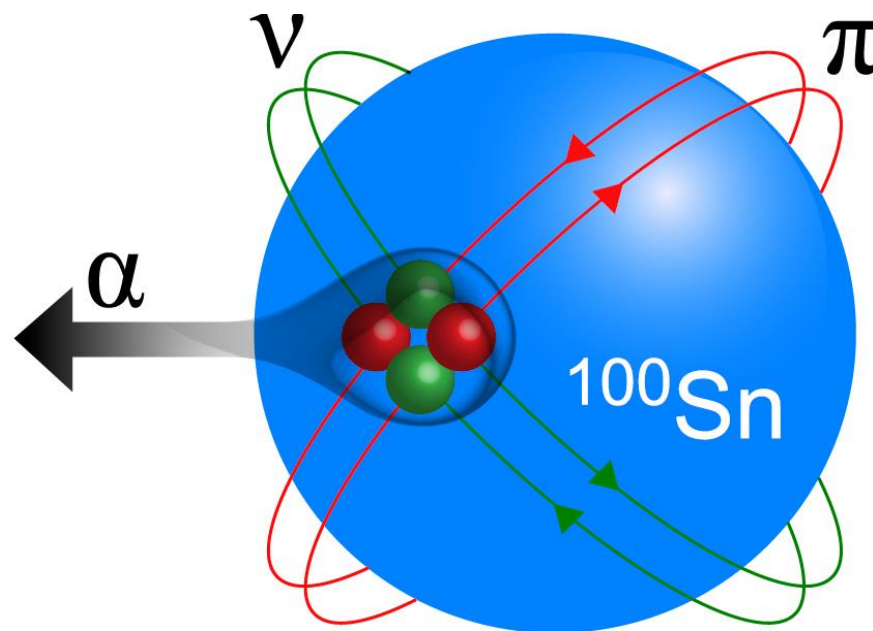


Superaligned α decay to doubly magic ^{100}Sn

Darek Seweryniak
Argonne National Laboratory



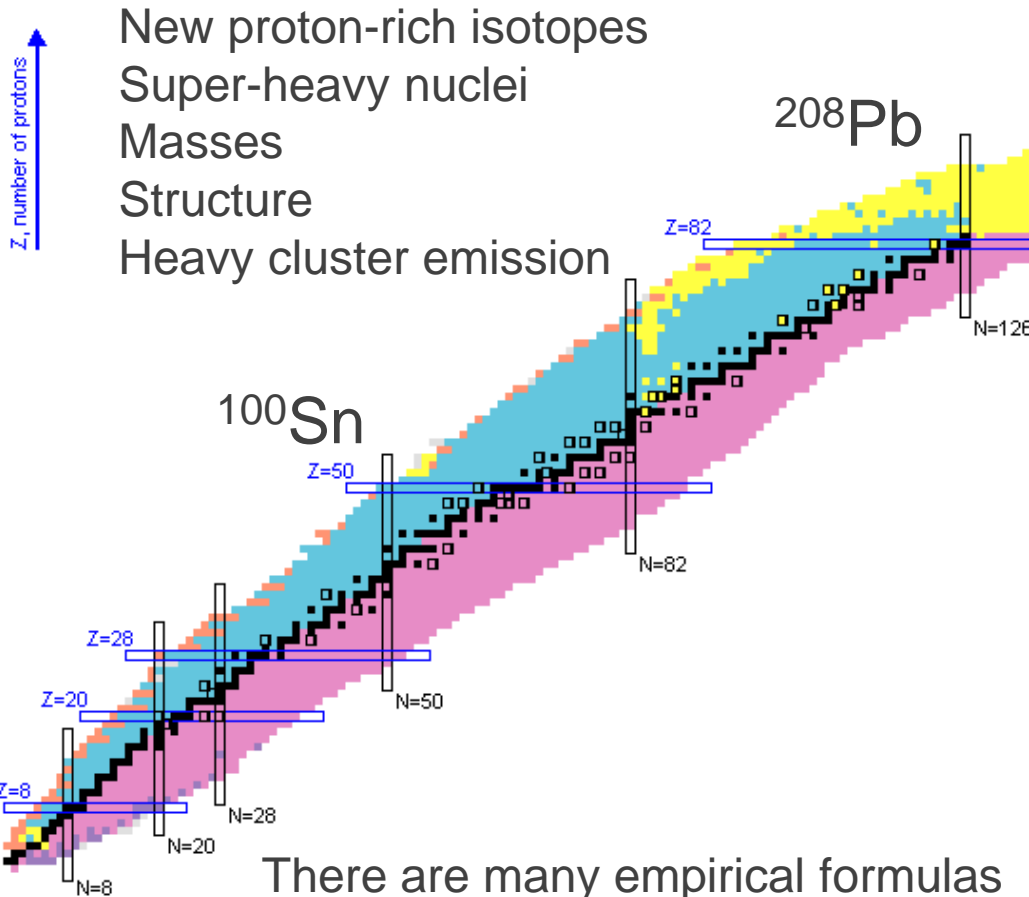
Symposium “Developments of Physics of Unstable Nuclei”

Yukawa International Seminar YKIS2022b

Outline

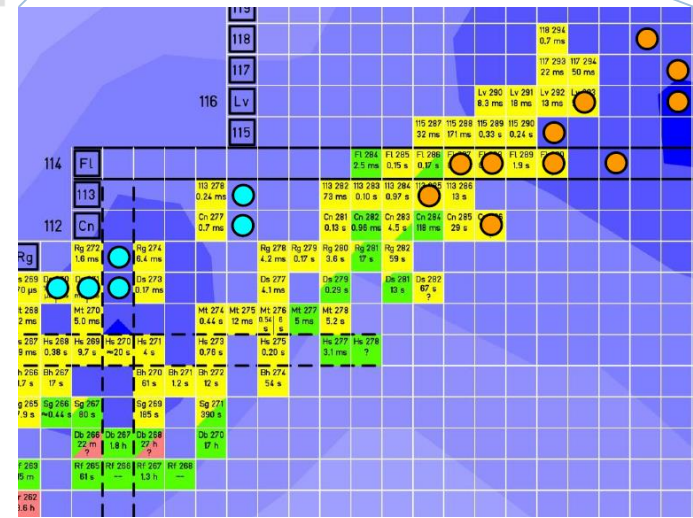
- Alpha decay landscape
- Microscopic description of α decay
- Doubly-magic ^{100}Sn
- Observation of the ^{108}Xe - ^{104}Te - ^{100}Sn α decay chain
- Discussion of α -decay reduced widths
 - ^{208}Pb region vs ^{100}Sn region
 - theoretical calculations for ^{104}Te
- Summary and Outlook

Alpha-decay landscape



There are many empirical formulas for calculating α -decay widths but **microscopic description of α decay remains challenging**

Super-Heavy Elements



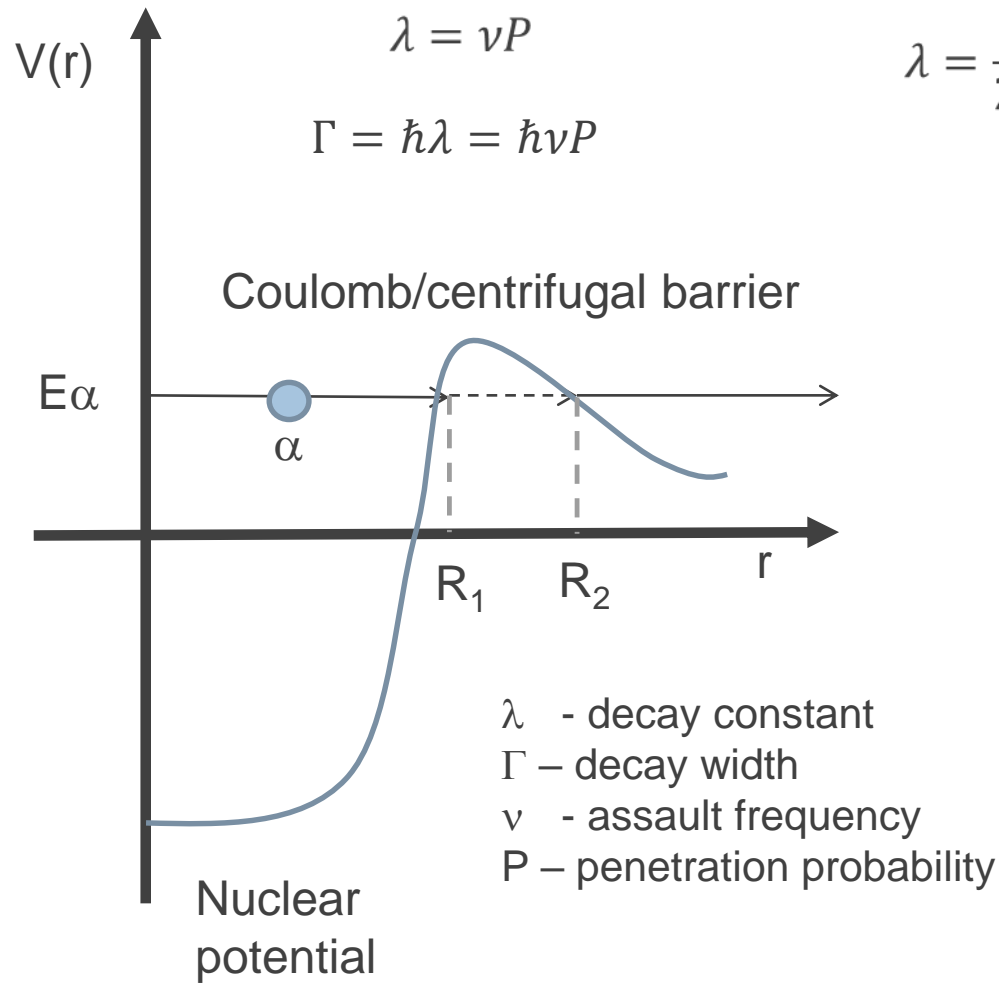
α decay of ^{212}Po to doubly magic ^{208}Pb is the simplest case and serves as a benchmark

α decay of ^{104}Te to doubly magic ^{100}Sn , enhanced α preformation factor due to strong π - ν interaction, **superallowed α decay**

Gamow alpha-decay model

G. Gamow, *Z. Phys.* 51, 204 (1928)

Probabilistic interpretation of quantum mechanics



$$\lambda = \frac{\nu}{2R_0} \exp \left[-2 \int_{R_1}^{R_2} \sqrt{\frac{2\mu}{\hbar} |Q_\alpha - V(r)|} dr \right]$$

Can be readily calculated
Very steep function of Q-value

$$\Gamma = \delta^2 P$$

$$\delta^2 = \frac{\Gamma_{exp}}{P_{calc}}$$

δ^2 - reduced α decay width,
often normalized to ^{212}Po
($\sim \alpha$ preformation factor)

R-matrix expression of the alpha-decay width

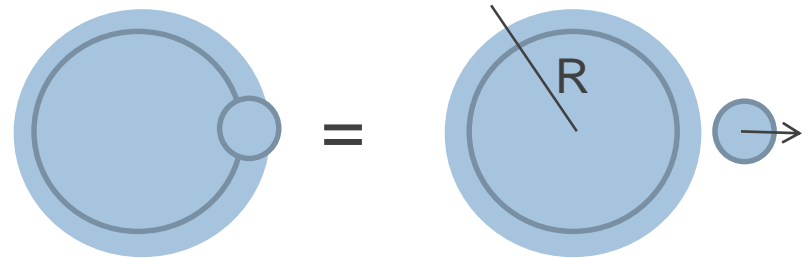
Decay width:

$$\Gamma_L(R) = 2\gamma_L^2(R)P_L(R)$$

$P_L(R)$ - penetrability
 R - channel radius (outside of the nucleus)

Reduced width amplitude:

$$\gamma_L(R) = \left(\frac{\hbar^2 R}{2\mu} \right)^{1/2} F_L(R)$$



Formation amplitude:

$$F_L(R) = \int d\xi_\alpha d\xi_D d\hat{R} [\underbrace{\phi_\alpha(\xi_\alpha)}_{\text{parent nucleus}} \underbrace{\psi_D(\xi_D)}_{\text{daughter nucleus}} Y_L(\hat{R})]_{\alpha_4, \nu_4}^* \underbrace{\psi_P(\xi_\alpha \xi_D; R)}_{\text{alpha}}$$

Overlap between **parent nucleus** and **daughter nucleus** + **alpha**
at a distance R outside of nuclear interactions

Spectroscopic factor formulation

Decay width:

$$\Gamma_L = S_L \Gamma_L^{sp}$$

Γ_L^{sp} – single-particle decay width
 L – α angular momentum

Spectroscopic factor:

$$S_L = \left| \langle \mathcal{A}[\phi_\alpha(\xi_\alpha) \Psi_j^D(\xi_D) \psi_L(\mathbf{R})]_{JM} | \Phi_{JM}^P \rangle \right|^2$$

↑ ↑ ↓ ↓

α particle daughter nucleus relative motion parent nucleus

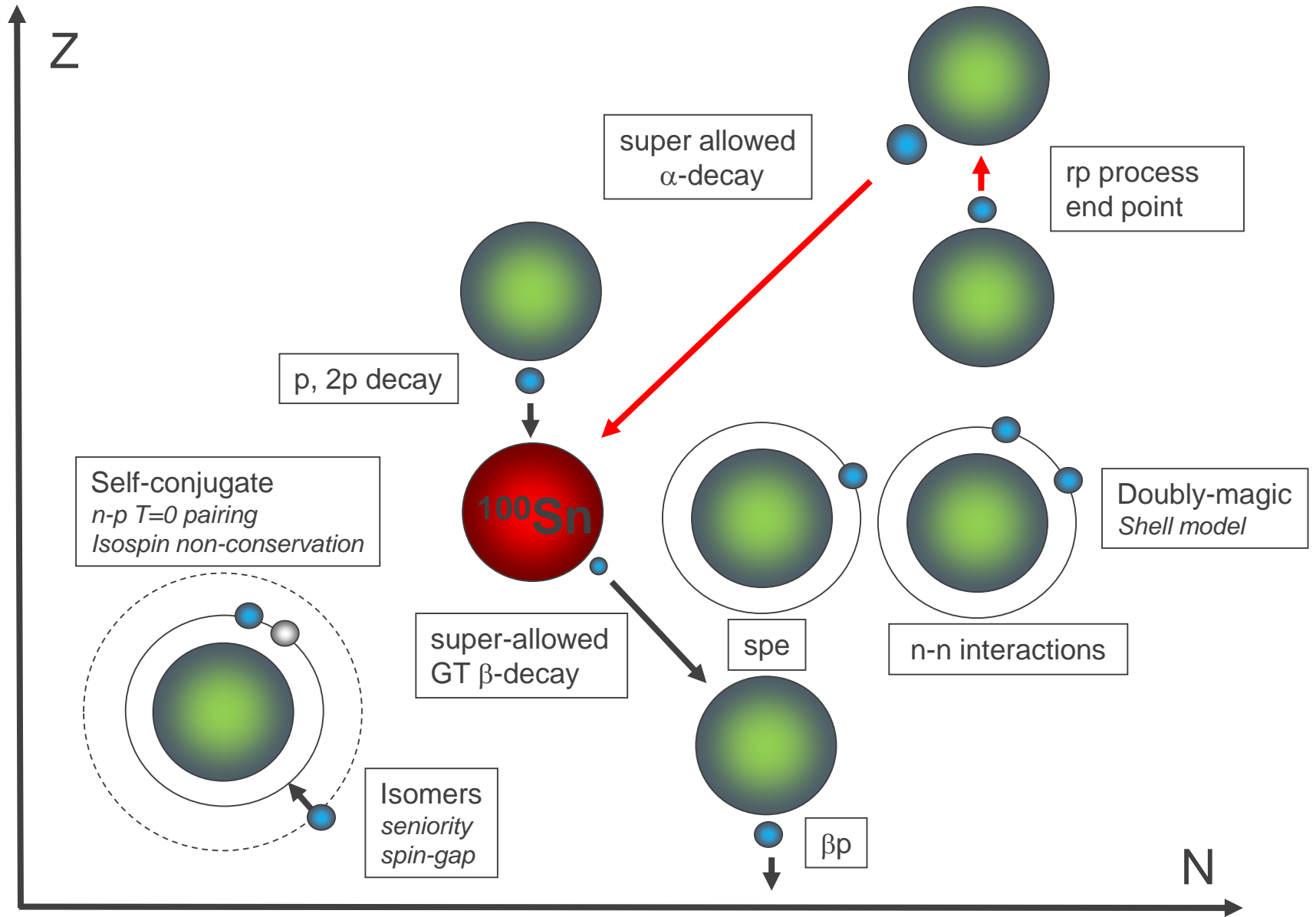
probability of finding α particle in the parent nucleus



Challenges of microscopic α -decay width calculations

- Microscopic description using **the shell model**
 - Underestimates experimental values by about **2 orders of magnitude**
 - Only Shell Model+Cluster Model reproduces ^{212}Po
K. Varga et al., PRL 69, 37 (1992)
- Large configuration space
- Antisymmetrization, Normalization
- Configuration mixing (nucleon-nucleon residual interaction)
 - pairing, **proton-neutron interaction**
- Contribution from the continuum

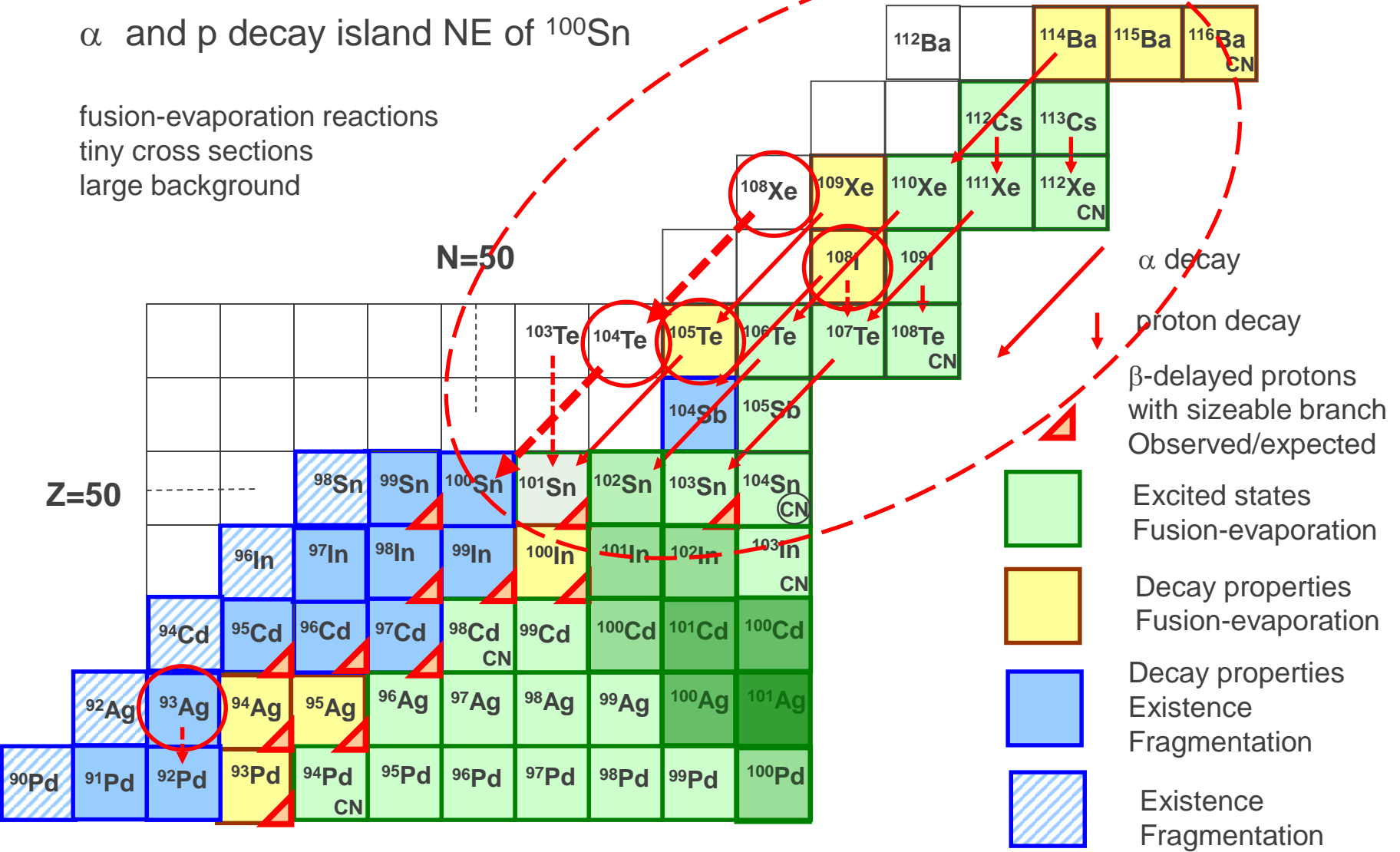
^{100}Sn physics



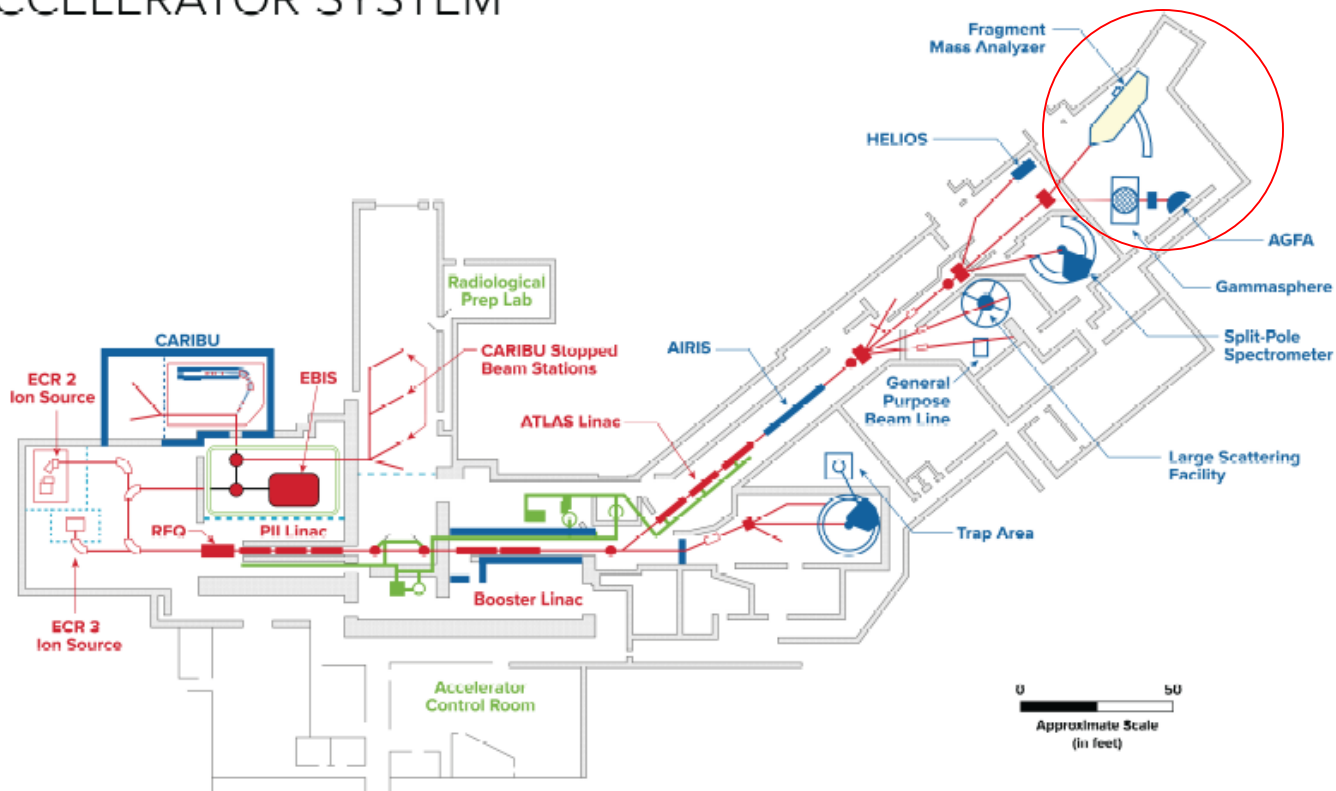
Charged-particle emitters near ^{100}Sn

α and p decay island NE of ^{100}Sn

fusion-evaporation reactions
 tiny cross sections
 large background

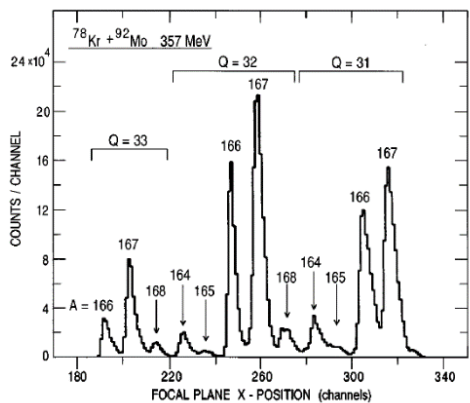
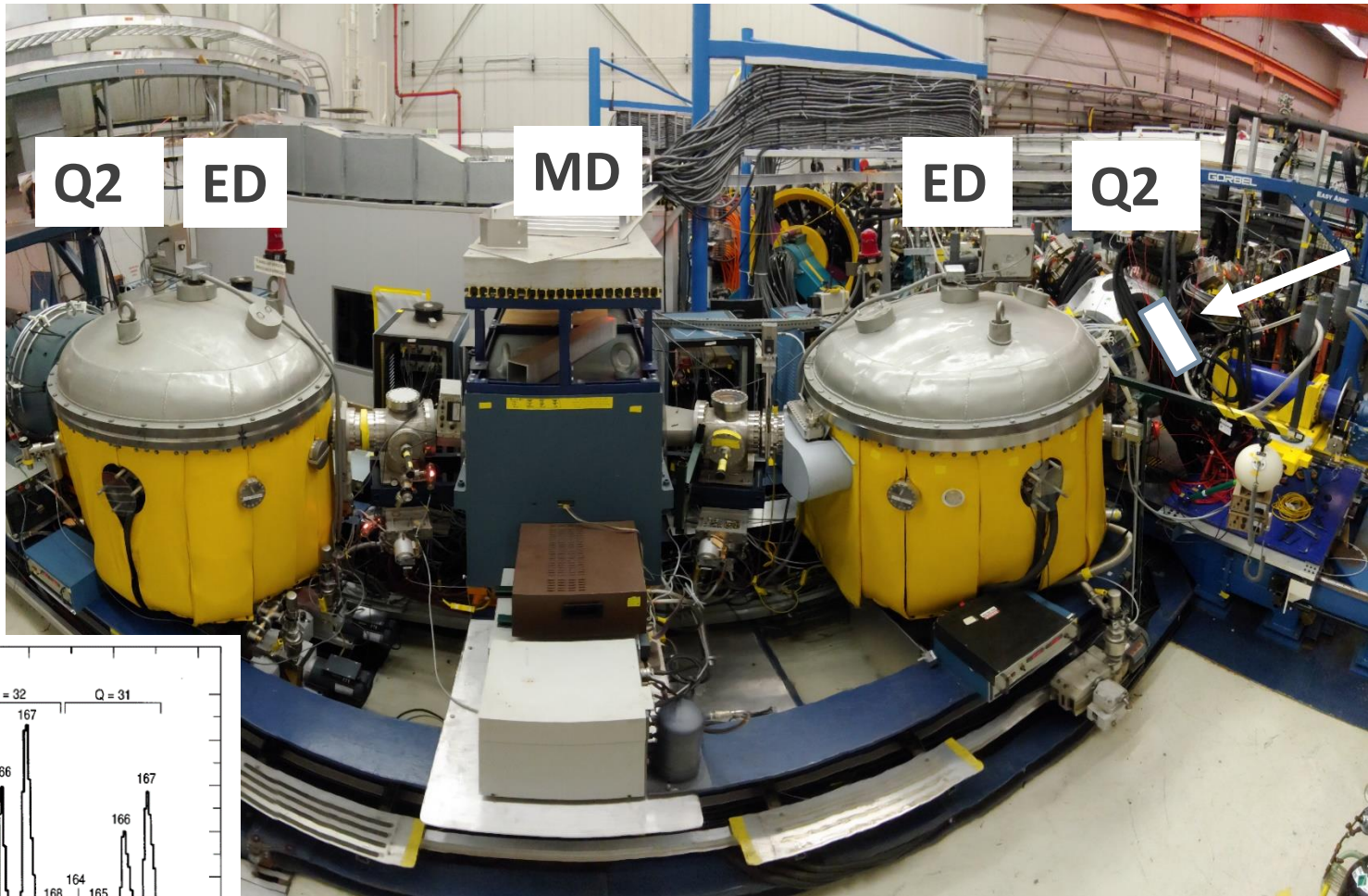


ATLAS ARGONNE TANDEM LINEAR ACCELERATOR SYSTEM



Triple LINAC: beam from H to U up to 10 MeV/n
In-flight radioactive beams with RASOR separator
CARIBU - ^{252}Cf fission fragment beams

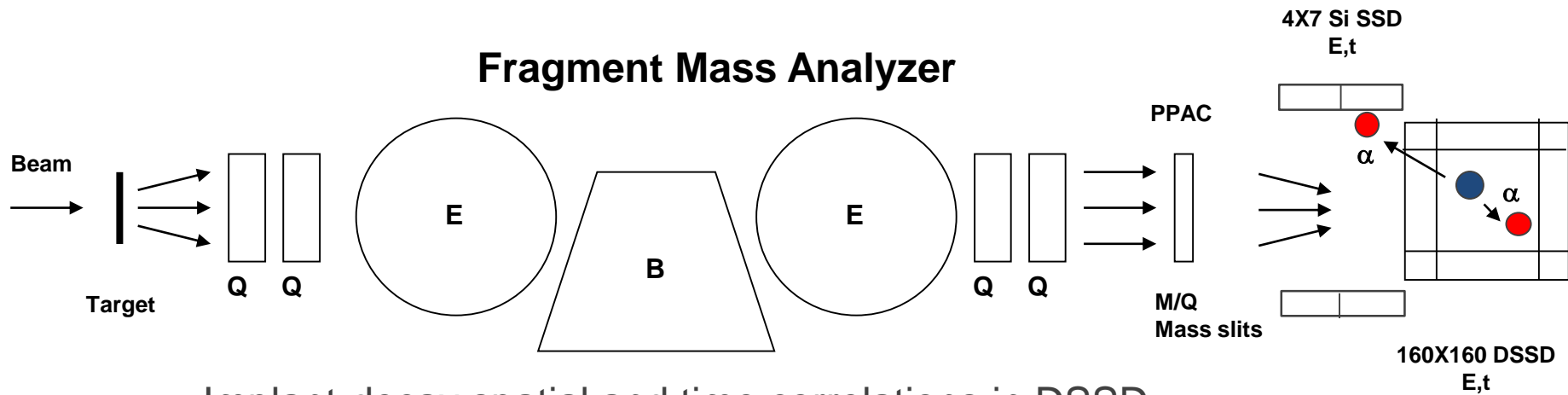
Argonne Fragment Mass Analyzer



YKIS2022b

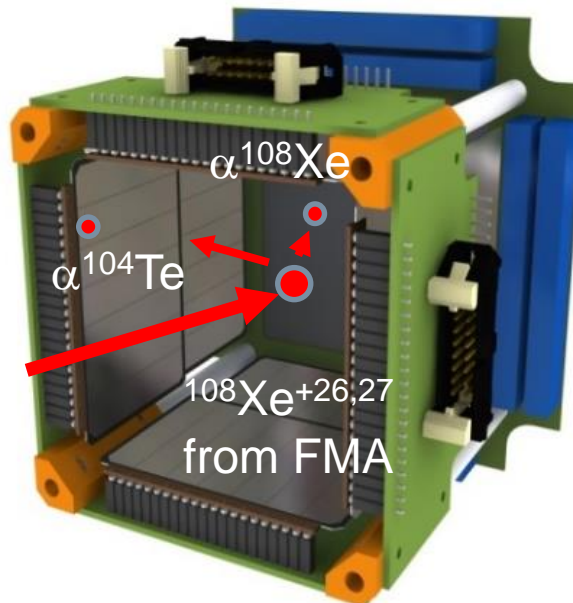
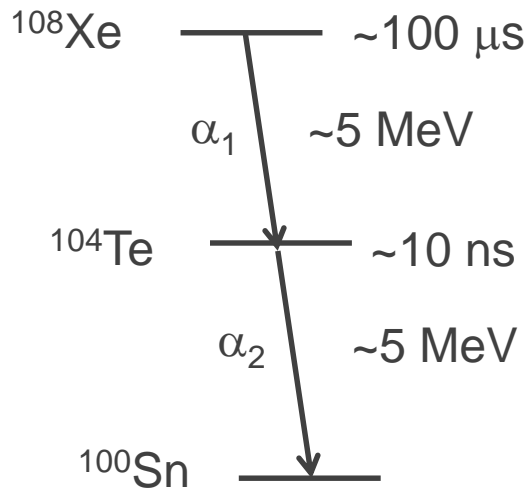
Super-allowed α decay ^{108}Xe - ^{104}Te - ^{100}Sn

Fragment Mass Analyzer



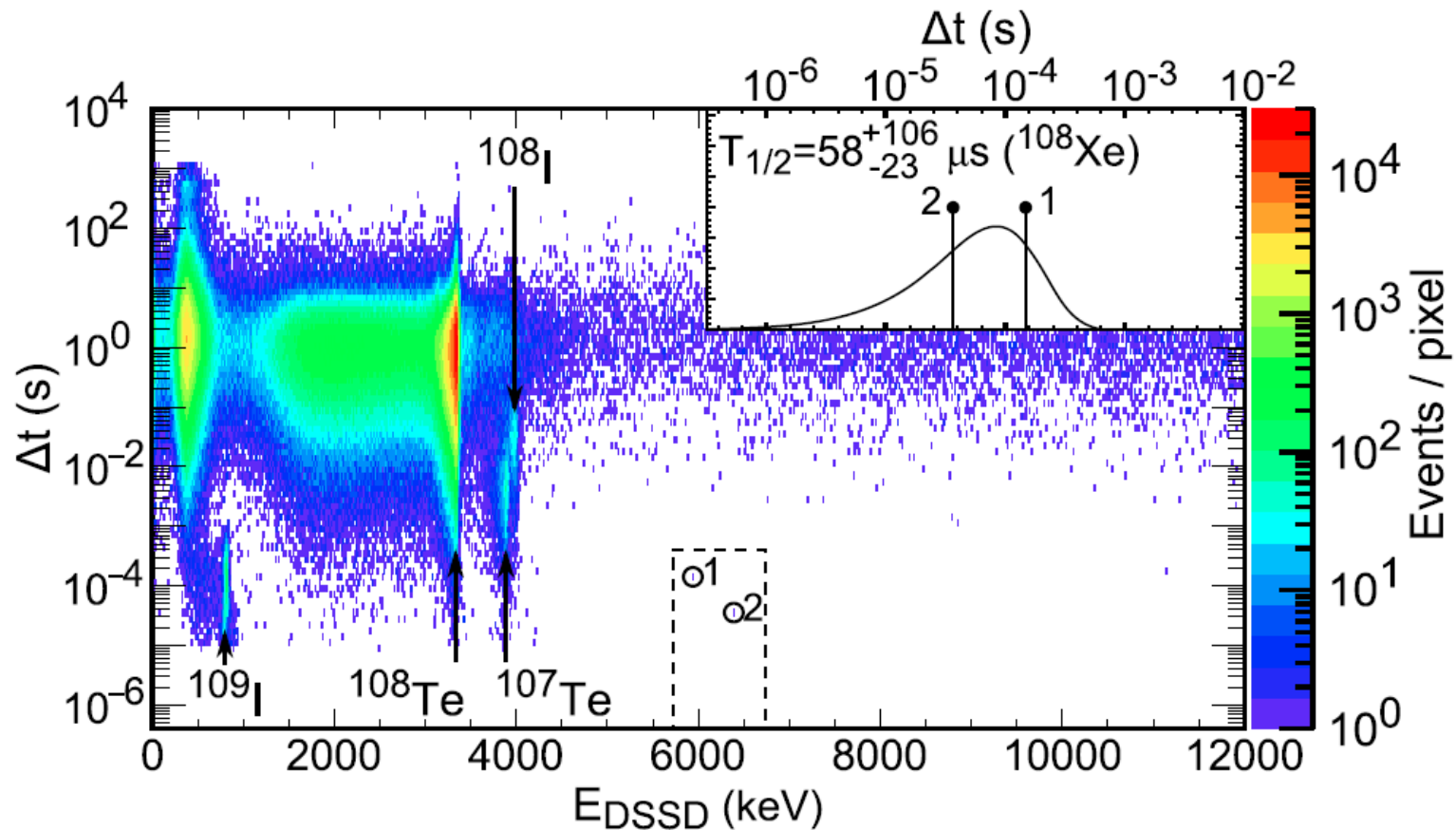
Implant-decay spatial and time correlations in DSSD

$^{58}\text{Ni}(^{54}\text{Fe},4n)^{108}\text{Xe}$ reaction, **cross section 100 pb out of 1b!**



Digital DAQ to detect $\alpha^{108}\text{Xe}$ - $\alpha^{104}\text{Te}$ pileup
Si box to catch escaping alphas

Recoil-decay correlations

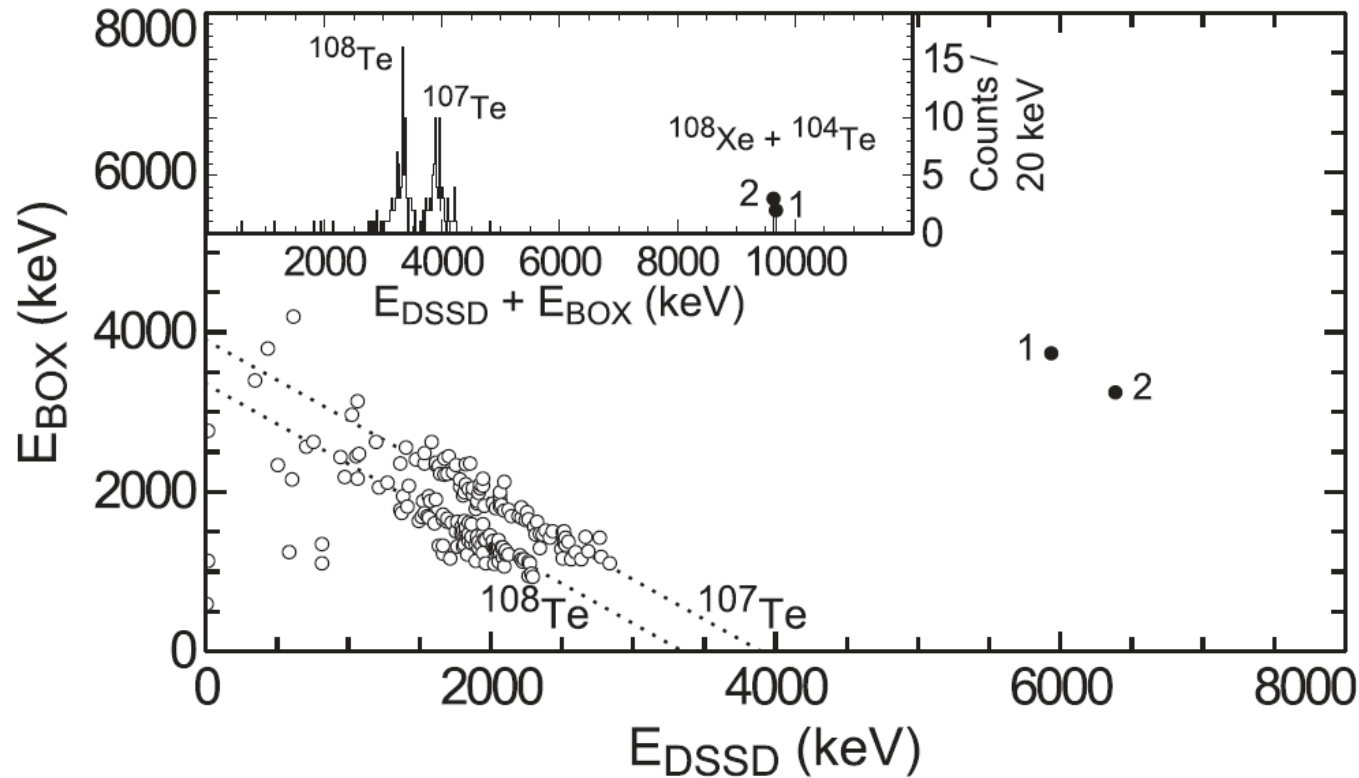


TWO fast high energy decay events

Expected **0.09** random events

BOTH events were in coincidence with the Si box (1 out of 400)

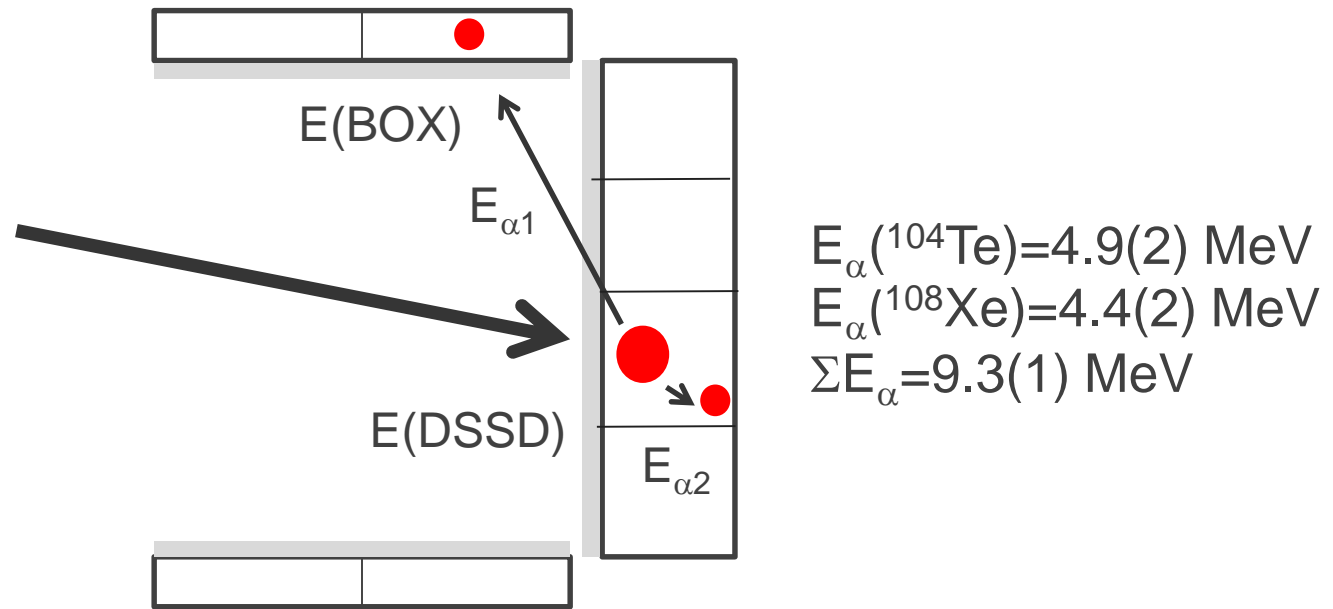
DSSD-Si box coincidences



- The same total energy for both events - $\Sigma E_{\alpha} = 9.3(1)$ MeV
- Compared to α emitters different energy split

$$E_{\alpha}(^{104}\text{Te}) = 4.9(2) \text{ MeV}, E_{\alpha}(^{108}\text{Xe}) = 4.4(2) \text{ MeV}$$

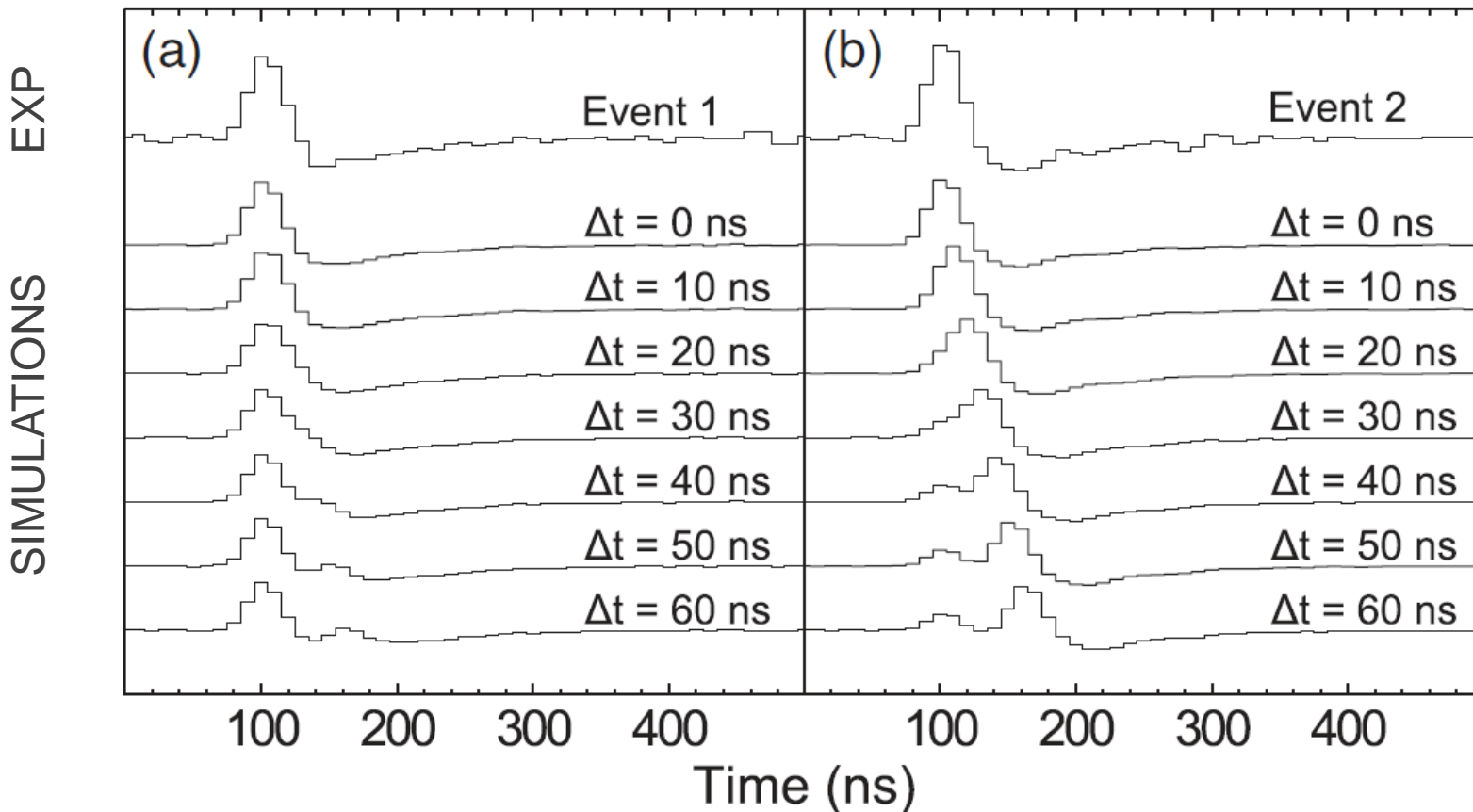
$^{108}\text{Xe}/^{104}\text{Te}$ α -particle energy determination



1. Implantation depth known from other α escapes
2. Emission angle known from detector geometry
3. Path in the DSSD calculated
4. Path in the box calculated from $E(\text{BOX})$
5. Path in dead layers calculated
6. Total $E_{\alpha 1}$ calculated from total path
7. $E_{\alpha 1}(\text{DSSD})$ deposited in DSSD calculated
8. Total $E_{\alpha 2}$ calculated as $E(\text{DSSD})-E_{\alpha 1}(\text{DSSD})$

DSSD traces for the ^{108}Xe - ^{104}Te pile-up events

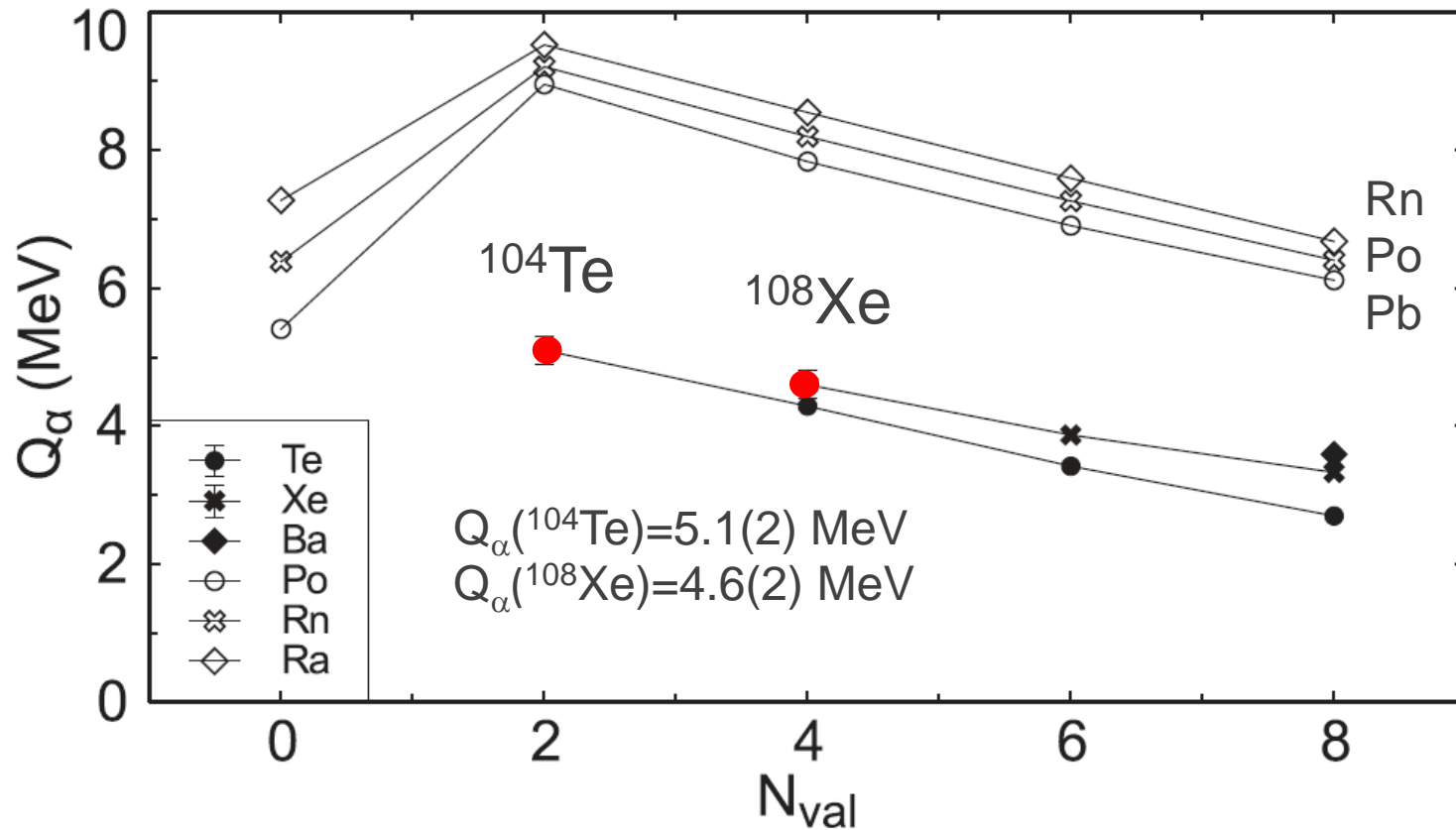
Doubly differentiated traces



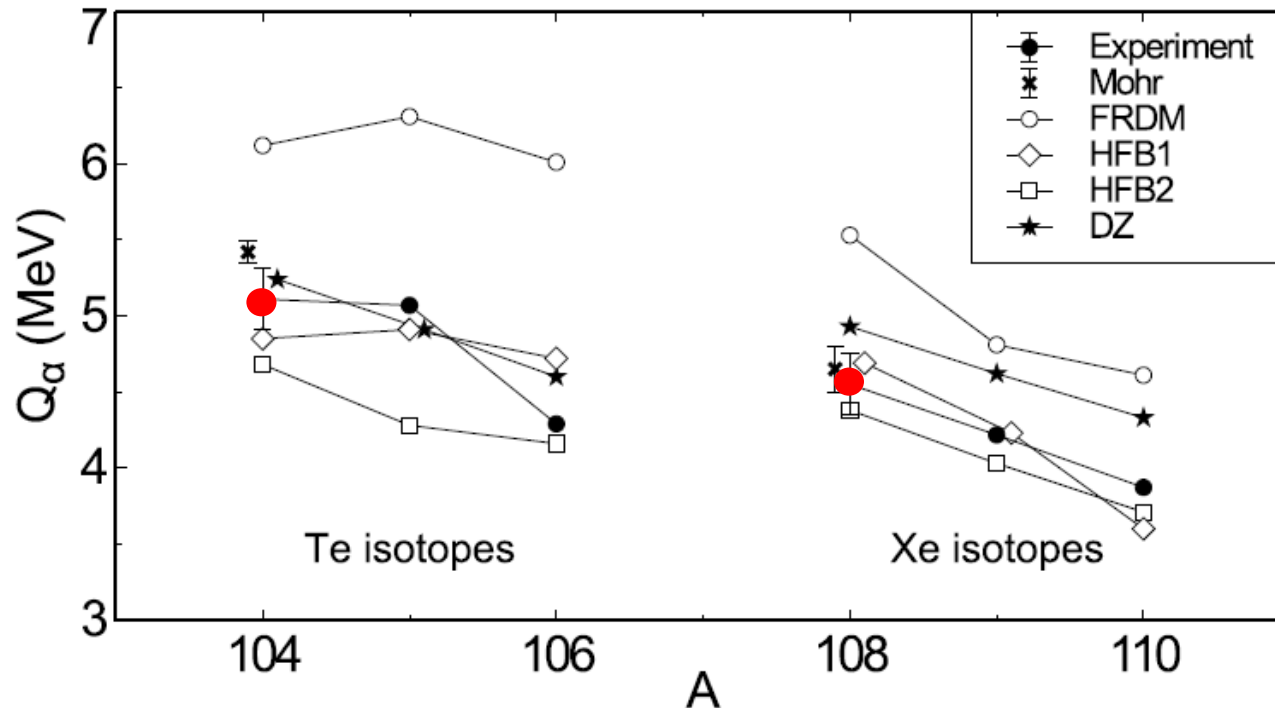
NO noticeable delay

TWO decays faster than 20 ns each imply $T_{1/2} < 18$ ns

Alpha-decay Q value systematics



Comparison with mass models



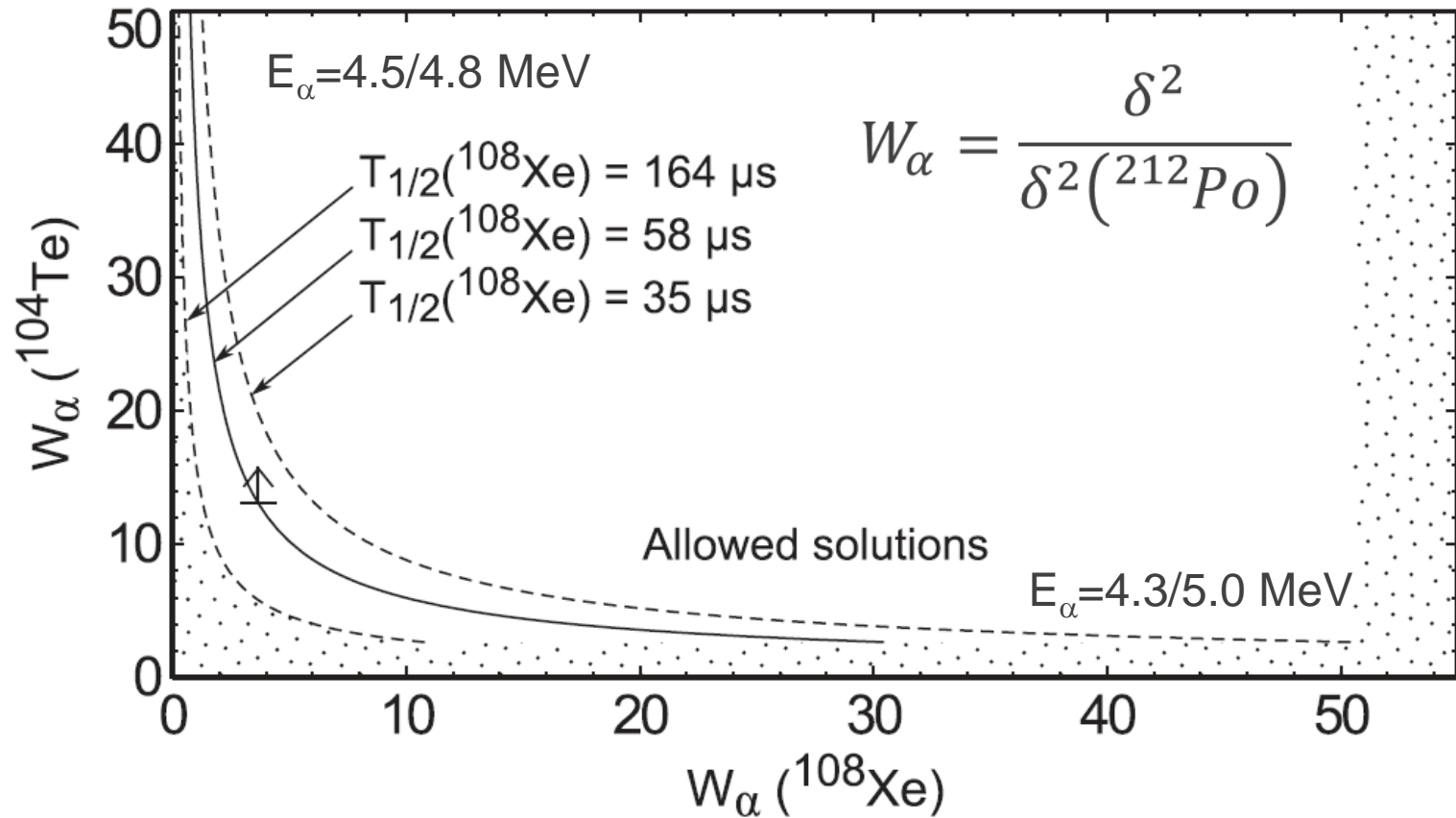
Locally adjusted double-folding potential

P. Mohr, Eur. Phys. J. A 31, 23 (2007)

$Q_\alpha(^{104}\text{Te})=5.42(0.07)$ MeV, $T_{1/2} (^{104}\text{Te}) = 5$ ns (assumed $\delta^2=10\%$)

$Q_\alpha(^{108}\text{Xe})=4.65(0.15)$, $T_{1/2} (^{108}\text{Xe}) = 60$ μs (assumed $\delta^2=5\%$)

^{108}Xe - ^{104}Te reduced width limits



$$W_\alpha(^{104}\text{Te}) > 13.1$$

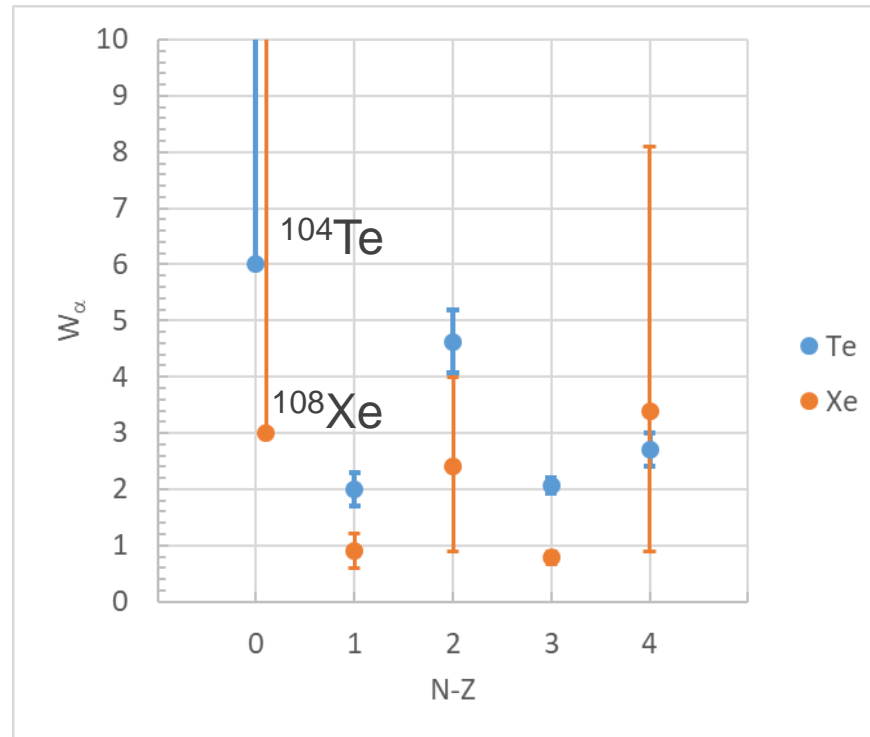
$$W_\alpha(^{104}\text{Te}) > 1.9$$

$$W_\alpha(^{108}\text{Xe}) = 3.7(-3.5 + 41.6)$$

$$W_\alpha(^{104}\text{Te})W_\alpha(^{108}\text{Xe}) > 25$$

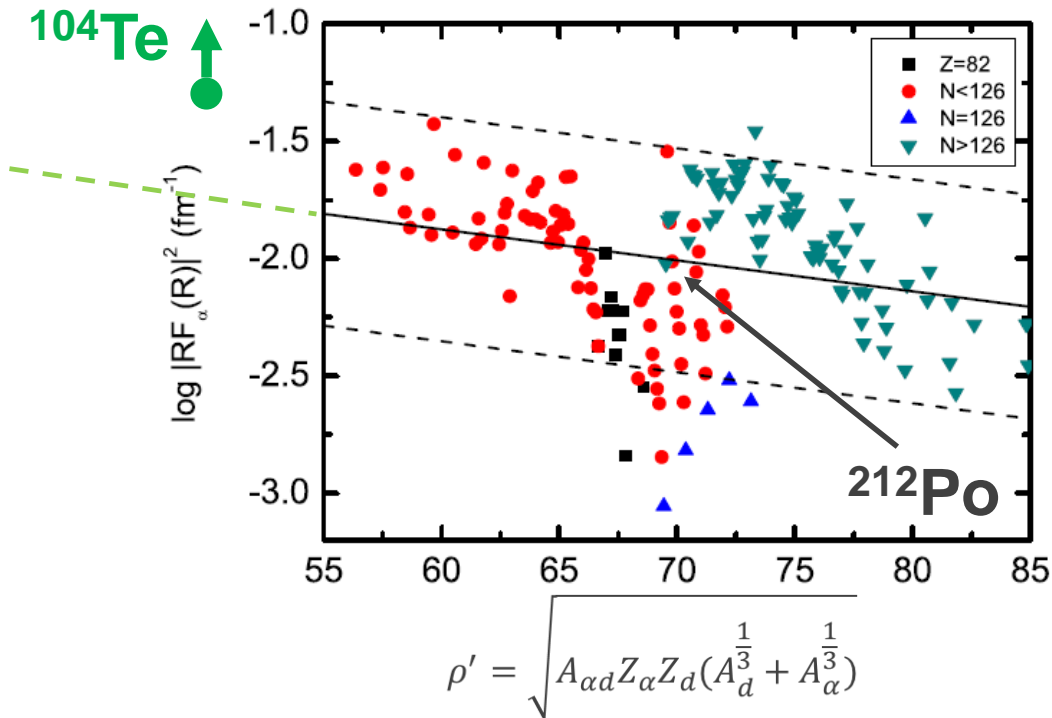
At least one W_α greater than 5

Reduced α -decay widths near ^{100}Sn

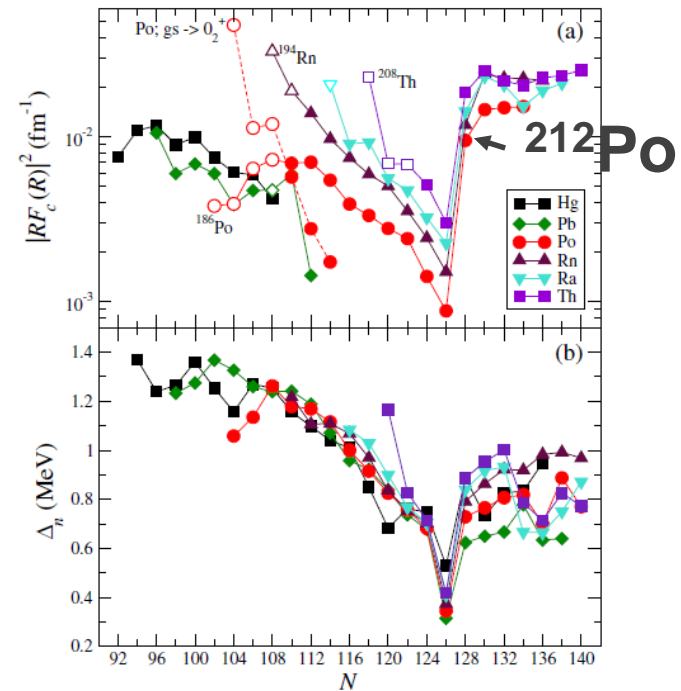


If $W_\alpha(\text{Te})/W_\alpha(\text{Xe}) \sim 2$ as for $^{106}\text{Te}/^{110}\text{Xe}$ pair: $W_\alpha(^{104}\text{Te}) > 6$, $W_\alpha(^{108}\text{Xe}) > 3$

Reduced α -decay width global systematics



C. Qi et al., *Phys. Rev. C* **81**, 064319 (2010)



A.N. Andreev et al., *PRL* **110**, 242502 (2013)

^{104}Te , $\rho' \sim 50$, $\log_{10} |R F(R)|^2 > -1.3$ ($W_{\alpha} > 5$)

Enhancement beyond pairing after size taken into account

Complex-energy shell model

R. Id Betan and W. Nazarewicz, Phys. Rev. C 86, 034338 (2012)

Spectroscopic factor, R-matrix

^{212}Po

Half life reproduced (spectroscopic factor formulation)

Calculated $T_{1/2}$ is 36 times too long (R-matrix)

Too small configuration space

^{104}Te

No convergence

$T_{1/2} < 500$ ns, assuming $Q_{\alpha} = 5.15$ MeV

Need to add proton-neutron interaction

Better treatment of continuum



Multistep shell model

Monika Patial, R. J. Liotta, and R. Wyss, *Phys. Rev. C* 93, 054326 (2016)

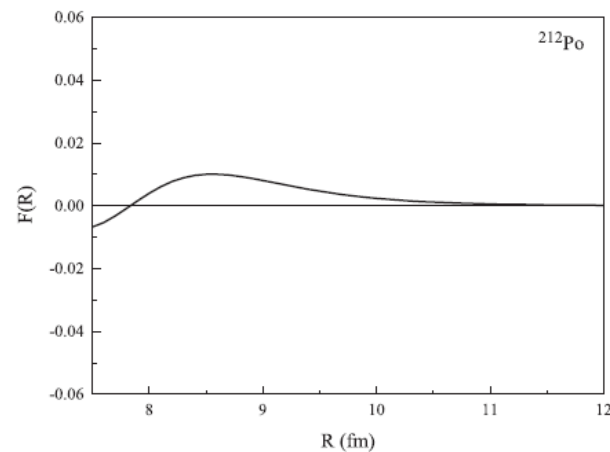
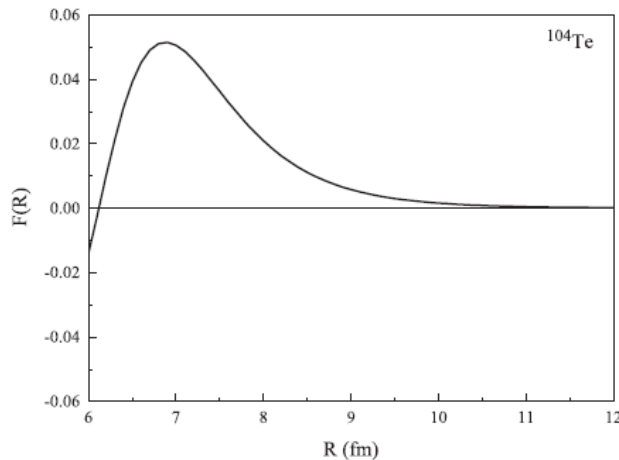
R-matrix

^{212}Po

Calculated $T_{1/2}=15\ \mu\text{s}$, experiment 298 ns

^{104}Te

Calculated $T_{1/2}=1.5\ \mu\text{s}$, assuming $Q_\alpha=5.06\ \text{MeV}$, experiment $T_{1/2}<15\ \text{ns}$
 α -particle formation probability **4.85 times larger** in ^{104}Te compared to ^{212}Po



Recent calculations for $^{104}\text{Te}/^{108}\text{Xe}$

Density-dependent cluster model plus two-potential approach

Dong Bai and Zhongzhou Ren, Eur. Phys. J. A 54, 220 (2018)

Relativistic Density Functional DD-PC1+separable pairing of finite range

Dynamical least-action paths from equilibrium deformation to scission

Assault frequency assumed

$T_{1/2}(^{108}\text{Xe})=50 \mu\text{s}$, $T_{1/2}(^{104}\text{Te})=197\text{ns}$

F. Mercier et al., Phys. Rev. C **102**, 011301(R) (2020)

Quartetting wave function approach

$T_{1/2}(^{104}\text{Te})=15 \text{ ns}$, $P_{\alpha} = 0.72$

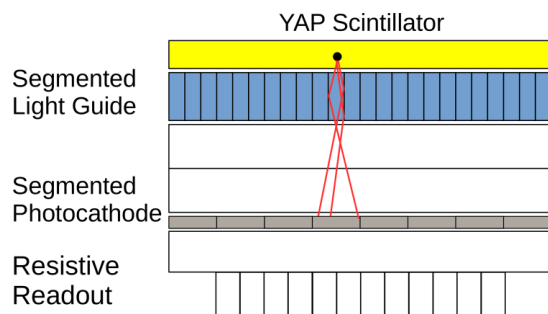
$T_{1/2}(^{212}\text{Po})=340 \text{ ns}$, $P_{\alpha} = 0.10$

Shuo Yang et al., Phys. Rev. C 101, 024316 (2020)

Cluster-Formation Model+DDCM

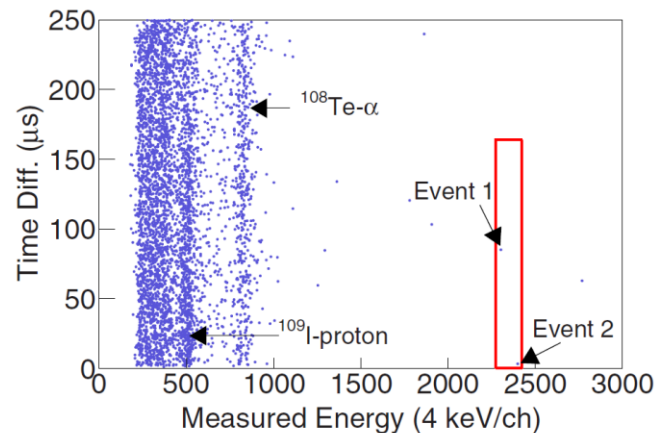
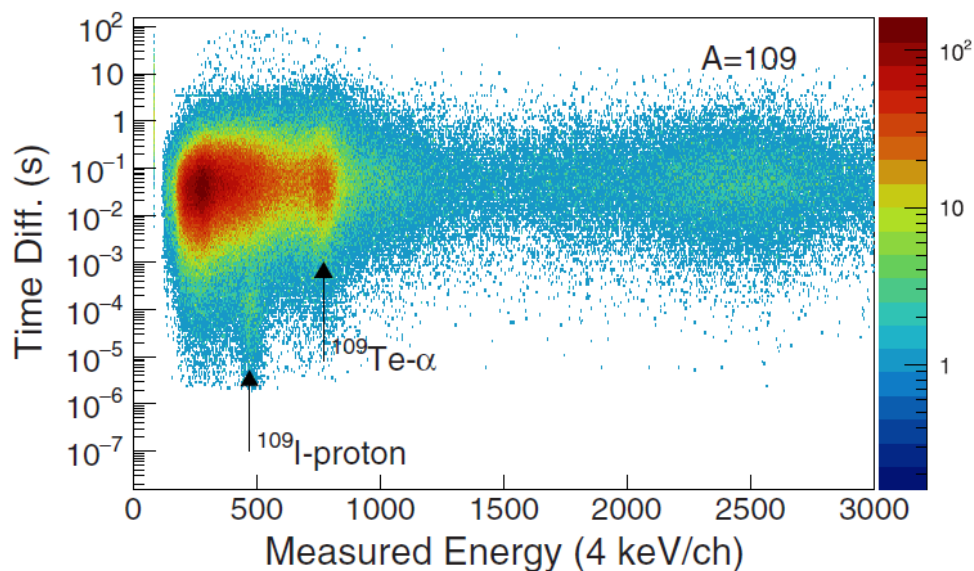
Niu Fan and Jingya Fan, Phys. Rev. C **104**, 064320 (2021)

Search for ^{108}Xe - ^{104}Te - ^{100}Sn with segmented scintillator Tandem, Tokai



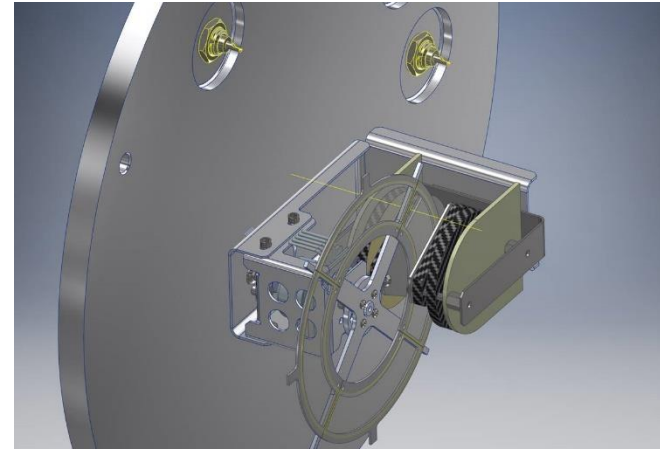
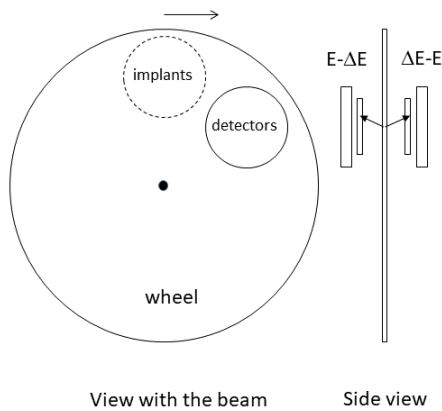
- YAP scintillator
 - faster detector
 - worse energy resolution
 - larger random background

- Two candidate events, if real:
 - $T_{1/2}(^{108}\text{Xe})=30(+57-12) \mu\text{s}$
 - $T_{1/2}(^{104}\text{Te}) < 4 \text{ ns}$, no pileup observed
 - Consistent with our paper
 - Even faster α decays?



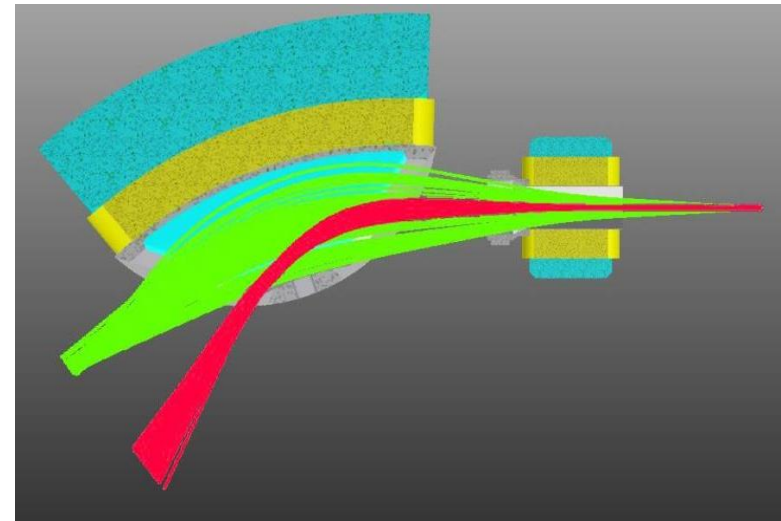
Y. Xiao et al., Phys. Rev. C **100**, 034315 (2019)

New approach - rotating stopper with AGFA



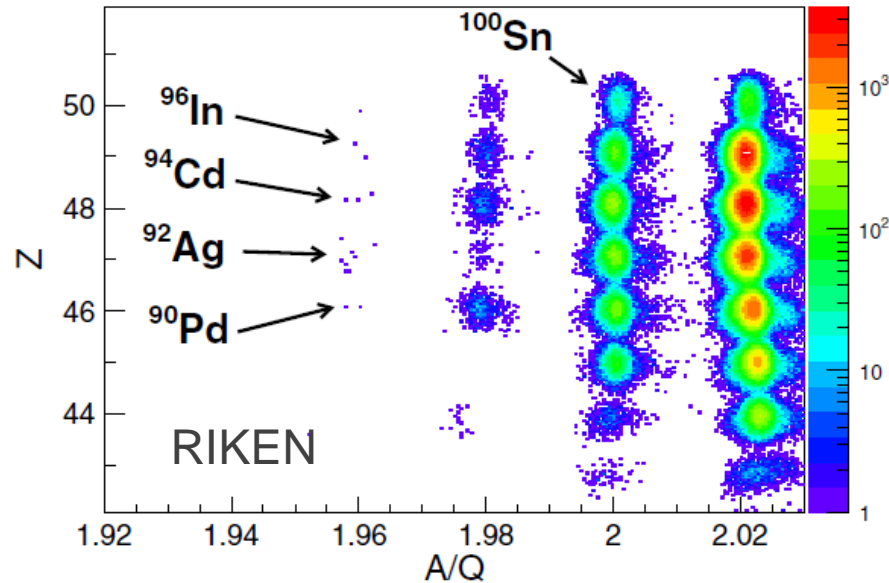
Search for triple α chains:
 ^{112}Ba - ^{108}Xe - ^{104}Te

- Higher efficiency
- Higher beam intensity
- $\Delta t \sim 1$ ns
- No mass dispersion
- Beam suppression?
- Background?



Argonne Gas-Filled Analyzer (AGFA)

Possible α -decay studies “north-east” of ^{100}Sn at fragmentation facilities



No results reported “north-east” of ^{100}Sn so far

- ^{108}Xe
- ^{112}Ba
- ^{103}Te 2p emitter candidate (*E. Olsen et al., PRL 111, 139903 (2013)*)
- ...

Summary and Outlook

- First observation of the ^{108}Xe - ^{104}Te - ^{100}Sn chain
 - Enhanced alpha preformation compared to ^{212}Po
- Microscopic description still work in progress
 - ^{104}Te , ^{108}Xe important for the role of neutron-proton interaction
- Future measurements
 - **More statistics**
 - ^{104}Te lifetime (~ 1 ns)
 - more precise α -decay widths for ^{108}Xe and other $N \sim Z$ α emitters
 - ^{112}Ba $N=Z$ α emitter

Superaligned α Decay to Doubly Magic ^{100}Sn

K. Auranen,^{1,*} D. Seweryniak,¹ M. Albers,¹ A. D. Ayangeakaa,^{1,†} S. Bottoni,^{1,‡} M. P. Carpenter,¹ C. J. Chiara,^{1,2,§} P. Copp,^{1,3} H. M. David,^{1,||} D. T. Doherty,^{4,¶} J. Harker,^{1,2} C. R. Hoffman,¹ R. V. F. Janssens,^{5,6} T. L. Khoo,¹ S. A. Kuvin,^{1,7} T. Lauritsen,¹ G. Lotay,⁸ A. M. Rogers,^{1,**} J. Sethi,^{1,2} C. Scholey,⁹ R. Talwar,¹ W. B. Walters,² P. J. Woods,⁴ and S. Zhu¹

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⁷Department of Physics, University of Connecticut, Storrs, Connecticut 06269, USA

⁸University of Surrey, Guildford GU2 7XH, United Kingdom

⁹Department of Physics, University of Jyväskylä, P.O. Box 35, FI-40014 University of Jyväskylä, Finland



Summary and Outlook

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Thank you for your attention!