

On “**Exotic**” Brane Dynamics, Gravity and Gauge Theories

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Contents

- A motivation
- Aspects of exotic branes
- GLSM
- Brane configurations
- Questions (which I would like to ask you)

Compactify 11-dim. supergravity (SUGRA) on a torus T^d

→ (11 - d)-dim maximal SUGRA has global symmetry G_0

$$G_0 = E_{d(d)} : \text{U-duality symmetry}$$

Deformation of maximal SUGRA by...

1. introducing mass parameter Romans' mass
2. gauging subgroup of G_0 embedding tensors

their String/M-theory origin ?

embedding tensors and M-theory: de Wit, Nicolai, Samtleben arXiv:0801.1294, etc.

1. Romans' mass \rightarrow D8-brane
2. embedding tensors \rightarrow domain walls

String/M-theory origins of many domain walls are not so common, but
referred to as

Exotic Branes

from standard branes via string dualities
in lower dimensions

Exotic branes should be ingredients of black hole microstructures,
as discussed in this workshop.

Exotic brane is ...

1. brane of codim $T \leq 2$

$T = 2$: defect brane (vortex), $T = 1$: domain wall (kink)

2. mass/tension $\propto 1/g_s^n$ (often with $n \geq 2$)

3. monodromy on background spacetime geometry by string dualities

| | D | domain walls | defect branes |
|-----|-----|-------------------|---------------------------------|
| ex) | 10 | D8 | D7, 7_3 |
| | 8 | D6, $6_3^{(1,1)}$ | D5, NS5, KK5, 5_2^2 , 5_3^2 |

Nomenclature $b_n^{(d,c)}$:

$$\text{mass} = \frac{R_1 R_2 \cdots R_b (R_{b+1} \cdots R_{b+c})^2 (R_{b+c+1} \cdots R_{b+c+d})^3}{g_s^n \ell_s^{b+2c+3d+1}}$$

String dualities

$$T_y : R_y \rightarrow \frac{\ell_s^2}{R_y}, \quad g_s \rightarrow \frac{\ell_s}{R_y} g_s$$

$$S : g_s \rightarrow \frac{1}{g_s}, \quad \ell_s \rightarrow g_s^{1/2} \ell_s$$

$$D5(12345) \xrightarrow{S} NS5(12345) \xrightarrow{T_9} KK5(12345,9) \xrightarrow{T_8} 5_2^2(12345,89) \xrightarrow{S} 5_3^2(12345,89)$$

5_1 5_2 5_2^1

SUGRA solutions corresponding to exotic branes are not well-defined :

1. asymptotically non-flat logarithmic ($T = 2$), linear ($T = 1$)
2. multi-valued monodromy by string dualities

SUGRA solution of exotic 5_2^2 -brane (string frame)

$$ds^2 = dx_{012345}^2 + H \{ (d\rho)^2 + \rho^2 (d\vartheta)^2 \} + \frac{H}{K} \{ (dx^8)^2 + (dx^9)^2 \}$$

$$B_{89} = -\frac{\vartheta}{K}, \quad e^{2\phi} = \frac{H}{K}$$

$$H = 1 + \log \frac{\Lambda}{\rho}, \quad K = H^2 + \vartheta^2$$

Brane charge is not given in a conventional way.

defect branes: Bergshoeff, Ortín, Riccioni arXiv:1109.4484

various setup: de Boer, Shigemori arXiv:1209.6056

Alice strings: Okada, Sakatani arXiv:1411.1043

Exotic branes are subject to string dualities.

It is important to study them in string theory beyond SUGRA.

2D Gauged Linear Sigma Model (GLSM)

Witten hep-th/9301042; Hori, Vafa hep-th/0002222

We studied 2D $\mathcal{N} = (4, 4)$ GLSM

TK, Sasaki arXiv:1304.4061

whose IR effective theory is string sigma model on exotic 5_2^2 -brane geometry.

This is a generalization of Tong's GLSM for ALF space.

ALE : Douglas, Moore hep-th/9603167; Johnson, Myers hep-th/9610140

ALF : Tong hep-th/0204186; Harvey, Jensen hep-th/0507204; Okuyama hep-th/0508097

GLSM for five-branes

- 2D $\mathcal{N} = (4, 4)$ gauge theory
with charged/adjoint hypermultiplets
- Higgs branch moduli space
= background geometry of five-branes
- Those for NS5-branes/KK5-branes (codim 3)
are well known
- Those for exotic five-branes (codim ≤ 2) are not so

... skip details

T-duality is realized by Roček-Verlinde (Hori-Vafa) transformation,
and its generalization :

$$\begin{aligned}
 & -\frac{1}{g^2} \int d^4\theta |\Theta|^2 - \left(\int d^2\tilde{\theta} \Theta \Sigma + (\text{h.c.}) \right) \rightarrow +g^2 \int d^4\theta (\Gamma + \bar{\Gamma} + V)^2 \\
 & +\frac{1}{g^2} \int d^4\theta |\Psi|^2 - \left(\int d^2\theta \Psi \Phi + (\text{h.c.}) \right) \rightarrow -g^2 \int d^4\theta \left\{ (\Xi + \bar{\Xi} - (C + \bar{C}))^2 - (\Psi - \bar{\Psi})(C - \bar{C}) \right\} \\
 & \hspace{25em} (\Phi = \bar{D}_+ \bar{D}_- C)
 \end{aligned}$$

The second transformation is performed without shift symmetry $\Psi \rightarrow \Psi + i\alpha$

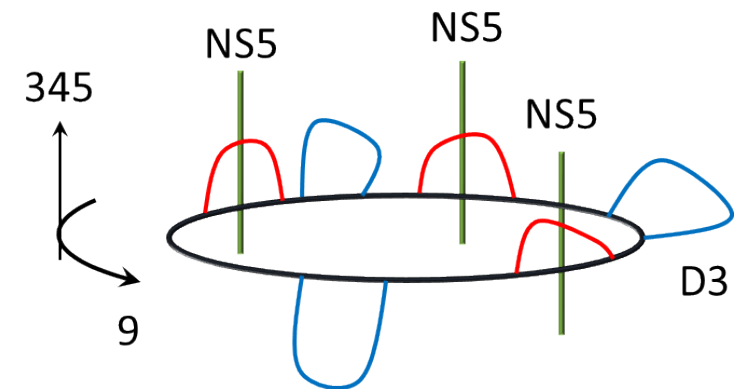
co-existence of the scalar parts of $\text{Im } \Xi$ and $\text{Im } \Psi$
implies the **doubled space** by T-duality

GLSM is also powerful when we embed the system into string theory.

Brane configurations

Tong's 2D GLSM is realized by intersection of k NS5 and a D3 wrapped on circles :

| type IIB | 0 | 1 | ② | ③ | ④ | ⑤ | 6 | 7 | 8 | ⑨ |
|----------|-------|---|---------|---|---|-------|---|---|---|---|
| k NS5 | — | — | — | — | — | — | | | | |
| D3 | — | — | — | | | | | | | — |
| | gauge | | Coulomb | | | Higgs | | | | |

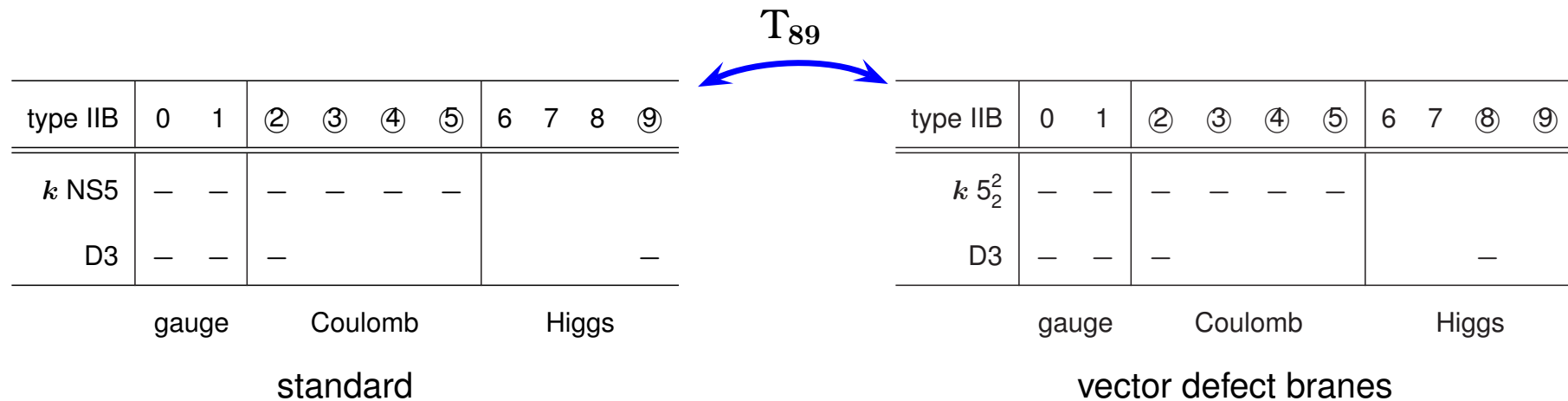


gauged flavor $U(1)_F$ as fluctuation on NS5-branes

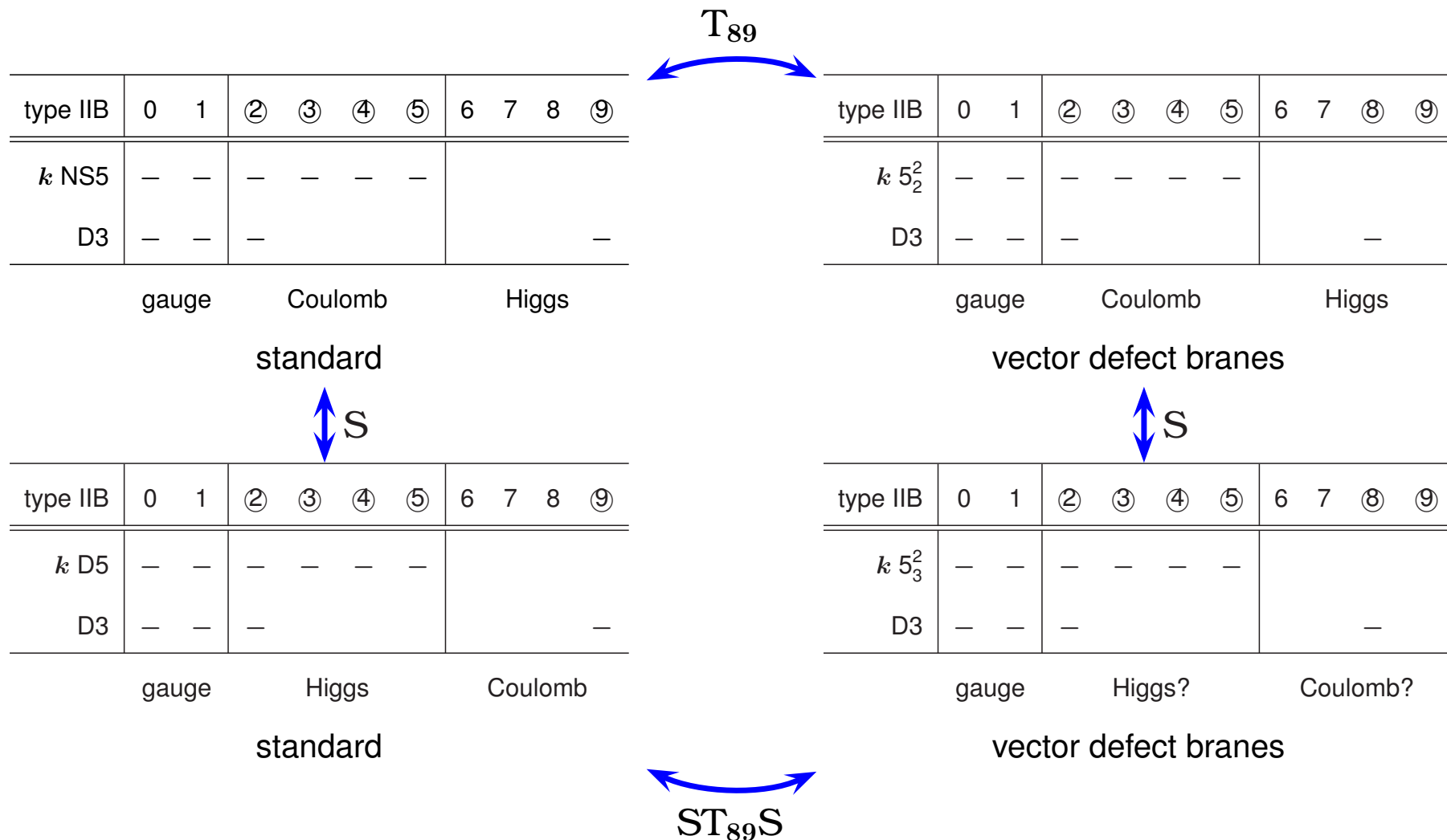
→ type IIB setup (6D $\mathcal{N} = (1, 1)$ vector) ~~type IIA setup (6D $\mathcal{N} = (2, 0)$ tensor)~~

2345-directions are compactified

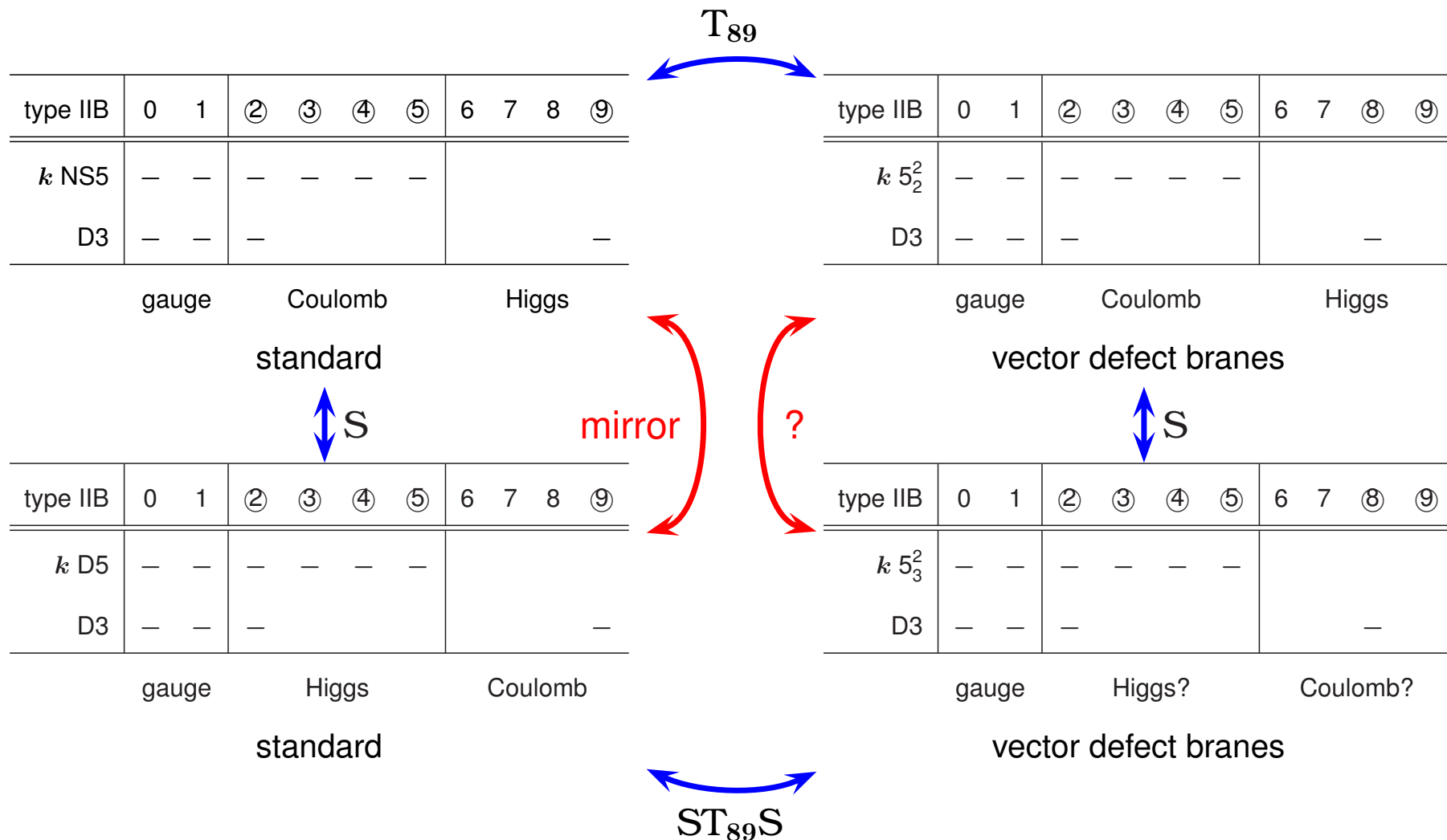
Once the GLSM is embedded into brane configurations,
we can apply string dualities to it and investigate various aspects of gauge theories :



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| type IIB | 0 | 1 | ② | ③ | ④ | ⑤ | 6 | 7 | ⑧ | ⑨ |
|-----------|---|---|---|---|---|---|---|---|---|---|
| $k 5_2^2$ | — | — | — | — | — | — | | | | |
| D3 | — | — | — | | | | | | — | |

gauge Coulomb Higgs

vector defect branes

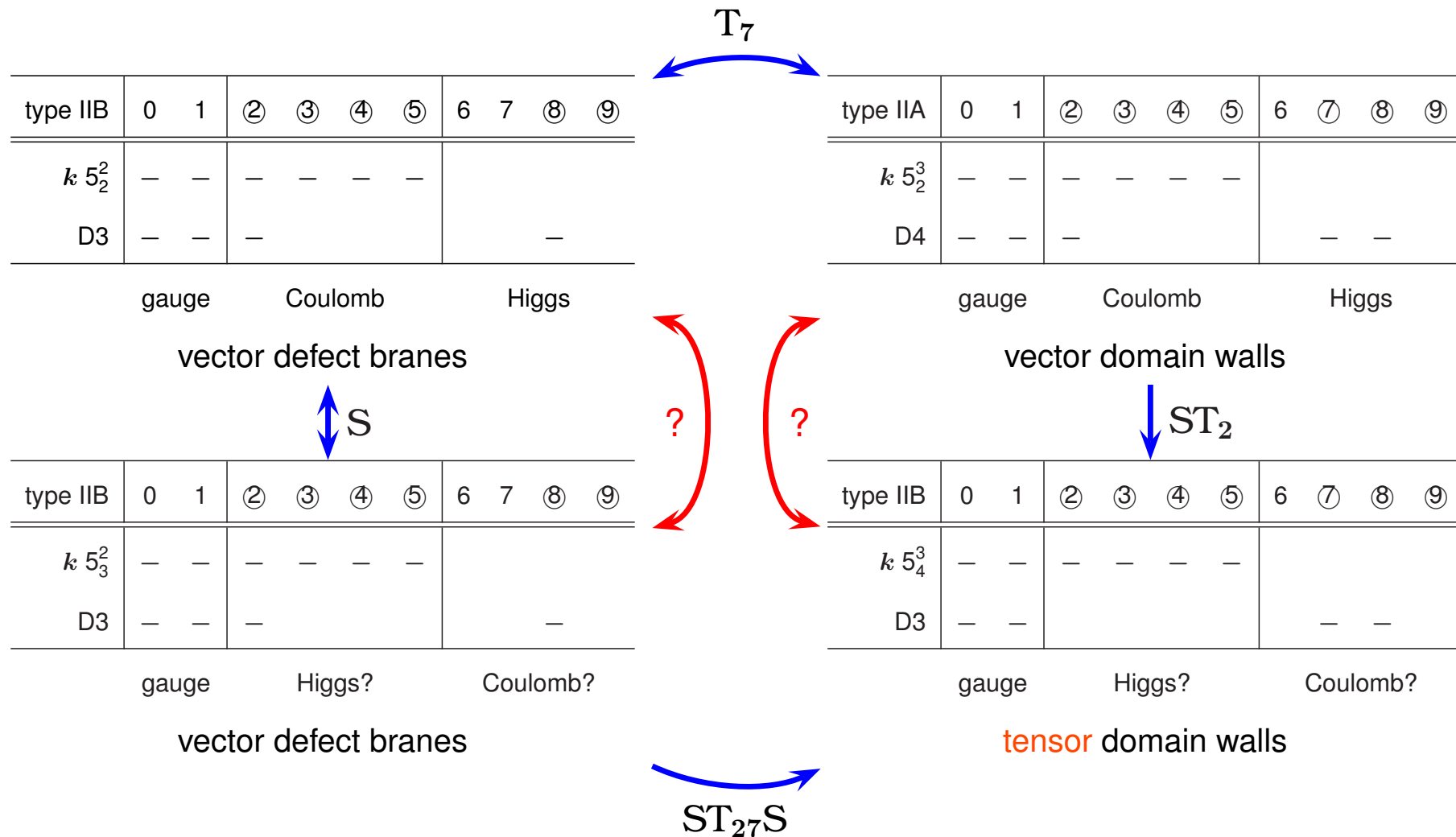


| type IIB | 0 | 1 | ② | ③ | ④ | ⑤ | 6 | 7 | ⑧ | ⑨ |
|-----------|---|---|---|---|---|---|---|---|---|---|
| $k 5_3^2$ | — | — | — | — | — | — | | | | |
| D3 | — | — | — | | | | | | — | |

gauge Higgs? Coulomb?

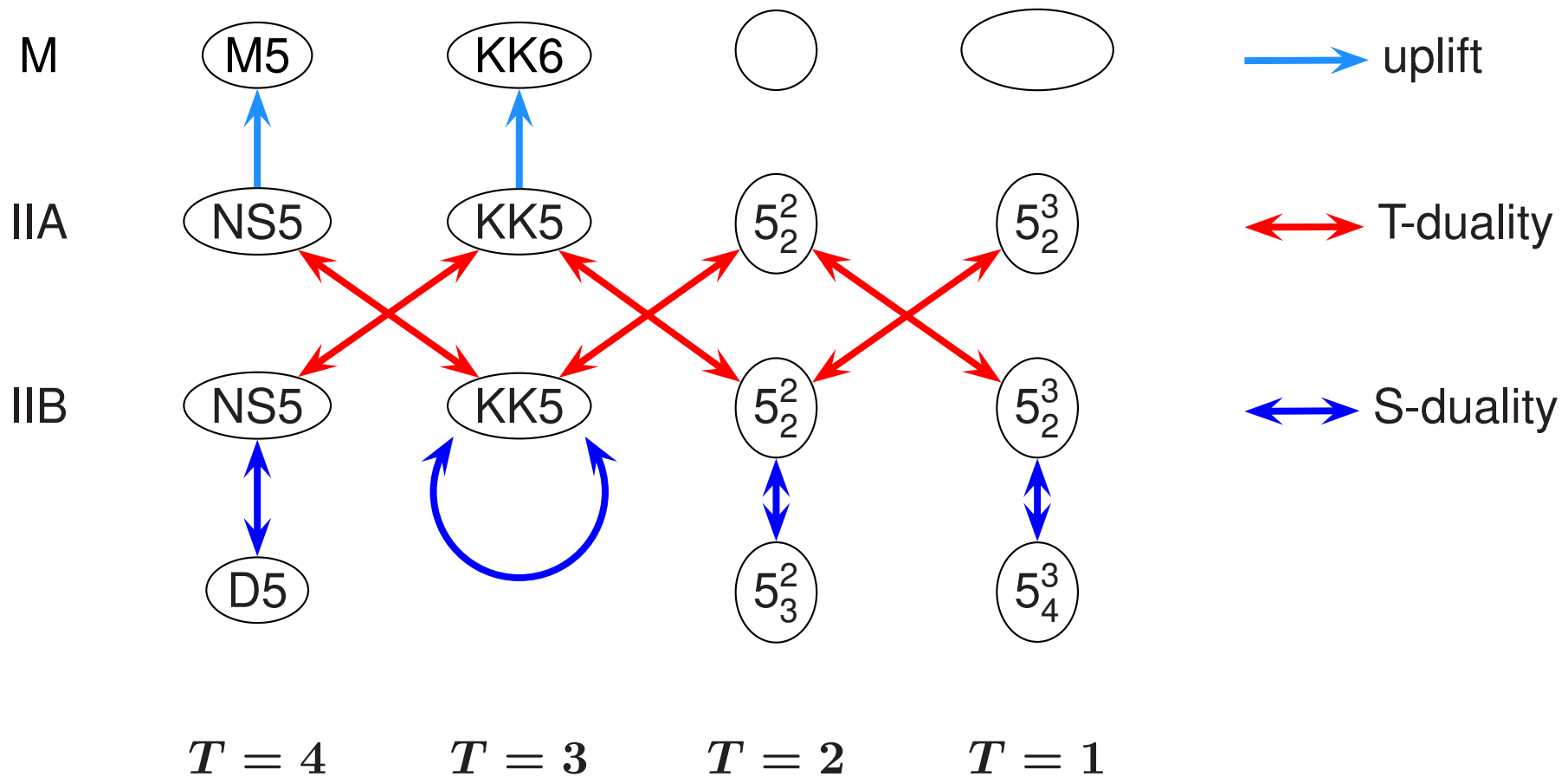
vector defect branes

Once the GLSM is embedded into brane configurations, we can apply string dualities to it and investigate various aspects of gauge theories :



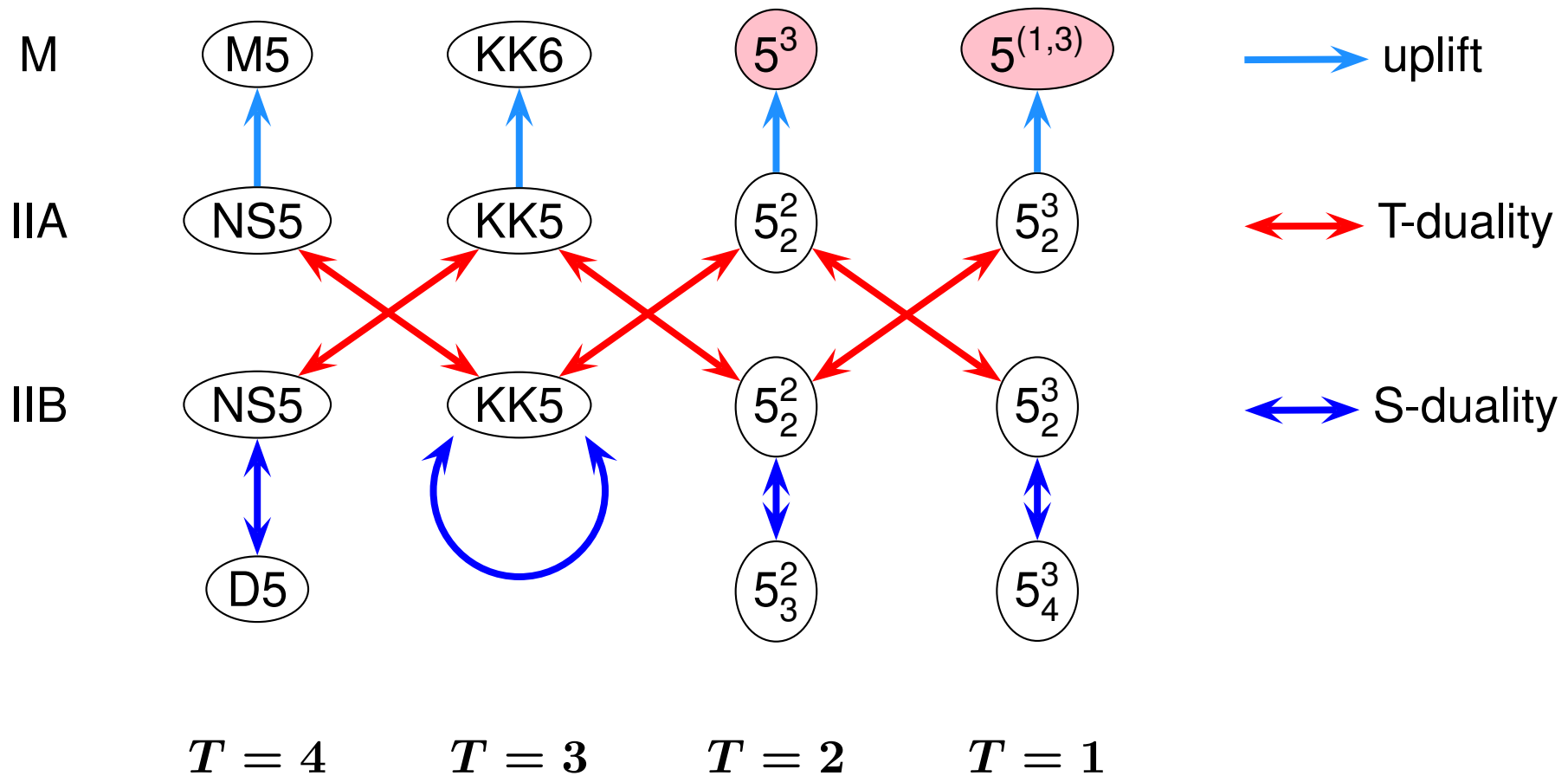
Questions which I would like to ask you are ...

1. Uplift **exotic** branes to M-theory
2. **Exotic** brane ending on **exotic** brane
3. Hanany-Witten effect across **exotic** branes



How can we use exotic M-branes (and other exotic branes)?

Worldvolume action : TK, Sasaki, Yata arXiv:1404.5442
 TK, Sasaki to appear

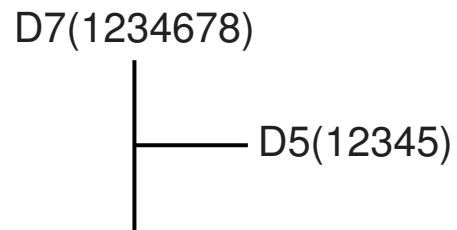


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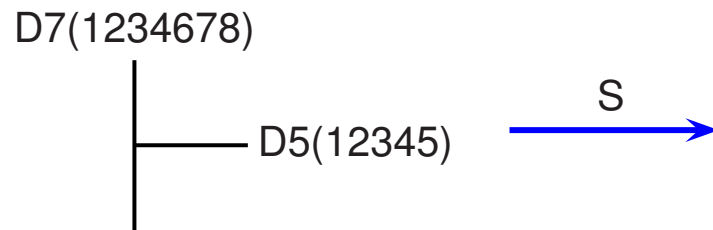
We start from D5-D7 system such as

| | 0 | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | 9 |
|----|---|---|---|---|---|---|---|---|---|---|
| D7 | - | - | - | - | - | | - | - | - | |
| D5 | - | - | - | - | - | - | | | | |



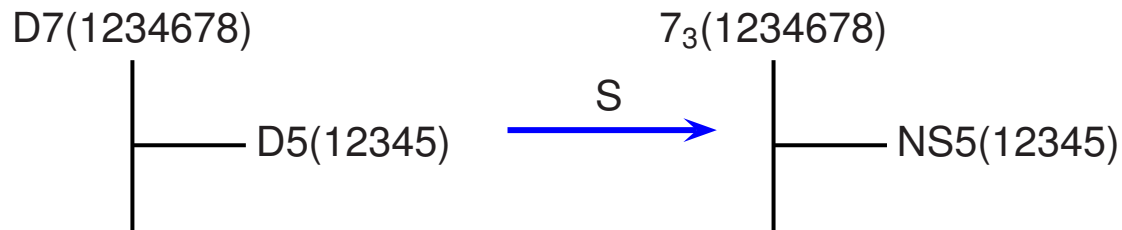
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|----|---|---|---|---|---|---|---|---|---|---|
| D7 | - | - | - | - | - | | - | - | - | |
| D5 | - | - | - | - | - | - | | | | |



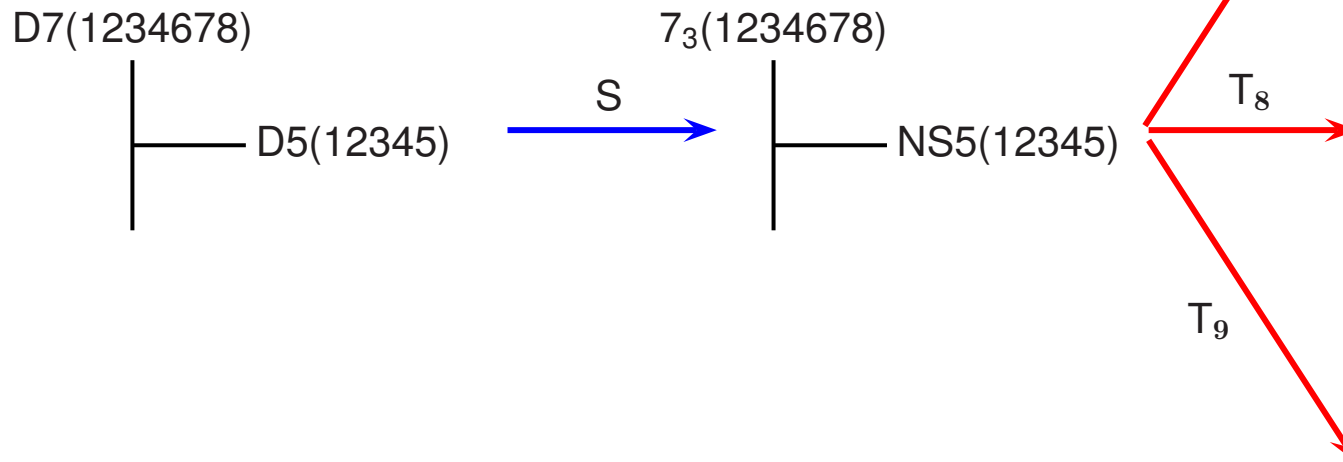
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|----|---|---|---|---|---|---|---|---|---|---|
| D7 | - | - | - | - | - | | - | - | - | |
| D5 | - | - | - | - | - | - | | | | |



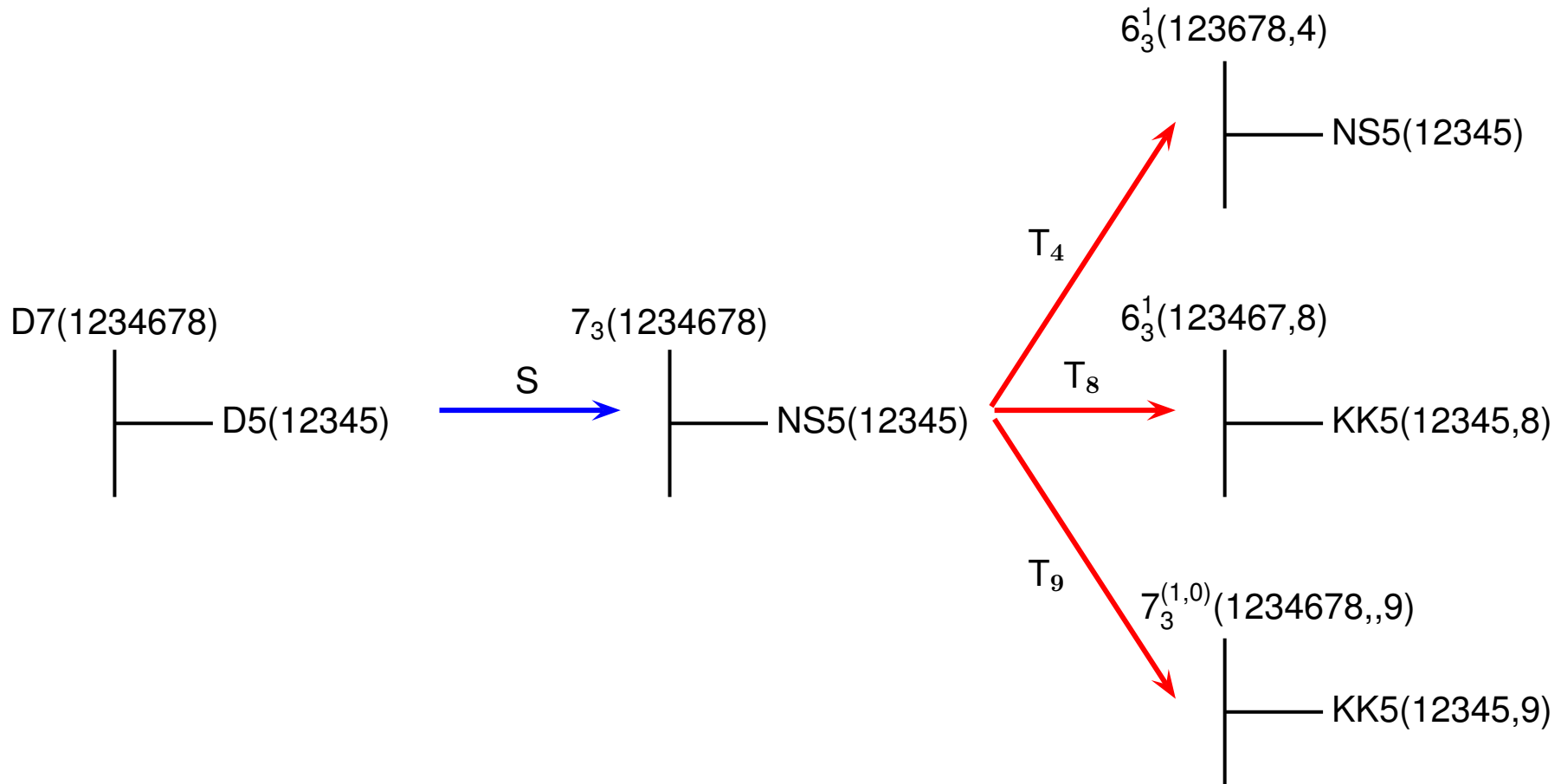
We start from D5-D7 system such as

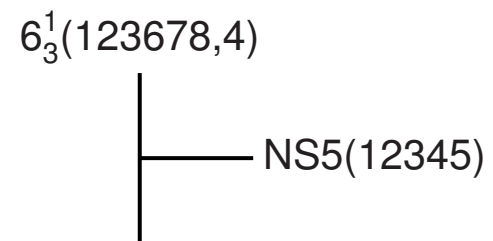
| | 0 | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | 9 |
|----|---|---|---|---|---|---|---|---|---|---|
| D7 | - | - | - | - | - | | - | - | - | |
| D5 | - | - | - | - | - | - | | | | |

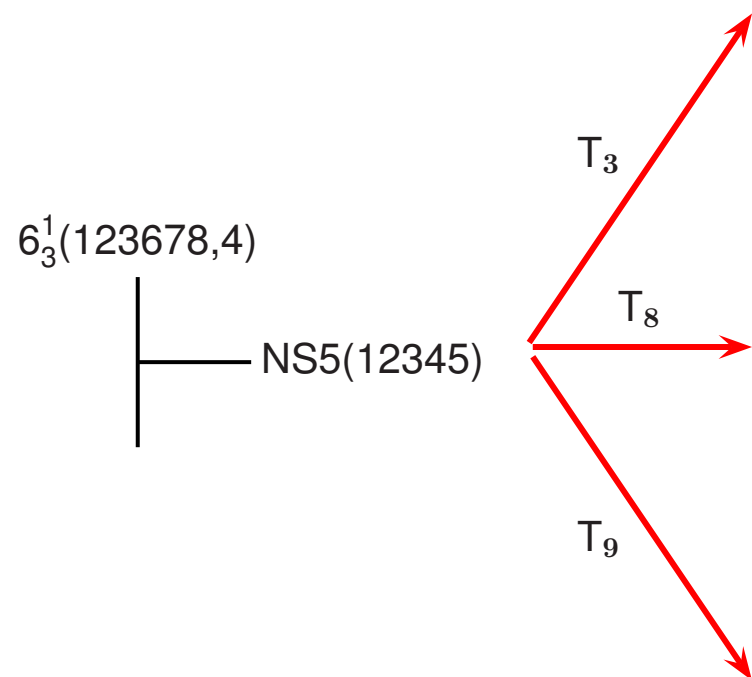


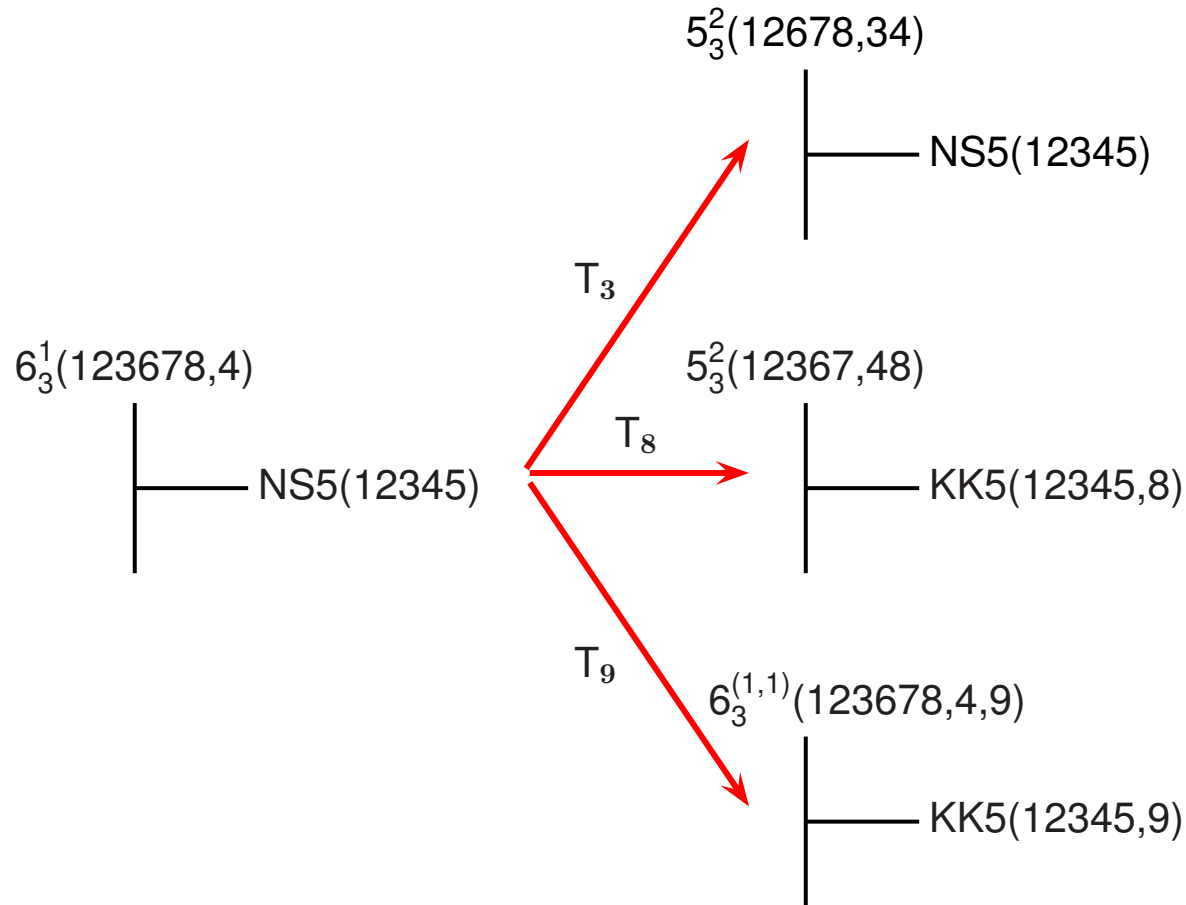
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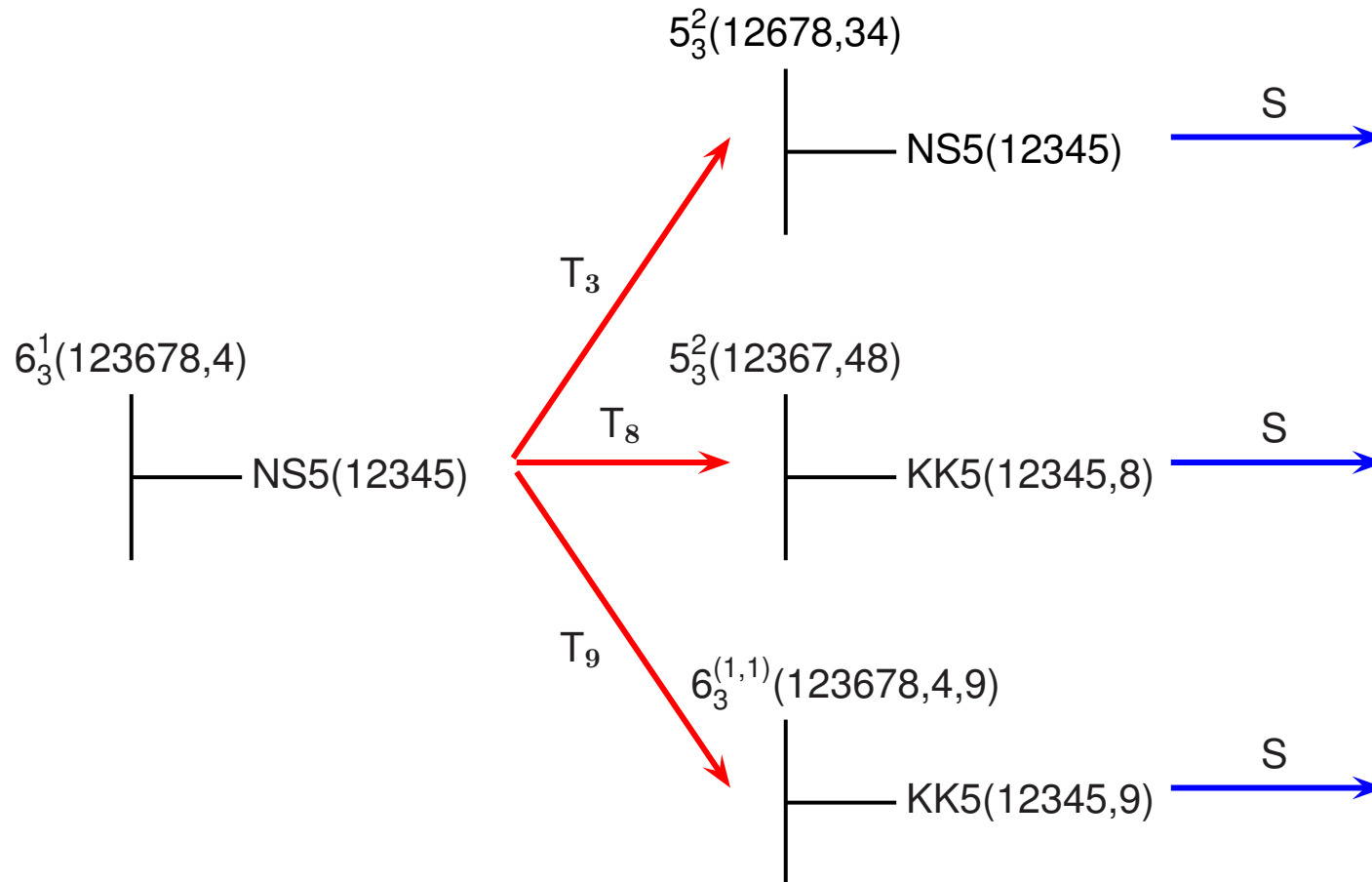
| | 0 | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | 9 |
|----|---|---|---|---|---|---|---|---|---|---|
| D7 | - | - | - | - | - | | - | - | - | |
| D5 | - | - | - | - | - | - | | | | |

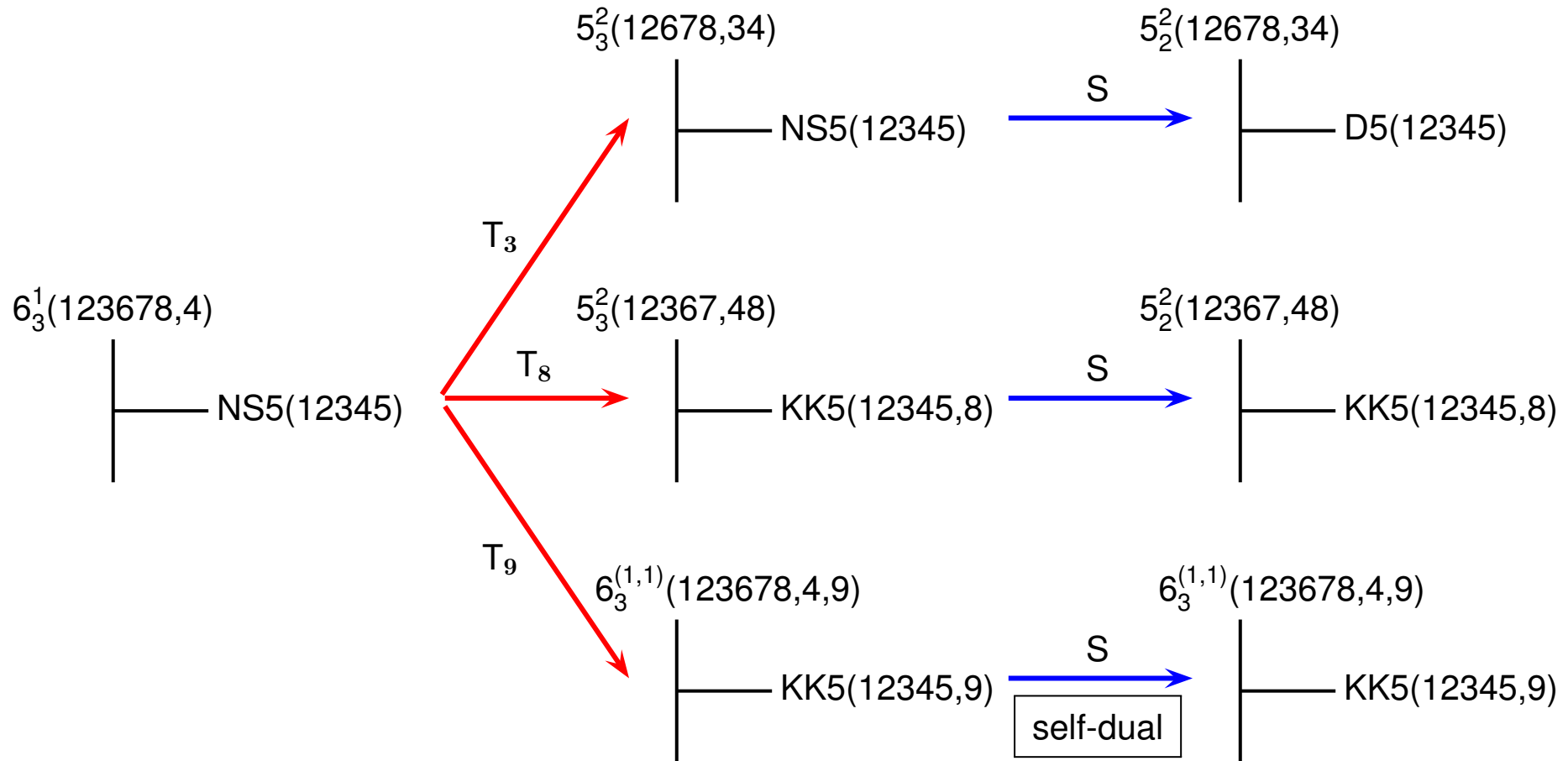




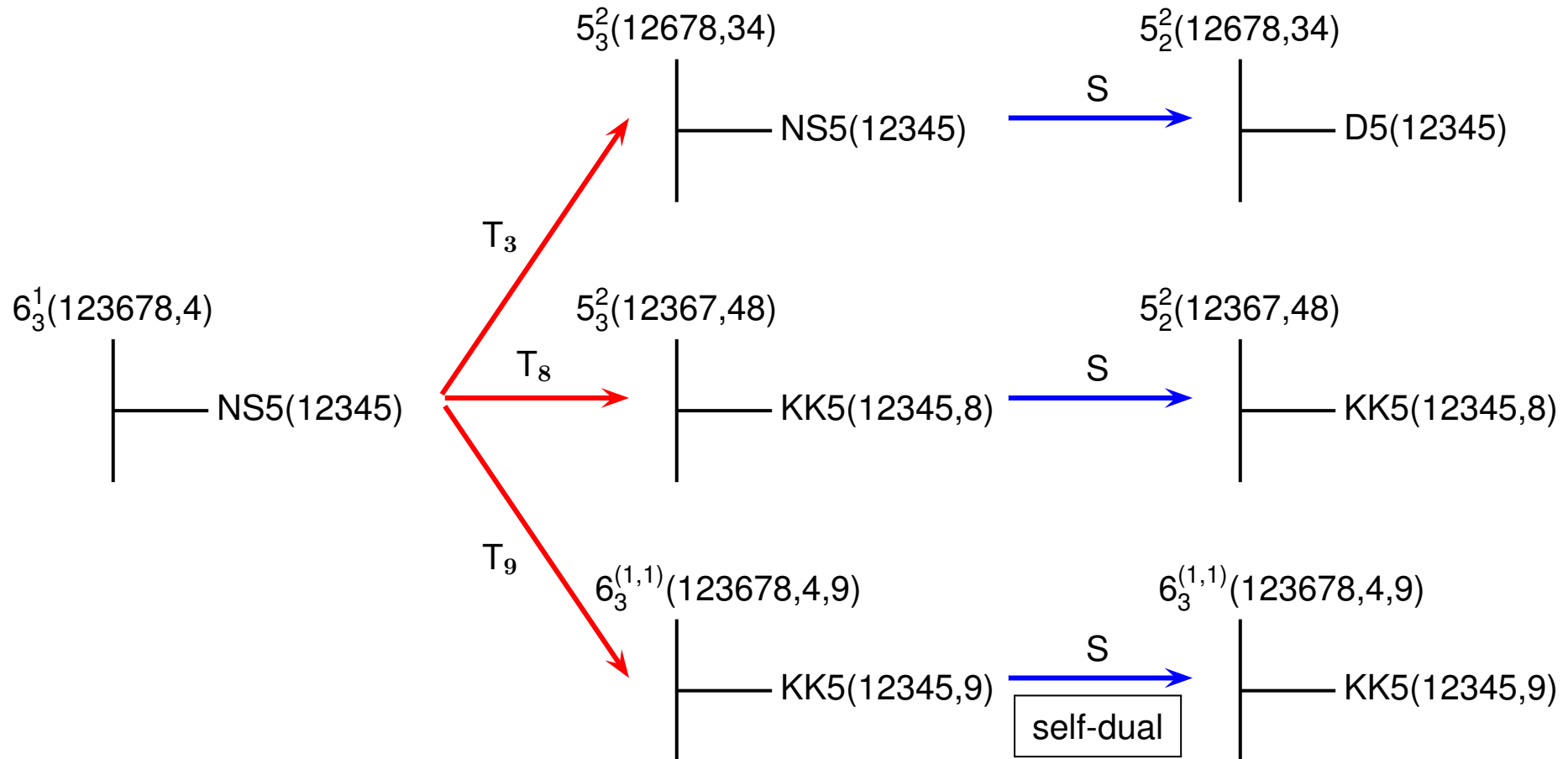




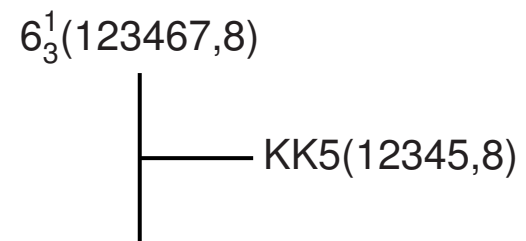


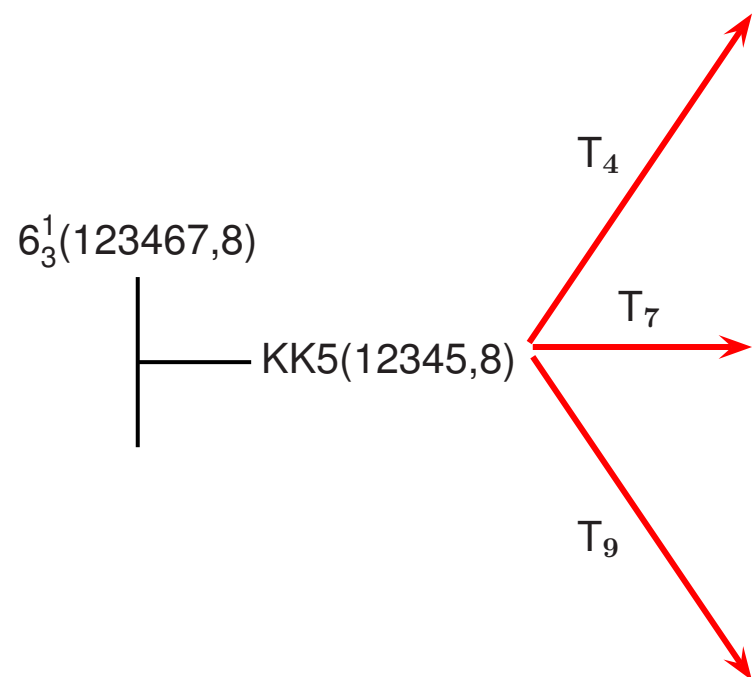


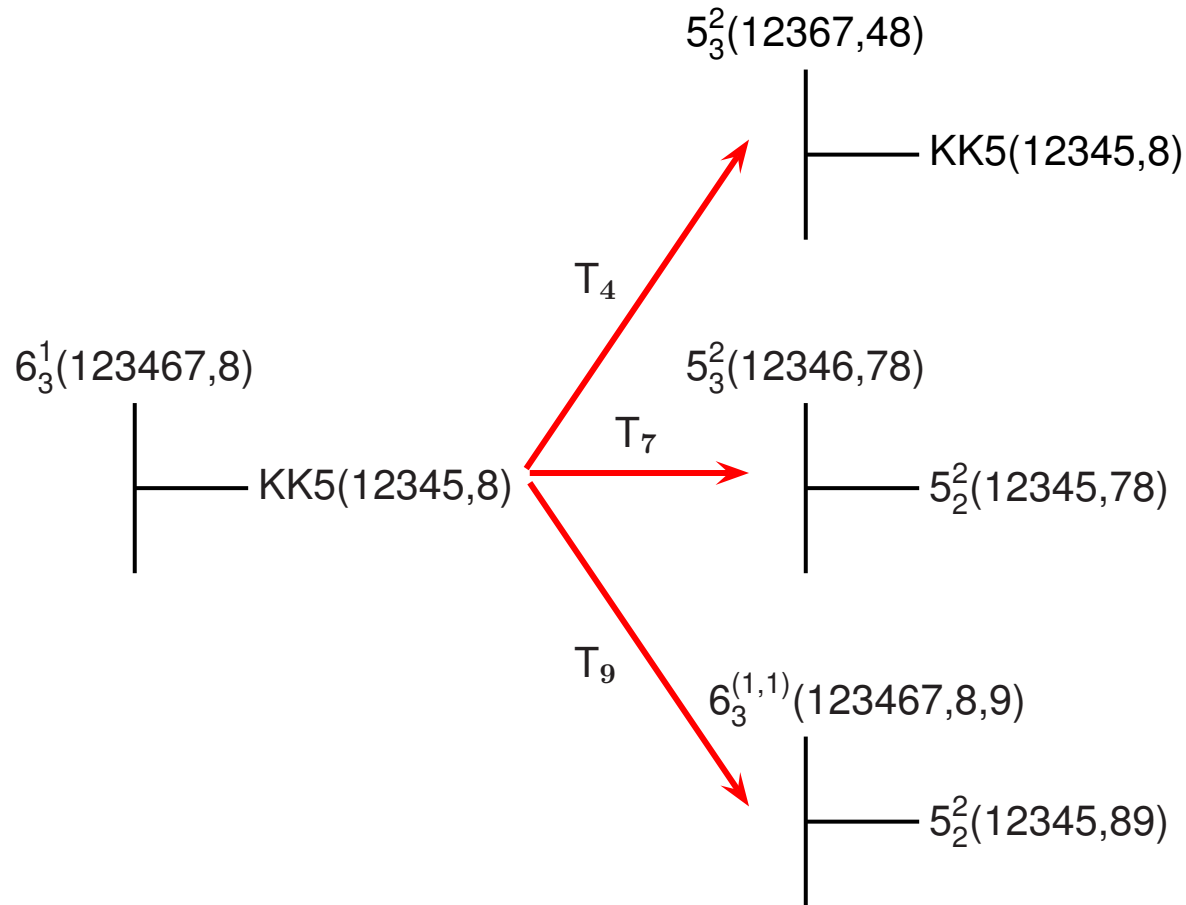
Which modes are localized at intersection?

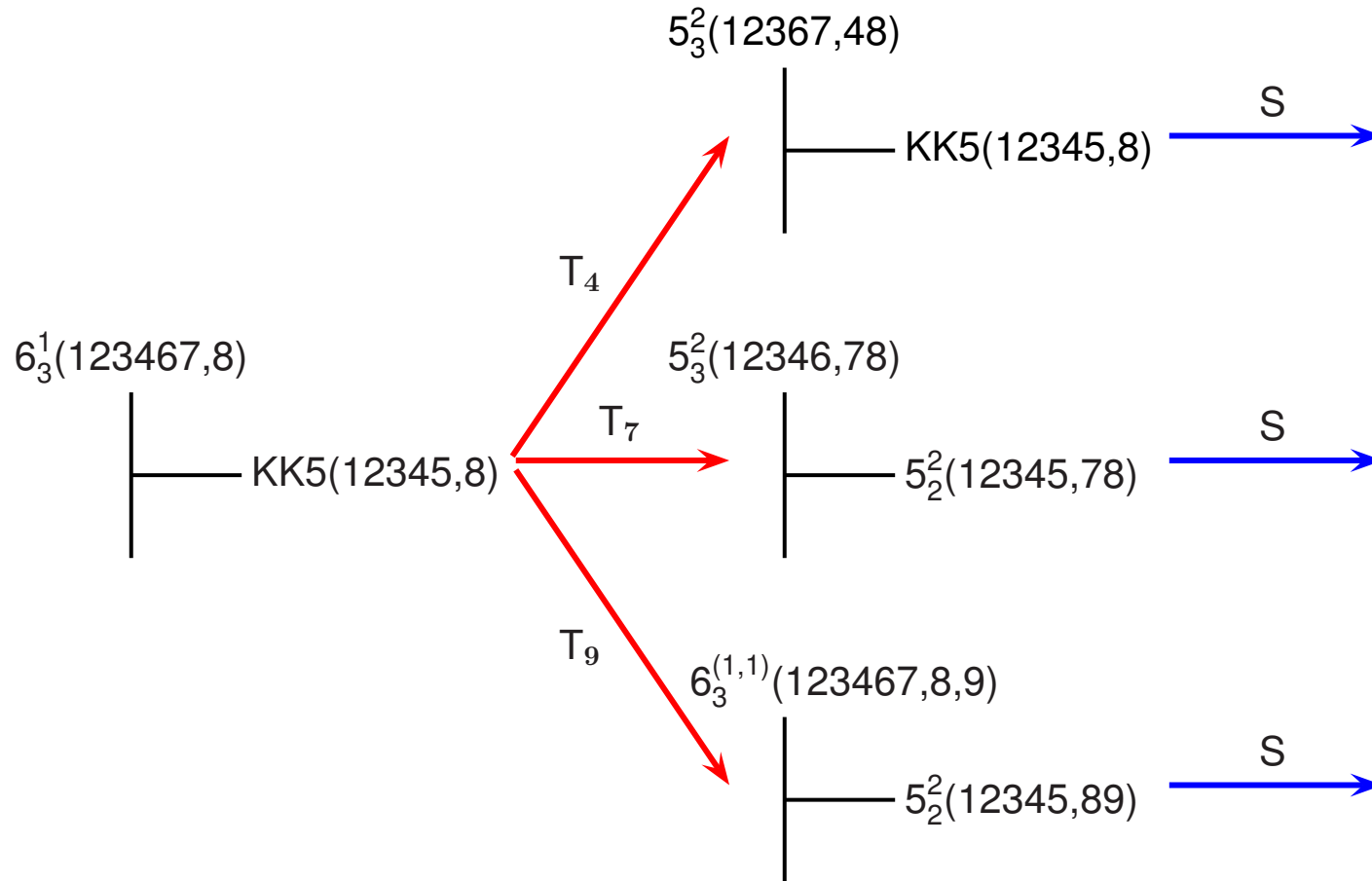


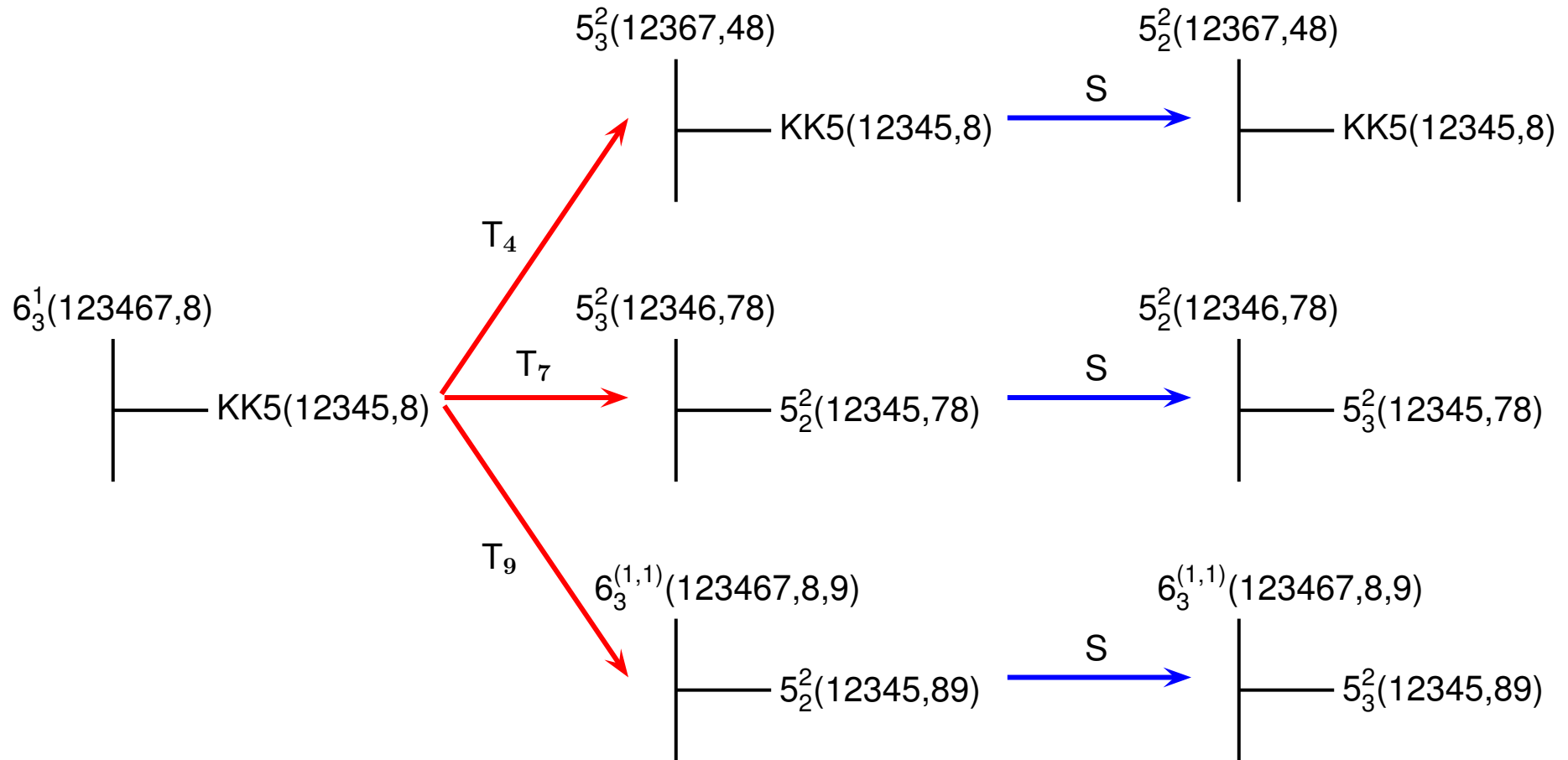
common 0123-directions: 4D exotic brane world scenario?



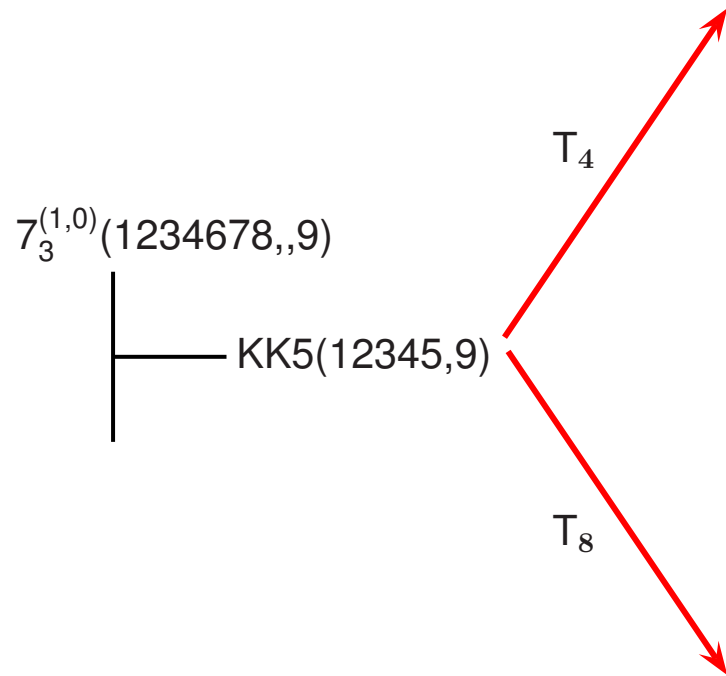


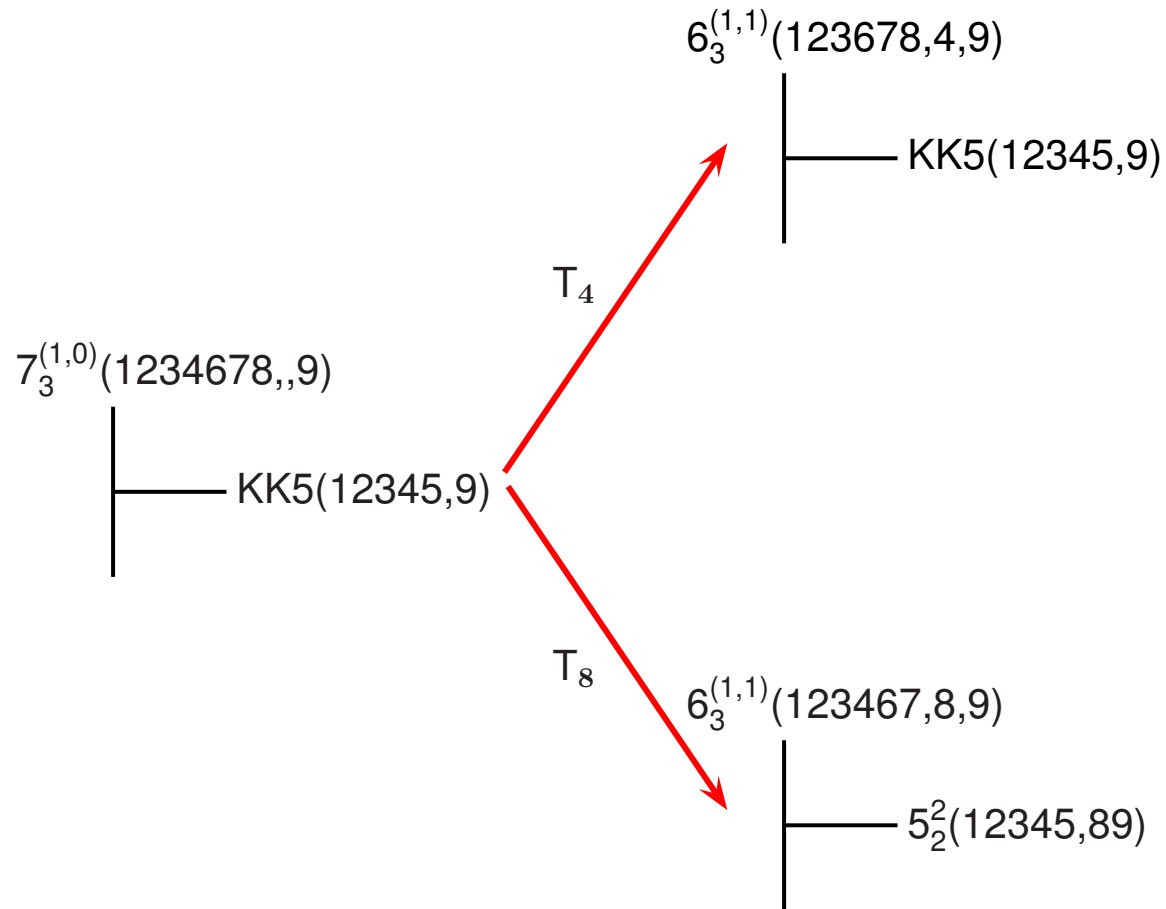


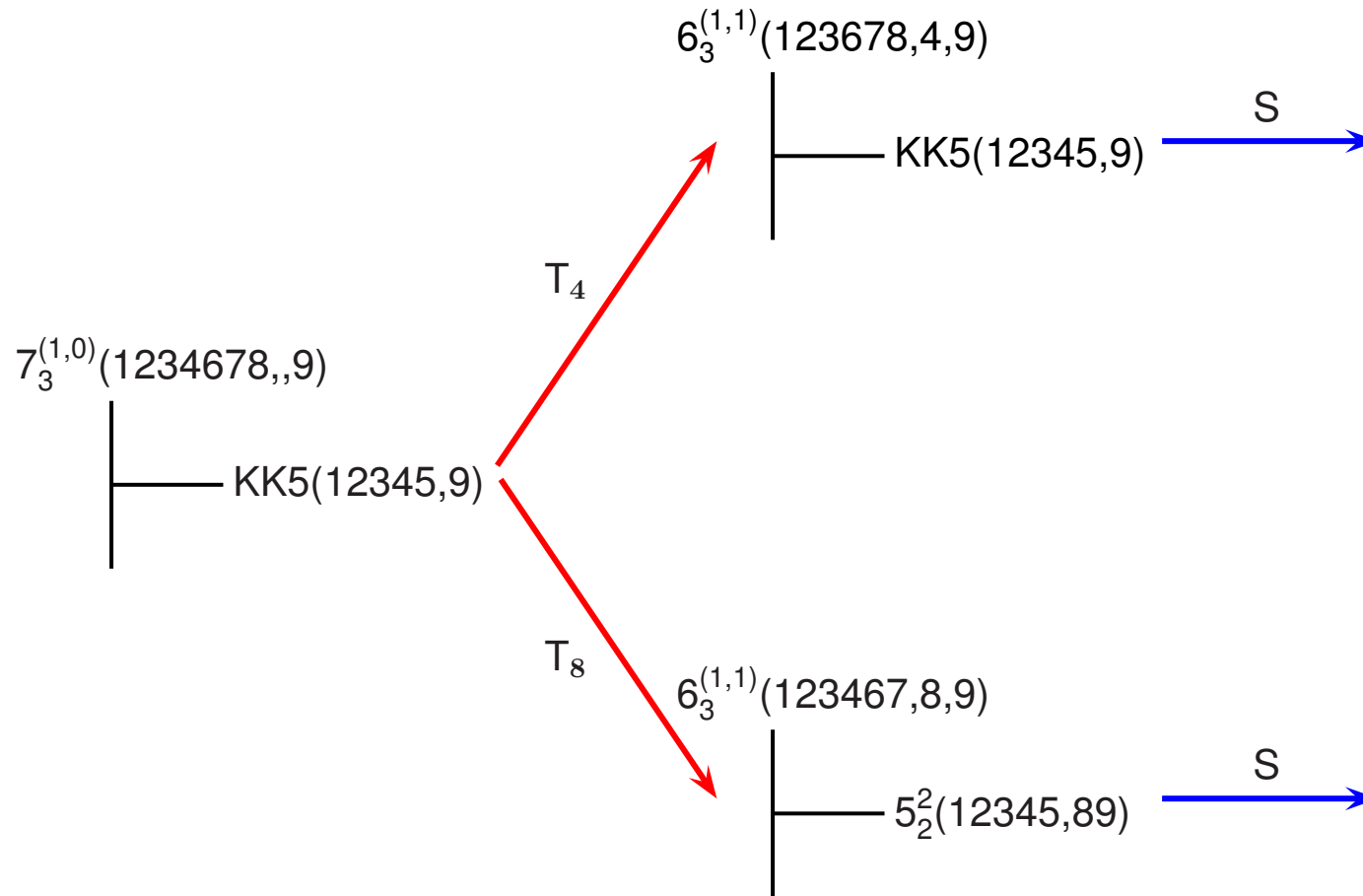


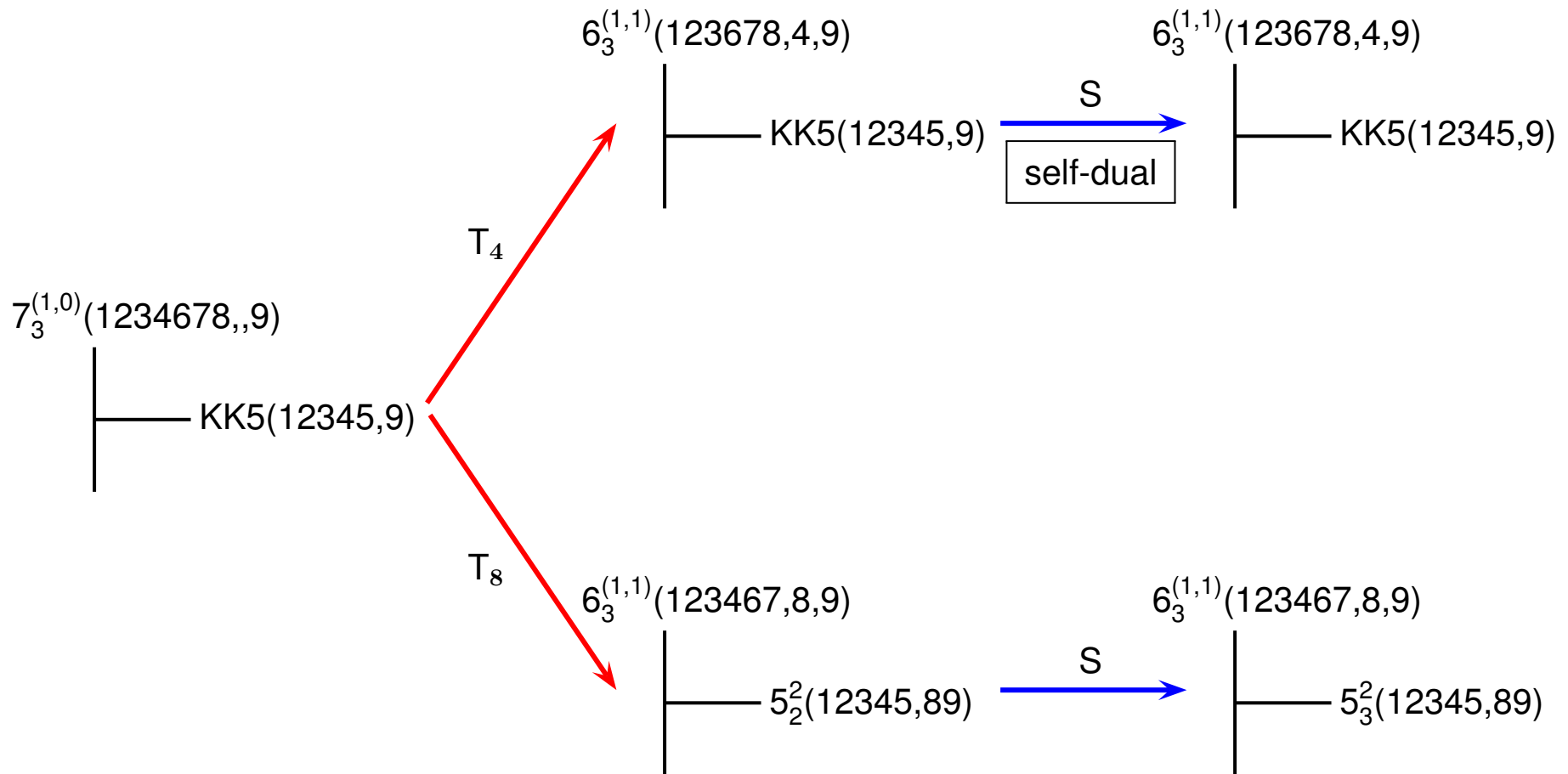


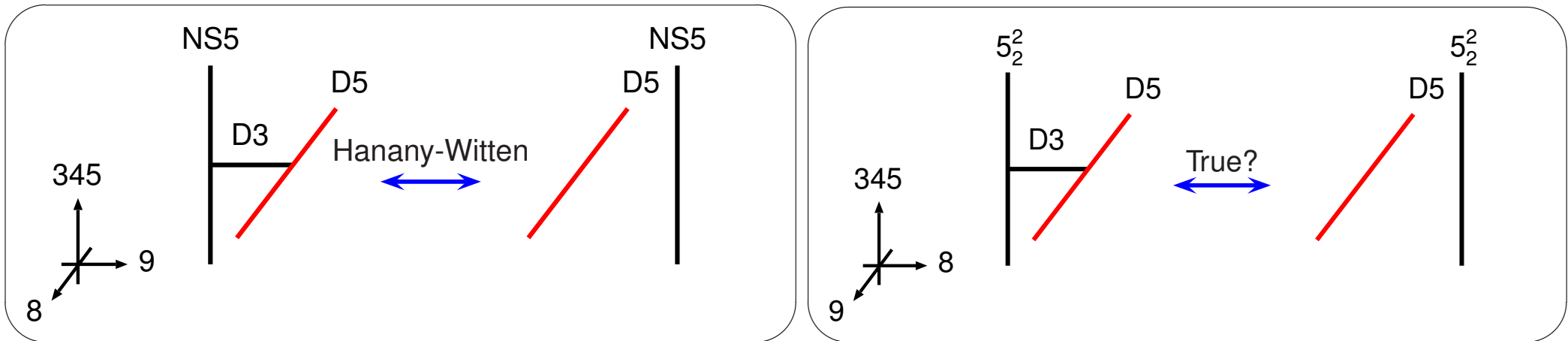
$$\begin{array}{c} 7_3^{(1,0)}(1234678,,9) \\ | \\ \text{---} \text{KK5}(12345,9) \end{array}$$











Linking number is defined as a sum of brane charges :

$$L_B = \frac{1}{2}(r - l) + L - R$$

However, can we define meaningful Dp -brane charges in the 5_2^2 background?

branch cut as D7 background?

instructive works : de Boer, Shigemori arXiv:1209.6056
Okada, Sakatani arXiv:1411.1043

Summary

- ✓ Exotic brane configurations will give us new examples of mirror symmetry between two gauge theories.

- ✓ Exotic brane configurations will tell us novel features of gravity and gauge theories if we understand the followings :
 1. exotic M-branes
 2. exotic brane ending on exotic brane
 3. Hanany-Witten effect in the presence of exotic brane

Orientifold planes are also reconsidered in exotic backgrounds.

- ✓ Exotic branes should be ingredients of BH microstructures (we hope).

Thanks

Appendix

| D | fundamental $n = 0$ | Dirichlet $n = 1$ | solitonic $n = 2$ | S_D -dual of (Dirichlet) $n = 3$ | S_D -dual of (fundamental) $n = 4$ |
|-----|------------------------|----------------------|-------------------------------------------------------------|---------------------------------------|-----------------------------------------|
| IIB | | D7 [C_8] | | 7_3 [$E_8 = S_{10}(C_8)$] | |
| 9 | | D6 [C_7] | | 6_3^1 [$E_{8,1} = S_9(C_7)$] | |
| 8 | | D5 [C_6] | NS5 [D_6] KK5 [$D_{7,1}$] 5_2^2 [$D_{8,2}$] | 5_3^2 [$E_{8,2} = S_8(C_7)$] | |
| 7 | | D4 [C_5] | | 4_3^3 [$E_{8,3} = S_7(C_5)$] | |
| 6 | | D3 [C_4] | | 3_3^4 [$E_{8,4} = S_6(C_4)$] | |
| 5 | | D2 [C_3] | | 2_3^5 [$E_{8,5} = S_5(C_3)$] | |
| 4 | F1 [B_2] | D1 [C_2] | | 1_3^6 [$E_{8,6} = S_4(C_2)$] | 1_4^6 [$F_{8,6} = S_4(B_2)$] |
| 3 | P | D0 [C_1] | | 0_3^7 [$E_{8,7} = S_3(C_1)$] | $0_4^{(6,1)}$ [$F_{8,7,1}$] |

E.A. Bergshoeff et al arXiv:1009.4657, 1102.0934, 1108.5067, etc.

| | fundamental | Dirichlet | solitonic | | |
|-----|--------------|--------------|-------------------------------------------------------------|-------------------------------|-------------------------------|
| D | $n = 0$ | $n = 1$ | $n = 2$ | $n = 3$ | $n = 4$ |
| IIA | | D8 [C_9] | | | |
| 9 | | D7 [C_8] | | $7_3^{(1,0)}$ [$E_{9,1,1}$] | |
| 8 | | D6 [C_7] | | $6_3^{(1,1)}$ [$E_{9,2,1}$] | |
| 7 | | D5 [C_6] | NS5 [D_6] KK5 [$D_{7,1}$] 5_2^2 [$D_{8,2}$] | $5_3^{(1,2)}$ [$E_{9,3,1}$] | 5_4^3 [$F_{9,3}$] |
| 6 | | D4 [C_5] | | $4_3^{(1,3)}$ [$E_{9,4,1}$] | $4_4^{(1,3)}$ [$F_{9,4,1}$] |
| 5 | | D3 [C_4] | | $3_3^{(1,4)}$ [$E_{9,5,1}$] | $3_4^{(2,3)}$ [$F_{9,5,2}$] |
| 4 | | D2 [C_3] | | $2_3^{(1,5)}$ [$E_{9,6,1}$] | $2_4^{(3,3)}$ [$F_{9,6,3}$] |
| 3 | F1 [B_2] | D1 [C_2] | | $1_3^{(1,6)}$ [$E_{9,7,1}$] | $1_4^{(4,3)}$ [$F_{9,7,4}$] |

E.A. Bergshoeff et al arXiv:1009.4657, 1102.0934, 1108.5067, etc.

$$\begin{aligned}
\mathcal{L}_{\text{NS5}} = & \sum_{a=1}^k \int d^4\theta \frac{1}{e_a^2} \left\{ -|\Sigma_a|^2 + |\Phi_a|^2 \right\} \\
& + \sum_a \int d^4\theta \left\{ |Q_a|^2 e^{V_a} + |\tilde{Q}_a|^2 e^{-V_a} \right\} + \sum_a \left[\int d^2\theta \tilde{Q}_a \Phi_a Q_a + (\text{h.c.}) \right] \\
& + \int d^4\theta \frac{1}{g^2} \left\{ -|\Theta|^2 + |\Psi|^2 \right\} \\
& + \sum_a \left[\int d^2\theta (s_a - \Psi) \Phi_a + \int d^2\tilde{\theta} (t_a - \Theta) \Sigma_a + (\text{h.c.}) \right]
\end{aligned}$$

(Σ_a, Φ_a) : $\mathcal{N} = (4, 4)$ vector multiplets

(Q_a, \tilde{Q}_a) : $\mathcal{N} = (4, 4)$ charged hypermultiplets

(Ψ, Θ) : $\mathcal{N} = (4, 4)$ adjoint hypermultiplet

(s_a, t_a) : FI parameters, positions of NS5-branes

$$\begin{aligned}
\mathcal{L}_{\text{KK5}} = & \sum_{a=1}^k \int d^4\theta \frac{1}{e_a^2} \left\{ -|\Sigma_a|^2 + |\Phi_a|^2 \right\} \\
& + \sum_a \int d^4\theta \left\{ |Q_a|^2 e^{V_a} + |\tilde{Q}_a|^2 e^{-V_a} \right\} + \sum_a \left[\int d^2\theta \tilde{Q}_a \Phi_a Q_a + (\text{h.c.}) \right] \\
& + \int d^4\theta \left\{ \frac{1}{g^2} |\Psi|^2 + \frac{g^2}{2} \left(\Gamma + \bar{\Gamma} + \sum_a V_a \right)^2 \right\} \\
& + \sum_a \left[\int d^2\theta (s_a - \Psi) \Phi_a + \int d^2\tilde{\theta} t_a \Sigma_a + (\text{h.c.}) \right]
\end{aligned}$$

$$\frac{1}{g^2} (\Theta + \bar{\Theta}) \doteq -(\Gamma + \bar{\Gamma} + \sum_a V_a)$$

$$\begin{aligned}
\mathcal{L}_{5_2^2} = & \sum_{a=1}^k \int d^4\theta \frac{1}{e_a^2} \left\{ -|\Sigma_a|^2 + |\Phi_a|^2 \right\} \\
& + \sum_a \int d^4\theta \left\{ |Q_a|^2 e^{V_a} + |\tilde{Q}_a|^2 e^{-V_a} \right\} + \sum_a \left[\int d^2\theta \tilde{Q}_a \Phi_a Q_a + (\text{h.c.}) \right] \\
& + \int d^4\theta \frac{g^2}{2} \left\{ \left(\Gamma + \bar{\Gamma} + \sum_a V_a \right)^2 - \left(\Xi + \bar{\Xi} - \sum_a (C_a + \bar{C}_a) \right)^2 \right\} \\
& + \int d^4\theta \left\{ -(\Psi - \bar{\Psi}) \sum_a (C_a - \bar{C}_a) \right\} \\
& + \sum_a \left[\int d^2\theta s_a \Phi_a + \int d^2\tilde{\theta} t_a \Sigma_a + (\text{h.c.}) \right]
\end{aligned}$$

$$\Phi_a = \bar{D}_+ \bar{D}_- C_a$$

$$\frac{1}{g^2} (\Psi + \bar{\Psi}) \doteq - \left(\Xi + \bar{\Xi} - \sum_a (C_a + \bar{C}_a) \right)$$