非可換ソリトンと無限個の保存量

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Ref: MH, ``Commuting Flows and Conservation Laws for NC Lax Hierarchies,"[hep-th/0311206]

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1. Introduction

Successful points on NC theories

- Appearance of new physical objects
- Description of real physics
- Various successful applications to D-brane dynamics etc.

NC Solitons play important roles (Integrable!)

Final goal: NC extension of all soliton theroies

Ward's observation: Almost all integrable equations are reductions of the ASDYM eqs.

R.Ward, Phil.Trans.Roy.Soc.Lond.A315(85)451

Anti-Self-Dual Yang-Mills eq.

Reductions

KP eq. BCS eq.
KdV eq. Boussinesq eq.
NLS eq. mKdV eq.
sine-Gordon eq. Burgers eq. ...
(Almost all!)

e.g. [Mason&Woodhouse]

NC Ward's observation: Almost all NC integrable equations are reductions of the NC ASDYM eqs.

MH&K.Toda, PLA316(03)77[hep-th/0211148]

NC ASDYM eq.

Successful

Reductions

Reductions

NC KP eq. NC BCS eq.

NC KdV eq. NC Boussinesq eq.

NC NLS eq. NC mKdV eq.

NC sine-Gordon eq. NC Burgers eq.

(Almost all !?)

Successful?

A general framework is needed

2. NC Sato's Theory

- Sato's Theory: one of the most beautiful theory of solitons
 - Based on the exsitence of hierarchies and tau-functions
- Sato's theory reveals essential aspects of solitons:
 - Construction of exact solutions
 - Structure of solution spaces
 - Infinite conserved quantities
 - Hidden infinite-dim. symmetry
 - Let's discuss NC extension of Sato's theory

Derivation of soliton equations

 Give a Lax operator which is a pseudodifferential operator

$$L := \partial_x + u_2 \partial_x^{-1} + u_3 \partial_x^{-2} + u_4 \partial_x^{-3} + \cdots$$

Define a differential operator

$$B_m := (L * \cdots * L)_{\geq 0}$$

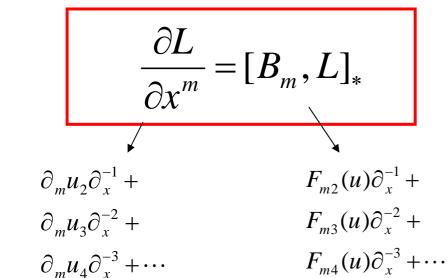
m times

$$u_k = u_k(x^1, x^2, x^3, \cdots)$$

Noncommutativity is introduced here:

$$[x^i, x^j] = i\theta^{ij}$$

• Introduce NC (KP) hierarchy equation:



Star product:

$$f(x) * g(x) := f(x) \exp\left(\frac{i}{2}\theta^{ij}\overleftarrow{\partial}_{i}\overrightarrow{\partial}_{j}\right)g(x)$$

which realizes the noncommutativity

Each coefficient yields a differential equation.

Closer look at NC (KP) hierarchy

For m=2

$$\begin{array}{ll} \partial_{x}^{-1} & \partial_{2}u_{2} = \underline{2u_{3}'} + u_{2}'' \\ \partial_{x}^{-2} & \partial_{2}u_{3} = \underline{2u_{4}'} + u_{3}'' + 2u_{2} * u_{2}' + 2[u_{2}, u_{3}]_{*} \\ \partial_{x}^{-3} & \partial_{2}u_{4} = \underline{2u_{5}'} + u_{4}'' + 4u_{3} * u_{2}' - 2u_{2} * u_{2}'' + 2[u_{2}, u_{4}]_{*} \\ \vdots & \vdots \end{array}$$

Infinite kind of fields are represented in terms of one kind of field $u_2 \equiv u$ MH&K.Toda, [hep-th/0309265]

$$u_{x} := \frac{\partial u}{\partial x}$$
$$\partial_{x}^{-1} := \int_{x}^{x} dx'$$

etc.

For m=3

$$\partial_{x}^{-1}$$
) $\partial_{3}u_{2} = u_{2}''' + 3u_{3}'' + 3u_{4}'' + 3u_{2}' * u_{2} + 3u_{2} * u_{2}'$
 \vdots

$$u_{t} = \frac{1}{4}u_{xxx} + \frac{3}{4}(u_{x} * u + u * u_{x}) + \frac{3}{4}\partial_{x}^{-1}u_{yy} + \frac{3}{4}[u, \partial_{x}^{-1}u_{yy}]_{*}$$
 (2+1)-dim. NC KP equation

and other NC equations $u = u(x^1, x^2, x^3, \dots)$ (NC hierarchy equations)

$$u = u(x^{1}, x^{2}, x^{3}, \cdots)$$

$$\uparrow \qquad \uparrow \qquad \uparrow$$

$$x \qquad y \qquad t$$

(KP hierarchy) → (various hierarchies.)

(Ex.) KdV hierarchy

Reduction condition

$$L^2 = B_2 (=: \partial_x^2 + u)$$
 : 2-reduction

gives rise to NC KdV hierarchy

which includes (1+1)-dim. NC KdV eq.:

$$u_{t} = \frac{1}{4}u_{xxx} + \frac{3}{4}(u_{x} * u + u * u_{x})$$

Note $\frac{\partial u}{\partial x_{2N}} = 0$: dimensional reduction in x_{2N} directions

KP:
$$u(x^1, x^2, x^3, x^4, x^5, ...)$$

 $x y t$: (2+1)-dim.
KdV: $u(x^1, x^3, x^5, ...)$
 $x t$: (1+1)-dim.

/-reduction of NC KP hierarchy yields wide class of other NC hierarchies

- No-reduction \rightarrow NC KP $(x, y, t) = (x^1, x^2, x^3)$
- 2-reduction \rightarrow NC KdV $(x,t) = (x^1, x^3)$
- 3-reduction \rightarrow NC Boussinesq $(x,t) = (x^1, x^2)$
- 4-reduction → NC Coupled KdV
- 5-reduction → ...
- 3-reduction of BKP → NC Sawada-Kotera
- 2-reduction of mKP → NC mKdV
- Special 1-reduction of mKP → NC Burgers
- Noncommutativity should be introduced into space-time coords

4. Conservation Laws

• Conservation laws: $\partial_t \sigma = \partial_i J^i = \sigma$: Conserved density time space

Then $Q := \int_{space}^{dx} \sigma$ is a conserved quantity.

$$\therefore \partial_t Q = \int_{space} dx \partial_t \sigma = \int_{spatial \atop inf \ inity} dS_i J^i = 0$$

Conservation laws for the hierarchies

$$\lim_{n \to \infty} \frac{\partial_{m} res_{-1} L^{n} = \partial_{x} J + \theta^{ij} \partial_{j} \Xi_{i}}{\operatorname{space}}$$

I have found the explicit form!

 $res_{-r}L^n$: coefficient of ∂_x^{-r} in L^n

Noncommutativity should be introduced in space-time directions only. →

 $t \equiv x^m$

∂_j should be space or time derivative
 → ordinary conservation laws!

Hot Results

Infinite conserved densities for NC hierarchy eqs. (n=1,2,...)

$$\sigma = res_{-1}L^{n} + \theta^{mi} \sum_{k=1}^{m} \sum_{l=n+1}^{n+1+m-k} (-1)^{m+n-k-l} \binom{m-k}{l-n-1} a_{n-l} \Diamond \partial_{i} b_{m-k}^{(m+n-k-l+1)}$$

$$t \equiv x^m$$

$$L^n = \partial_x^n + \sum_{l=1}^\infty a_{n-l} \partial_x^{n-l}, \quad B_m = \partial_x^m + \sum_{k=1}^m b_{m-k} \partial_x^{m-k}$$

♦ : Strachan's product

$$f(x) \diamond g(x) := f(x) \left(\sum_{s=0}^{\infty} \frac{(-1)^s}{(2s+1)!} \left(\frac{1}{2} \theta^{ij} \overleftarrow{\partial}_i \overrightarrow{\partial}_j \right)^{2s} \right) g(x)$$

MH, [hep-th/0311206]

We can calculate the explicit forms of conserved densities for wide class of NC soliton equations.

- Space-Space noncommutativity: NC deformation is slight: $\sigma = res_{-1}L^n$
- Space-time noncommutativity
 - NC deformation is drastical:
 - Example: NC KP and KdV equations $([t,x]=i\theta)$

$$\sigma = res_{-1}L^n - 3\theta((res_{-1}L^n) \diamond u_3' + (res_{-2}L^n) \diamond u_2')$$

meaningful?

4. Conclusion and Discussion

- We proved the existence of infinite conserved quantities for wide class of NC Lax hierarchies.
- We gave the infinite conserved densities explicitly from the viewpoint of Sato's theory, which suggests that infinite-dim. symmetry would be hidden in the NC (soliton) equations.
- The results show that NC soliton eqs. have very special properties though they include non-linear terms and infinite (time-!) derivatives.
- Of course, there are still many things to be seen in order to confirm the integrability.

Further directions

- Completion of NC Sato's theory
 - Hirota's bilinearization and tau-functions $u = \partial_x^2 \log \tau$
 - → hidden symmetry (deformed affine Lie algebras?)
 - Cf. Dimakis&Mueller-Hoissen, Wang&Wadati...
 - Geometrical descriptions from NC extension of the theories of Krichever, Mulase and Segal-Wilson and so on.
- Confirmation of NC Ward's conjecture
 - NC twistor theory
 - e.g. Kapustin&Kuznetsov&Orlov, Hannabuss, Hannover group,...
 - D-brane interpretations → physical meanings
- Foundation of Hamiltonian formalism with space-time noncommutativity
 - Initial value problems, Liouville's theorem, Noether's thm,...