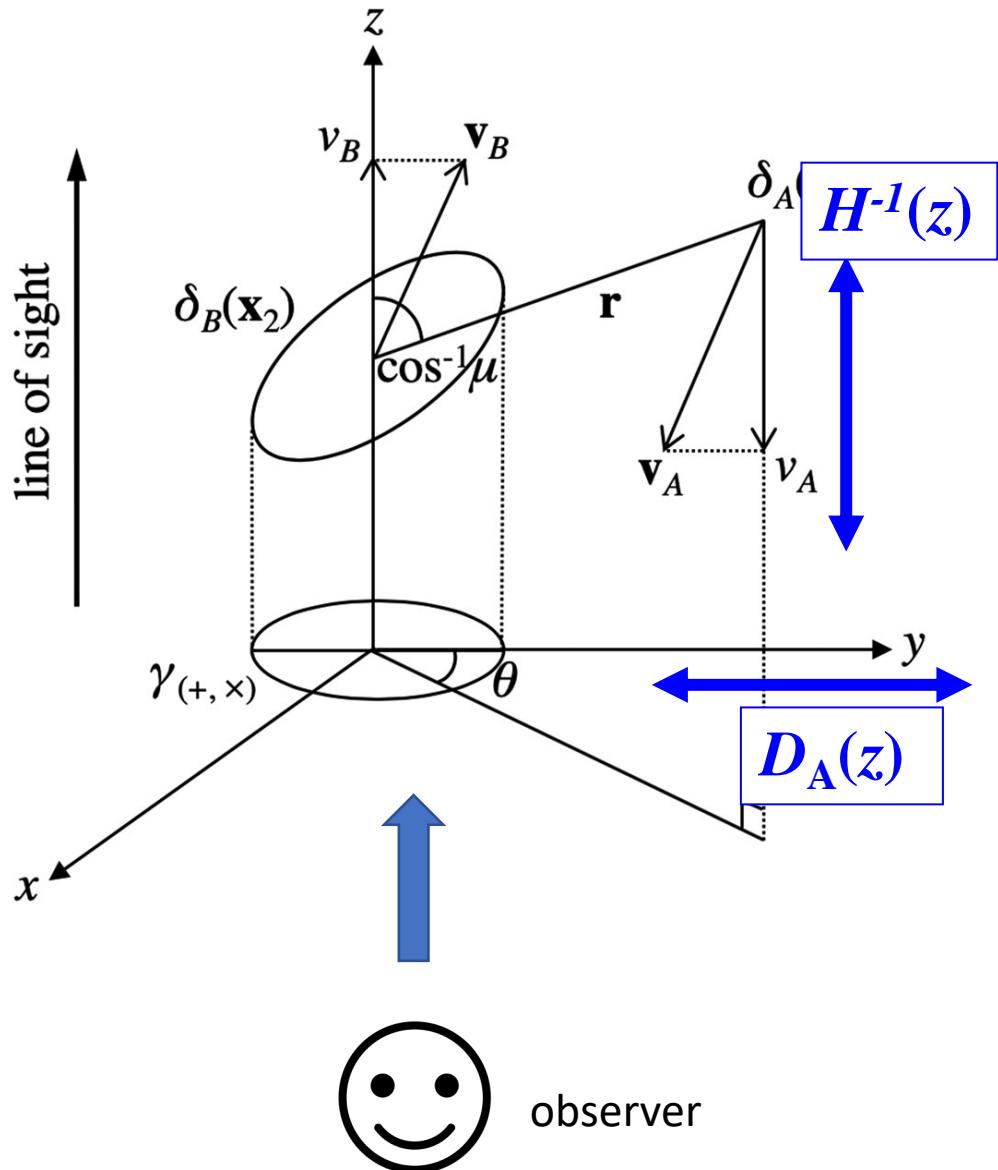
Testing galaxy-ellipticity correlation in LA model with observations

Blazek, McQuinn & Seljak (2011)



- LA model predicts the measurement of IA of the SDSS DR6 Luminous red galaxies by Okumura & Jing (2009)
- But this is the projected correlation function, not full 3D correlation.

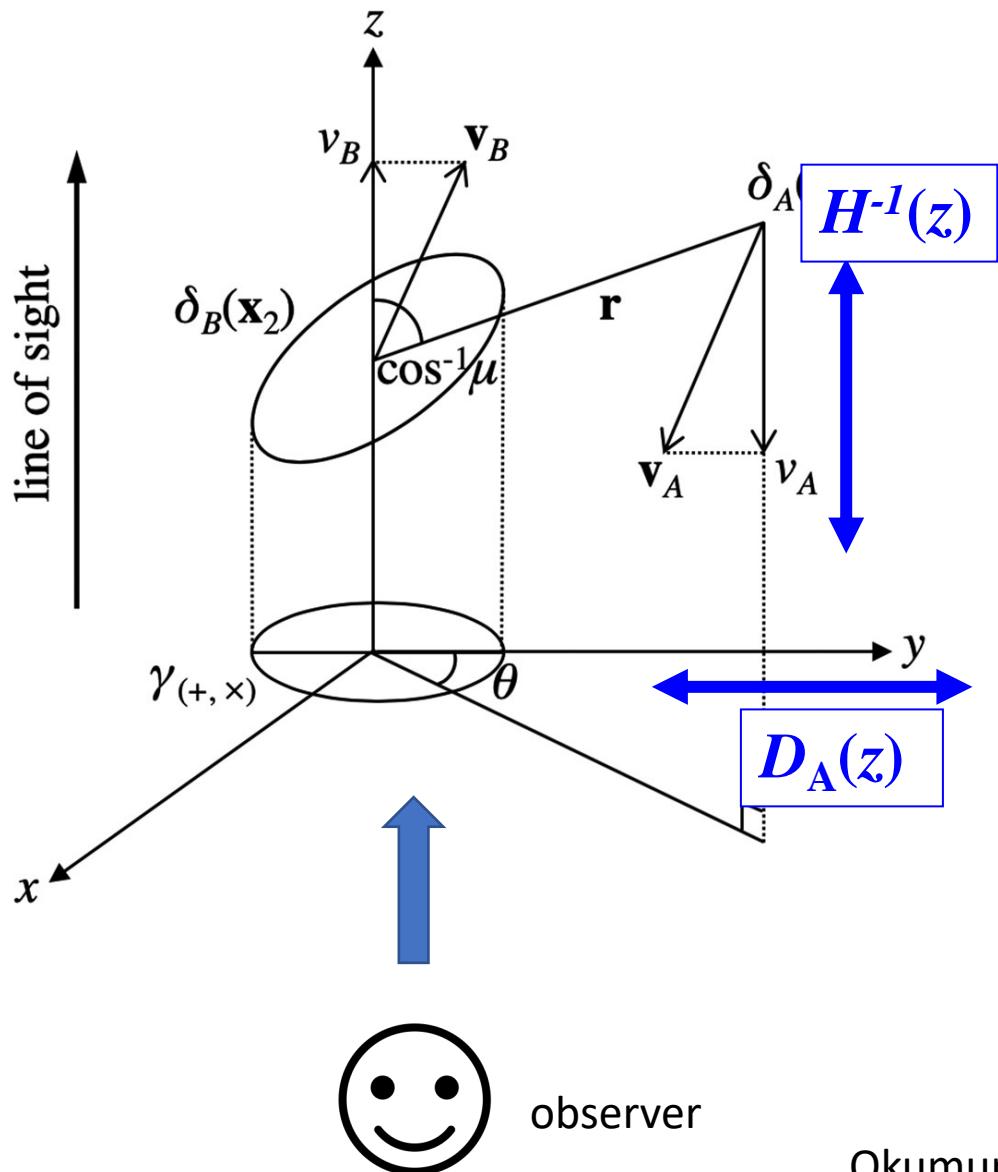
Formulating the IA statistics in redshift space



- Galaxy shapes/orientations are observed on the sky as 2-d projection (θ)
- The line-of-sight component of velocity (1-d) modulates galaxy positions (RSD)
- On the other hand, separation between two galaxies is described as 3-d, $\mathbf{r} = (r, \mu)$
- Original formula for the galaxy-intrinsic ellipticity (GI) correlation function (Hirata & Seljak 2004, Blazek et al 2011)

$$\xi_{g+}(\mathbf{r}) = \tilde{C}_1 b_g \cos(2\phi) \int_0^\infty \frac{k_\perp dk_\perp}{2\pi^2} J_2(k_\perp r_\perp) \times \int_0^\infty dk_\parallel \frac{k_\perp^2}{k_\parallel^2} P_{\delta\delta}(k) \cos(k_\parallel r_\parallel),$$

Formulating the IA statistics in redshift space



- Original formula for real-space GI correlation

$$\xi_{g+}(\mathbf{r}) = \tilde{C}_1 b_g \cos(2\phi) \int_0^\infty \frac{k_\perp dk_\perp}{2\pi^2} J_2(k_\perp r_\perp) \times \int_0^\infty dk_\parallel \frac{k_\perp^2}{k^2} P_{\delta\delta}(k) \cos(k_\parallel r_\parallel),$$

- New, equivalent formula

$$\xi_{g+}^R(\mathbf{r}) = \tilde{C}_1 b_g \cos(2\phi) (1 - \mu^2) \Xi_{\delta\delta,2}^{(0)}(r)$$

$$\xi_{g+,0}^R(r) = -\xi_{g+,2}^R(r) = \frac{2}{3} \tilde{C}_1 b_g \Xi_{\delta\delta,2}^{(0)}(r)$$

With RSD

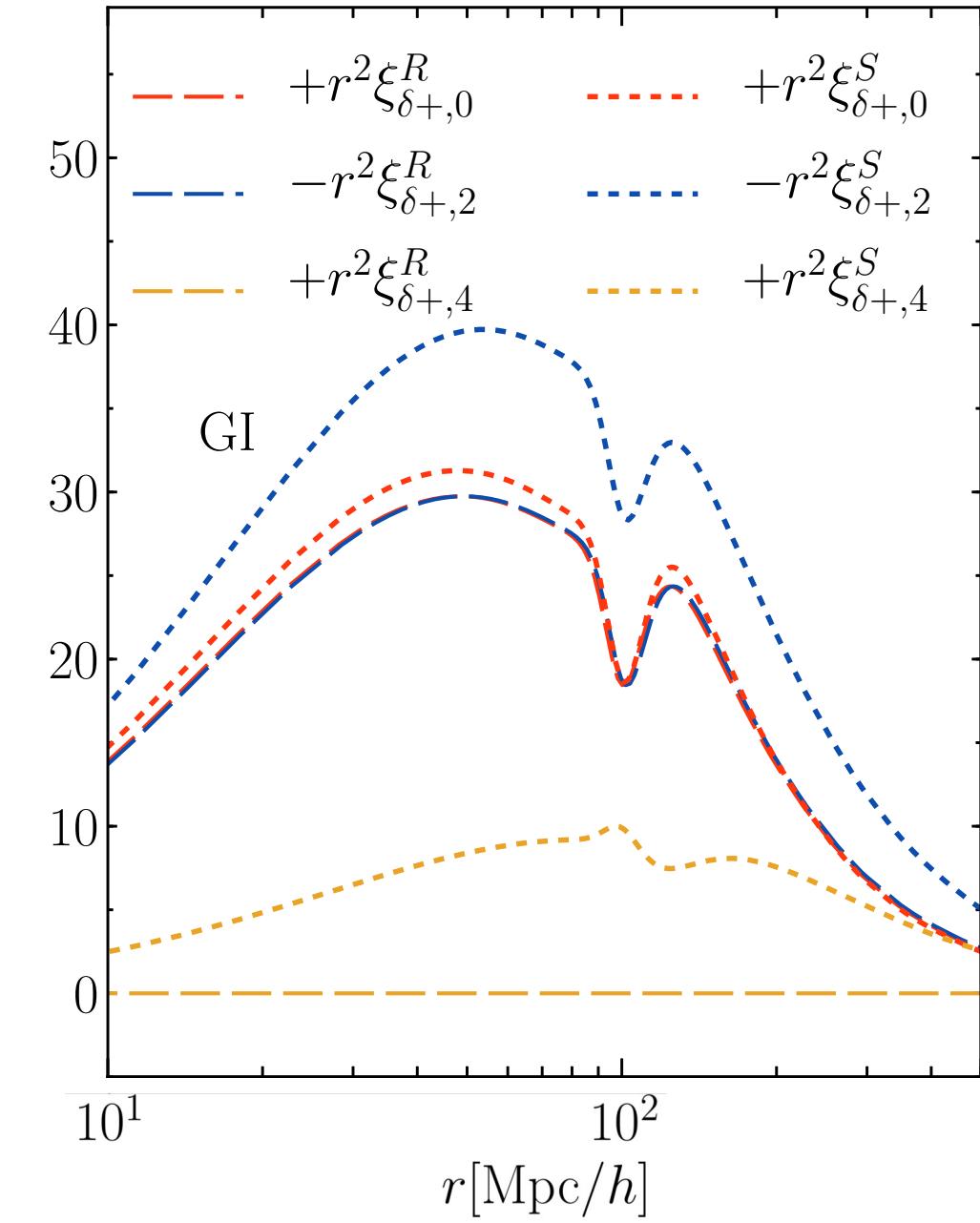
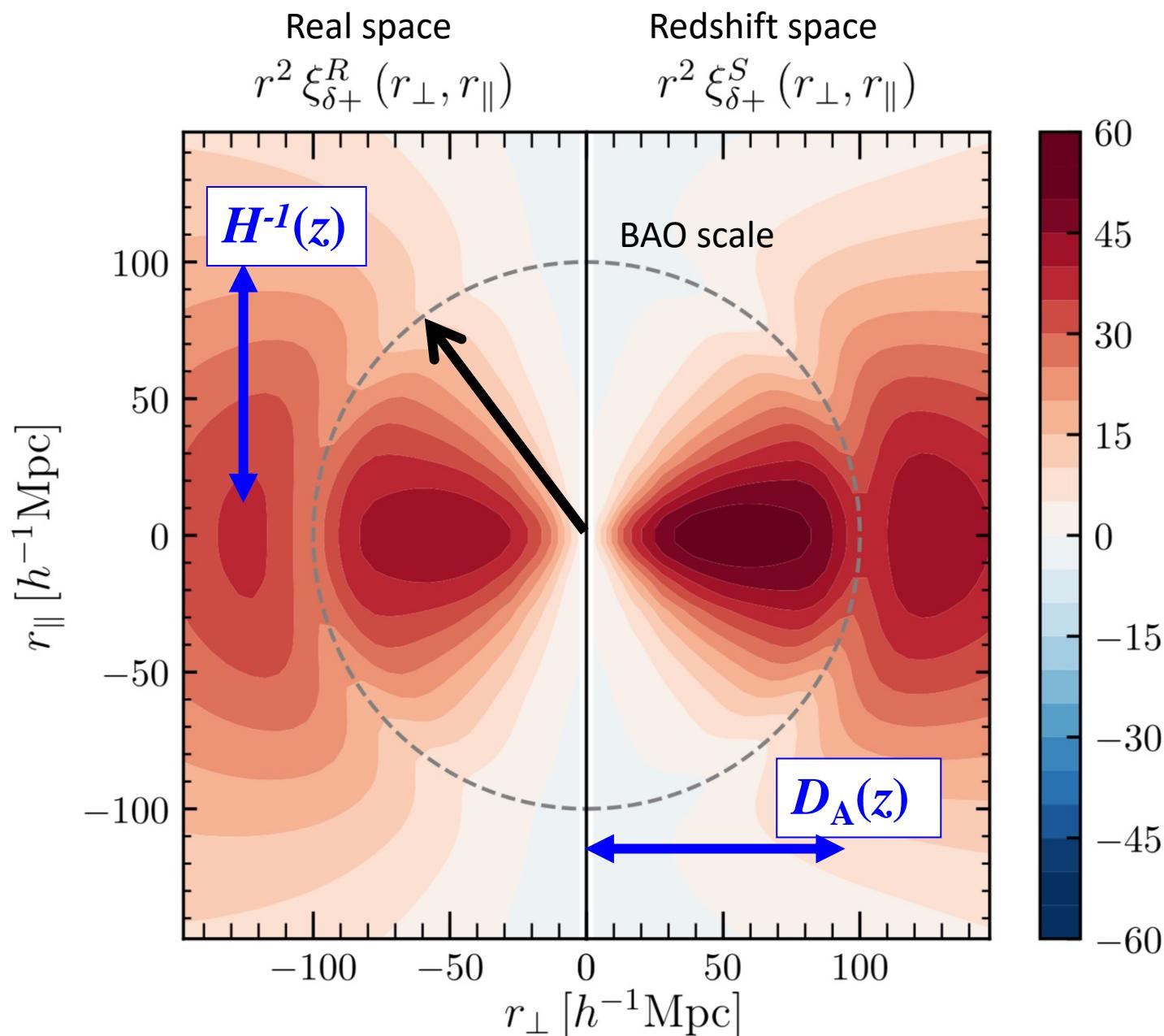
$$\xi_{g+,0}^S(r) = \xi_{g+,0}^R(r) + \frac{2}{105} \tilde{C}_1 f \left[5 \Xi_{\delta\Theta,2}^{(0)}(r) - 2 \Xi_{\delta\Theta,4}^{(0)}(r) \right],$$

$$\xi_{g+,2}^S(r) = \xi_{g+,2}^R(r) - \frac{2}{21} \tilde{C}_1 f \left[\Xi_{\delta\Theta,2}^{(0)}(r) + 2 \Xi_{\delta\Theta,4}^{(0)}(r) \right],$$

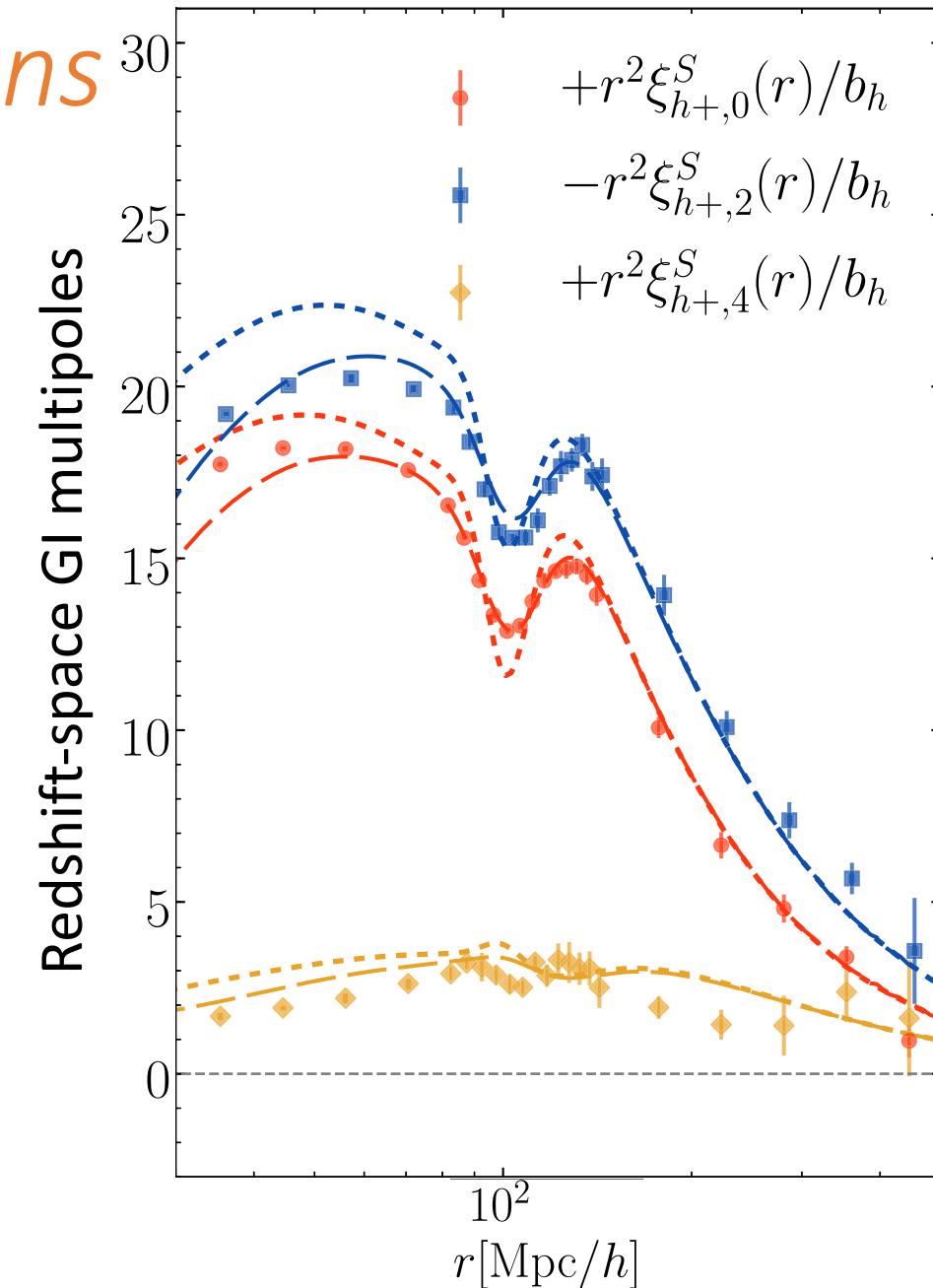
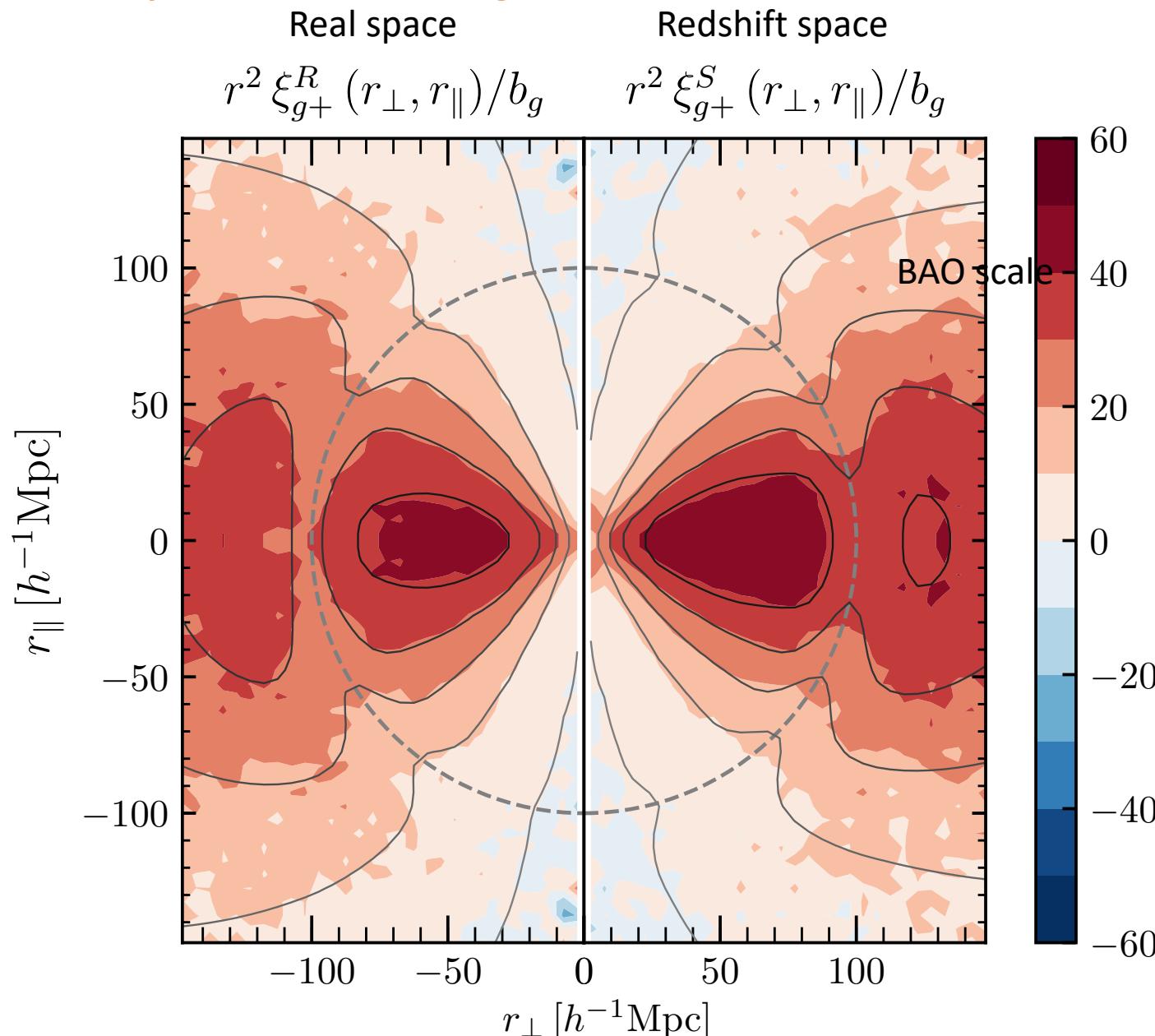
$$\xi_{g+,4}^S(r) = \frac{8}{35} \tilde{C}_1 f \Xi_{\delta\Theta,4}^{(0)}(r).$$

$$\Xi_{XY,\ell}^{(n)}(r) = (aHf)^n \int_0^\infty \frac{k^{2-n} dk}{2\pi^2} P_{XY}(k) j_\ell(kr)$$

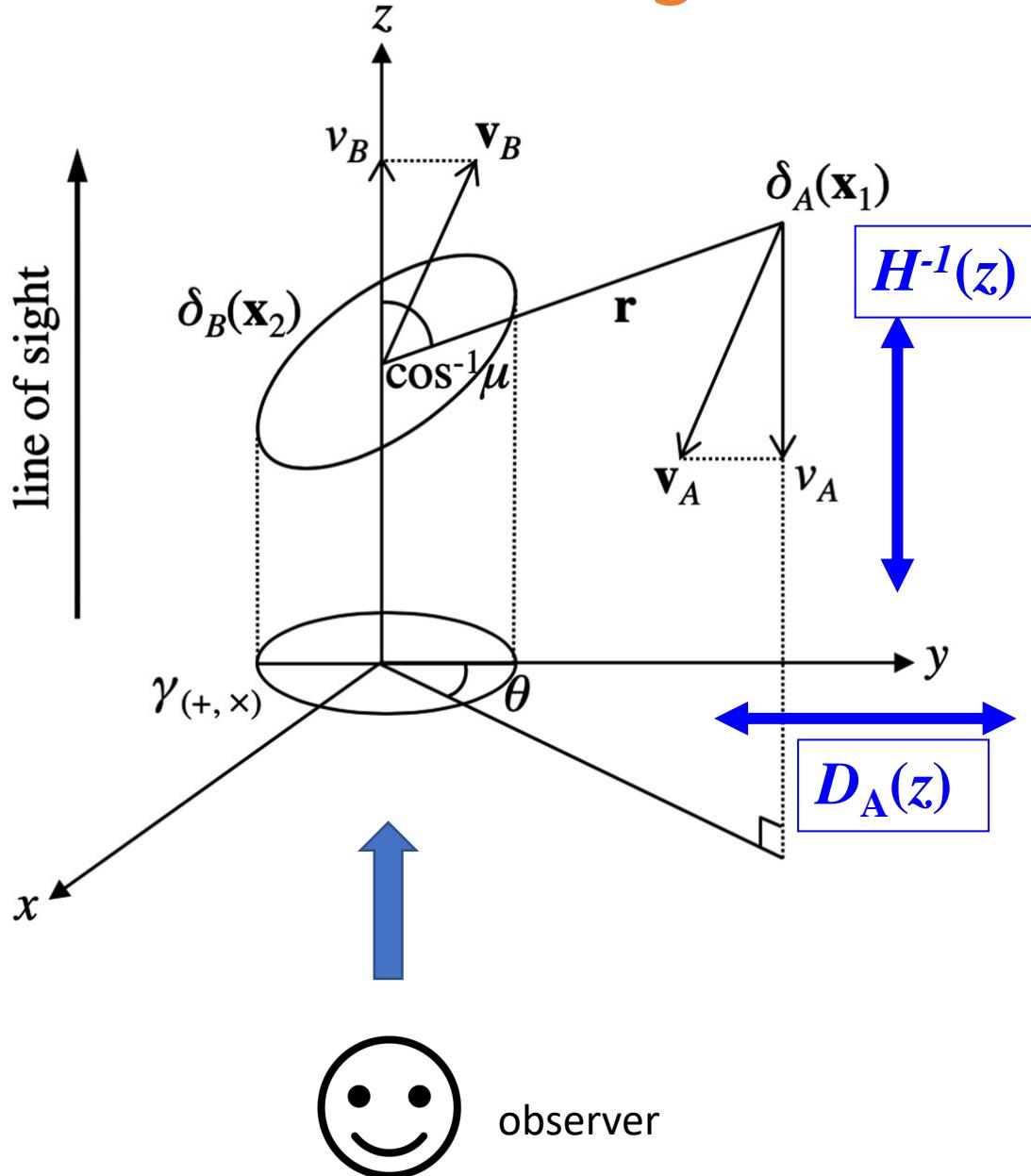
GI correlation in linear theory



Comparison of 2-D GI to simulations



Formulating the IA statistics in redshift space



- New formula for II correlation
 - See Xia+ (2017) for a similar expression for the isotropic moment

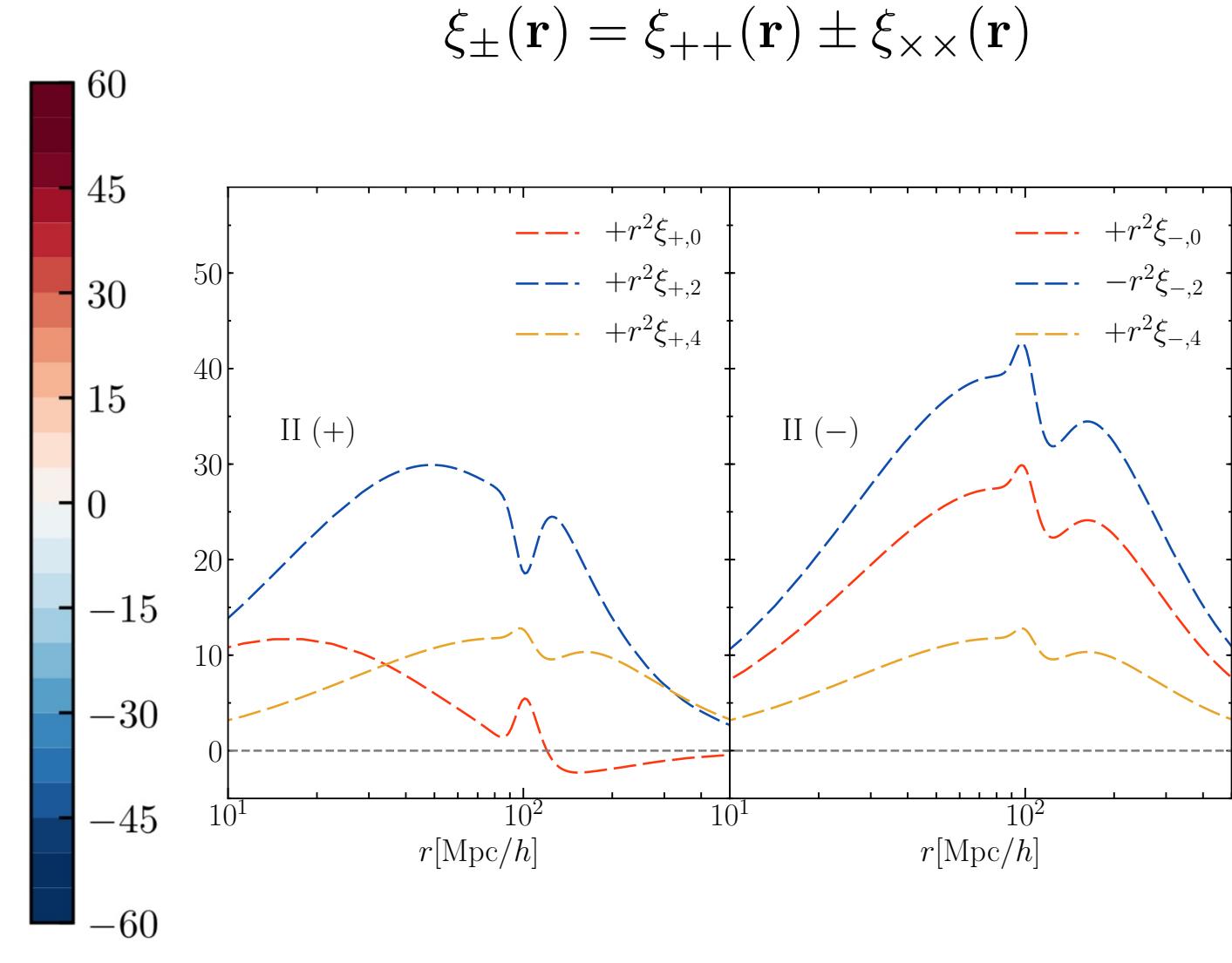
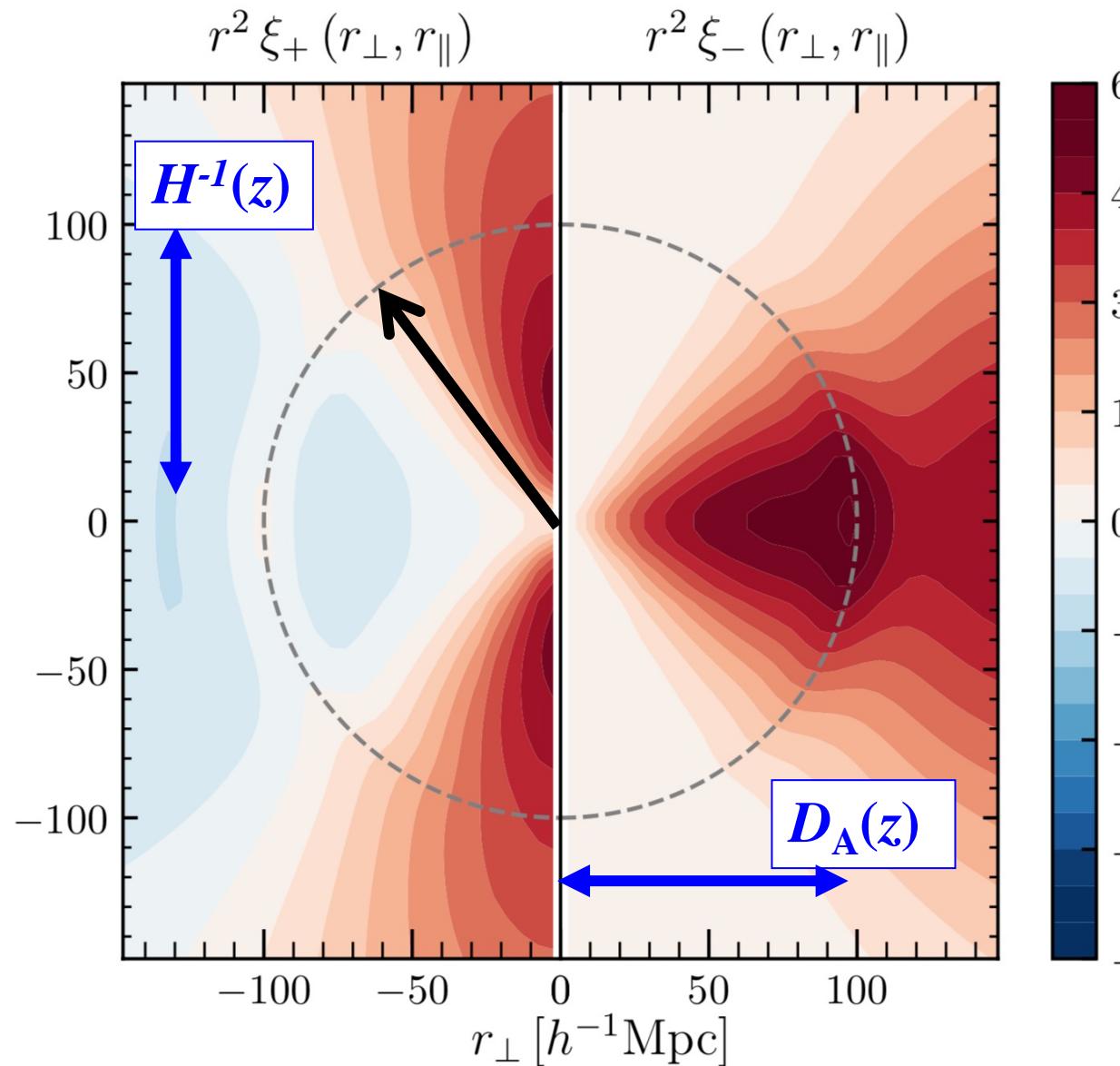
$$\xi_{\pm}(\mathbf{r}) = \xi_{++}(\mathbf{r}) \pm \xi_{\times\times}(\mathbf{r})$$

$$\begin{aligned} \xi_{+}(\mathbf{r}) = & \frac{8}{105} \tilde{C}_1^2 \left[7 \mathcal{P}_0(\mu) \Xi_{\delta\delta,0}^{(0)}(r) + 10 \mathcal{P}_2(\mu) \Xi_{\delta\delta,2}^{(0)}(r) \right. \\ & \left. + 3 \mathcal{P}_4(\mu) \Xi_{\delta\delta,4}^{(0)}(r) \right], \end{aligned}$$

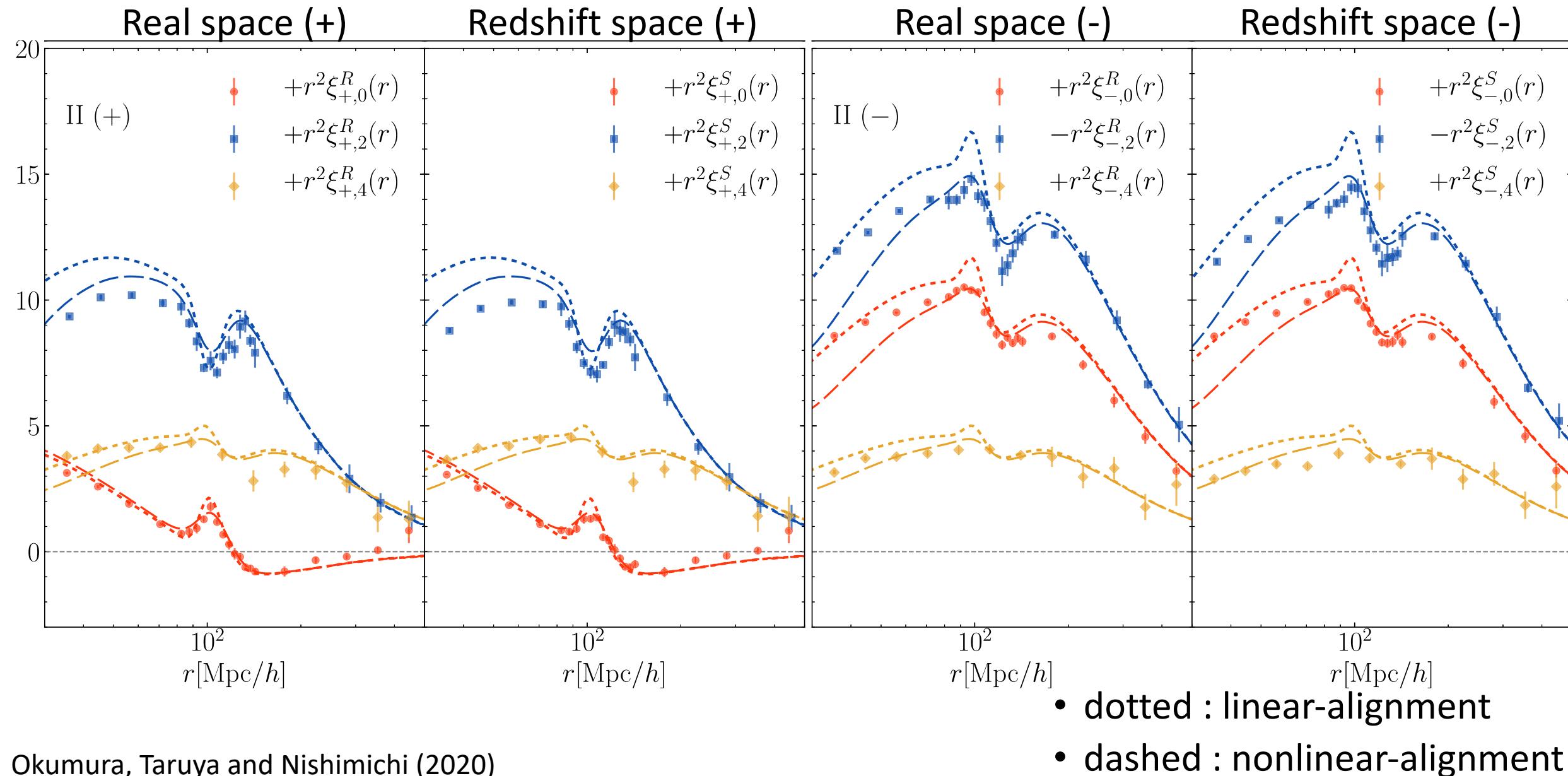
$$\begin{aligned} \xi_{-}(\mathbf{r}) = & \tilde{C}_1^2 \cos(4\phi) (1 - \mu^2)^2 \Xi_{\delta\delta,4}^{(0)}(r) \\ = & \frac{8}{105} \tilde{C}_1^2 \cos(4\phi) \\ & \times [7 \mathcal{P}_0(\mu) + 10 \mathcal{P}_2(\mu) + 3 \mathcal{P}_4(\mu)] \Xi_{\delta\delta,4}^{(0)}(r). \end{aligned}$$

$$\Xi_{XY,\ell}^{(n)}(r) = (aHf)^n \int_0^\infty \frac{k^{2-n} dk}{2\pi^2} P_{XY}(k) j_\ell(kr)$$

$\text{II correlation in } 2D$

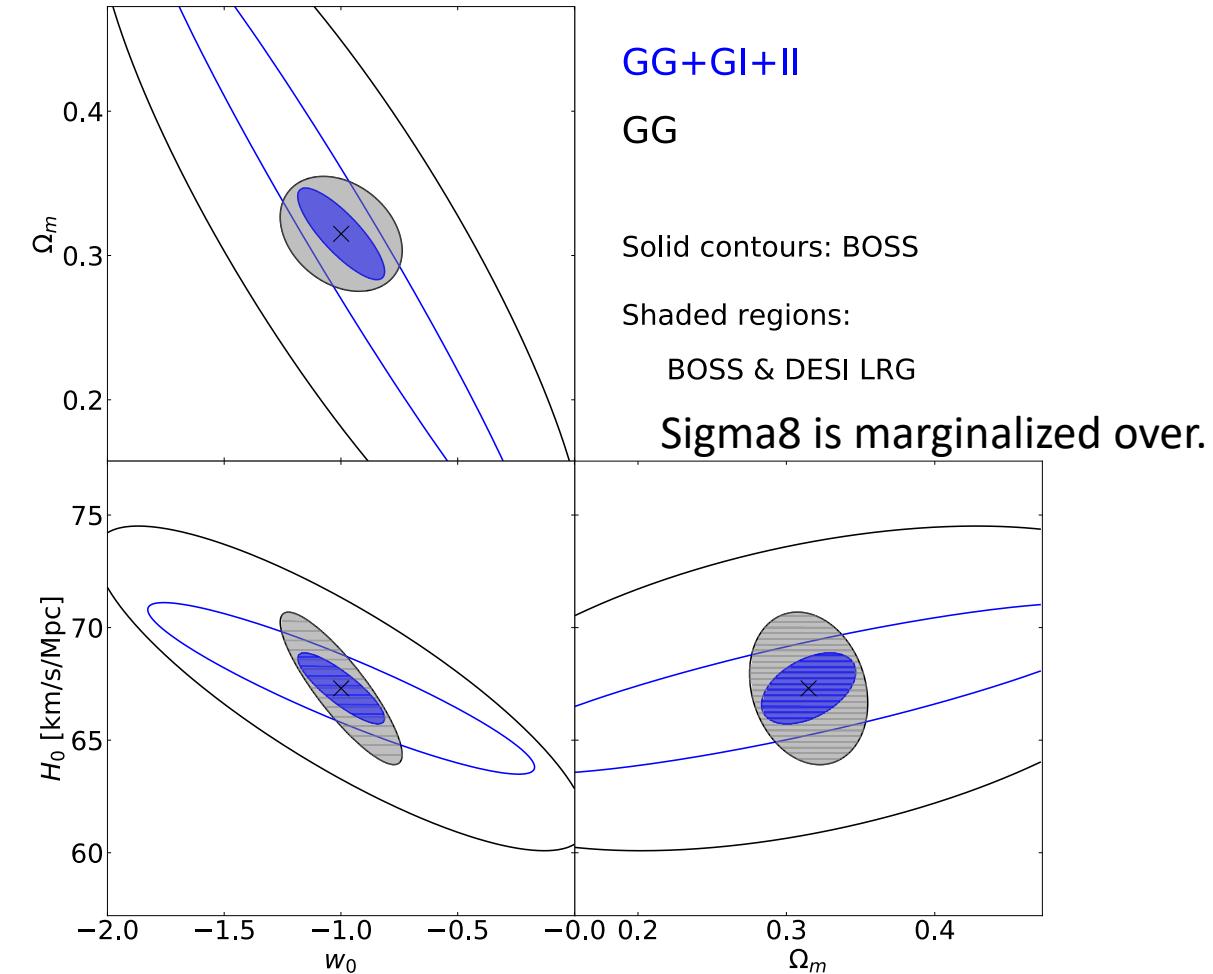
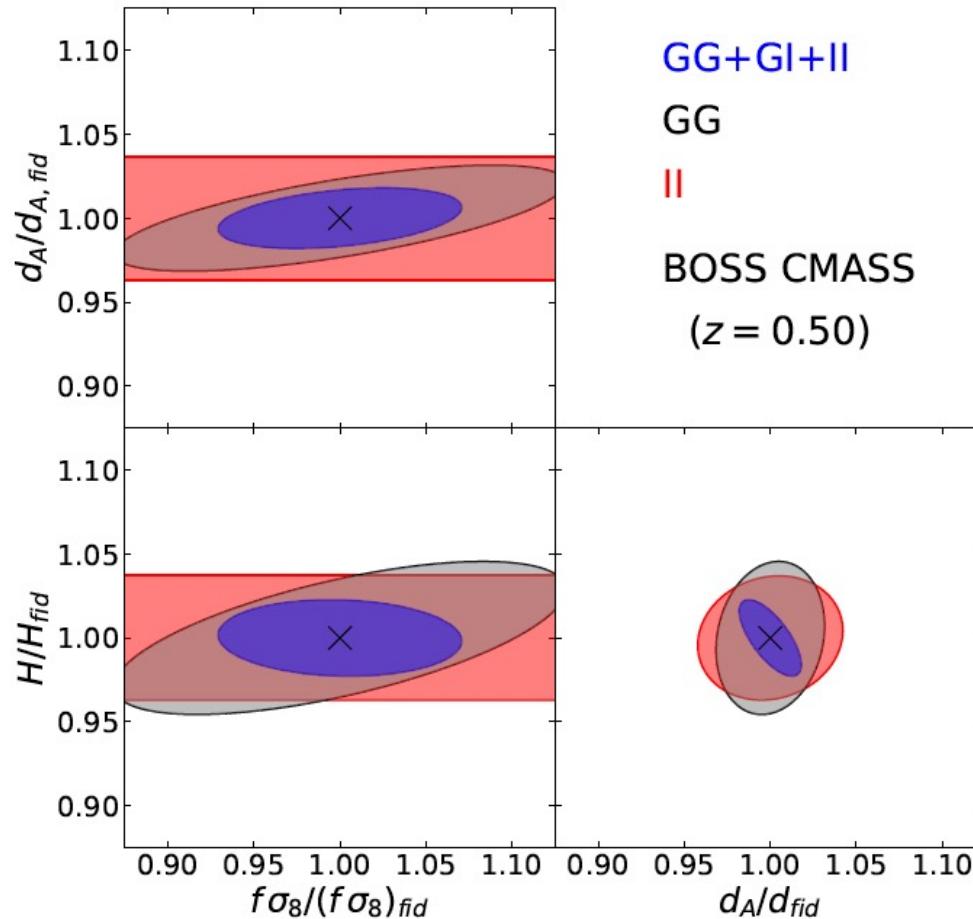


Comparison of $l=2$ multipoles to simulations



IA measurements enhance the science return from galaxy redshift surveys

- Under the assumption that the linear alignment model describes the IA perfectly,



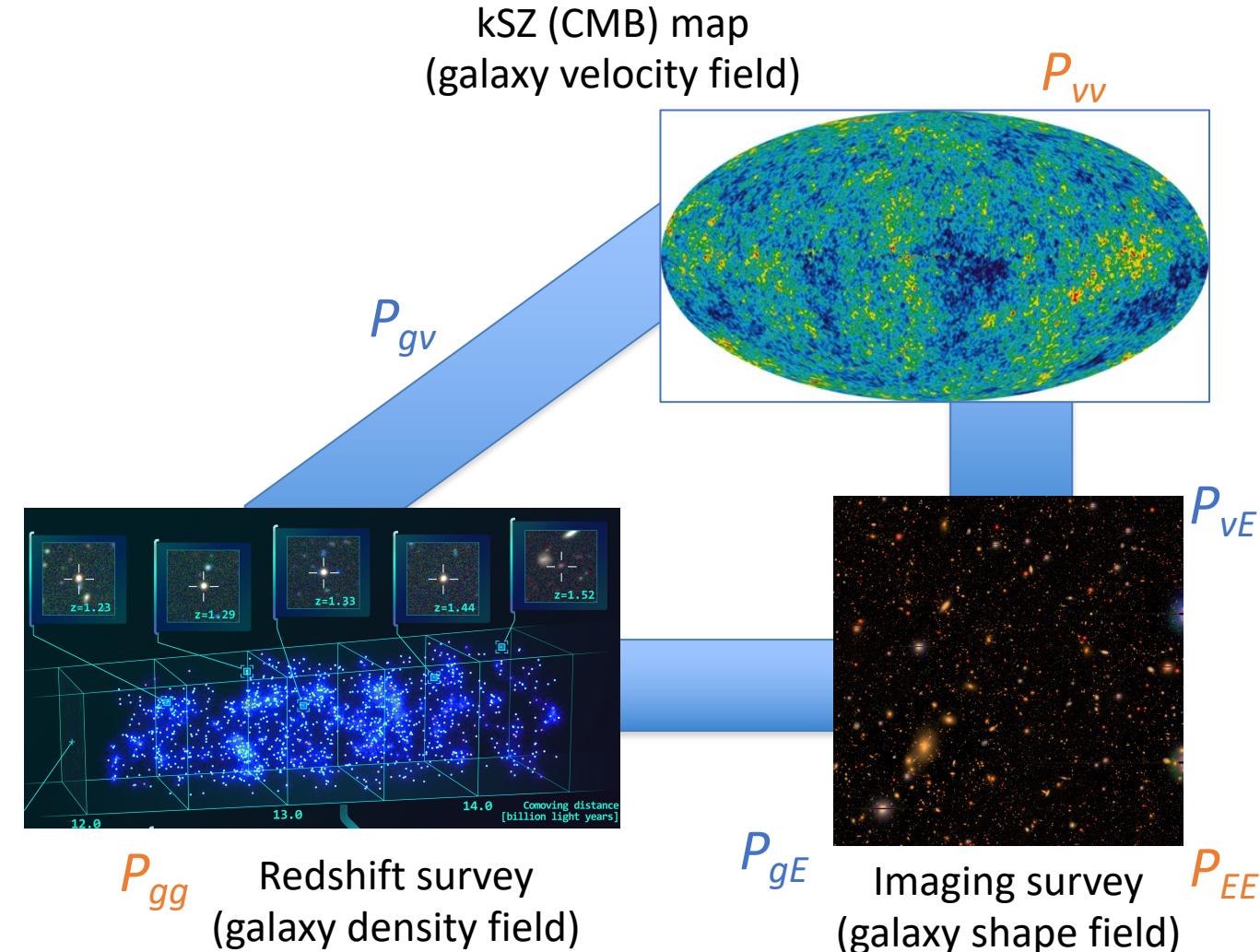
Clustering σ_8 and IA amplitude A_{IA} are marginalized over.

Taruya & Okumura (2020)

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Synergies between galaxy redshift surveys and CMB



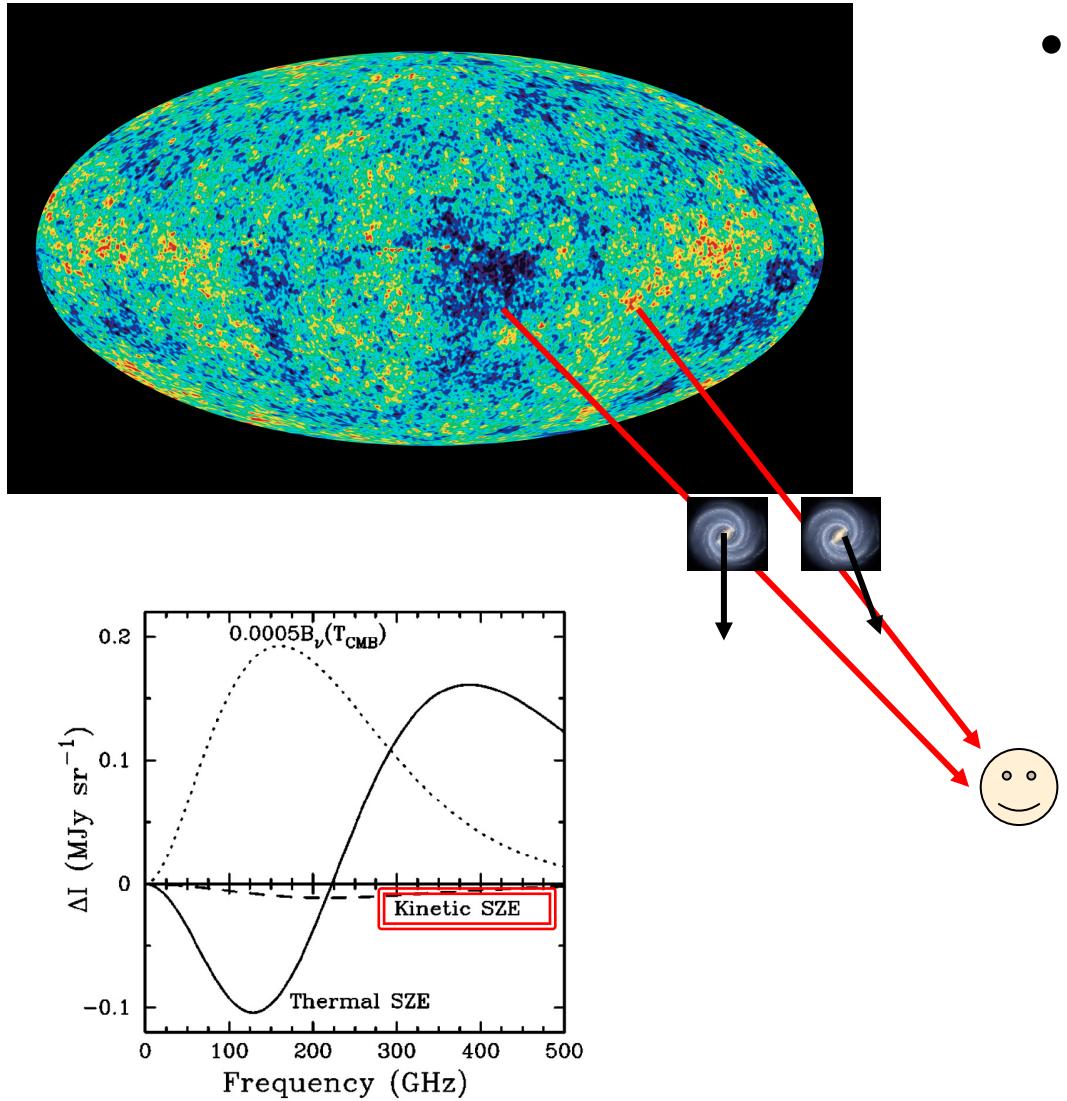
- Okumura, Nishimichi, Umetsu & Osato (2018)
- Okumura, Taruya & Nishimichi (2019, 2020)
- Okumura & Taruya (2020)
- van Gemenen & Chisari (2020)

$$\xi_{v+}(\mathbf{r}) = \tilde{C}_1 \cos(2\phi) \mu(1 - \mu^2) \Xi_{\delta\Theta, 3}^{(1)}(r)$$

$$\Xi_{XY,\ell}^{(n)}(r) = (aHf)^n \int_0^\infty \frac{k^{2-n} dk}{2\pi^2} P_{XY}(k) j_\ell(kr)$$

Direct measurement of velocities: Kinetic Sunyaev-Zeld'ovich (kSZ) effect

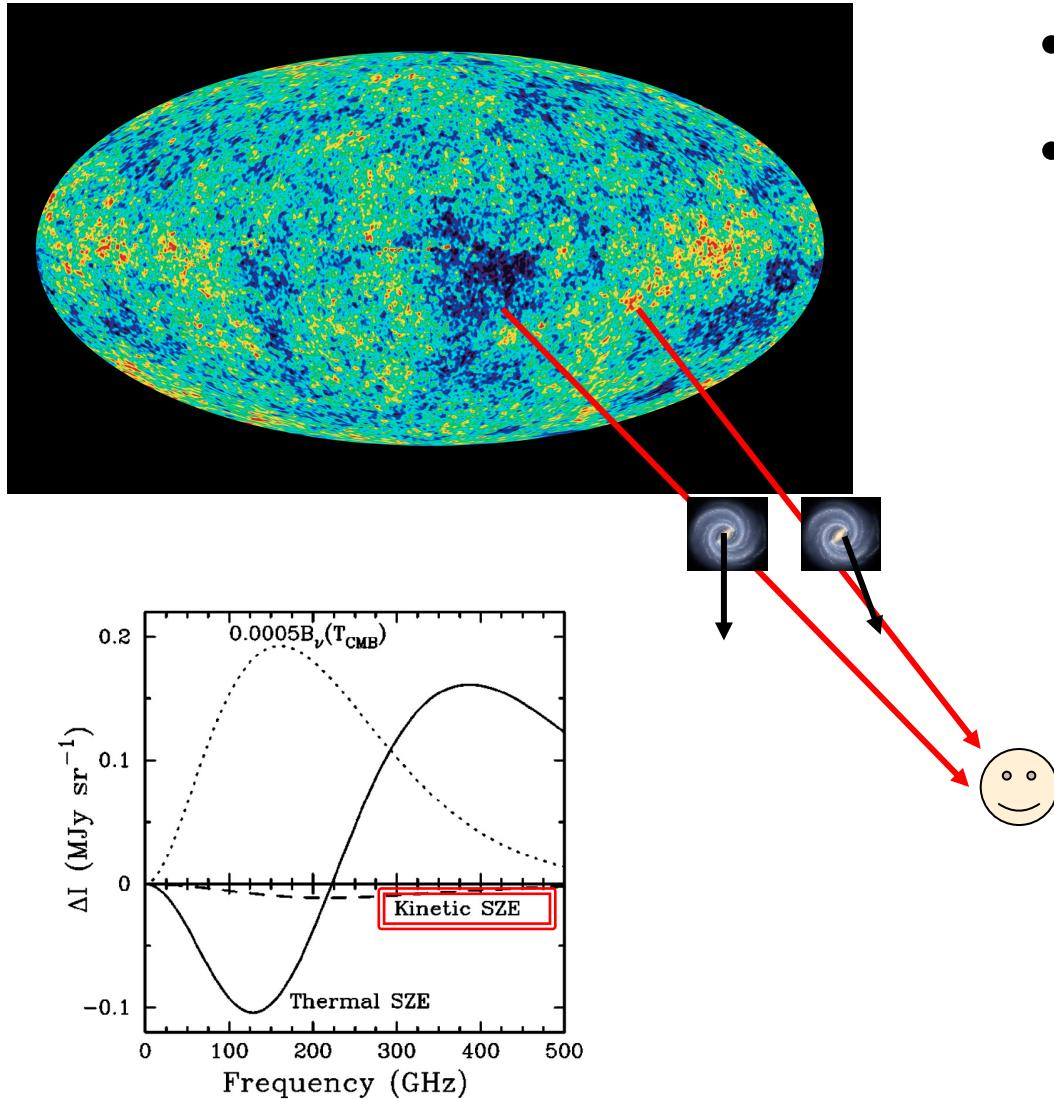
Cosmic microwave background



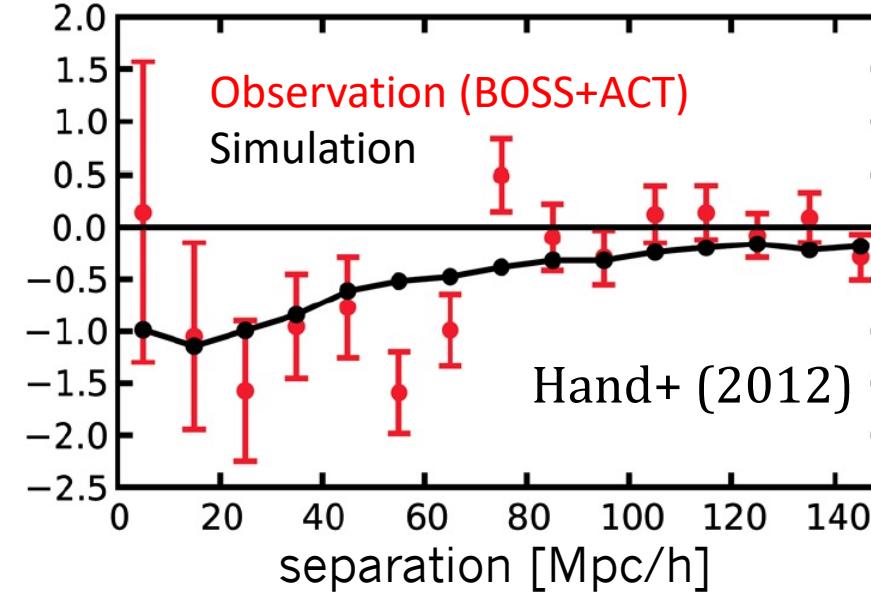
- Kinetic SZ (kSZ) effect (1980)
 - Doppler effect of cluster bulk velocity w.r.t. CMB rest frame
$$\Delta T_{kSZ} / T_{CMB} = -\tau_e v_{\parallel} \quad (v_{\parallel}: \text{line-of-sight velocity})$$
 - By measuring the temperature distortion, one can directly measure the velocity field of galaxy clusters, so it is a powerful observable to test modified gravity theories.
 - However, this effect is very tiny and hard to measure in observation.

Direct measurement of velocities: Kinetic Sunyaev-Zeld'ovich (kSZ) effect

Cosmic microwave background



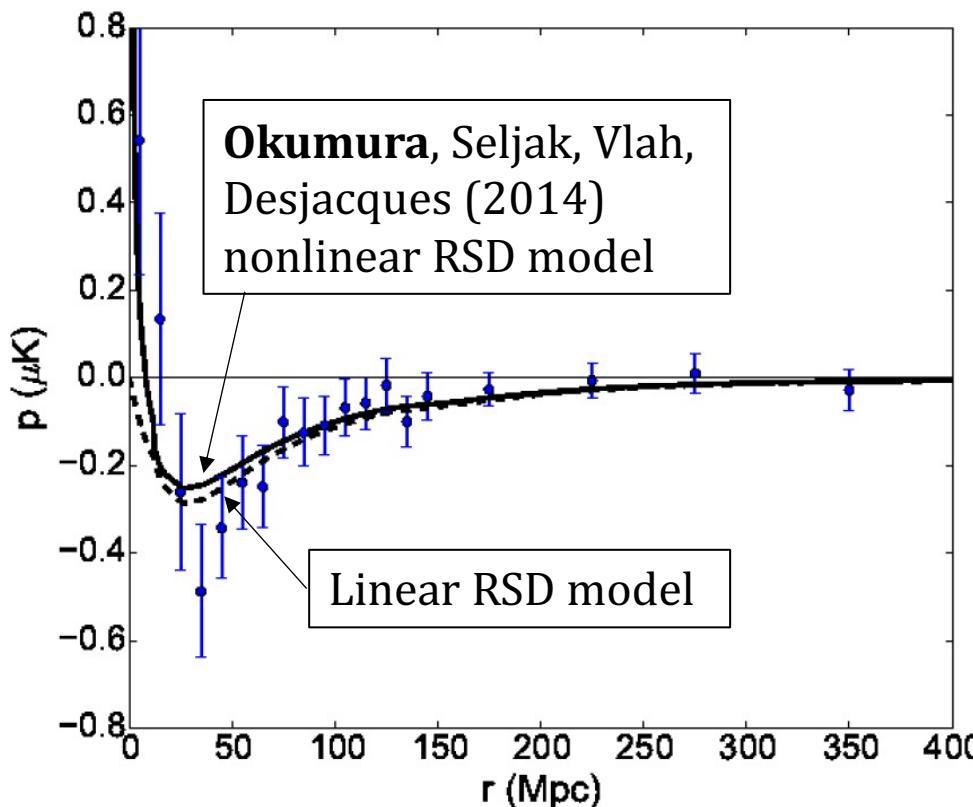
- Kinetic SZ (kSZ) effect (1980)
- First detection of kSZ effect (2012)



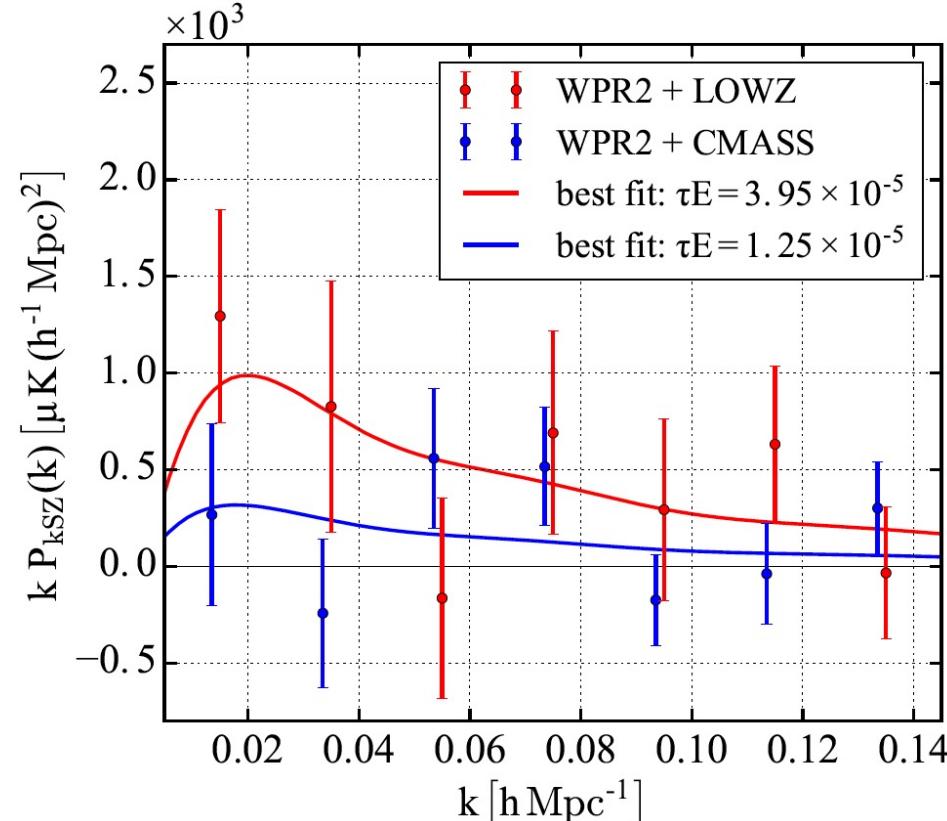
Large-scale velocity field probed via kSZ effect

$$P_{k\text{SZ}}(k) \propto P_{\delta\theta}(k) \propto P_{\delta\delta}(k)/k$$

- De Bernardis et al. (2016)
- Configuration space
- ACT x BOSS

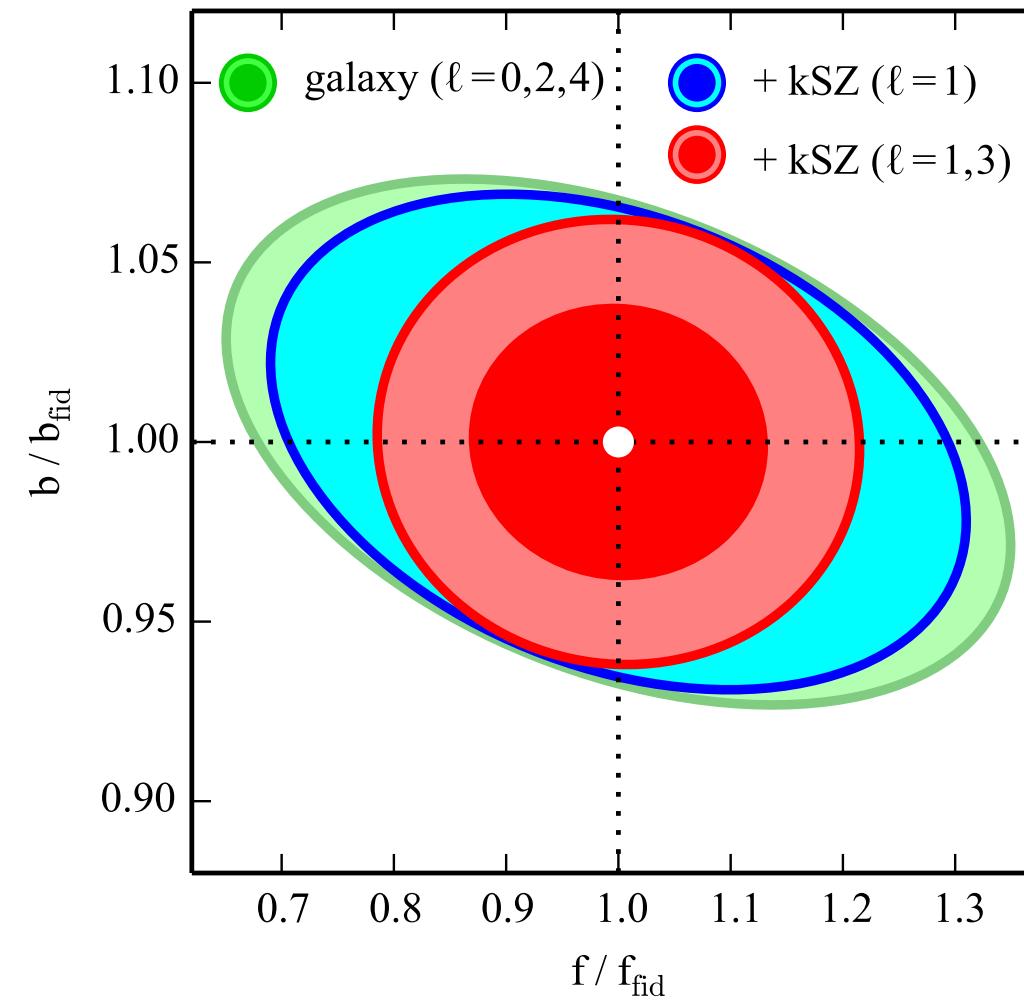
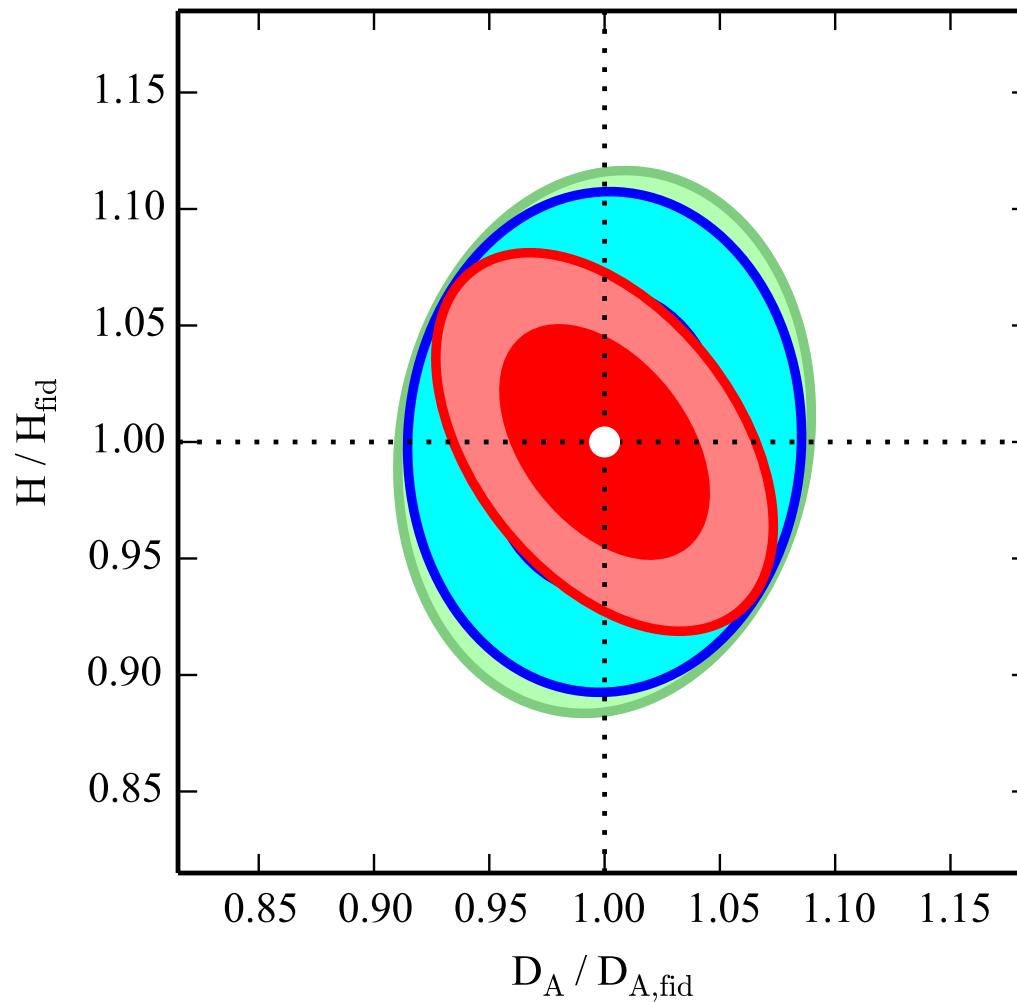


- Sugiyama, **Okumura**, Spergel (2018)
- Fourier space
- Planck x BOSS



kSZ measurements enhance the science return from galaxy redshift surveys

- Sugiyama, Okumura, Spergel (2017)
- CMB-S4 x DESI



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- **Fisher matrix forecast with galaxy clustering + IA + kSZ**
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Galaxy density, velocity and ellipticity power spectra in linear theory

- Galaxy clustering: $\delta_g^S(\mathbf{k}; z) = K_g(\mu; z)\delta_m(\mathbf{k}; z)$

$$K_g(\mu; z) = b_g(z) + f(z)\mu^2$$

- kSZ: $\delta T(\mathbf{k}; z) = (T_0\tau/c)v_{\parallel}(\mathbf{k}; z) = K_v(\mathbf{k}; z)\delta_m(\mathbf{k}; z)$

$$K_v(k, \mu; z) = i \frac{T_0\tau}{c} \frac{f(z)\mu aH(z)}{k}$$

- IA: $\gamma_E(\mathbf{k}; z) = K_E(\mu; z)\delta_m(\mathbf{k}; z)$

$$K_E(\mu; z) = b_K(z)(1 - \mu^2)$$

$$b_K(z) = 0.01344 A_{\text{IA}}(z)\Omega_m/D(z)$$

- Power spectra (**6 in total**)

$$P_{ij}(k, \mu; z) = K_i(k, \mu; z)K_j(k, \mu; z)P_{\text{lin}}(k; z)$$

- Geometric distortions

$$P_{ij}^{\text{obs}}(k_{\perp}^{\text{fid}}, k_{\parallel}^{\text{fid}}; z) = \frac{H(z)}{H^{\text{fid}}(z)} \left\{ \frac{D_A^{\text{fid}}(z)}{D_A(z)} \right\}^2 P_{ij}(k_{\perp}, k_{\parallel}; z)$$

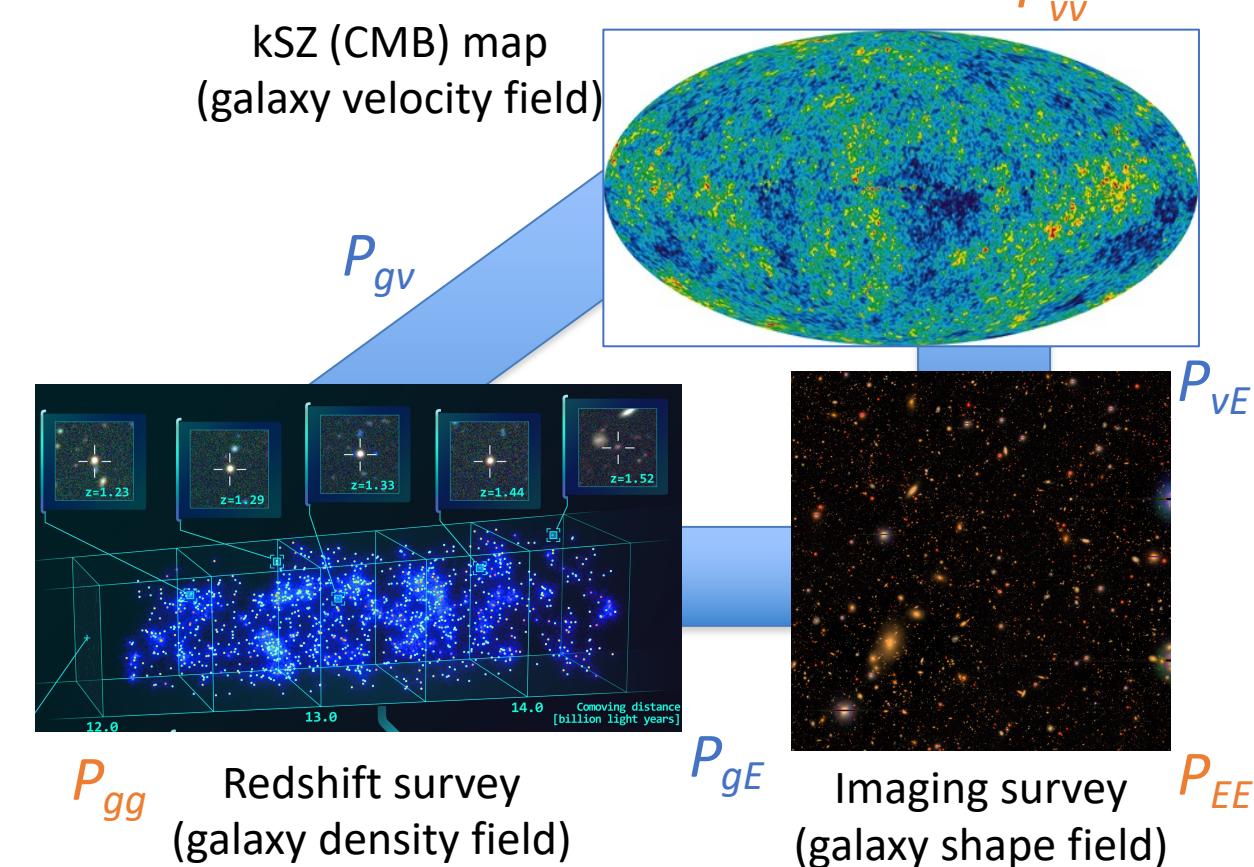
$$k_{\parallel}^{\text{fid}} = k_{\parallel} H^{\text{fid}}(z)/H(z) \quad \text{and} \quad k_{\perp}^{\text{fid}} = D_A(z)/D_A^{\text{fid}}(z)$$

- Fitting parameters

$$\theta_{\alpha} = (b\sigma_8, A_{\text{IA}}\sigma_8, \tau, f\sigma_8, H, D_A)$$

Amplitude (nuisance)
parameters

Dynamical and
geometric quantities



Fisher matrix formalism

$$F_{\alpha\beta} = \frac{V_s}{4\pi^2} \int_{k_{\min}}^{k_{\max}} dk k^2 \int_{-1}^1 d\mu \sum_{a,b=1}^{N_P} \frac{\partial P_a(k, \mu)}{\partial \theta_\alpha} [\text{Cov}^{-1}]_{ab} \frac{\partial P_b(k, \mu)}{\partial \theta_\beta}$$

- 6 x 6 Gaussian covariance matrix

$$\text{Cov}_{ab}(k, \mu) =$$

<i>Auto</i>	$P_{gg} \rightarrow$	$2\{\tilde{P}_{gg}\}^2 \quad 2\{P_{gE}\}^2 \quad 2\{P_{gv}\}^2 \quad 2\tilde{P}_{gg}P_{gE} \quad 2\tilde{P}_{gg}P_{gv} \quad 2P_{gv}P_{gE}$
	$P_{EE} \rightarrow$	$2\{P_{gE}\}^2 \quad 2\{\tilde{P}_{EE}\}^2 \quad 2\{P_{vE}\}^2 \quad 2P_{gE}\tilde{P}_{EE} \quad 2P_{gE}P_{vE} \quad 2\tilde{P}_{EE}P_{vE}$
<i>power</i>	$P_{vv} \rightarrow$	$2\{P_{gv}\}^2 \quad 2\{P_{vE}\}^2 \quad 2\{\tilde{P}_{vv}\}^2 \quad 2P_{gv}P_{vE} \quad 2\tilde{P}_{vv}P_{gv} \quad 2P_{vE}\tilde{P}_{vv}$
<i>Cross</i>	$P_{gE} \rightarrow$	$2\tilde{P}_{gg}P_{gE} \quad 2P_{gE}\tilde{P}_{EE} \quad 2P_{gv}P_{vE} \quad \tilde{P}_{gg}\tilde{P}_{EE} + \{P_{gE}\}^2 \quad \tilde{P}_{gg}P_{vE} + P_{gE}P_{gv} \quad P_{gv}\tilde{P}_{EE} + P_{gE}P_{vE}$
<i>power</i>	$P_{gv} \rightarrow$	$2\tilde{P}_{gg}P_{gv} \quad 2P_{gE}P_{vE} \quad 2\tilde{P}_{vv}P_{gv} \quad \tilde{P}_{gg}P_{vE} + P_{gE}P_{gv} \quad \tilde{P}_{gg}\tilde{P}_{vv} + \{P_{gv}\}^2 \quad P_{gE}\tilde{P}_{vv} + P_{gv}P_{vE}$
	$P_{vE} \rightarrow$	$2P_{gv}P_{gE} \quad 2\tilde{P}_{EE}P_{vE} \quad 2P_{vE}\tilde{P}_{vv} \quad P_{gv}\tilde{P}_{EE} + P_{gE}P_{vE} \quad P_{gE}\tilde{P}_{vv} + P_{gv}P_{vE} \quad \tilde{P}_{EE}\tilde{P}_{vv} + \{P_{vE}\}^2$

- Poisson shot noise

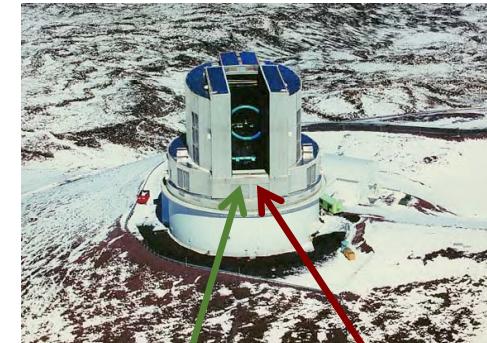
$$\tilde{P}_{gg} = P_{gg} + \frac{1}{n_g}, \quad \tilde{P}_{vv} = P_{vv} + (1 + R_N^2) \left(\frac{T_0 \tau}{c} \right)^2 \frac{(f\sigma_v)^2}{n_v}, \quad \tilde{P}_{EE} = P_{EE} + \frac{\sigma_\gamma^2}{n_\gamma}$$

Survey setup

- Assume a PFS-like emission line galaxy (ELG) survey

- Parameters from PFS white paper (Takada et al 2014)

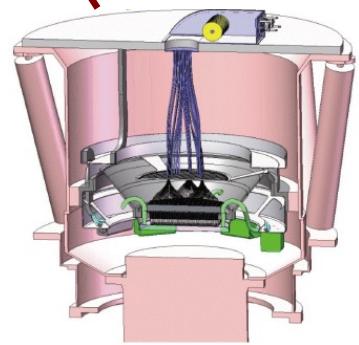
Redshift		Volume V_s ($h^{-3}\text{Gpc}^3$)	$10^4 n$ ($h^3\text{Mpc}^{-3}$)	b_g
z_{\min}	z_{\max}			
0.6	0.8	0.59	1.9	1.18
0.8	1.0	0.79	6.0	1.26
1.0	1.2	0.96	5.8	1.34
1.2	1.4	1.09	7.8	1.42
1.4	1.6	1.19	5.5	1.50
1.6	2.0	2.58	3.1	1.62
2.0	2.4	2.71	2.7	1.78



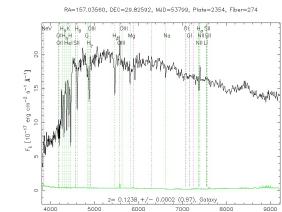
Subaru (NAOJ)



HSC



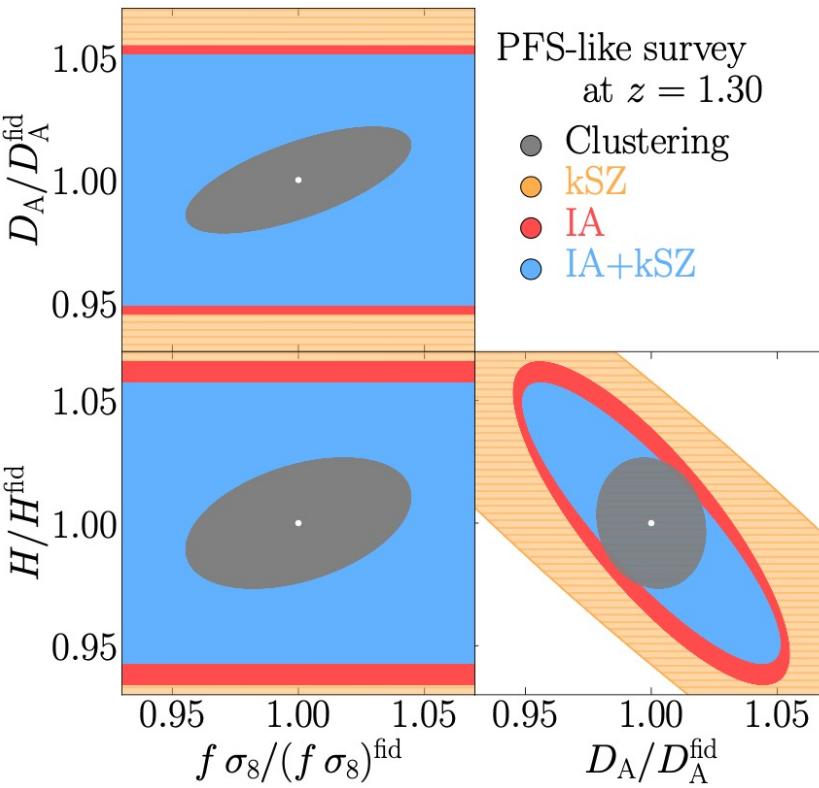
PFS



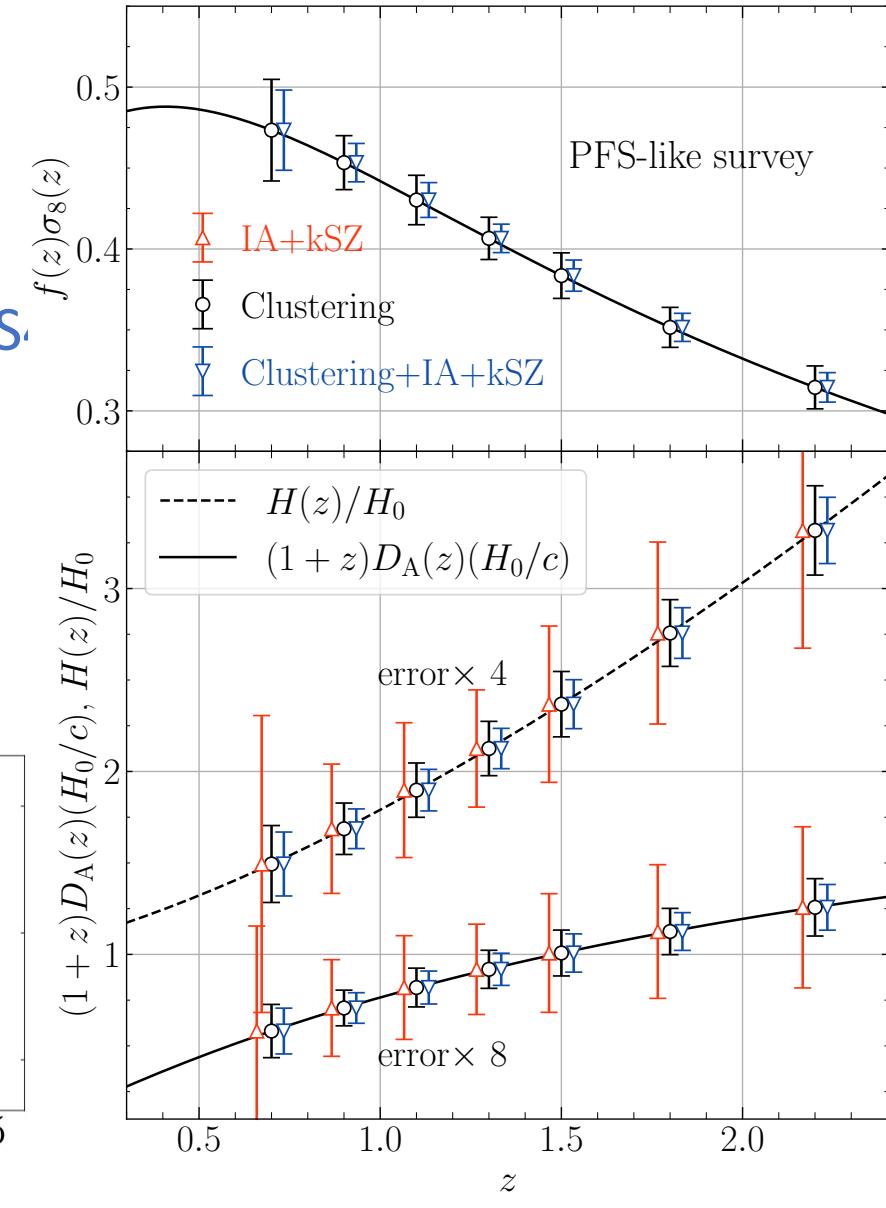
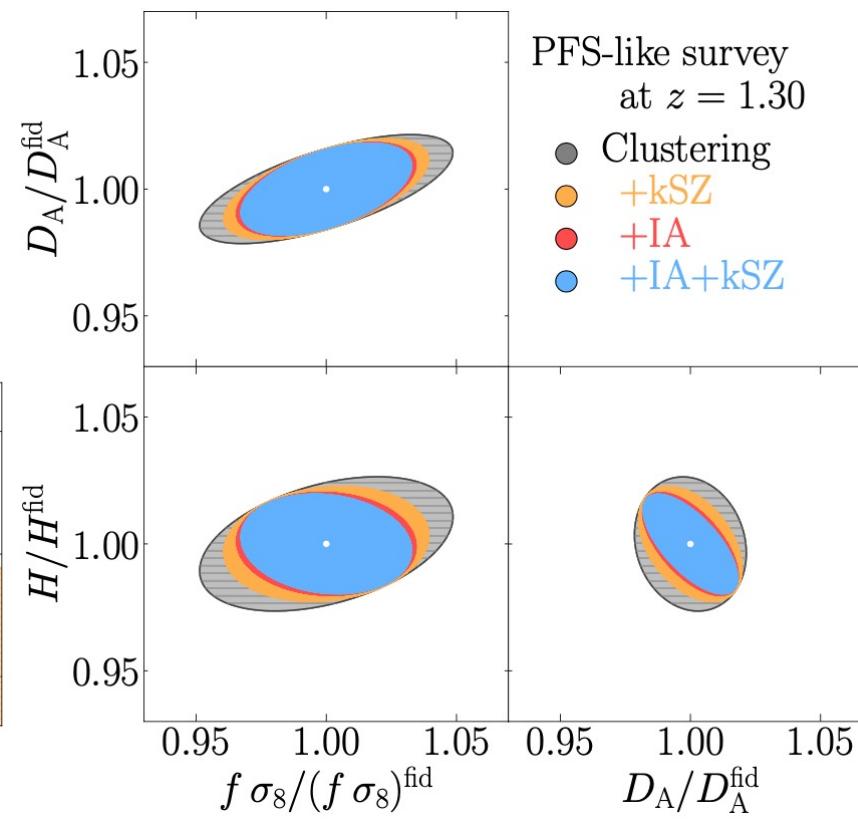
- Intrinsic alignment:
 - Beautiful galaxy images are obtained thanks to Hyper Suprime-Cam (HSC), $\sigma_y=0.2$.
 - Shi et al (2021) proposed an estimator to directly detect IA of host halos using the observation of ELGs, $A_{IA} = 18$.
- kSZ:
 - CMB-S4, which is completely overlapped with the area of the PFS
 - Fiducial values: linear theory for σ_ν , and the inverse S/N of the kSZ temperature fluctuations $R_N = 10$ (Sugiyama et al 2017).

Geometric and dynamical constraints

- Clustering (PFS)
- kSZ (PFS + CMB-S4)
- IA (PFS + HSC)
- IA+kSZ (PFS + HSC + CMB-S4)



- Clustering (PFS)
- + kSZ (PFS + CMB-S4)
- + IA (PFS + HSC)
- + IA+kSZ (PFS + HSC + CMB-S4)



Cosmological constraints (1)

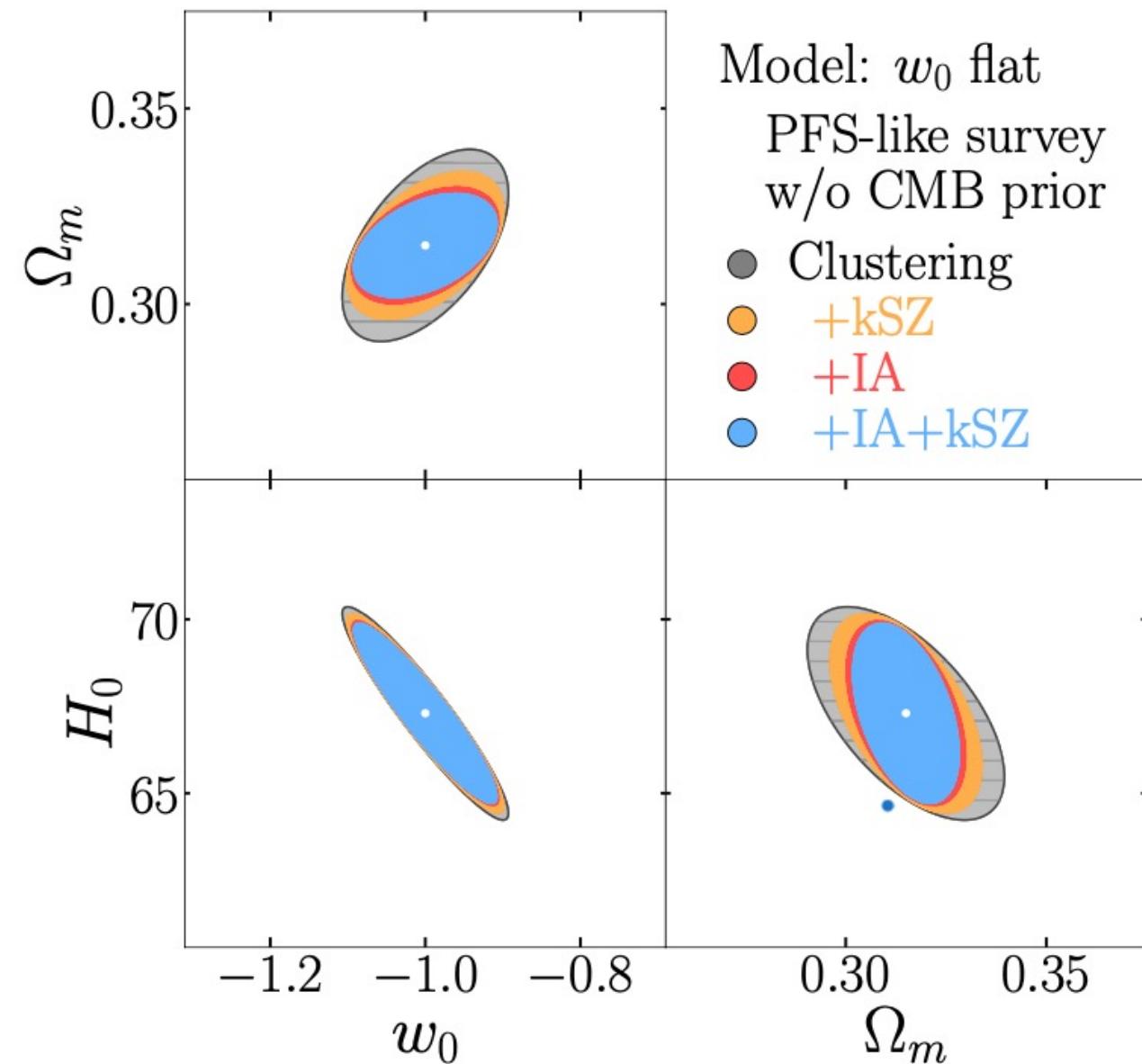
- Projection of the Fisher matrix to the cosmological parameter space:

$$S_{AB} = \sum_{\alpha, \beta} \frac{\partial \theta_\alpha}{\partial p_A} F_{\alpha\beta} \frac{\partial \theta_\beta}{\partial p_B}$$

$$\theta_\alpha = (b\sigma_8, A_{\text{IA}}\sigma_8, \tau, f\sigma_8, H, D_A)$$

$$\rightarrow p_A = (\Omega_m, w_0, H_0, \sigma_8)$$

constant w , flat ($\Omega_K = 0$) model



Cosmological constraints (2)

- Projection of the Fisher matrix to the cosmological parameter space:

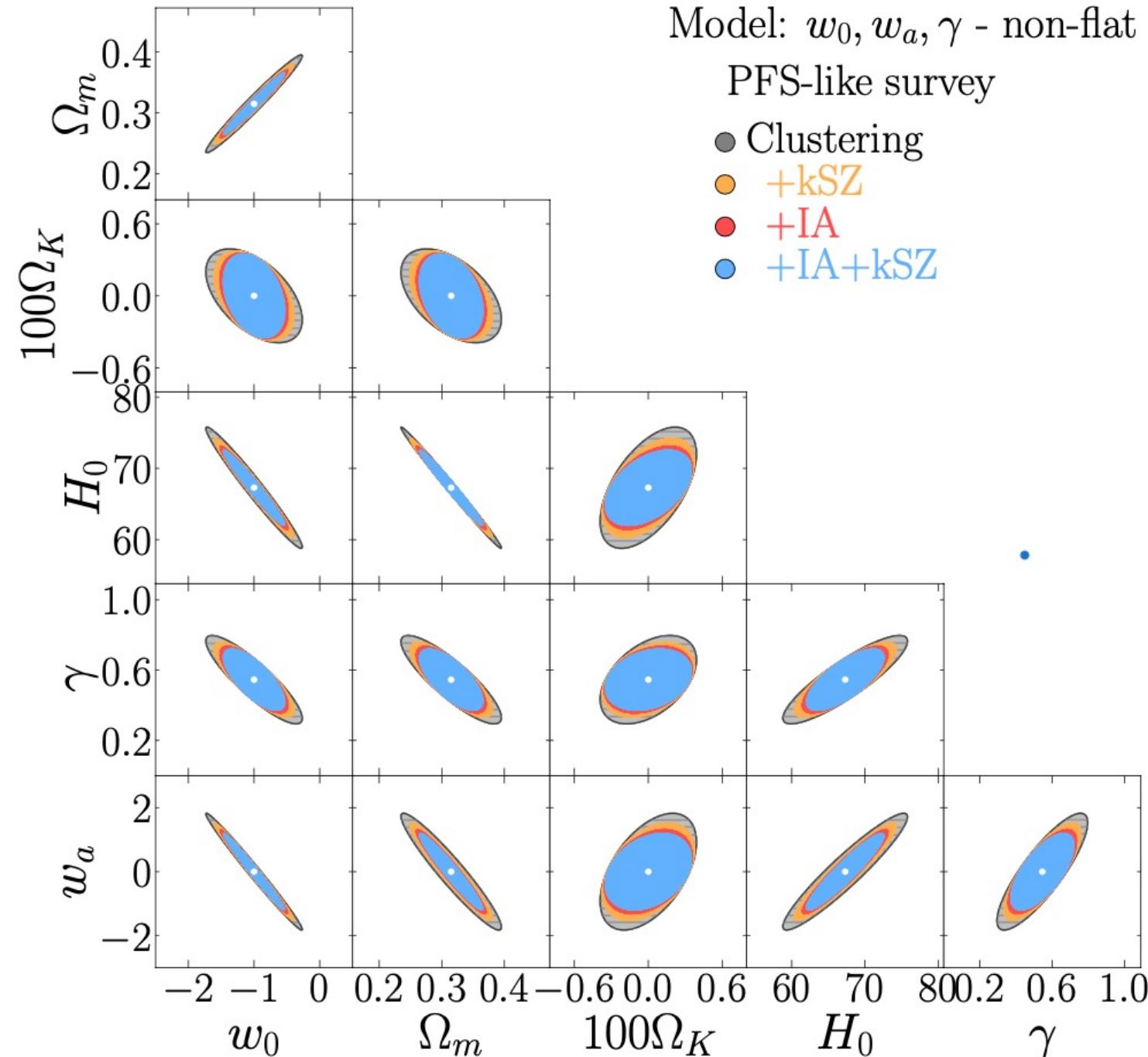
$$S_{AB} = \sum_{\alpha, \beta} \frac{\partial \theta_\alpha}{\partial p_A} F_{\alpha\beta} \frac{\partial \theta_\beta}{\partial p_B}$$

$$\theta_\alpha = (b\sigma_8, A_{\text{IA}}\sigma_8, \tau, f\sigma_8, H, D_A)$$

$$\rightarrow p_A = (\Omega_m, \Omega_K, w_0, w_a, H_0, \gamma, \sigma_8)$$

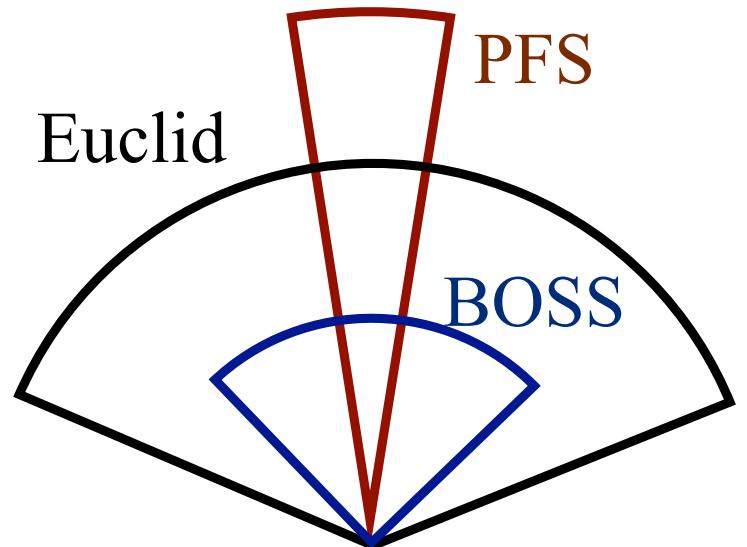
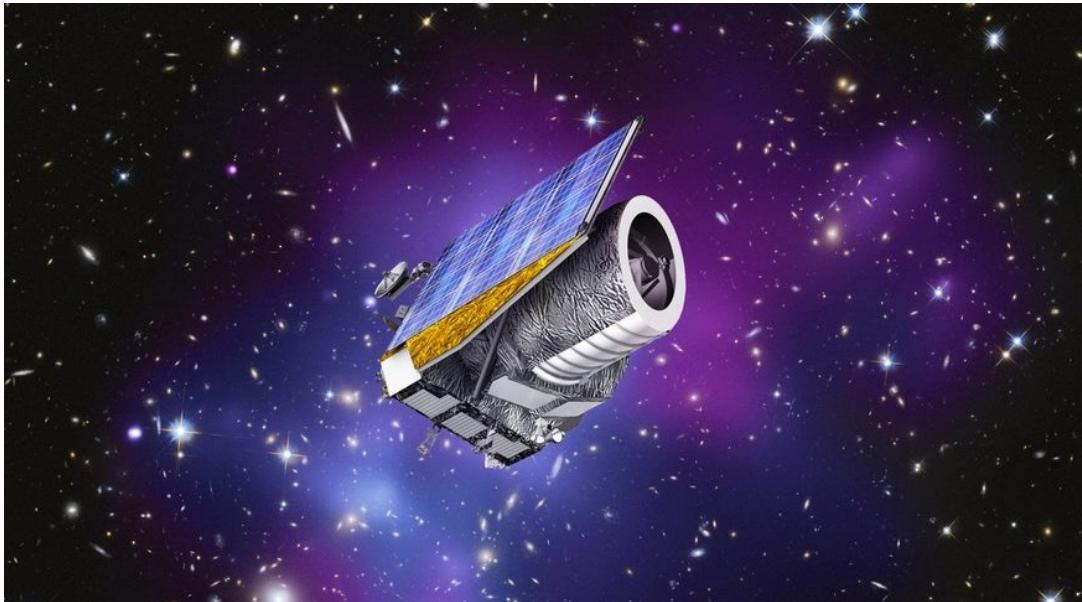
time-varying $w(a) = w_0 + (1-a)w_a$,
 non-flat ($\Omega_K \neq 0$) model with
 modified gravity parameter γ

$$S_{AB} = S_{AB}^{\text{LSS}} + S_{AB}^{\text{CMB}}$$



Deep vs wide surveys

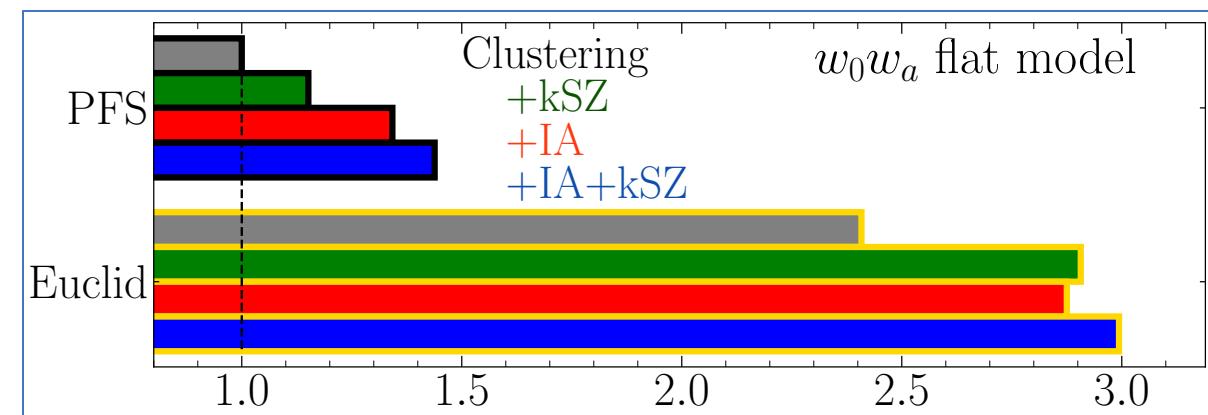
- Euclid satellite mission for galaxy surveys



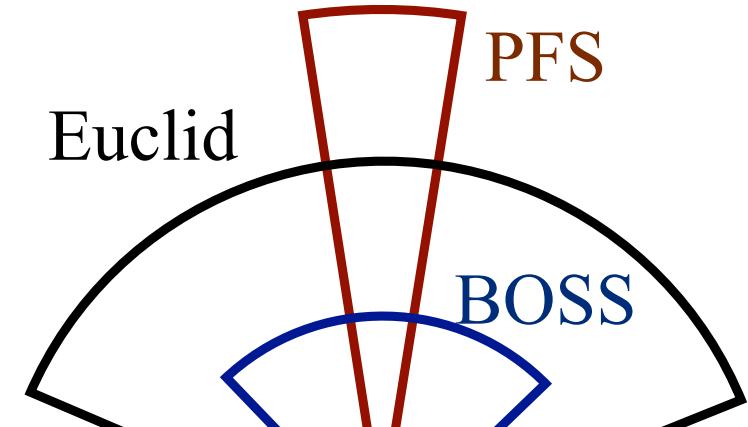
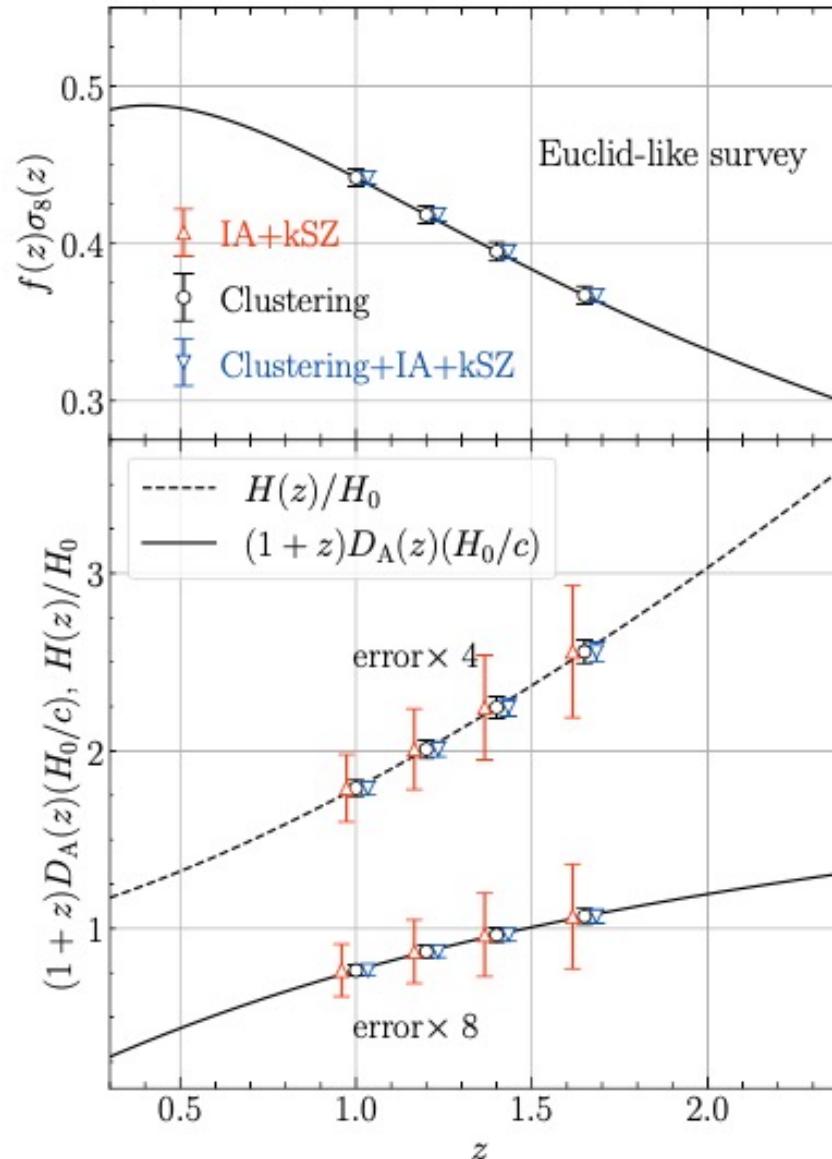
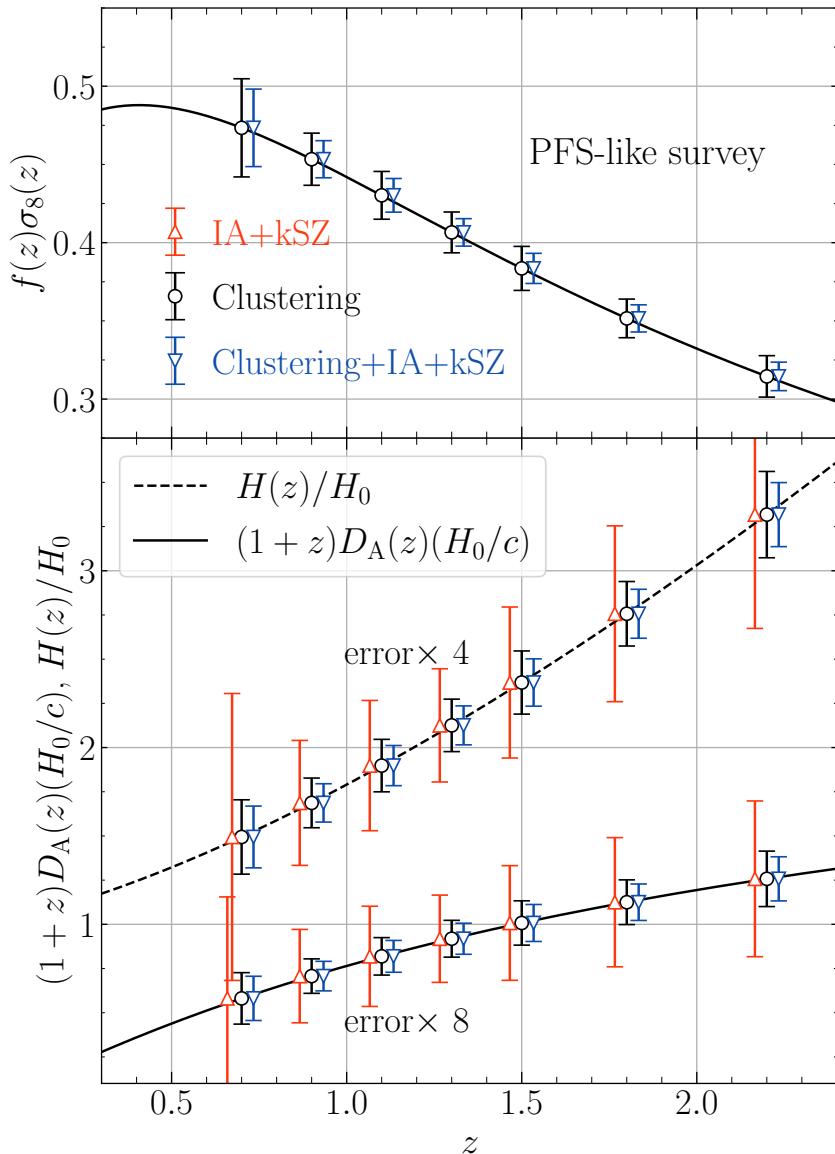
- Figure-of-Merit (FoM)

$$\text{FoM} = \left\{ \det(\bar{S}_{AB}) \right\}^{1/N_p}$$

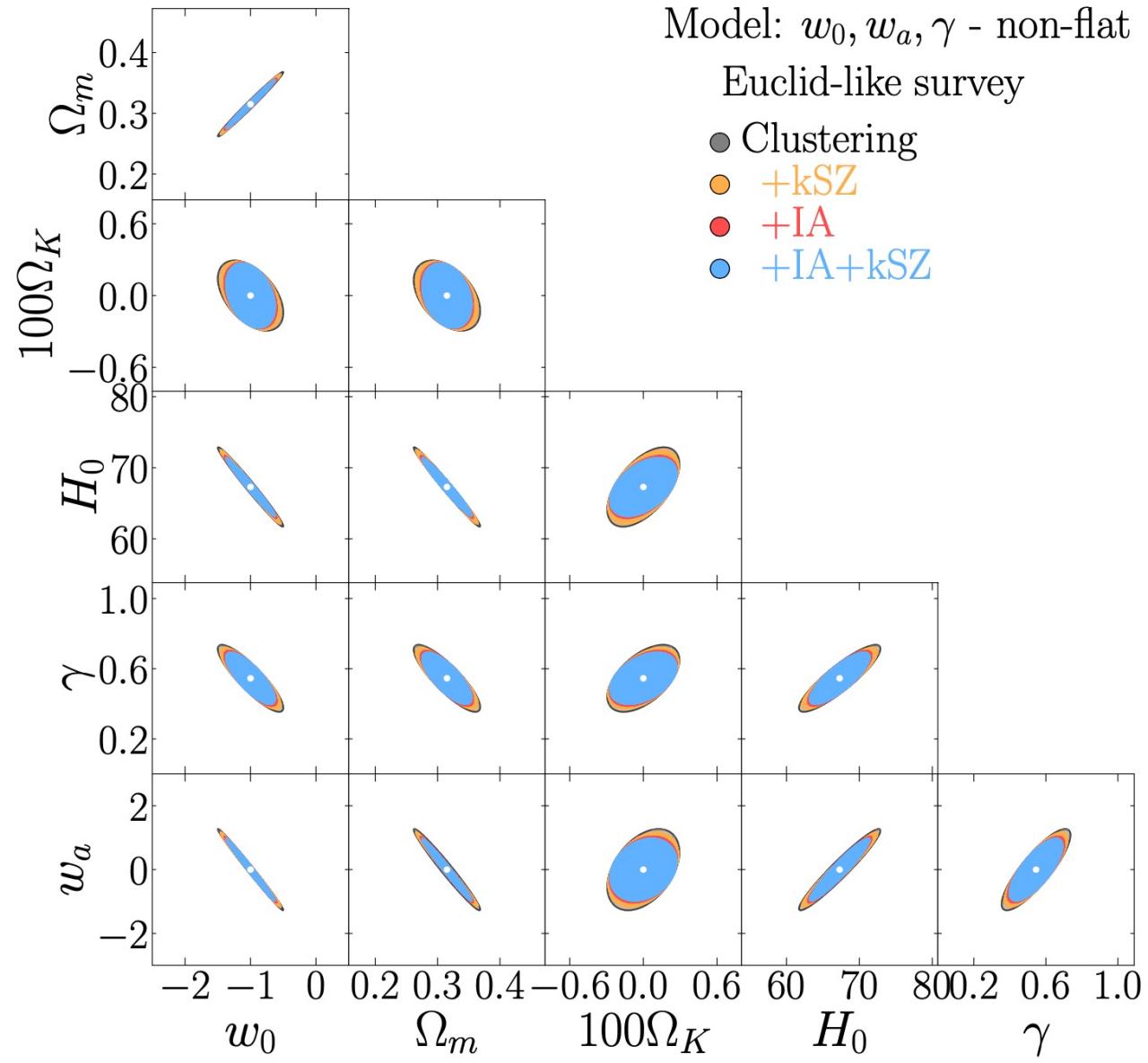
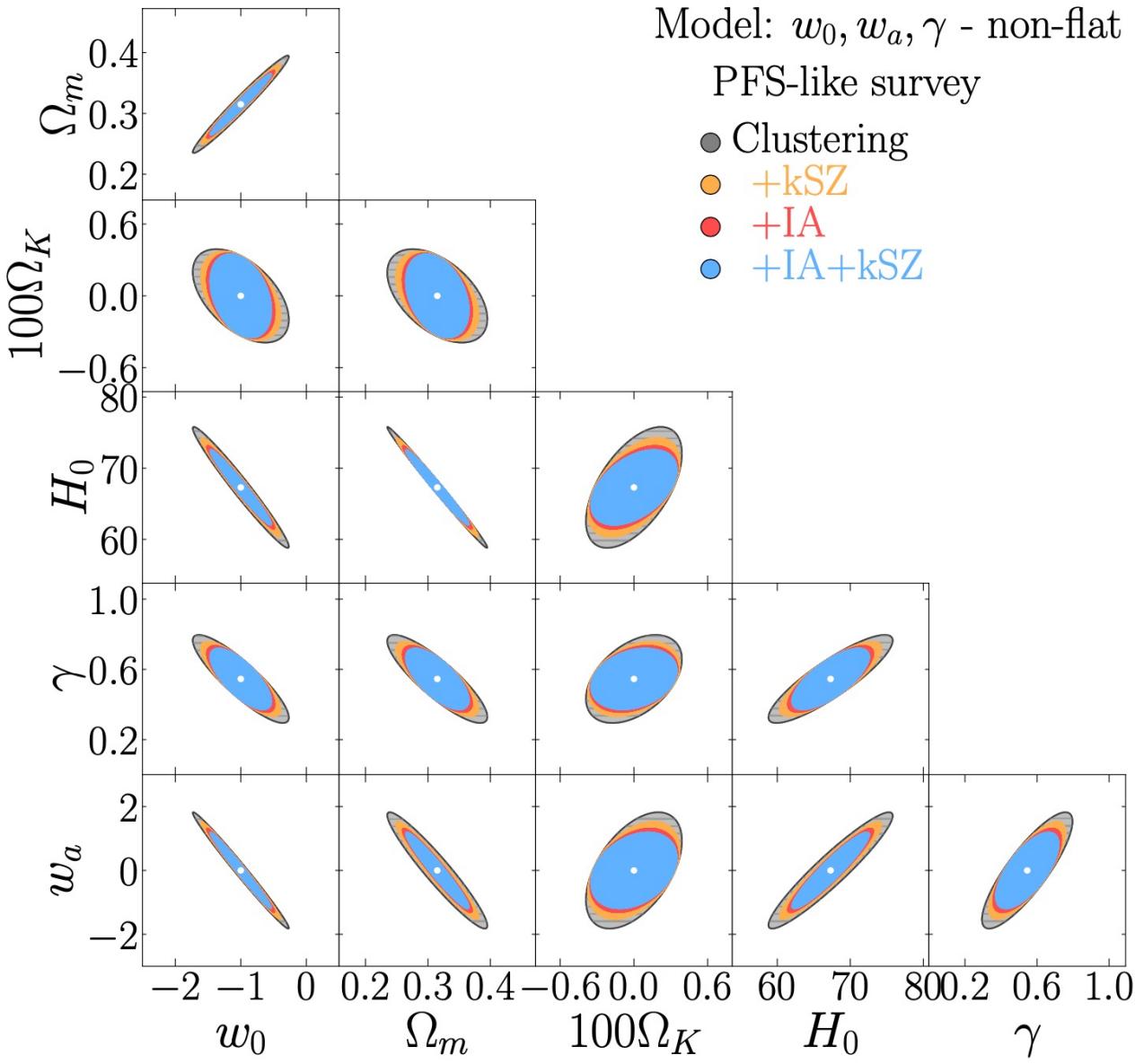
- Normalize by the clustering-only case with PFS



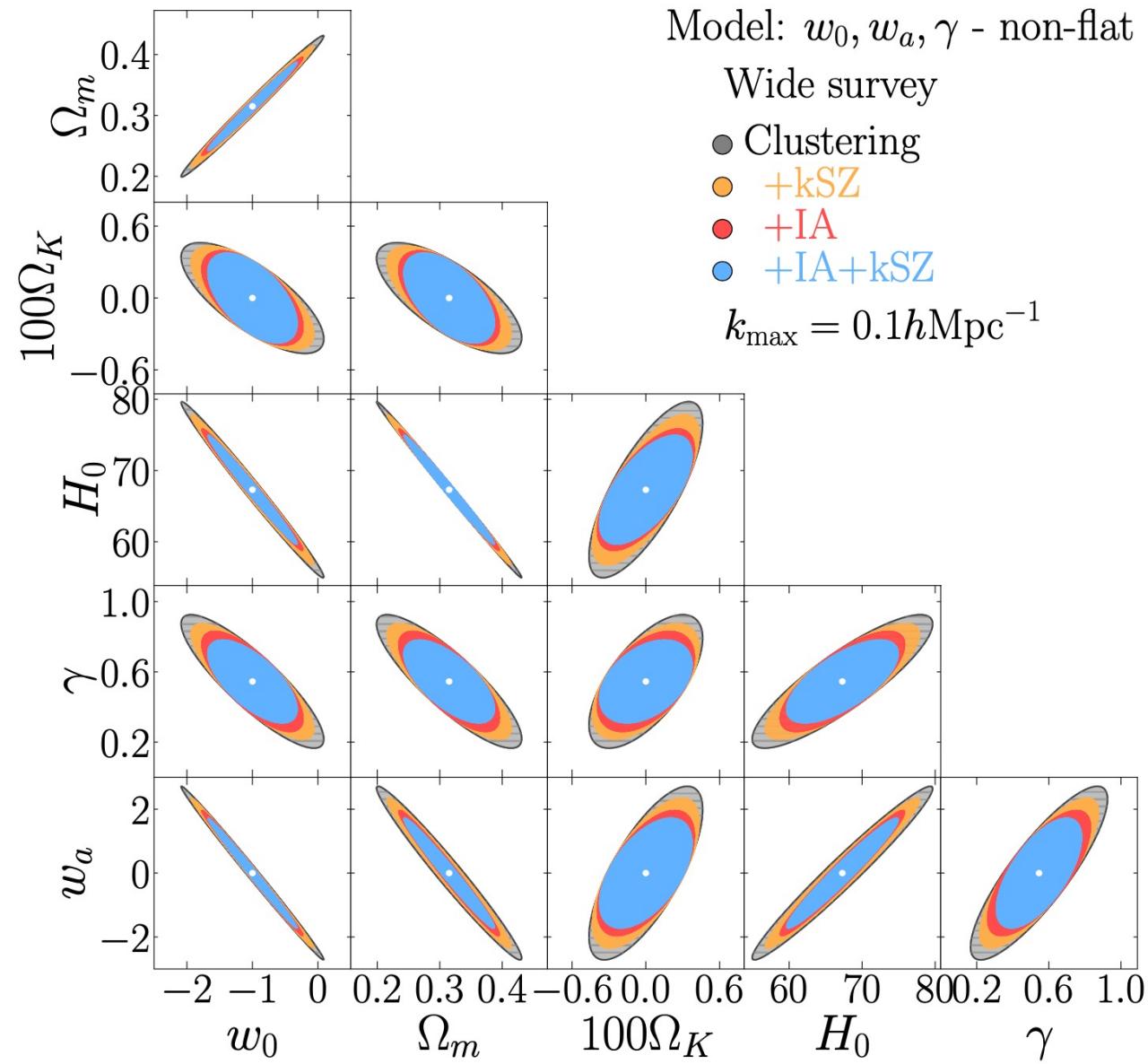
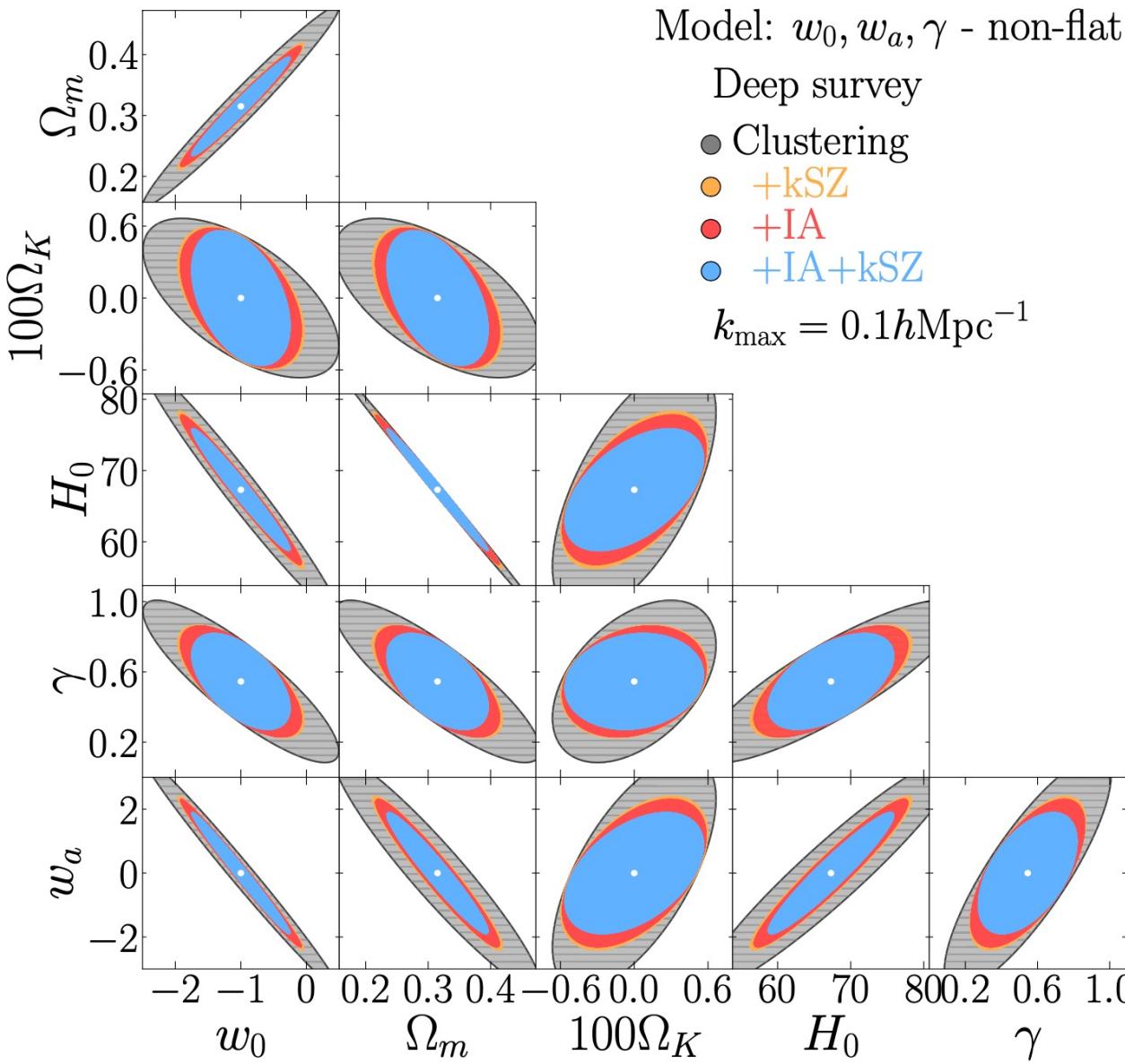
Deep vs wide surveys



Deep vs wide surveys



Conservative analysis with cutoff of $k_{\text{max}} = 0.10 h/\text{Mpc}$



Outline

- Galaxy intrinsic alignment (IA) as a dynamical and geometric probe
- Kinetic Sunyaev-Zel'dovich (kSZ) effect as a dynamical and geometric probe
- Fisher matrix forecast with galaxy clustering + IA + kSZ
 - Geometric and dynamical constraints
 - Cosmological parameter constraints
 - Deep vs wide galaxy surveys
- **Measurement of IA of red galaxies at $z > 1$**

First Evidence of Intrinsic Alignments of Red Galaxies at $z > 1$: Cross-correlation between CFHTLenS and FastSound Samples

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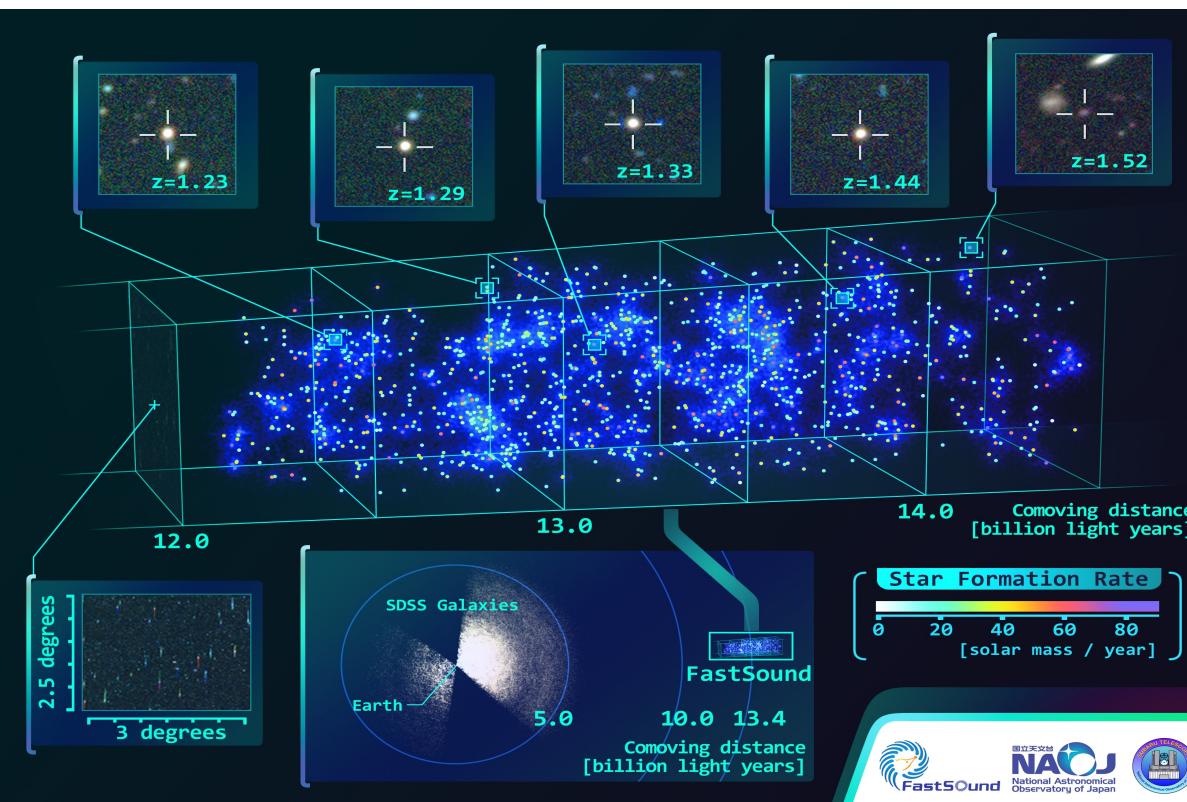
ABSTRACT

We report the first evidence of intrinsic alignment (IA) of red galaxies at $z > 1$. We measure the gravitational shear-intrinsic ellipticity (GI) cross-correlation function at $z \sim 1.3$ using galaxy positions from the FastSound spectroscopic survey and galaxy shapes from Canada-Hawaii-France telescope lensing survey data. Adopting the non-linear alignment model, we obtain a 2.4σ level detection of the IA amplitude, $A^{\text{LA}} = 27.48_{-11.54}^{+11.53}$, larger than the value extrapolated from the constraints obtained at lower redshifts. Our measured IA is translated into a $\sim 20\%$ contamination to the weak lensing power spectrum for the red galaxies. This marginal detection of IA for red galaxies at $z > 1$ motivates the continuing investigation of the nature of IA for weak lensing studies. Furthermore, our result provides the first step to utilize IA measurements in future high- z surveys as a cosmological probe, complementary to galaxy clustering and lensing.

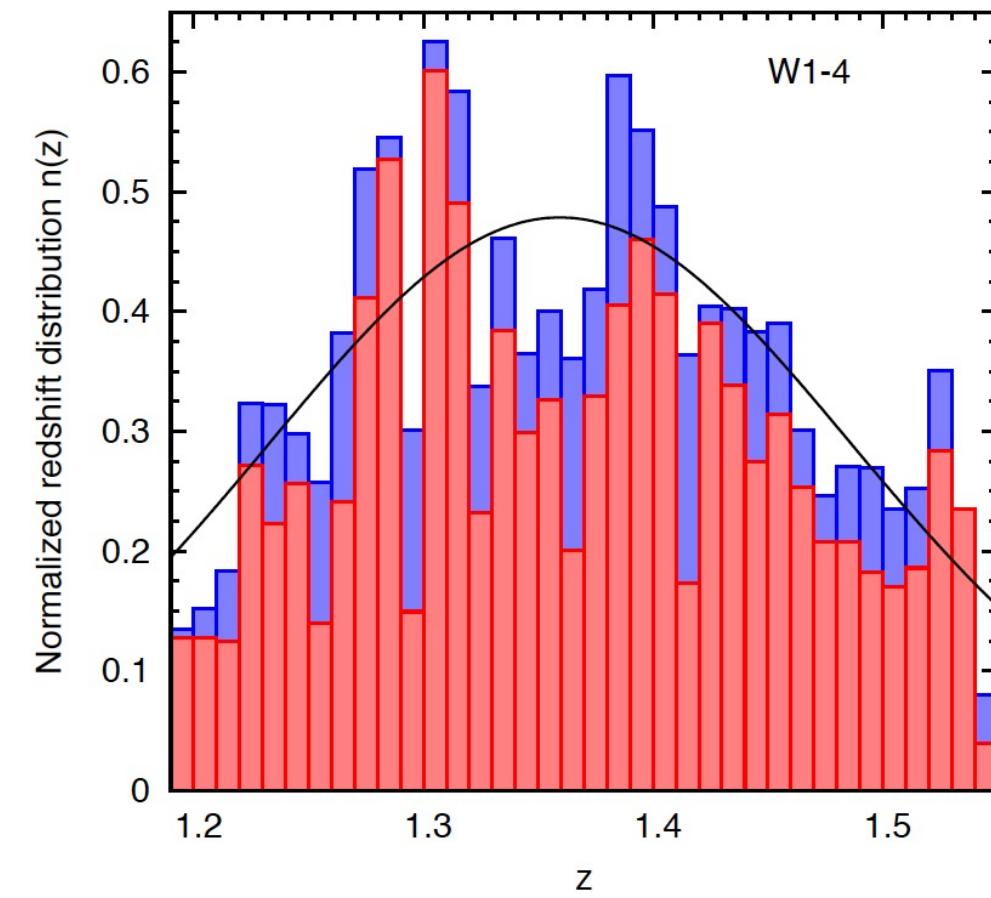


M. Tonegawa
(postdoc at
APCTP, Korea)

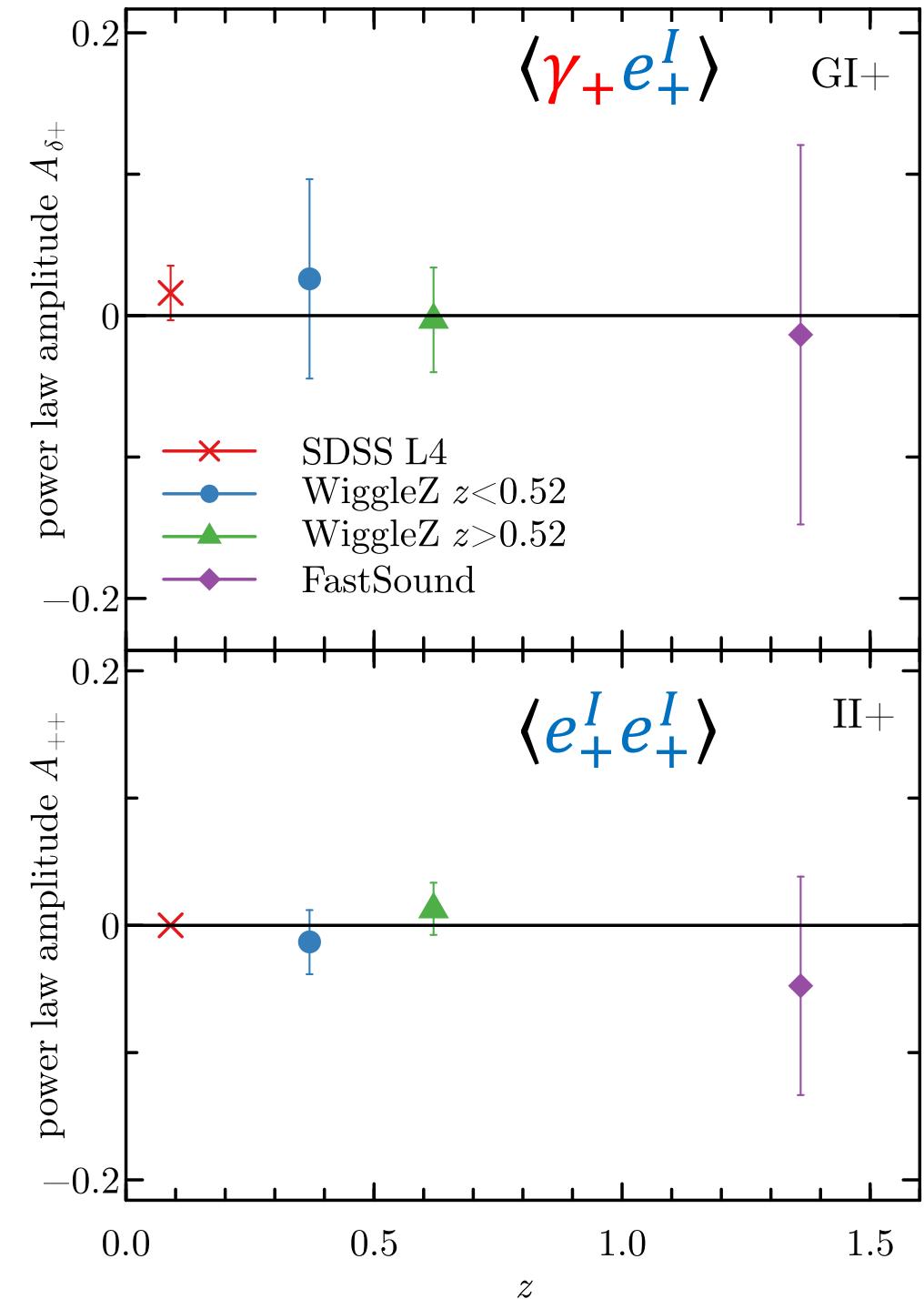
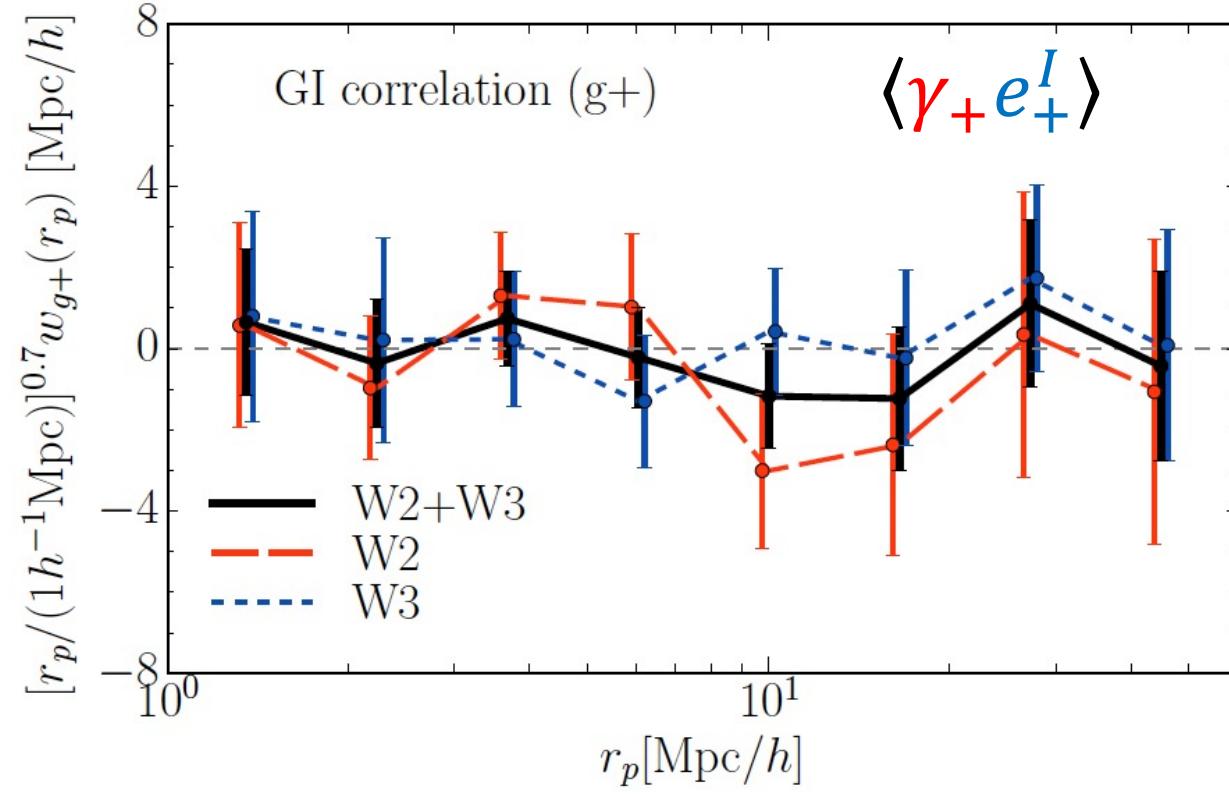
GI correlation measurement from CFHTLens-FastSound



- $z \sim 1.2-1.5$
- Shape sample: ~ 12000 LRGs from CFHTLenS
- Density sample: ~ 3000 H α from FastSound

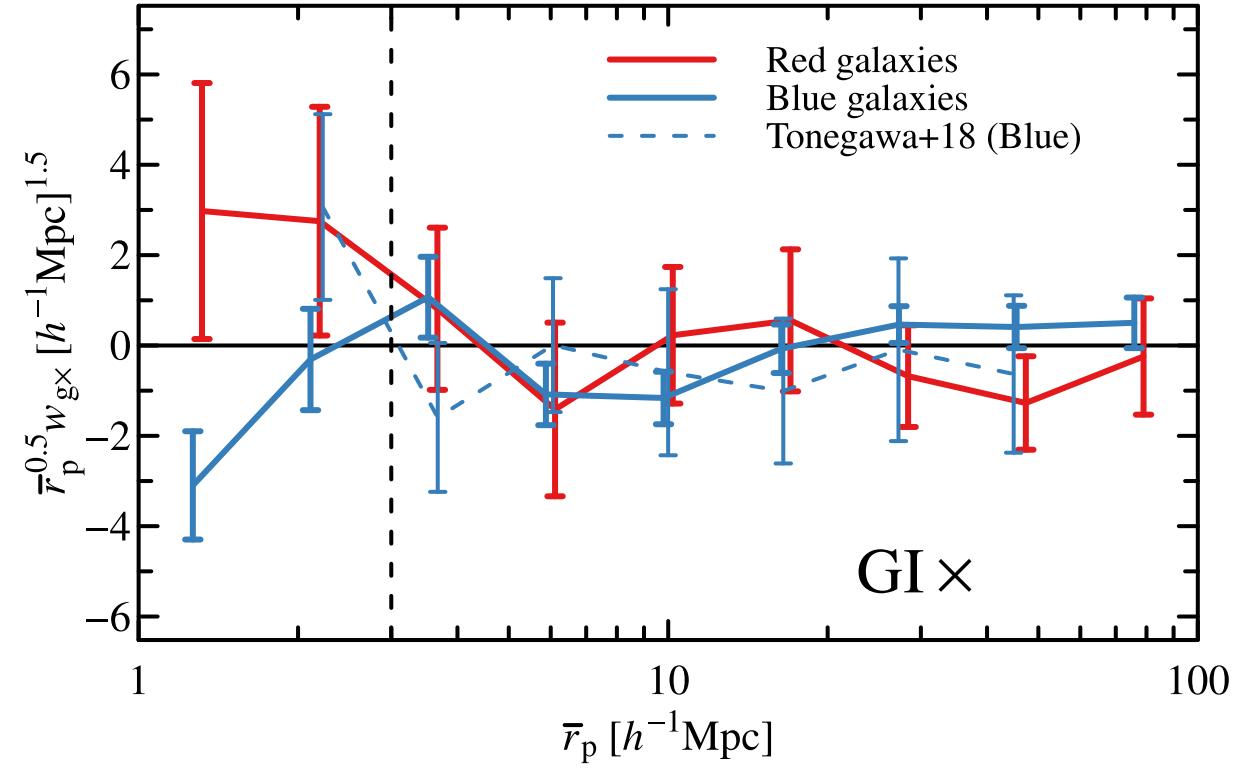
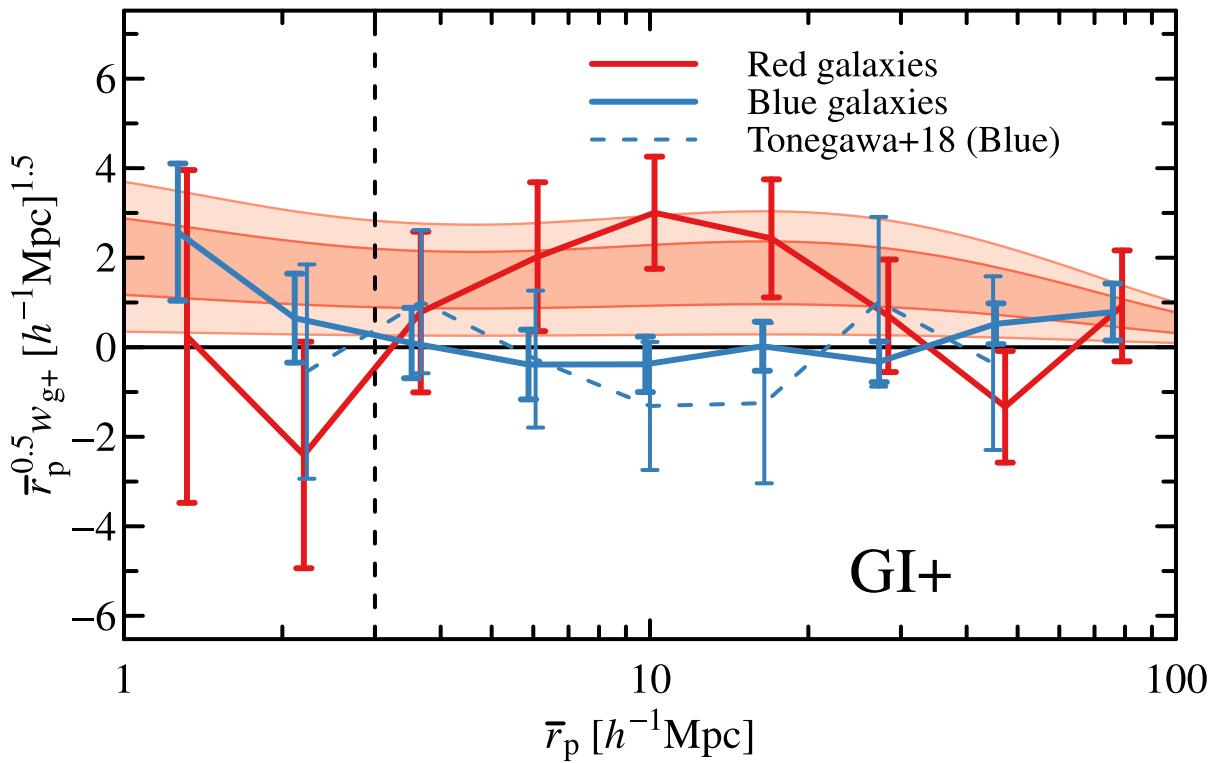


*Previous result:
No IA correlation for H α*

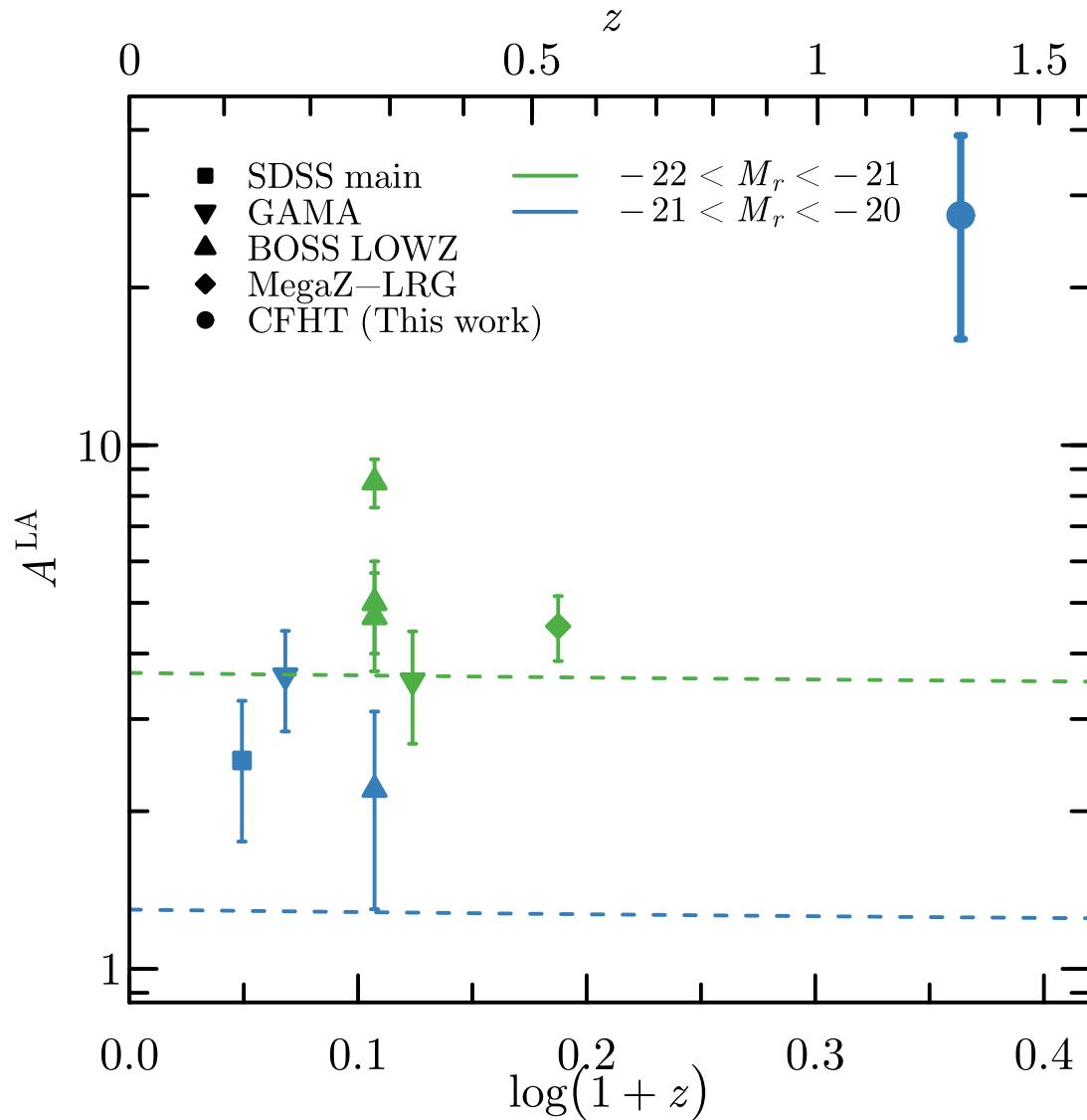


2.4- σ detection of IA for LRGs at $z>1$

- GI correlation function
 - Galaxy shape: CFHTLenS photo-z LRG
 - Galaxy position: FastSound spec-z H α emitters



2.4- σ detection of IA for LRGs at $z > 1$



- 2.4- σ (w/o shear and magnification)
 $A^{\text{LA}} = 27.48^{+11.53}_{-11.54}$
- With gG and mG, the significance gets larger.
- It is straightforward to extend the analysis to HSC LRG shape – PFS [OII] emitter position cross-correlation.

Conclusions

- Conventionally, cosmological constraints on the growth and expansion history of the universe have been obtained from the measurements of RSD and BAO embedded in the galaxy distribution.
- We studied how well one can improve the cosmological constraints from the combination of the galaxy density field with velocity (kSZ) and ellipticity (IA) fields.
- For illustration, we consider the Subaru PFS whose survey footprint perfectly overlaps with the HSC and CMB-S4 experiment.
- We found adding the kSZ and IA effects significantly improves cosmological constraints.
- We measured the GI cross-correlation at $z>1$ and found the 2.4σ evidence of IA of elliptical galaxies, which can be greatly improved by upcoming surveys like DESI.