

Irreducible Decomposition of a Superconducting Gap Function

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Consider the mean-field many-body Hamiltonian for a superconducting gap function,

$$\hat{\Delta}(\mathbf{k}) = \frac{1}{2} \sum_{\mathbf{k}, i, j} \psi_i^\dagger(\mathbf{k}) \Delta_{ij}(\mathbf{k}) \psi_j^\dagger(-\mathbf{k}) + h.c. \quad (1)$$

Assume that the normal state has unitary symmetries of a group G :

$$\hat{g} \psi_i^\dagger(\mathbf{k}) \hat{g}^{-1} = \psi_j^\dagger(p_g \mathbf{k}) [u_g]_{ji}, \quad u_g u_h = z_{g,h} u_{gh}, \quad \hat{g} i \hat{g}^{-1} = i, \quad g, h \in G. \quad (2)$$

Then

$$\hat{g} \hat{\Delta} \hat{g}^{-1} = \frac{1}{2} \sum_{\mathbf{k}, i, j} \psi_i^\dagger(p_g \mathbf{k}) [u_g \Delta(\mathbf{k}) u_g^T]_{ij} \psi_j^\dagger(-p_g \mathbf{k}) + h.c. \quad (3)$$

$$= \frac{1}{2} \sum_{\mathbf{k}, i, j} \psi_i^\dagger(\mathbf{k}) [u_g \Delta(p_g^{-1} \mathbf{k}) u_g^T]_{ij} \psi_j^\dagger(-\mathbf{k}) + h.c. \quad (4)$$

Thus we introduce the left action of G on the gap function $\Delta(\mathbf{k})$ by

$$(D_g \Delta)(\mathbf{k}) := u_g \Delta(p_g^{-1} \mathbf{k}) u_g^T. \quad (5)$$

It satisfies

$$(D_g(D_h \Delta))(\mathbf{k}) = u_g(D_h \Delta)(p_g^{-1} \mathbf{k}) u_g^T \quad (6)$$

$$= u_g u_h \Delta(p_h^{-1} p_g^{-1} \mathbf{k}) u_h^T u_g^T \quad (7)$$

$$= (z_{g,h})^2 u_{gh} \Delta(p_{gh}^{-1} \mathbf{k}) u_{gh}^T \quad (8)$$

$$= (z_{g,h})^2 (D_{gh} \Delta)(\mathbf{k}). \quad (9)$$

Therefore, if the factor system obeys $(z_{g,h})^2 \equiv 1$, then $D_g D_h = D_{gh}$.

The irreducible decomposition of a given superconducting gap function under G is expected to be obtained from the projection formula onto an irreducible representation α of G ,

$$P_\alpha = \frac{\dim \alpha}{|G|} \sum_{g \in G} (\chi_g^\alpha)^* D_g. \quad (10)$$

Namely,

$$\Delta(\mathbf{k}) = \sum_{\alpha \in \text{irreps}} \Delta^{(\alpha)}(\mathbf{k}), \quad (11)$$

$$\Delta^{(\alpha)}(\mathbf{k}) = \frac{\dim \alpha}{|G|} \sum_{g \in G} (\chi_g^\alpha)^* (D_g \Delta)(\mathbf{k}). \quad (12)$$

Here χ_g^α denotes the irreducible character.