

# A Note on Redefining the Second Stiefel–Whitney Class

Ken Shiozaki

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The Whitney sum formula

$$w(E + F) = w(E)w(F) \quad (1)$$

implies

$$w_1(E + F) = w_1(E) + w_1(F), \quad (2)$$

$$w_2(E + F) = w_2(E) + w_2(F) + w_1(E)w_1(F). \quad (3)$$

We would like to ask whether  $w_2$  can be redefined so that

$$w'_2(E + F) = w'_2(E) + w'_2(F) \quad (4)$$

holds. This is similar to a quadratic refinement.

In general, for a symmetric bilinear form  $b(x, y)$  with values in  $\mathbb{Z}_2$ , there exists a  $\mathbb{Z}_2$ -valued function  $q(x)$  satisfying

$$q(x + y) = q(x) + q(y) + b(x, y). \quad (5)$$

Such a  $q$  is called a quadratic refinement. It is not unique. In the present situation,

$$b(E, F) := w_1(E)w_1(F) \quad (6)$$

is a symmetric bilinear form. Thus there exists a quadratic refinement  $q(E)$ . Using it, define

$$w'_2(E) := w_2(E) + q(E). \quad (7)$$

Then

$$w'_2(E + F) = w_2(E + F) + q(E + F) \quad (8)$$

$$= w_2(E) + w_2(F) + w_1(E)w_1(F) + q(E) + q(F) + w_1(E)w_1(F) \quad (9)$$

$$= w'_2(E) + w'_2(F). \quad (10)$$

How should one choose  $q(E)$ ? As a general principle, after restricting the base space on which the principal bundle is defined, a quadratic refinement should be defined in terms of topological invariants on lower-dimensional base spaces [1]. Let us explicitly construct it when the base space is the two-dimensional torus  $T^2$ .

Write

$$w_1(E) = \nu_{1x}(E)dx + \nu_{1y}(E)dy, \quad (11)$$

and

$$w_2(E) = \nu_2(E)dxdy. \quad (12)$$

Then the sum rule is

$$\nu_2(E + F) = \nu_2(E) + \nu_2(F) + \nu_{1x}(E)\nu_{1y}(F) + \nu_{1y}(E)\nu_{1x}(F). \quad (13)$$

For the symmetric bilinear form

$$b(E, F) = \nu_{1x}(E)\nu_{1y}(F) + \nu_{1y}(E)\nu_{1x}(F), \quad (14)$$

one can take the quadratic refinement

$$q(E) = \nu_{1x}(E)\nu_{1y}(E). \quad (15)$$

Indeed,

$$q(E + F) = (\nu_{1x}(E) + \nu_{1x}(F))(\nu_{1y}(E) + \nu_{1y}(F)) \quad (16)$$

$$= \nu_{1x}(E)\nu_{1y}(E) + \nu_{1x}(F)\nu_{1y}(F) + \nu_{1x}(E)\nu_{1y}(F) + \nu_{1x}(F)\nu_{1y}(E) \quad (17)$$

$$= q(E) + q(F) + b(E, F). \quad (18)$$

Therefore, redefining

$$\tilde{\nu}_2(E) := \nu_2(E) - \nu_{1x}(E)\nu_{1y}(E), \quad (19)$$

we obtain the additive rule

$$\tilde{\nu}_2(E + F) = \tilde{\nu}_2(E) + \tilde{\nu}_2(F). \quad (20)$$

Such a redefinition is also discussed in [2].

To see the viewpoint based on topological invariants of subspaces, it is useful to write the construction more formally. Let  $S_x^1$  and  $S_y^1$  be the subcycles of  $T^2$  in the  $k_x$  and  $k_y$  directions, respectively. The restriction maps to these circles define homomorphisms

$$r_\mu : H^*(S_\mu^1; \mathbb{Z}/2) \rightarrow H^*(T^2; \mathbb{Z}/2). \quad (21)$$

The restriction of  $E$  to  $S_\mu^1$  determines an element

$$r_\mu w_1(E|_{S_\mu^1}) \in H^1(T^2; \mathbb{Z}/2). \quad (22)$$

The quadratic refinement is then defined by

$$q(E) = r_x w_1(E|_{S_x^1}) \cup r_y w_1(E|_{S_y^1}). \quad (23)$$

## References

- [1] Seishiro Ono and Ken Shiozaki, *Towards complete characterization of topological insulators and superconductors: A systematic construction of topological invariants based on Atiyah-Hirzebruch spectral sequence*, arXiv:2311.15814.
- [2] Yanzhu Chen, Sheng-Jie Huang, Yi-Ting Hsu, and Tzu-Chieh Wei, *Topological invariants beyond symmetry indicators: Boundary diagnostics for twofold rotationally symmetric superconductors*, arXiv:2109.06959.