

~~Deformation and~~

# Pairing correlation in neutron-rich halo nuclei

Kouichi Hagino  
Kyoto University, Kyoto, Japan

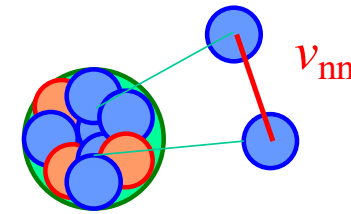


1. Introduction: Borromean nuclei and di-neutron correlations
2. Pairing in two-neutron halo: Coulomb breakup of  $^{22}\text{C}$
3. Role of di-neutron correlation in E1 excitations
4. Summary

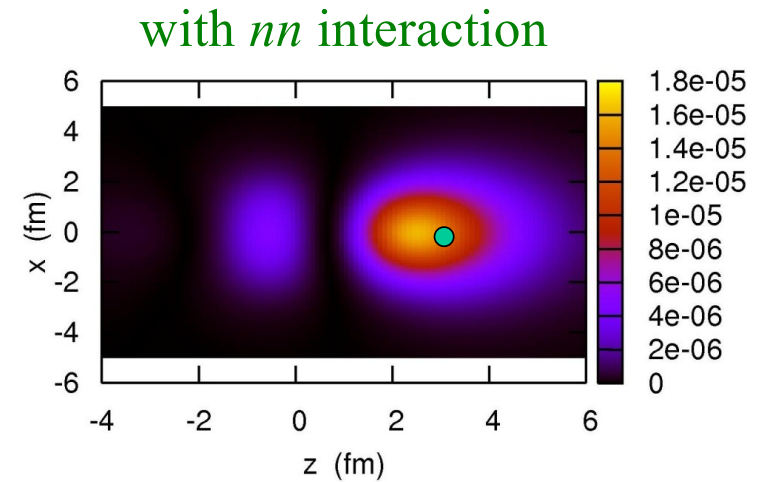
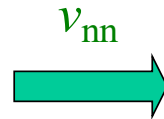
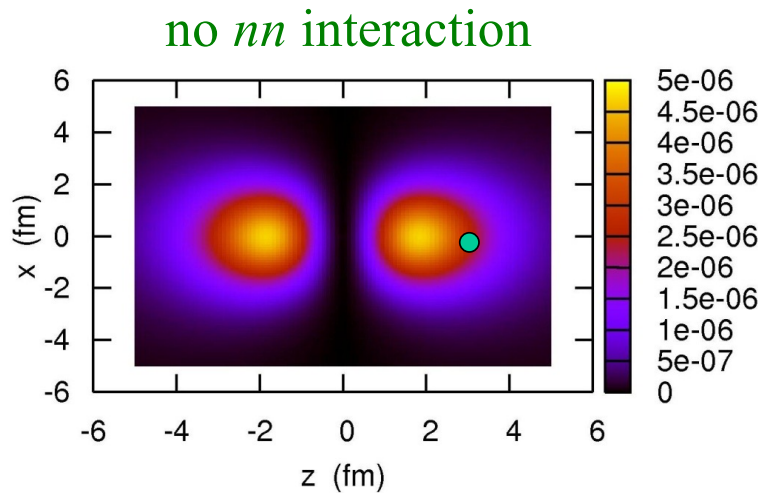
# Borromean nuclei and di-neutron correlation

## Di-neutron correlation

spatially compact configurations of two neutrons



3-body model calculation for  $^{11}\text{Li}$

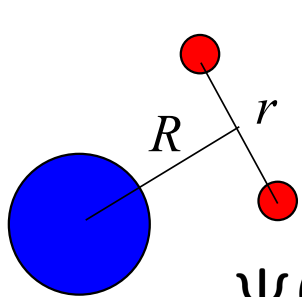
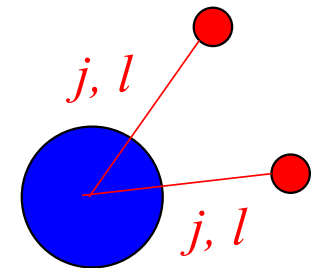
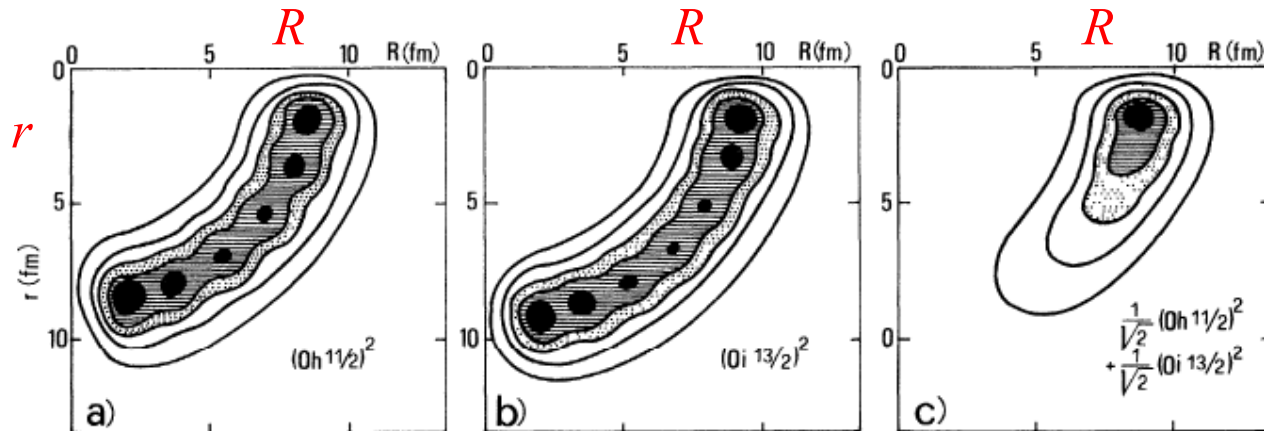


K. Hagino and H. Sagawa, PRC72 (2005) 044321

## Borromean nuclei and di-neutron correlation

the essential ingredient:  
admixture of configurations with  
opposite parities

$$\Psi(\mathbf{r}_1, \mathbf{r}_2) = \sum_{j,l} C_{jl} [\phi_{jl}(\mathbf{r}_1)\phi_{jl}(\mathbf{r}_2)]^{(00)}$$



$$\Psi(\mathbf{r}_1, \mathbf{r}_2) \rightarrow \Psi(r, R)$$

$$(0h_{11/2})^2$$

$$l = 5$$

$$(0i_{13/2})^2$$

$$l = 6$$

$$\frac{1}{\sqrt{2}}(0h_{11/2})^2$$

$$+ \frac{1}{\sqrt{2}}(0i_{13/2})^2$$

F. Catara, A. Insolia, E. Maglione,  
and A. Vitturi, PRC29 (1984) 1091

## Analysis of the Coulomb breakup of $^{22}\text{C}$

✓ Y. Togano et al., PLB761, 412 (2016)

$$\sigma_{\text{R}} \rightarrow \text{rms radius} = 3.44(8) \text{ fm}$$

$$S_{2\text{n}} = 0.56^{+0.27}_{-0.20} \text{ MeV}$$

✓ N. Kobayashi et al., PRC86, 054606 (2012).

momentum distribution  $\rightarrow$  a large s-d mixture

$$S_{2\text{n}} = 0.40\text{-}1.20 \text{ MeV}$$

✓ N. Nakatsuka, T. Nakamura et al. in preparation

exclusive Coulomb breakup cross sections  $\leftarrow$

### 3-body model calculation

G.F. Bertsch and H. Esbensen, Ann. of Phys. 209 (1991) 327.

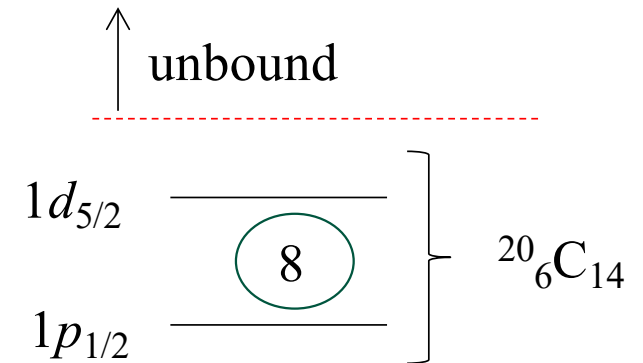
K. Hagino and H. Sagawa, PRC72 (2005) 044321.

$V_{\text{nn}}$  : a density-dependent contact interaction

### 3-body model calculation for $^{22}\text{C}$

$$^{22}_6\text{C}_{16} = ^{20}_6\text{C}_{14} + \text{n} + \text{n}$$

*spherical core*

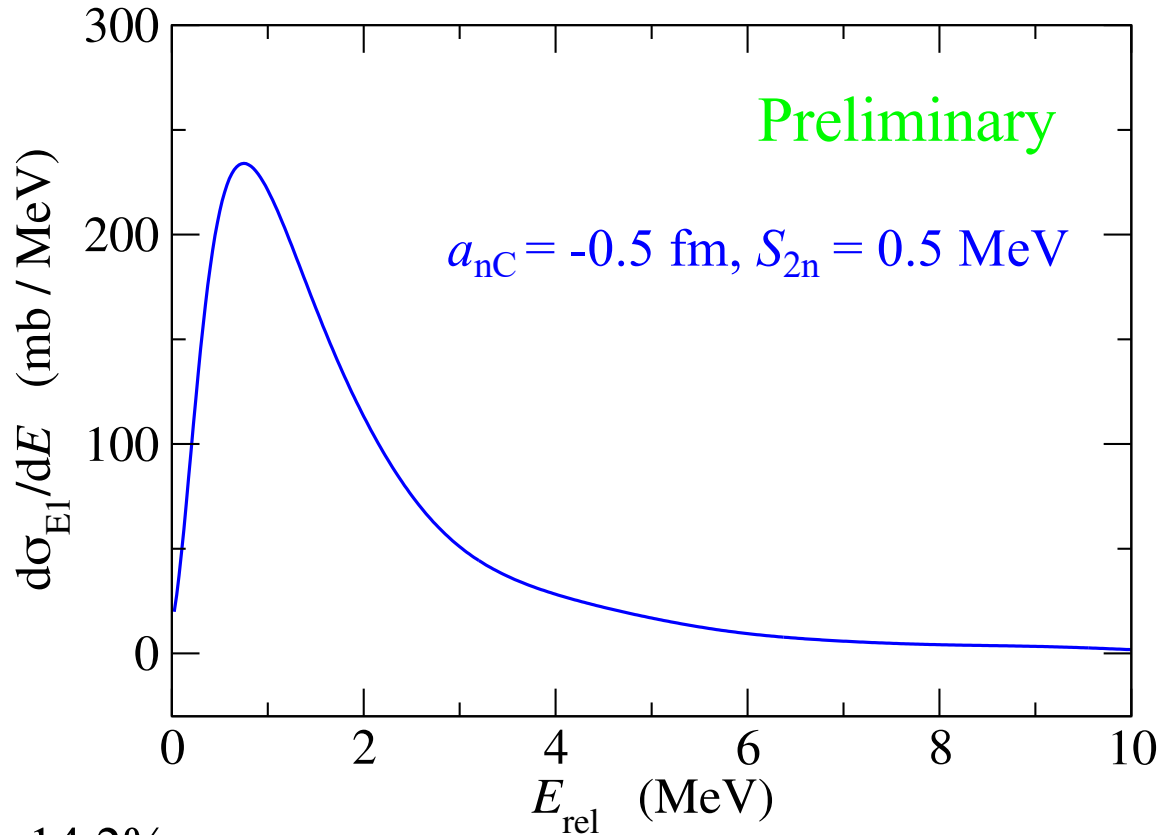


$V_{\text{nC}}$  : Woods-Saxon

change  $V_0$  for  $l = 0$  to  
adjust  $a_{\text{nC}}$

$\rightarrow$  adjustable parameters:  
 $a_{\text{nC}}$  and  $S_{2\text{n}}$

# Analysis of the Coulomb breakup of $^{22}\text{C}$



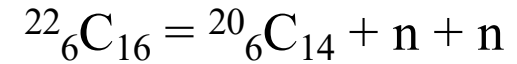
$P(s^2) = 14.2\%$

$P(d_{3/2}^2) = 63.9\%$

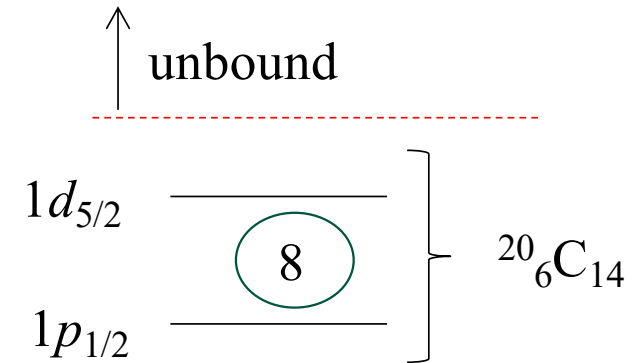
N. Nakatsuka, T. Nakamura et al. in preparation

rms radius = 3.20 fm [Exp.: 3.44 (8) fm]

## 3-body model calculation for $^{22}\text{C}$



*spherical core*



$V_{nC}$  : Woods-Saxon

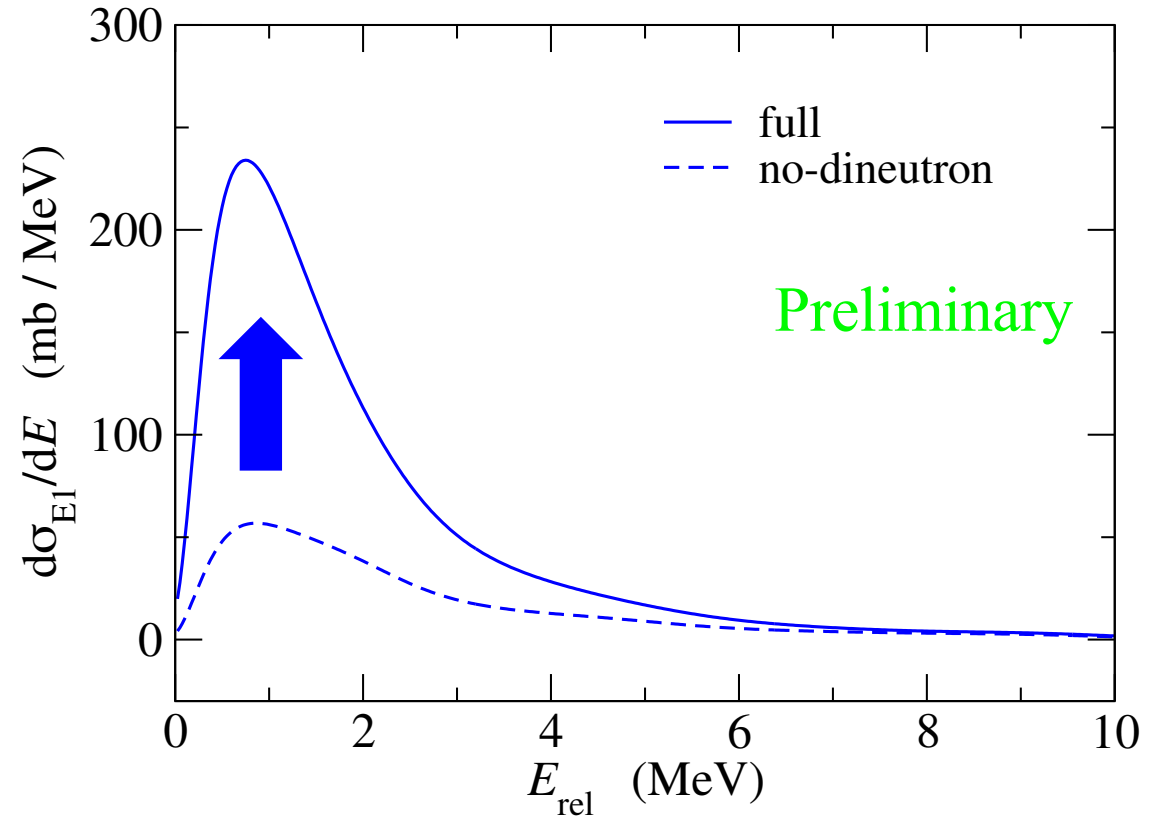
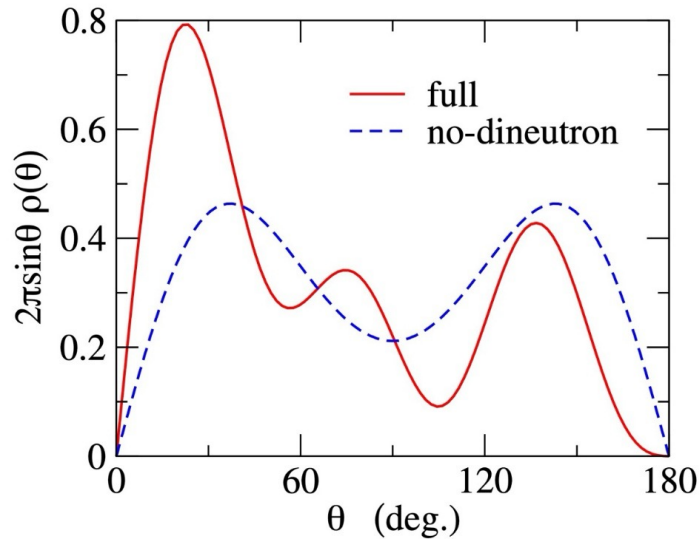
change  $V_0$  for  $l = 0$  to  
adjust  $a_{nC}$

→ adjustable parameters:  
 $a_{nC}$  and  $S_{2n}$

## Role of di-neutron correlation

the ground state  $0^+$

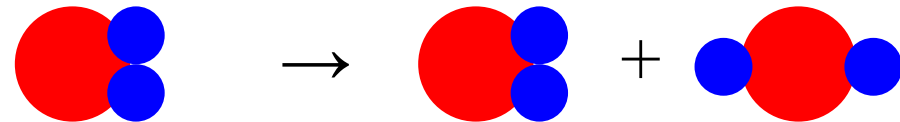
full correlations vs no di-neutron  
(even  $l$  only)



no dineutron in the ground state  $\rightarrow$  a large hindrance in the  $E1$  strengths

$\leftarrow$  due to a small  $R_{c-2n}$  ( $4.25 \rightarrow 3.01$  fm)

$$B_{\text{tot}}(E1) \propto \langle R_{c-2n}^2 \rangle$$



## Role of di-neutron correlation

di-neutron correlation:

$$|\Psi\rangle = \alpha_0|s^2\rangle + \alpha_1|p^2\rangle + \dots$$

E1

$|sp\rangle$

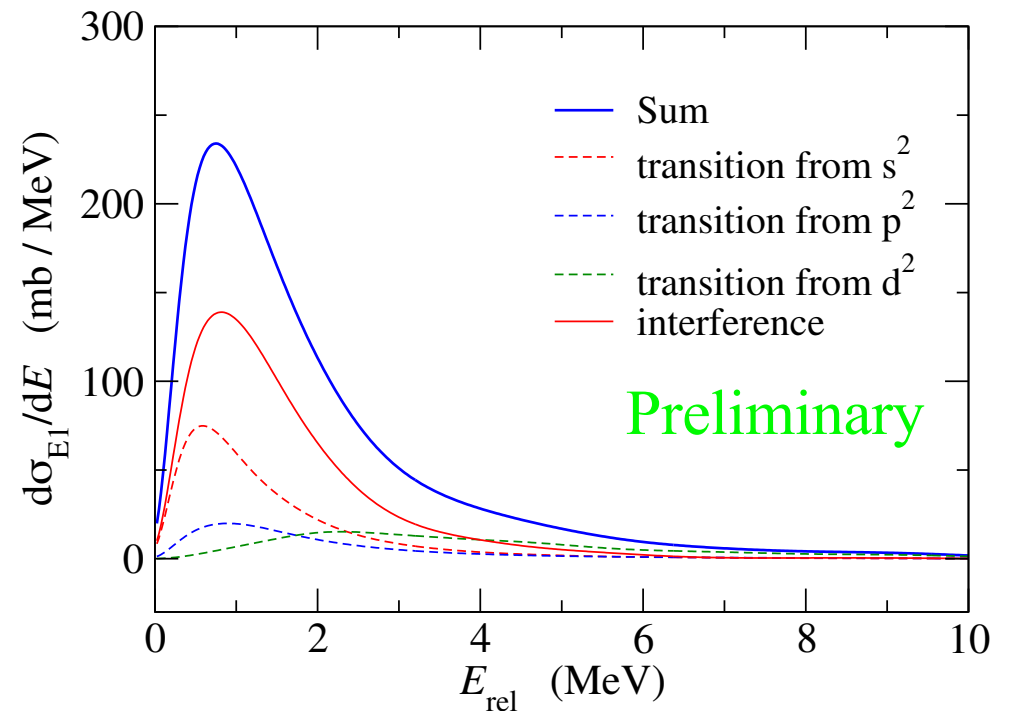
$$\begin{aligned} \rightarrow |\langle sp|T_{E1}|\Psi\rangle|^2 &= |\alpha_0|^2 |\langle sp|T_{E1}|s^2\rangle|^2 \\ &+ |\alpha_1|^2 |\langle sp|T_{E1}|p^2\rangle|^2 \\ &+ \alpha_0\alpha_1^* \langle sp|T_{E1}|s^2\rangle \langle p^2|T_{E1}|sp\rangle + c.c. + \dots \end{aligned}$$

← interferences

dineutron correlation: mixture of different parity states

→ the same final state for the E1 transitions

→ large interferences

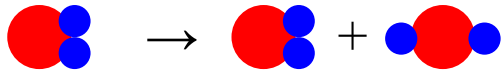


## Summary

### Coulomb breakup of $^{22}\text{C}$

- 3-body model calculation
- Enhancement of E1 strengths due to dineutron correlation

✓ cluster sum rule  $B_{\text{tot}}(E1) \propto \langle R_{\text{C}-2n}^2 \rangle$



✓ interference

### Future perspective

2n-halo nuclei with a deformed core  $\rightarrow$  a rotor + 2 neutrons?

maybe important in  $^{22}\text{C}$  (a deformed  $^{20}\text{C}$ )

