#### Unified description of nuclear structure and reactions employing the dispersive optical model (DOM)

YKIS2022b workshop 5/23/2022



- •Motivation: essential for studying unstable nuclei and link to neutron stars
- Optical potentials 2022
- Removal probabilities puzzle(s)
- •Green's functions/propagator method
  - vehicle for ab initio calculations --> matter & finite nuclei
  - as a framework to link data at positive and negative energy (and to generate predictions for unstable nuclei as well as neutron skins)
- -> dispersive optical model (DOM <- started by Claude Mahaux)
- DOM with non-local potentials <sup>12</sup>C, <sup>16-18</sup>O, <sup>40,48</sup>Ca, <sup>58,64</sup>Ni, <sup>112,124</sup>Sn, <sup>208</sup>Pb
- Revisit (e,e'p) data from NIKHEF <sup>40</sup>Ca and <sup>48</sup>Ca --> N-Z dependence
- Neutron skin in <sup>48</sup>Ca and <sup>208</sup>Pb -> PREX II
- Conclusion and outlook
- Ground-state energy and EOS —> saturation properties

# Optical potentials 2022

#### Reviews

IOP Publishing

Journal of Physics G: Nuclear and Particle Physics J. Phys. G: Nucl. Part. Phys. 44 (2017) 033001 (57pp) https://doi.org/10.1088/1361-6471/44/3/033001

**Topical Review** 

#### Novel applications of the dispersive optical model

#### W H Dickhoff<sup>1,4</sup>, R J Charity<sup>2</sup> and M H Mahzoon<sup>1,3</sup>

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Progress in Particle and Nuclear Physics 105 (2019) 252-299



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Recent developments for the optical model of nuclei

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#### Progress in Particle and Nuclear Physics 118 (2021) 103847



#### Review

Quenching of single-particle strength from direct reactions with stable and rare-isotope beams

T. Aumann<sup>a,b</sup>, C. Barbieri<sup>c,d,e</sup>, D. Bazin<sup>f,g</sup>, C.A. Bertulani<sup>h</sup>, A. Bonaccorso<sup>i</sup>, W.H. Dickhoff<sup>j</sup>, A. Gade<sup>f,g</sup>, M. Gómez-Ramos<sup>a,k</sup>, B.P. Kay<sup>1</sup>, A.M. Moro<sup>k,m</sup>, T. Nakamura<sup>n</sup>, A. Obertelli<sup>a,\*</sup>, K. Ogata<sup>o,p</sup>, S. Paschalis<sup>q</sup>, T. Uesaka<sup>r</sup>



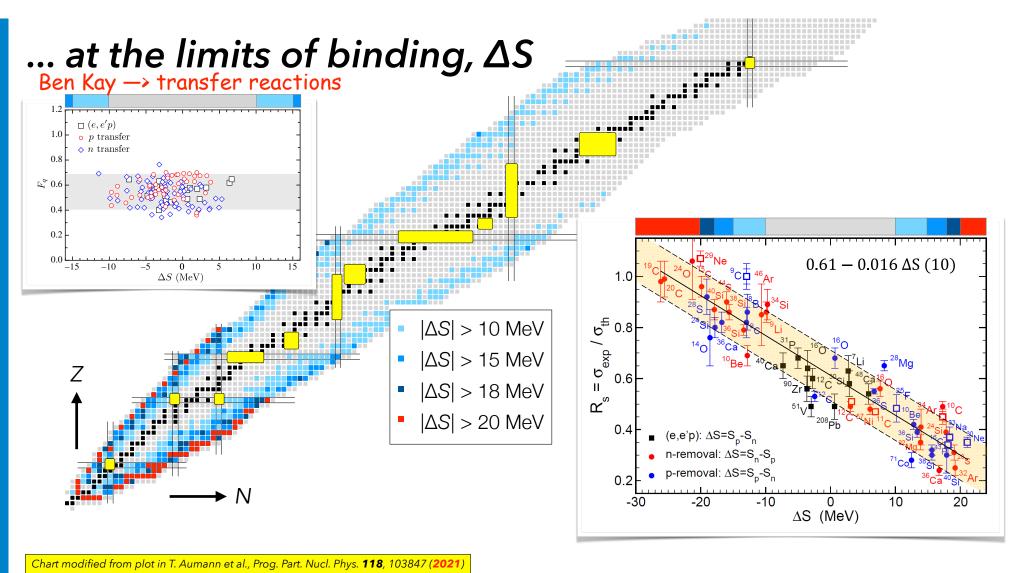
# Unstable Nuclei and Optical Potentials

- Masses and electromagnetic properties -> unambiguous!
- All other properties depend on description of strongly interacting probe with target —> ambiguous! (or in other words model dependent)
- Optical potentials are an important ingredient in analyzing e.g. transfer reactions, knockout reactions like (e,e'p) and (p,2p), the latter available in inverse kinematics!

#### • Status 2022

- Global potentials still local, non-dispersive and not constrained by scattering data for exotic nuclei
- Ab initio methods have as yet very limited relevance e.g. for FRIB physics
- DOM potentials available for a reasonable set of nuclei and can be extrapolated
- DOM potentials are nonlocal, dispersive, and describe structure and reaction domains simultaneously

**Optical Potential** 

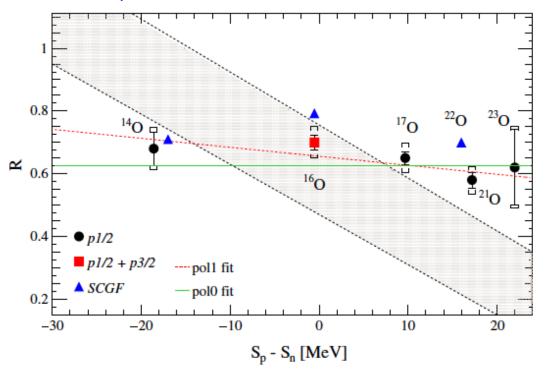


J. A. Tostevin and A. Gade, Phys. Rev. C **103**, 054610 (**2021**) [an update to Phys. Rev. C **90**, 057602 (**2014**)

40

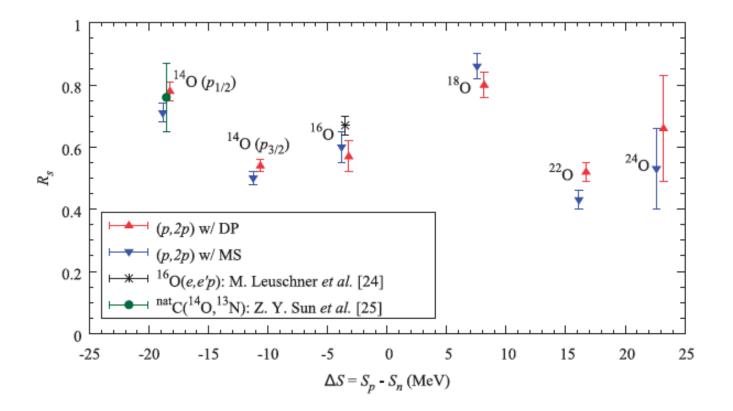
O(p,2p) L. Atar et al. Phys. Rev. Lett. 120, 052501 (2018)

- "Ab initio" interaction has "no" tensor force —> spectroscopic factors?
- Reaction model: distorted waves not constrained by experimental information as a function of nucleon asymmetry
- Inconsistent with np dominance observed in 2N knockout reactions (Or et al.)



# O(p,2p)

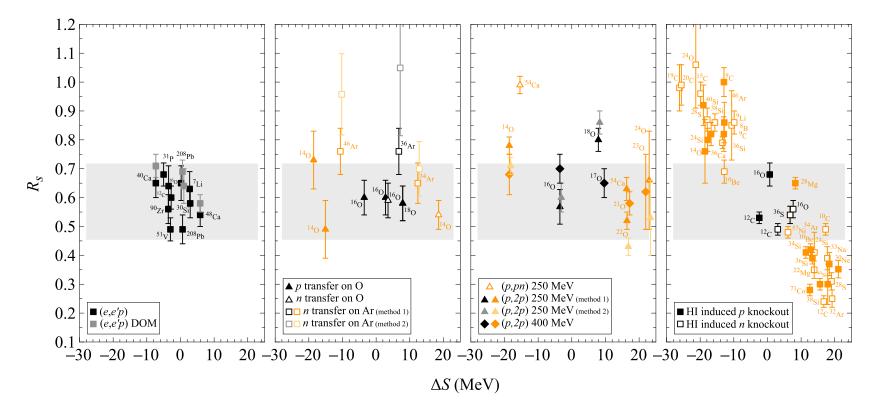
- S. Kawase et al. Prog. Theor. Exp. Phys. 2018, 021D01
- DWIA uses optical potentials not constrained by scattering data for unstable nuclei



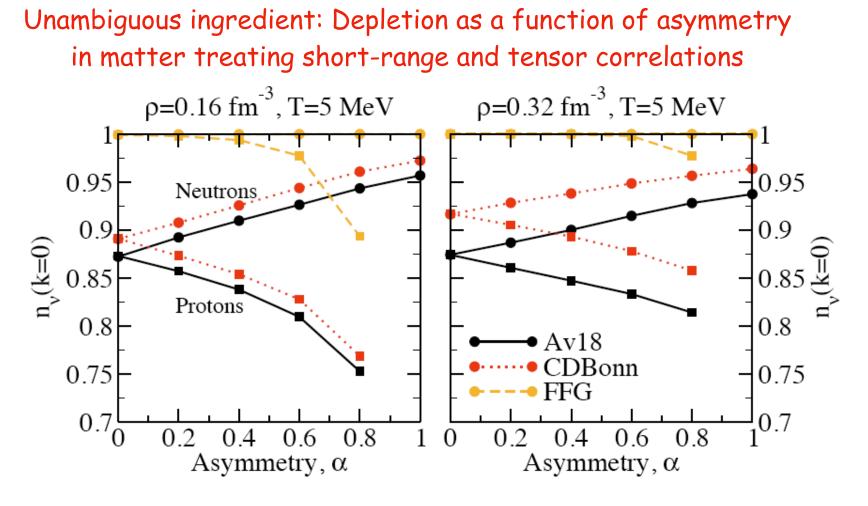
#### Status of "reduction" factors/spectroscopic factors

T. Aumann, C. Barbieri, D. Bazin et al.

Progress in Particle and Nuclear Physics 118 (2021) 103847



**Fig. 56.** The four panels of this plot show the quenching (reduction) factors for (a) electron-induced knockout reactions [87,172,237,376], (b) transfer reactions with radioactive ion beams [55,57,203], (c) quasifree (p, 2p) proton knockout on stable nuclei (from the compilation in [239]) and radioactive nuclei [58,59], and (d) the inclusive intermediate-energy knockout data [46]. The measurements are compared to predictions based on effective-interaction shell-model SFs while, in the case of (e, e'p), the integrated strength is compared to the independent-particle expectation.



A. Rios, A. Polls, and W. H. Dickhoff Depletion of the nuclear Fermi sea. <u>Phys. Rev. C79, 064308 (2009)</u>.

### Properties and relevance of optical potential or self-energy

- Describes all observables related to elastic scattering
- Generates wave functions inside the nucleus distorted waves used to describe other nuclear reactions aimed at extracting nuclear structure information
- Dispersive optical model describes additional observables related to the ground state
- · Can be extrapolated to exotic nuclei [but uncertain]
- Astrophysics: relevant for the description of rapid neutron capture
- · Lots of recent activity to generate the optical potential starting from the "NN" interaction
  - Nuclear matter approach [old and recent]
  - Multiple scattering using free NN interaction [T-matrix]
  - Green's function method for finite nuclei [old and recent]
  - Coupled cluster method
  - No core shell model [including symmetry-adapted version]
  - DFT + 2nd order/RPA contributions

### Some problems with "ab initio" approaches

- Multiple scattering approaches useful at higher energy but hard to improve over spectator approximation (problems with polarization data)
- Methods that generate the propagator directly and then generate optical potential by inverting the Dyson equation
  - Coupled cluster
  - No core shell model
  - Both methods struggle with the  $~\eta~$  infinitesimal

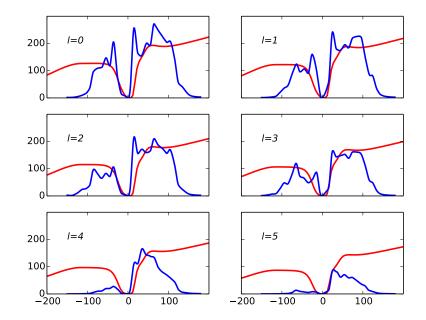
$$G_{\ell j}(r,r';E) = \sum_{m} \frac{\langle \Psi_{0}^{A} | a_{r\ell j} | \Psi_{m}^{A+1} \rangle \langle \Psi_{m}^{A+1} | a_{r'\ell j}^{\dagger} | \Psi_{0}^{A} \rangle}{E - (E_{m}^{A+1} - E_{0}^{A}) + i\eta} + \sum_{n} \frac{\langle \Psi_{0}^{A} | a_{r'\ell j}^{\dagger} | \Psi_{n}^{A-1} \rangle \langle \Psi_{n}^{A-1} | a_{r\ell j} | \Psi_{0}^{A} \rangle}{E - (E_{0}^{A} - E_{n}^{A-1}) - i\eta}$$
$$G_{\ell j}(r,r';E) = G_{\ell j}^{(0)}(r,r';E) + \int d\tilde{r} \; \tilde{r}^{2} \int d\tilde{r}' \; \tilde{r}'^{2} G_{\ell j}^{(0)}(r,\tilde{r};E) \Sigma_{\ell j}(\tilde{r},\tilde{r}';E) G_{\ell j}(\tilde{r}',r';E)$$

- Methods that use HO basis cannot describe density of states and higher angular momentum

**Optical Potential** 

Illustrate: comparison with ab initio FRPA calculation

 Volume integrals of imaginary part of nonlocal ab initio (FRPA) self-energy compared with DOM result for <sup>40</sup>Ca



• Ab initio S. J. Waldecker, C. Barbieri and W. H. Dickhoff Microscopic self-energy calculations and dispersive-optical-model potentials. Phys. Rev. C84, 034616 (2011), 1-11.

#### DOM initiated by Mahaux

St. Louis extensions (nonlocality, energy domain, isotope chains)

Nonlocal and dispersive optical potential

$$\operatorname{Re} \Sigma_{\ell j}(r, r'; E) = \operatorname{Re} \Sigma_{\ell j}(r, r'; \varepsilon_{F})$$

$$- \mathcal{P} \int_{\varepsilon_{T}^{+}}^{\infty} \frac{dE'}{\pi} \operatorname{Im} \Sigma_{\ell j}(r, r'; E') \left[ \frac{1}{E - E'} - \frac{1}{\varepsilon_{F} - E'} \right]$$

$$+ \mathcal{P} \int_{-\infty}^{\varepsilon_{T}^{-}} \frac{dE'}{\pi} \operatorname{Im} \Sigma_{\ell j}(r, r'; E') \left[ \frac{1}{E - E'} - \frac{1}{\varepsilon_{F} - E'} \right]$$

- Allows consideration of negative energy experimental information [charge density]
- Subtracted dispersion relation emphasizes influence of energies close to the Fermi energy
- Empirical information constrains binding potential at the Fermi energy as well as volume integrals of the imaginary part at positive energy

**Optical Potential** 

Dispersive Optical Model (St. Louis group)

2000 1800

1600

1400 [mp 1200

1000

400

200

10000

9000 8000

3000

2000

100 E<sub>lab</sub> [MeV]

100 E<sub>lab</sub> [MeV]

n+208PF

E<0 ->

ь 800 600

- Mahaux & Sartor 1991 -> Washington University group since 2006
- Use experimental data to constrain the nucleon self-energy while linking structure and reaction domain using dispersion relations

Indirectly:

 $\begin{array}{c} \mathrm{d}\sigma/\mathrm{d}\Omega \, \left[\mathrm{mb/sr}\right] \\ \mathrm{d}\sigma & \mathrm{d}\Omega \\ \mathrm{d} & \mathrm{d} \\ \mathrm{d} & \mathrm{d} \end{array}$ 

 $p+^{208}Ph$ 

 $E_{lab} > 100$ 

a bardana

 $\theta_{cm}$  [deg]

120 150 180

100

a ha a ha a ha a ha a ha

 $\theta_{cm}$  [deg]

30 60 90 120 150 180

Generates proton/neutron distorted waves

30

 $k_{-1} > 100$ 

 $40>E_{lab}>20$ 

60 90 120 150 180

 $\theta_{c.m.}$  [deg]

20

30

60

 $\theta_{c.m.}$  [deg]

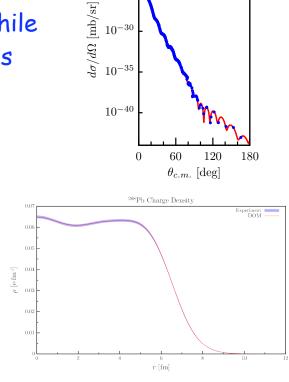
90 120 150 180

Overlap functions with their normalization (spectroscopic factors)

Mack Atkinson -> TRIUMF -> LLNL



DOM



 $10^{-20}$ 

 $10^{-25}$ 

DOM

Experiment •



Eur. Phys. J. A (2017) **53**: 178 DOI 10.1140/epja/i2017-12371-9

THE EUROPEAN PHYSICAL JOURNAL A

Regular Article – Theoretical Physics

# Toward a complete theory for predicting inclusive deuteron breakup away from stability

G. Potel<sup>1,a</sup>, G. Perdikakis<sup>1,2,3,b</sup>, B.V. Carlson<sup>4,c</sup>, M.C. Atkinson<sup>5</sup>, W.H. Dickhoff<sup>5</sup>, J.E. Escher<sup>6</sup>, M.S. Hussein<sup>4,7,8</sup>, J. Lei<sup>9,d</sup>, W. Li<sup>1</sup>, A.O. Macchiavelli<sup>10</sup>, A.M. Moro<sup>9</sup>, F.M. Nunes<sup>1,11</sup>, S.D. Pain<sup>12</sup>, and J. Rotureau<sup>1</sup>

- State of the art inclusive (d,p)
- Employs local DOM potentials constrained for <sup>40</sup>Ca and <sup>48</sup>Ca and extrapolated to <sup>60</sup>Ca
- Explores link with  $(n,\gamma)$  process

# Why DOM?

#### Compare standard optical potential with DOM

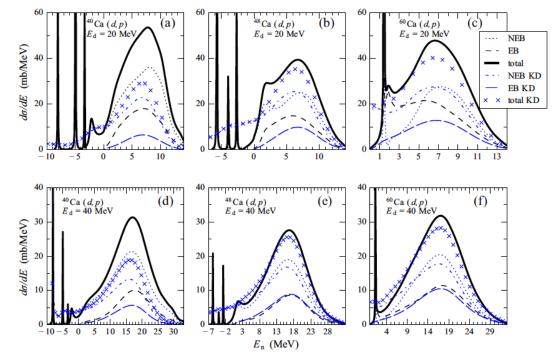
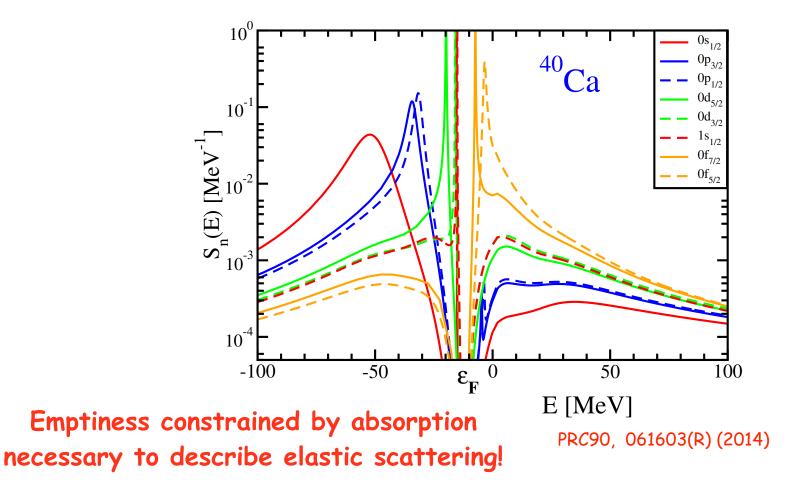


Fig. 8. Comparison of KD phenomenological optical potential and the DOM: elastic breakup (EB) and non-elastic breakup (NEB) proton spectra for the reactions  ${}^{40}\text{Ca}(d, p)$ ,  ${}^{48}\text{Ca}(d, p)$ , and  ${}^{60}\text{Ca}(d, p)$ , at  $E_d = 20$  MeV and  $E_d = 40$  MeV.

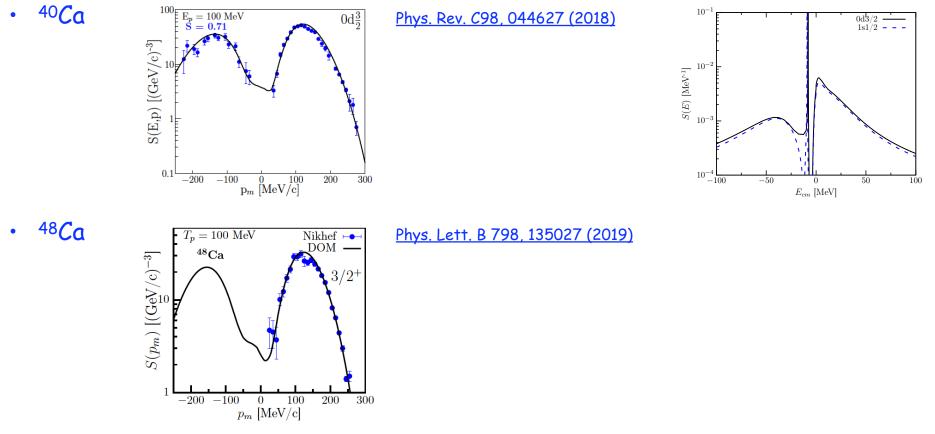
 Global potentials do not generate relevant information at negative energies —> dispersive approach essential for unstable nuclei

# Spectral function for bound states from DOM analysis

[0,200] MeV -> constrained by elastic scattering data



# Check with (e,e'p) cross sections (Mack Atkinson)

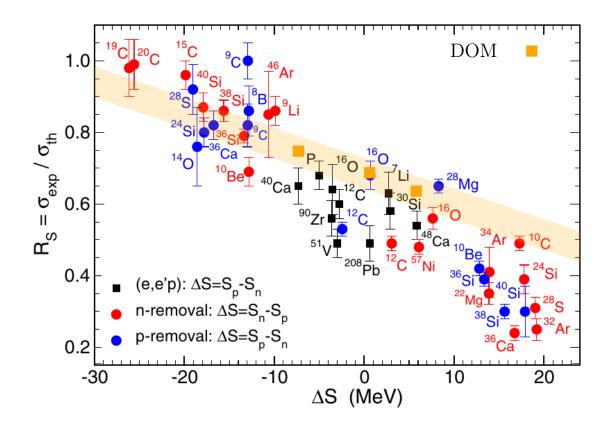


- No further adjustments!
- Both structure and reaction properties allowed to change when 8 n added

DOM

# Compare with Gade plot

Very near the Fermi energy in <sup>40</sup>Ca and <sup>48</sup>Ca from (e,e'p) -> error band

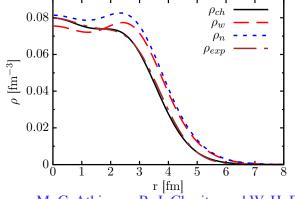


Quenching sp strength review: Aumann et al, Prog. Part. Nucl. Phys. 118, 103847 (2021)

DOM

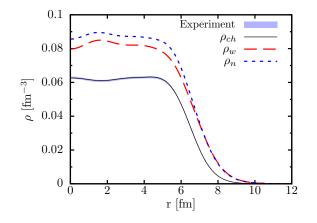
# Neutron skins in <sup>48</sup>Ca and <sup>208</sup>Pb from DOM predictions

• DOM 2017



M. H. Mahzoon, M. C. Atkinson, R. J. Charity, and W. H. Dickhoff Phys. Rev. Lett. **119**, 222503 (2017), 1-5.

• DOM 2020



M. C. Atkinson, M. H. Mahzoon, M. A. Keim, B. A. Bordelon, C. D. Pruitt, R. J. Charity, and W. H. Dickhoff Phys. Rev. C 101,044303 (2020), 1-15.

#### MCMC DOM prediction of neutron skins



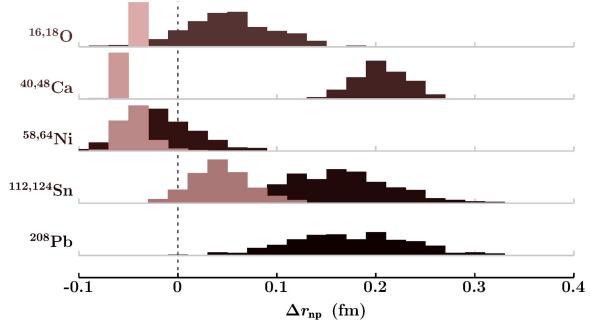


TABLE I. Neutron skins ( $\Delta r_{np}$ ), in fm, from this work. The 16th, 50th, and 84th percentile values of the skin distribution are reported as  $50_{16}^{84}$ .

<sup>16</sup> O	<sup>18</sup> O	<sup>40</sup> Ca	<sup>48</sup> Ca	<sup>58</sup> Ni	<sup>64</sup> Ni	<sup>112</sup> Sn	<sup>124</sup> Sn	<sup>208</sup> Pb
$-0.025^{-0.023}_{-0.027}$	$0.06_{0.02}^{0.11}$	$-0.051^{-0.048}_{-0.055}$	$0.22_{0.19}^{0.24}$	$-0.03\substack{-0.02\\-0.05}$	$-0.01^{0.03}_{-0.04}$	$0.05_{0.02}^{0.08}$	$0.17_{0.12}^{0.23}$	$0.18_{0.12}^{0.25}$

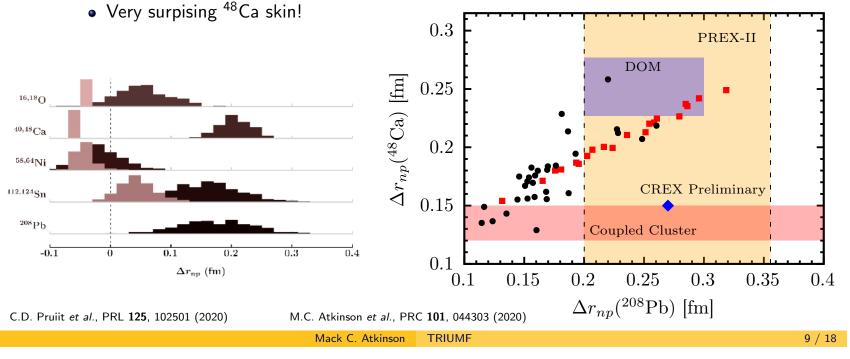
C. D. Pruitt, R. J. Charity, L. G. Sobotka, M. C. Atkinson, and W. H. Dickhoff Phys. Rev. Lett. 125, 102501 (2020), 1-6.

DOM

#### Neutron skin puzzle & DOM (slide Mack Atkinson now at LLNL)

#### Neutron Skin: $\Delta r_{np} = r_n - r_p$

- $r_n$  can be measured through parity-violating electron scattering (weak)
- PREX-II at Jefferson Lab measured <sup>208</sup>Pb skin
  - Preliminary CREX results for <sup>48</sup>Ca released at DNP meeting 2021



Optical Potentia

# Outlook FRIB era

- Accurate global optical potential requires a lot of theoretical effort
  - Topical Collaboration: Holt, Dickhoff, Elster, Lovell, Nunes, Potel, Stroberg (DOE proposal)
- · Accurate global optical potential requires a lot of experimental effort as well!
- Removal probability puzzle only slow progress
- Neutron skins very interesting

#### Conceptual paradigm shift: need to treat reactions and structure simultaneously

• (p,2p) good starting point

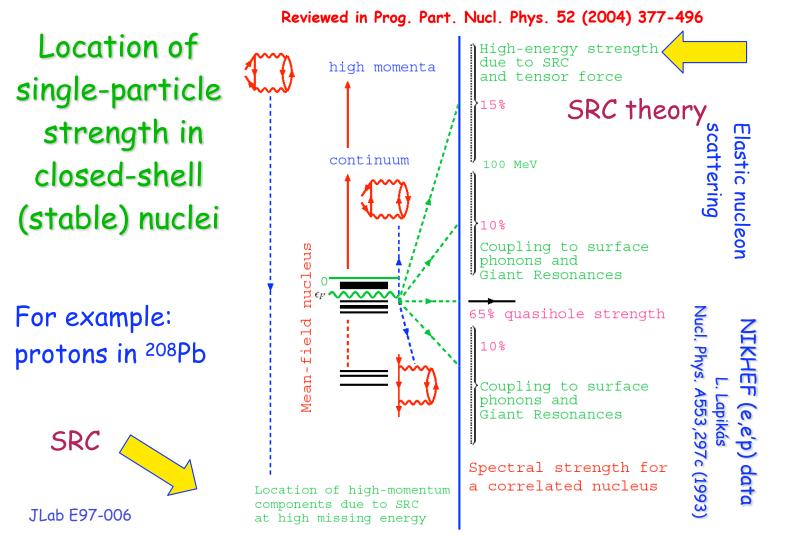
PHYSICAL REVIEW C 105, 014622 (2022)

First DOM analysis -->

First application of the dispersive optical model to (p, 2p) reaction analysis within the distorted-wave impulse approximation framework

K. Yoshida<sup>(D)</sup>,<sup>1,\*</sup> M. C. Atkinson<sup>(D)</sup>,<sup>2</sup> K. Ogata<sup>(D)</sup>,<sup>3,4,5</sup> and W. H. Dickhoff<sup>(D)</sup>

**Optical Potential** 



Phys. Rev. Lett. 93, 182501 (2004) D. Rohe et al.

High-momentum predictions & relation to ground-state energy Ground-state energy can be included in the DOM

$$E/A = \frac{1}{2A} \sum_{\ell j} (2j+1) \int_0^\infty \!\!\! dk k^2 \frac{k^2}{2m} n_{\ell j}(k) + \frac{1}{2A} \sum_{\ell j} (2j+1) \int_0^\infty \!\!\! dk k^2 \int_{-\infty}^{\varepsilon_F} dE \; E \; S_{\ell j}(k;E)$$

Succeeds	Α	DOM $E_0^A/A$	Mass equation	Expt. $E_0^A/A$
	<sup>12</sup> C	-7.85	-7.29	-7.68
	<sup>40</sup> Ca	-8.46	-8.50	-8.55
	<sup>48</sup> Ca	-8.66	-8.59	-8.66
	<sup>208</sup> Pb	-7.76	-7.81	-7.87

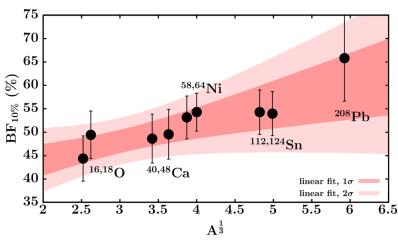
Phys. Rev. C 102, 044333 (2020)

Because fraction of binding energy from 10% most deeply bound nucleons includes the

high-momentum contribution

Phys. Rev. Lett. 125, 102501 (2020)

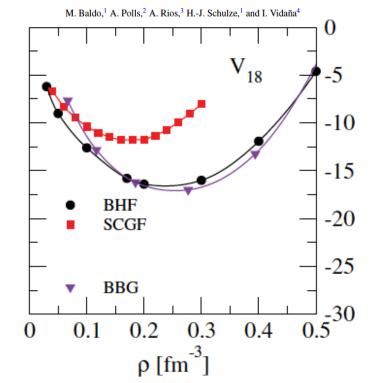
Predicted in Phys. Rev. C51, 3040 (1995)



#### SCGF & SRC compared to BHF and BBG

PHYSICAL REVIEW C 86, 064001 (2012)

Comparative study of neutron and nuclear matter with simplified Argonne nucleon-nucleon potentials

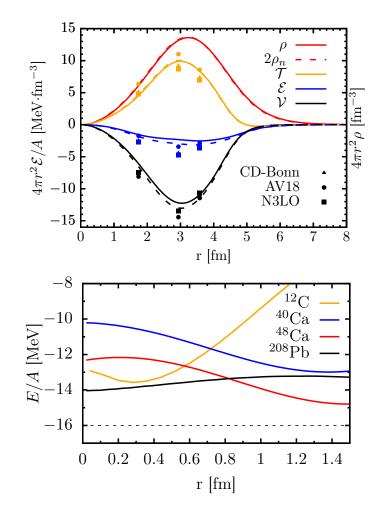


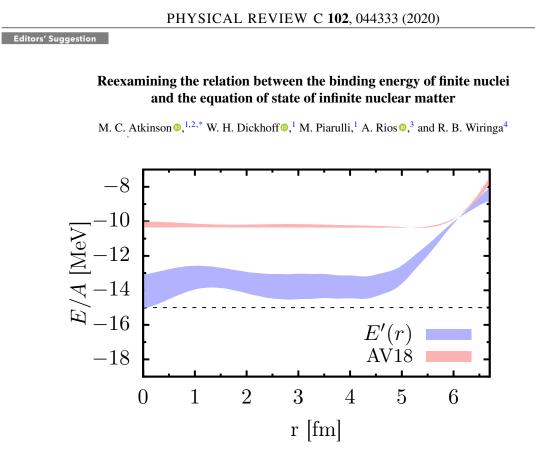
• BBG requires a repulsive NNN at high density to improve density

Neutron skins and EOS

#### Consequence

#### Maybe 16 MeV binding is not needed!





Neutron skins and EOS

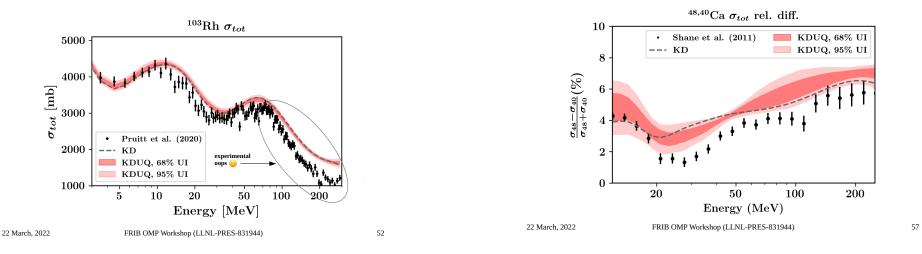
# Cole Pruitt (LLNL)

"OMP uncertainty - characterized empirically for KD and CH89 - is larger than assumed:

Global KD OMP has systematic over/underestimation for  $\sigma_{rxn}$ 

Roughly 20% std. dev. for (n,\*)/(p,\*) cross sections on stable targets KDUQ/CHUQ help pinpoint where data and model are grossly inconsistent"

• Total neutron cross sections using an iterative procedure using KD potential



• Problem at high energy but relative OK —> another look at the neutron skin of <sup>48</sup>Ca?

**Optical Potential** 

#### Propagator / Green's function and spectral functions & spectroscopic factors

Lehmann representation

$$\begin{aligned} G_{\ell j}(k,k';E) &= \sum_{m} \frac{\langle \Psi_{0}^{A} | \, a_{k\ell j} \, | \Psi_{m}^{A+1} \rangle \, \langle \Psi_{m}^{A+1} | \, a_{k'\ell j}^{\dagger} \, | \Psi_{0}^{A} \rangle}{E - (E_{m}^{A+1} - E_{0}^{A}) + i\eta} \\ &+ \sum_{n} \frac{\langle \Psi_{0}^{A} | \, a_{k'\ell j}^{\dagger} \, | \Psi_{n}^{A-1} \rangle \, \langle \Psi_{n}^{A-1} | \, a_{k\ell j} \, | \Psi_{0}^{A} \rangle}{E - (E_{0}^{A} - E_{n}^{A-1}) - i\eta} \end{aligned}$$

- Any other single-particle basis can be used & continuum integrals implied
- Overlap functions --> numerator
   Corresponding eigenvalues
   --> denominator
- Spectral function  $S_{\ell j}(k; E) = \frac{1}{\pi} \operatorname{Im} G_{\ell j}(k, k; E) \qquad E \leq \varepsilon_F^ = \sum_n \left| \langle \Psi_n^{A-1} | a_{k\ell j} | \Psi_0^A \rangle \right|^2 \delta(E - (E_0^A - E_n^{A-1}))$
- Discrete transitions

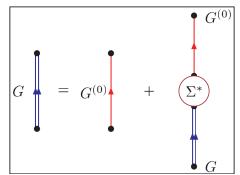
 $\sqrt{S_{\ell j}^n} \phi_{\ell j}^n(k) = \langle \Psi_n^{A-1} | a_{k\ell j} | \Psi_0^A \rangle$ 

- Momentum distribution: integrate spectral function to  $\varepsilon_F^-$
- Positive energy —> see later



DOM

# Propagator from Dyson Equation and "experiment"



#### Equivalent to ...

 $G = G^{(0)} + \Sigma^{*}$   $F = G^{(0)} + \Sigma^{*}$  $\Rightarrow$  as extracted from (e,e'p) reaction

$$\frac{k^2}{2m}\phi_{\ell j}^n(k) + \int dq \ q^2 \ \Sigma_{\ell j}^*(k,q;E_n^-) \ \phi_{\ell j}^n(q) = E_n^- \ \phi_{\ell j}^n(k)$$
  
Spectroscopic factor  $\ \mathcal{Z}_{\ell j}^n = \int dk \ k^2 \ \left| \langle \Psi_n^{A-1} | \ a_{k\ell j} \ | \Psi_0^A \rangle \right|^2 < 1$ 

Dyson equation also yields  $\left[\chi^{elE}_{\ell j}(r)\right]^* = \langle \Psi^{A+1}_{elE} | \, a^{\dagger}_{r\ell j} \, | \Psi^A_0 \rangle$  for positive energies

Elastic scattering wave function for protons or neutrons

Dyson equation therefore provides:

Link between scattering and structure data from dispersion relations