

Can accreting isolated neutron stars be detected?



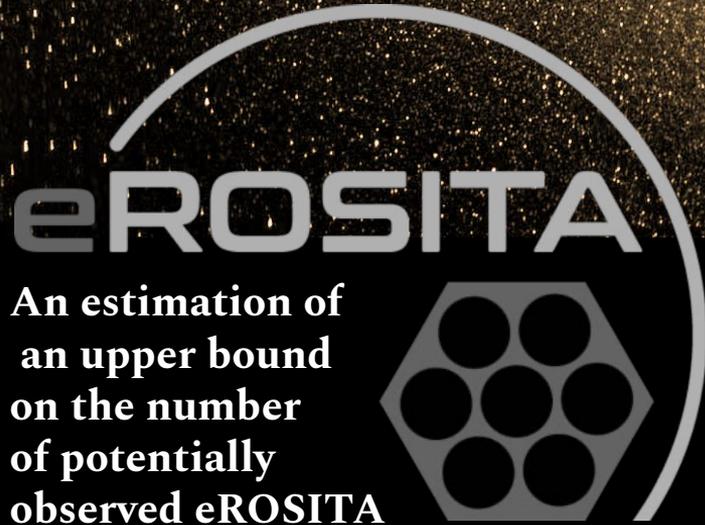
Marina Afonina,
Anton Biryukov, Sergei Popov
Sternberg Astronomical Institute

Kyoto, Japan
February 2026

2512.10666

arXiv

Results



An estimation of
an upper bound
on the number
of potentially
observed eROSITA

isolated accreting neutron stars

2512.10666

arXiv

One-phase ISM		
Propeller model	Constant field	Exponential decay
A	4700 ± 110	3200 ± 80
B	4700 ± 110	3100 ± 80
C	1300 ± 60	1600 ± 60
Two-phase ISM		
Propeller model	Constant field	Exponential decay
A	1300 ± 60	840 ± 40
B	1200 ± 60	760 ± 40
C	270 ± 30	280 ± 30

most realistic model

With 300 million NSs in the Galaxy, if the accretion efficiency is high and the accretor stage can be reached, hundreds to thousands low-velocity NSs are potentially observable soft X-ray sources

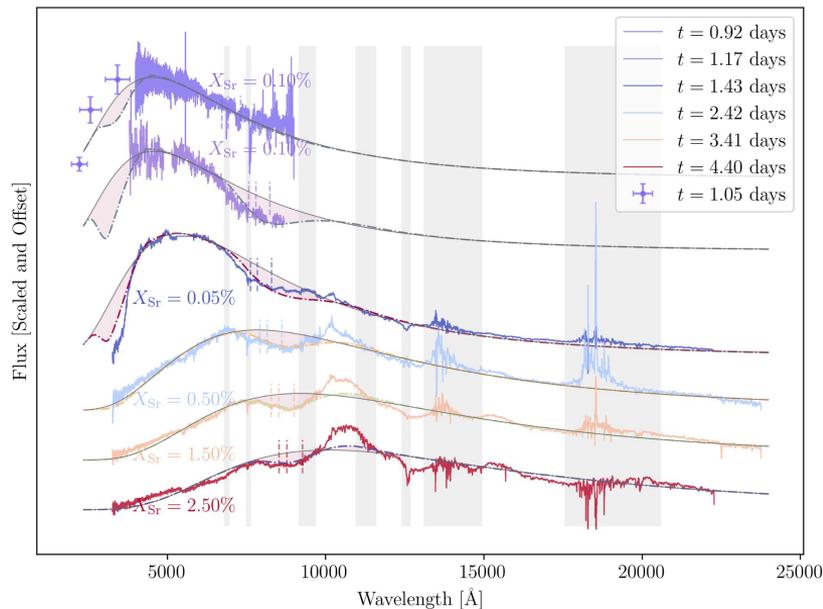
Strontium vs. helium spectral lines in the kilonova AT2017gfo

Aayush Arya

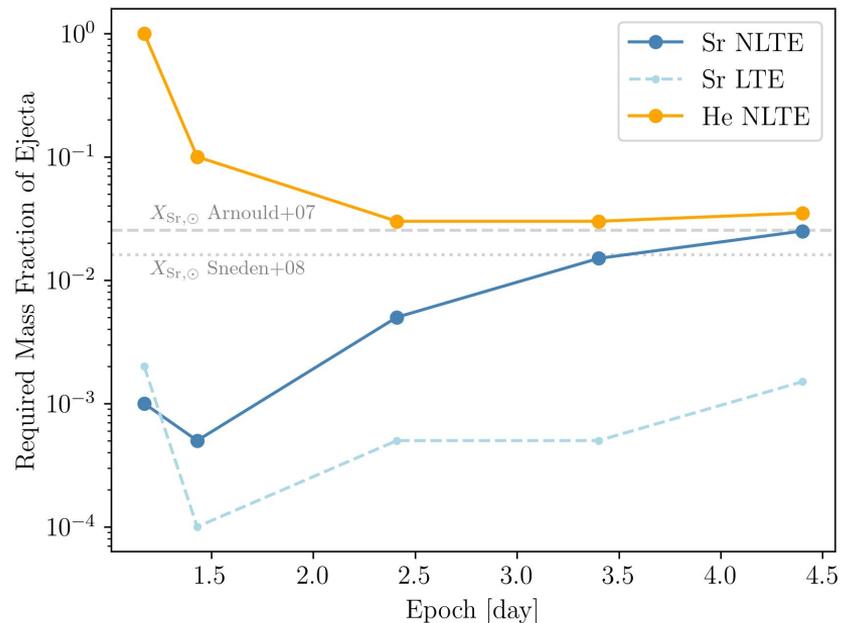
Niels Bohr Institute, University of Copenhagen, Denmark

with Rasmus Damgaard, Albert Sneppen, Stuart Sim, David Dougan, Connor Ballance, and Darach Watson

1 μ m feature with NLTE radiative transfer



Strontium vs. Helium: What is it?



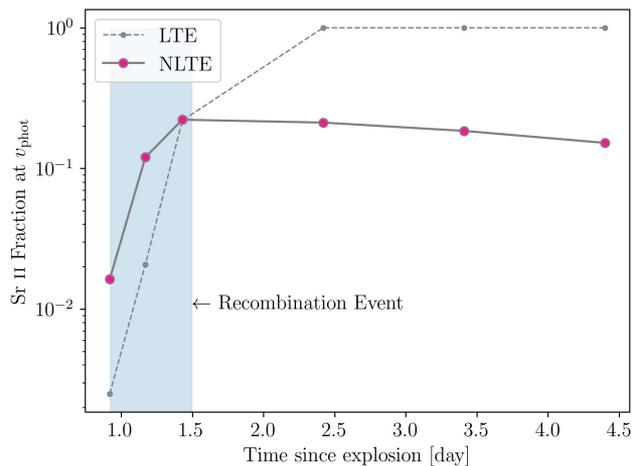
Strontium vs. helium spectral lines in the kilonova AT2017gfo

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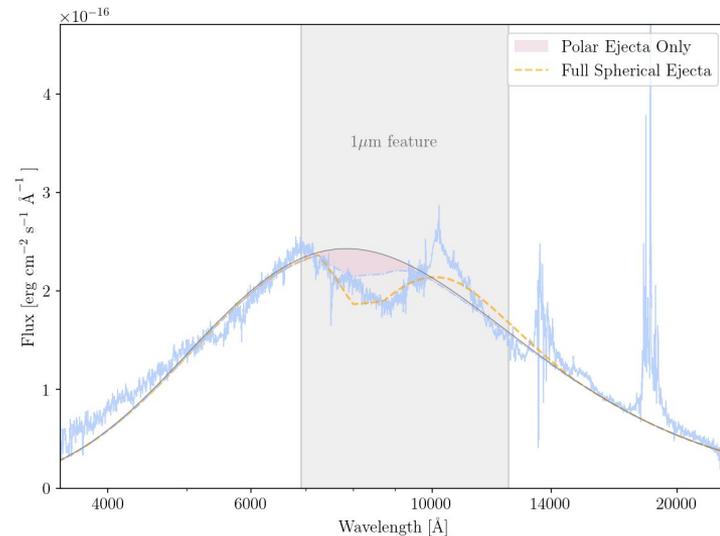
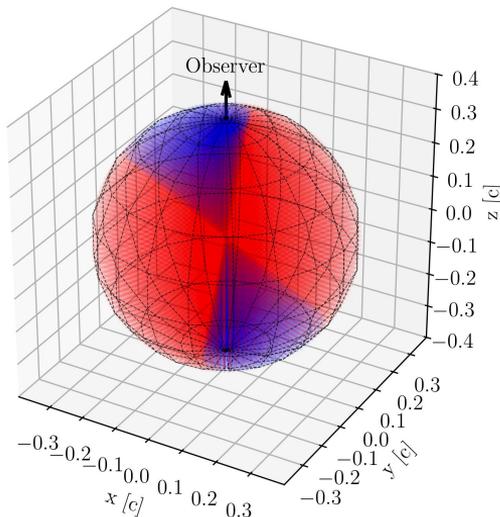
Niels Bohr Institute, University of Copenhagen, Denmark

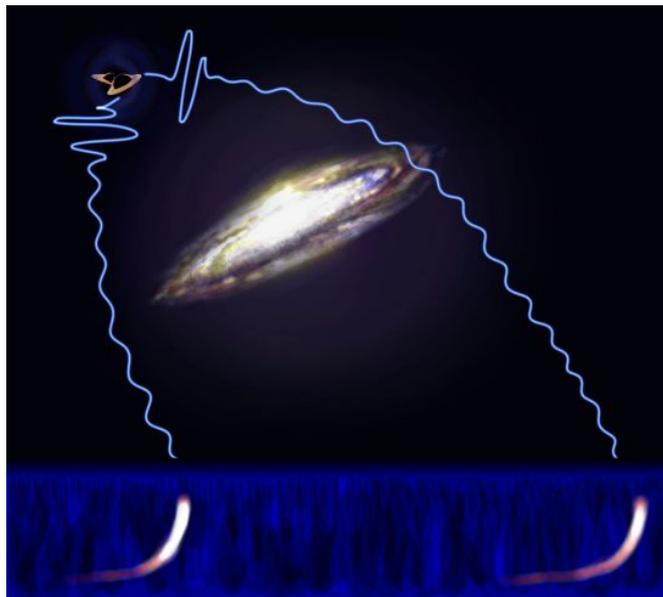
with Rasmus Damgaard, Albert Sneppen, Stuart Sim, David Dougan, Connor Ballance, and Darach Watson

Emergence of the feature: Recombination Event



Helium in polar ejecta cannot form a full P Cygni

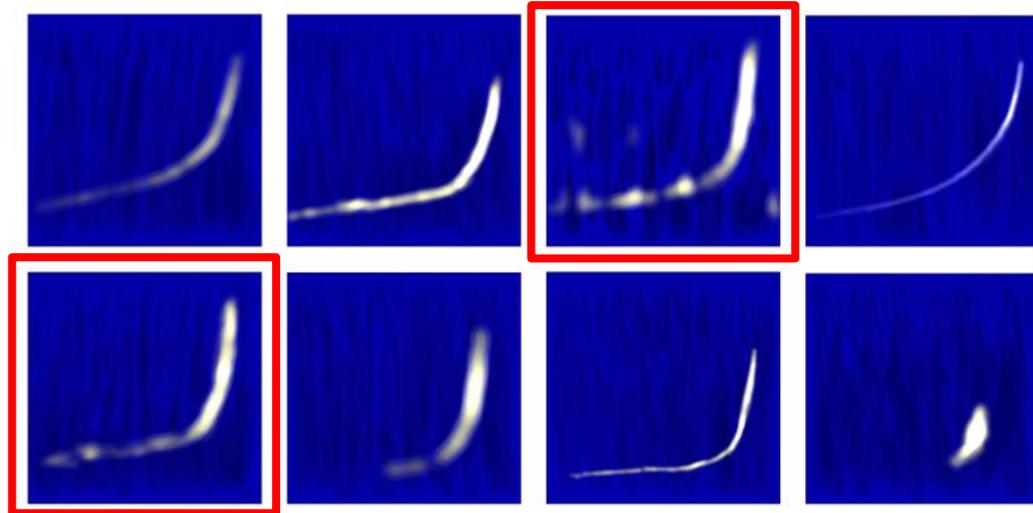




AB, Koustav

Strongly lensed gravitational waves

- Multiple copies of the same signal
- Useful probe of dark matter, cosmology, testing GR



Number of lensed signals $\sim 0.001 \times N$

Number of unlensed pairs $\sim N^2 / 2$

False positives grow faster!

Lensing, or luck?

Detection prospects of strongly lensed gravitational waves

Ankur Barsode*, Koustav N. Maity, Parameswaran Ajith

International Centre for Theoretical Sciences, Tata Institute of Fundamental Research, Bangalore

*ankur.barsode@icts.res.in

arXiv:2510.23238

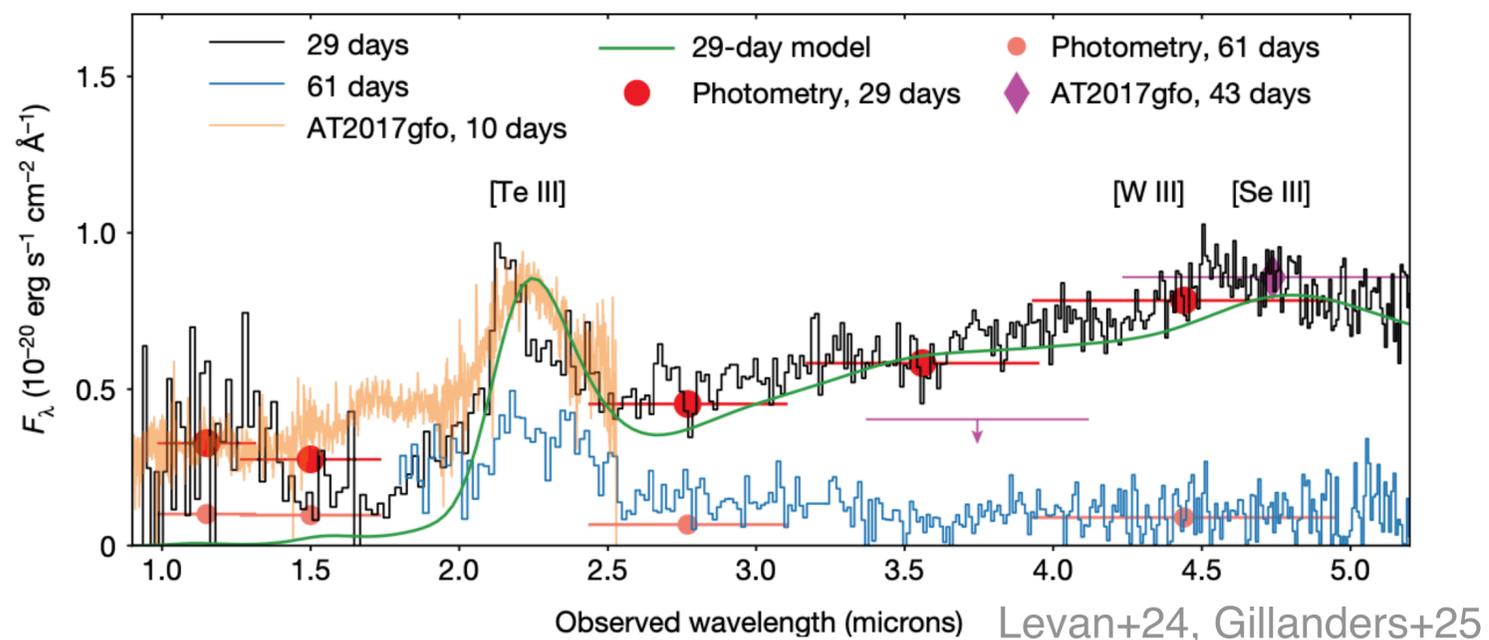
- First significant detection likely in LVK's O5!
- Routine detections thereafter

Lensing, not luck!

Dust formation in kilonovae

Nanae Domoto (U. Tokyo, RESCEU), Kenta Hotokezaka (RESCEU), Daniel Kasen (UC Barkley)

Late-time emission of observed kilonovae

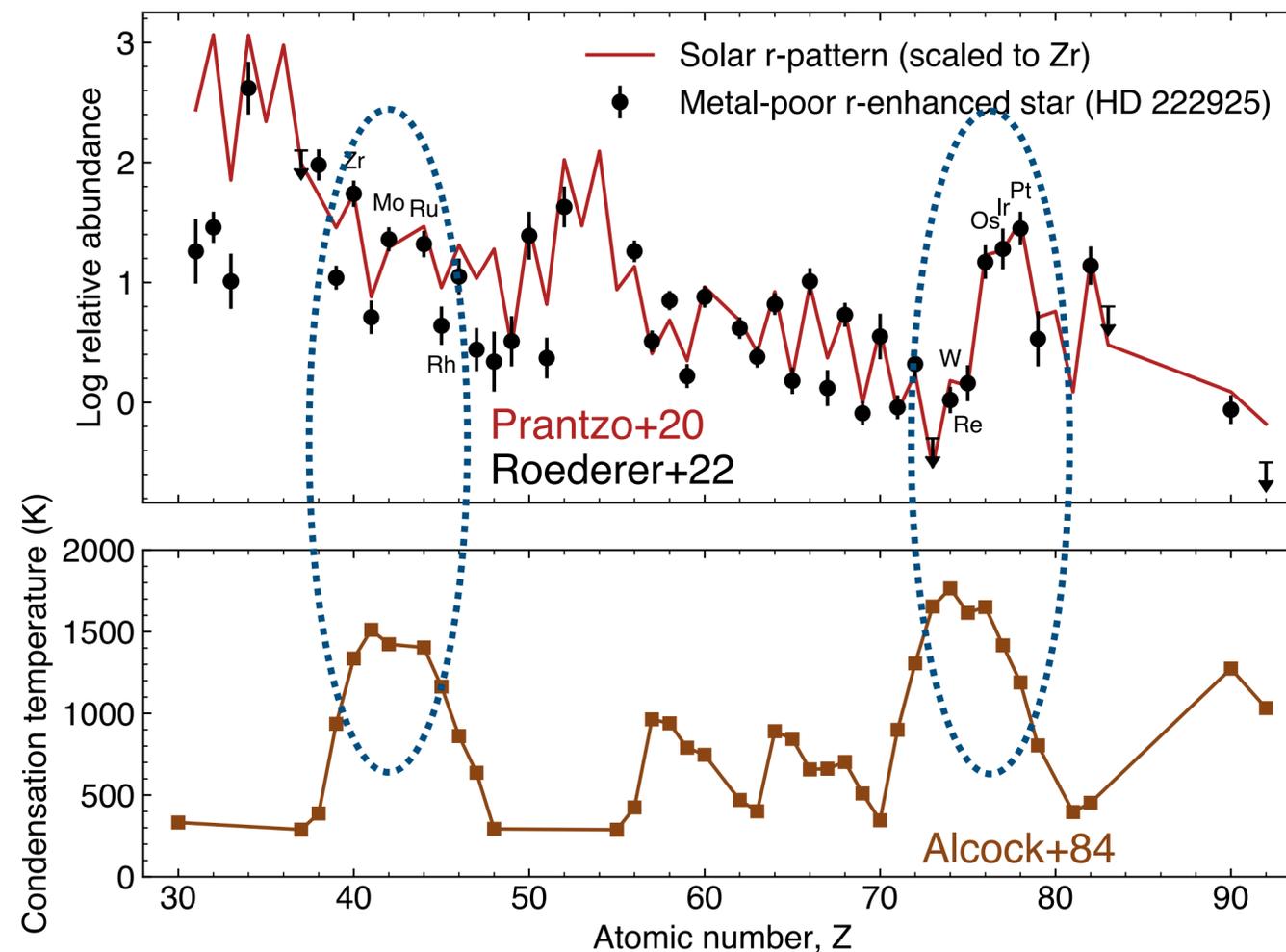


- AT2017gfo (GW170817) and AT2023vfi (GRB 230307A) revealed prominent red emission at late time
- Continuum-like emission is difficult to explain by atomic processes

→ **R-process dust grains?**

Refractory element

Palme & Wlotzka 1987, Wood et al. 2019

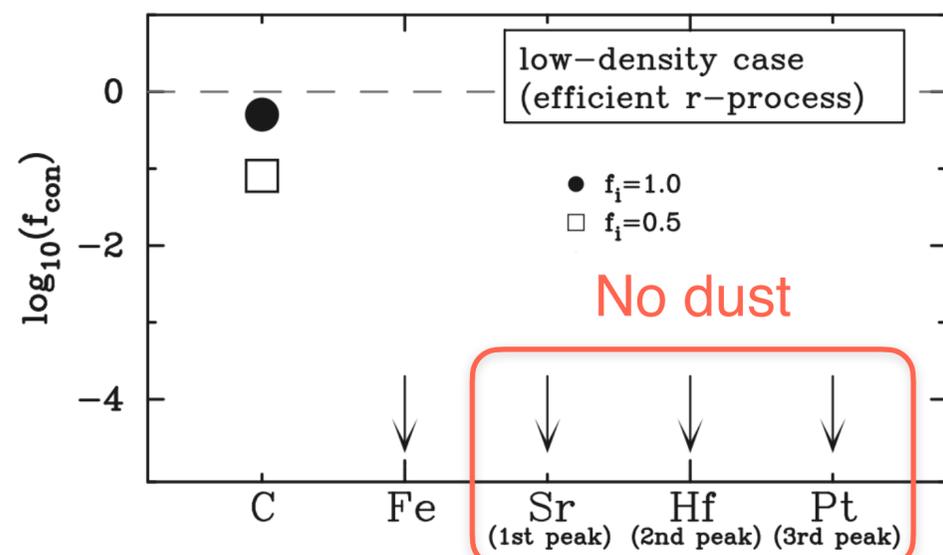


- Zr, W, Os, ... have high condensation temperatures ~1800 K
- Ejecta temperature drops below this at ~10 days after the merger
- “Refractory metal nuggets” has been observed in meteorites (considered to form in early solar system)

Dust formation in kilonovae

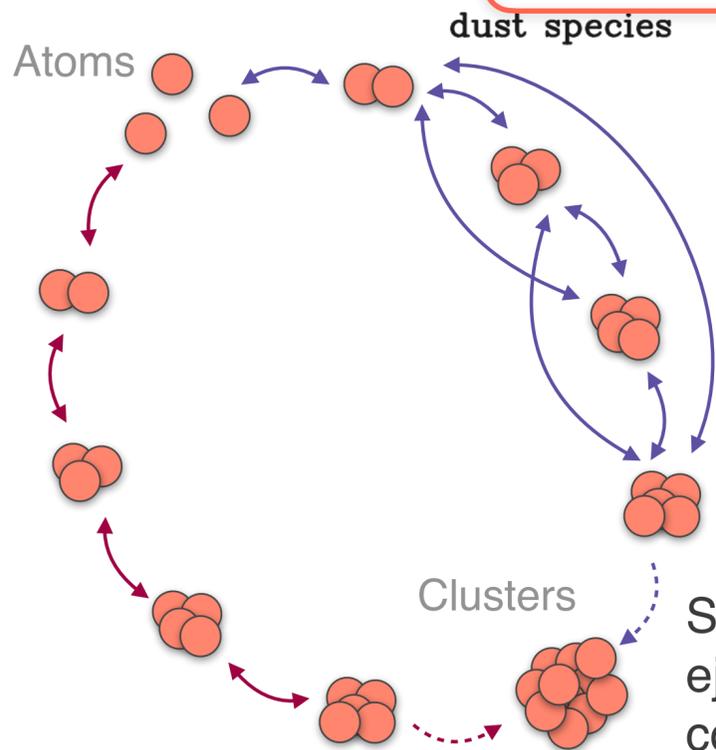
Nanae Domoto (U. Tokyo, RESCEU), Kenta Hotokezaka (RESCEU), Daniel Kasen (UC Barkley)

Methodology



Previous study has concluded that dust formation is unlikely (Takami+14)

However... **only low density ejecta was considered; the nucleation theory was used**



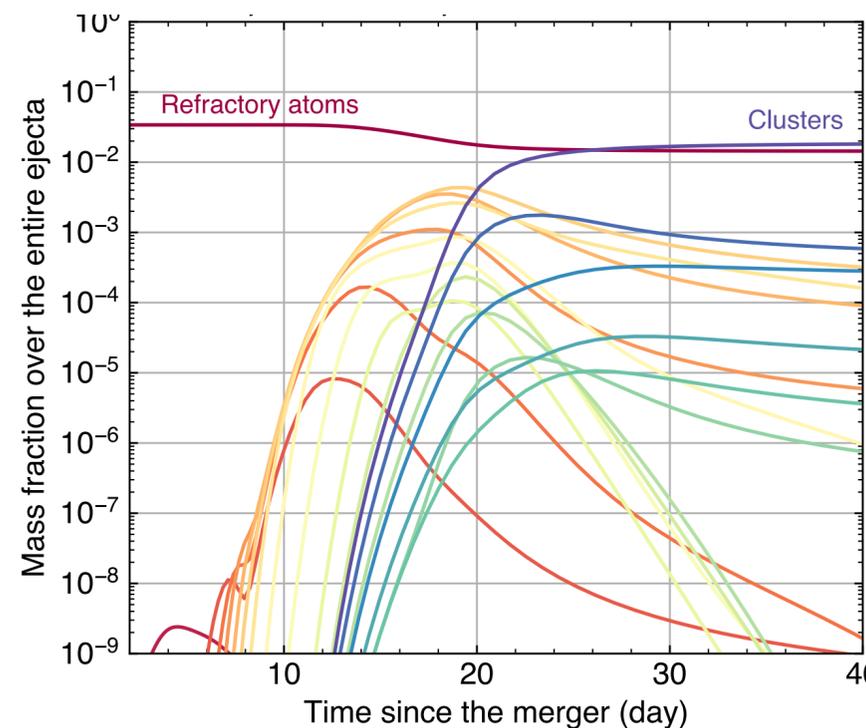
“molecular” picture

Cherchneff & Lilly 08,
Cherchneff & Dwek 09, 10,
Lazzati & Heger 16, Sluder+18, ...

Solving rate equations in expanding ejecta using the theoretical rate coefficients from Matúška+19

“classical” picture

Results & Discussions

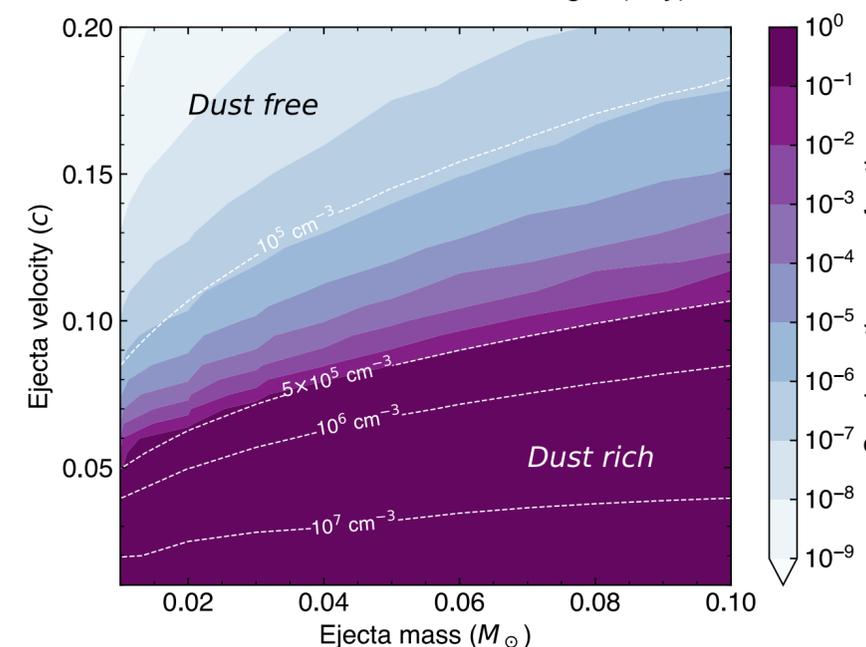


Refractory element grains may form at ~10 days after the merger.

Resulting grain mass in this case is $\sim 10^{-3} M_{\text{sun}}$.
This may produce thermal red emission.

Condensation mass fraction is strongly sensitive to the velocity (density) of ejecta.

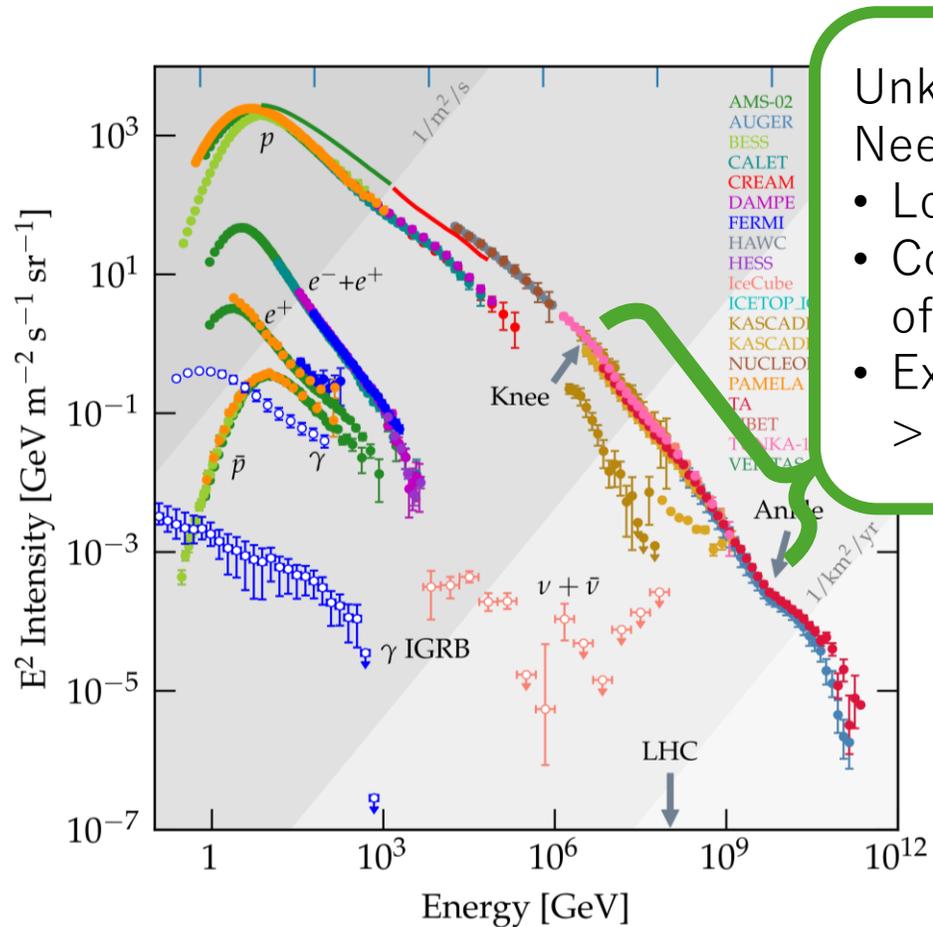
More data of molecular and cluster formation of r-process elements are needed.



Please visit the poster No.4 for more details!

Interacting SNe: super-PeV candidates that explain the composition trends

Nick Ekanger (Tohoku University), with Shigeo S. Kimura, Kazumi Kashiyama



Unknown!

Need:

- Local
- Composed of nuclei
- Explain flux $> 10^{15} \text{ eV}$

- Typical SNRs cannot reach PeV
- *Interacting SNe?*

$$E_{max} \sim r_L Z e B$$

$$B^2/8\pi \sim \epsilon_B \rho_w V_w^2/2$$

$$\rho_w \sim \dot{M}/4\pi R_0^2 V_w$$

$$* E_{max} \sim Z e (\epsilon_B \dot{M} V_w)^{1/2} \sim 5 \times 10^{16} \text{ eV}$$

$$L_{CR|GeV}/L_{CR|PeV} \sim 10^3$$

$$L_{CR} \sim \epsilon_{CR} M V^2 \mathcal{R}/2$$

$$L_{CR|SNR}/L_{CR|ISN} \sim \frac{M_{ej} \mathcal{R}_{SN}}{M_{CSM} \mathcal{R}_{ISN} / \ln(E_{max}/E_{min})} \sim 10^3$$

For parameters:

- $Z = 2$
- $\epsilon_B = 10^{-3}$
- $\dot{M} = 10^{-2} M_\odot \text{ yr}^{-1}$
- $V_w = 100 \text{ km s}^{-1}$
- $M_{ej} = 10 M_\odot$
- $M_{CSM} = 1 M_\odot$
- $\mathcal{R}_{SN}/\mathcal{R}_{ISN} \sim 10$
- $E_{min} \sim 1 \text{ GeV}$

* E_{max} , also:

- Spallation, pp
- Photodisintegration
- Adiabatic loss
- Escape

Magnetic field amplification

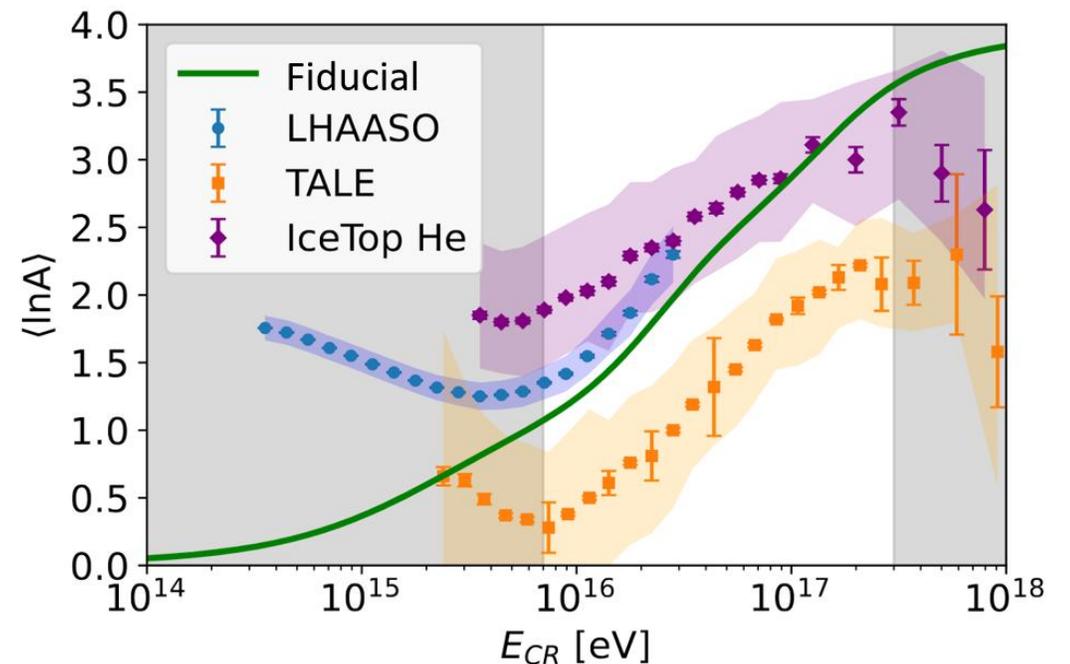
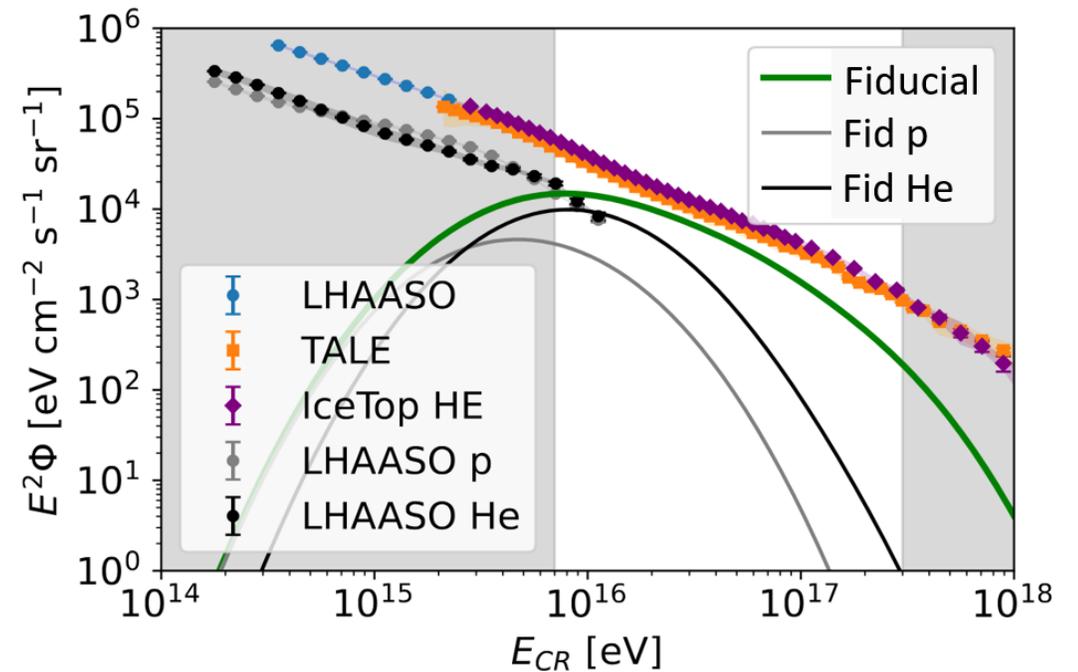
- CRs escape, producing a current
- Excite modes $< r_L$ (*non-resonant streaming* or *Bell instability*)
- $A \sim 3 - 5$ for In parameters

$$A = \max\left(\frac{B_{sat,NRS}}{B_w}, 1\right)$$

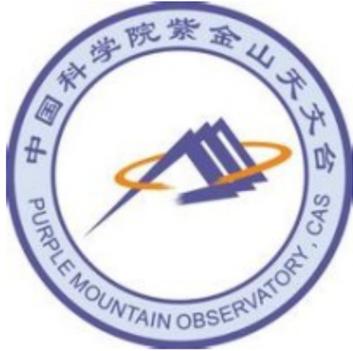
Heavy ion abundance enhancement

- Simulations show $f_{inj} \sim (A/Z_{ion})^{5/2}$
- Z_{ion} from photoionization code, $T \sim 15000$ K

		$f_{inj} \sim (A/Z_{ion})^{5/2}$
H	$\sim 90\%$	$\rightarrow 24\%$
He	$\sim 9\%$	$\rightarrow 53\%$
CNO	$\sim 0.1\%$	$\rightarrow 20\%$
Fe	$\sim 0.002\%$	$\rightarrow 2\%$



Dependence of postmerger properties on the thermal heating efficiency in neutron star mergers

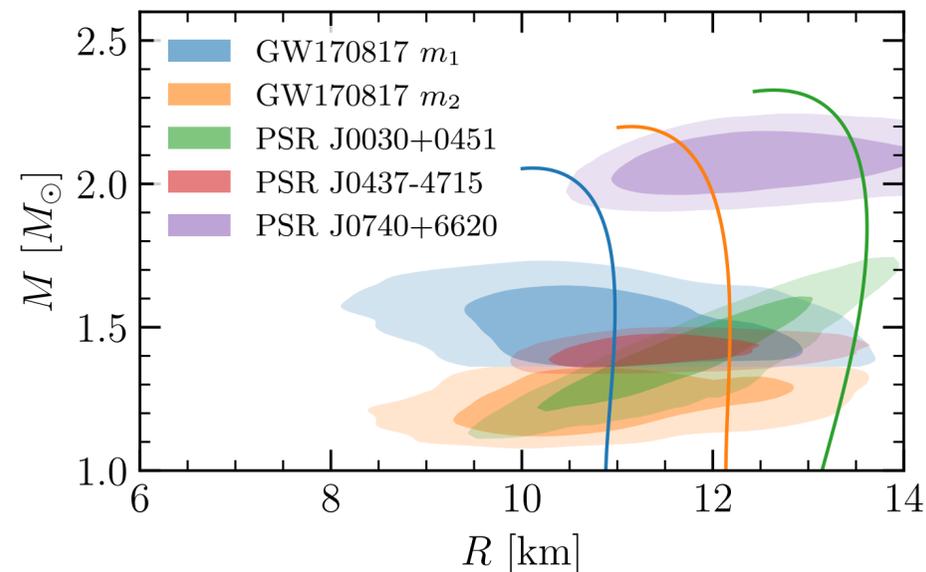
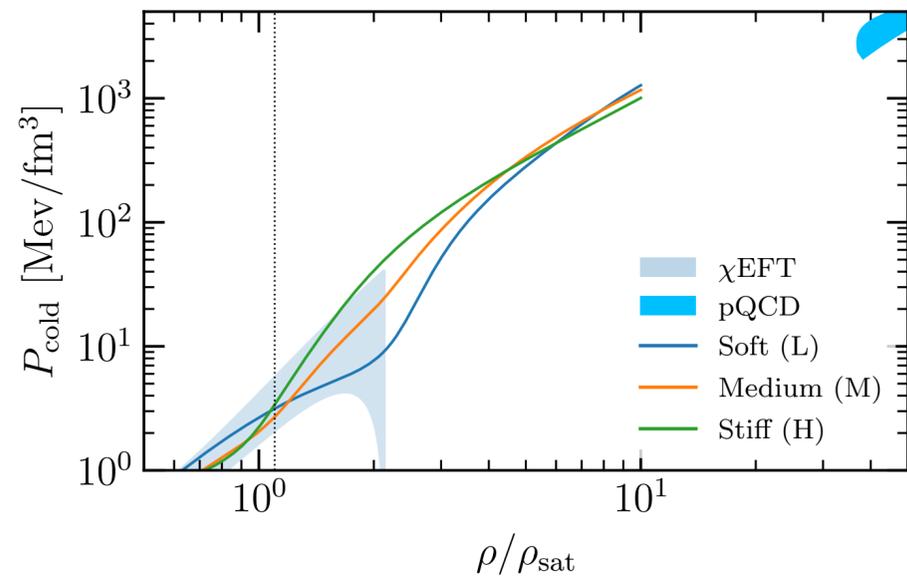


Ming-zhe Han, Yong Gao, Kenta Kiuchi, Masaru Shibata

DOI: [10.1103/p6s8-khkw](https://doi.org/10.1103/p6s8-khkw) arxiv: [2504.08514](https://arxiv.org/abs/2504.08514)



EOS



EOS constraints

χ EFT at low density (up to $\sim 1.1\rho_{\text{sat}}$)

pQCD at high density

M- Λ measurements of **GW170817**

M-R measurements of

PSR J0030+0451,

PSR J0740+6620,

and **PSRJ0437-4715**

Thermal effect

$$\mathcal{E} = \mathcal{E}_{\text{cold}} + \mathcal{E}_{\text{th}}$$

$$P = P_{\text{cold}} + P_{\text{th}}$$

$$P_{\text{th}} = (\Gamma_{\text{th}} - 1)\rho\mathcal{E}_{\text{th}}$$

Numerical Setups

ID solver FUKA

Evolution SACRA

EOS Soft, Medium, Stiff

Mass Equal mass with total mass $m_0 = 2.7, 2.9 M_{\odot}$

Thermal index $\Gamma_{\text{th}} = 1.1, 1.2, 1.3, \dots, 2.0$

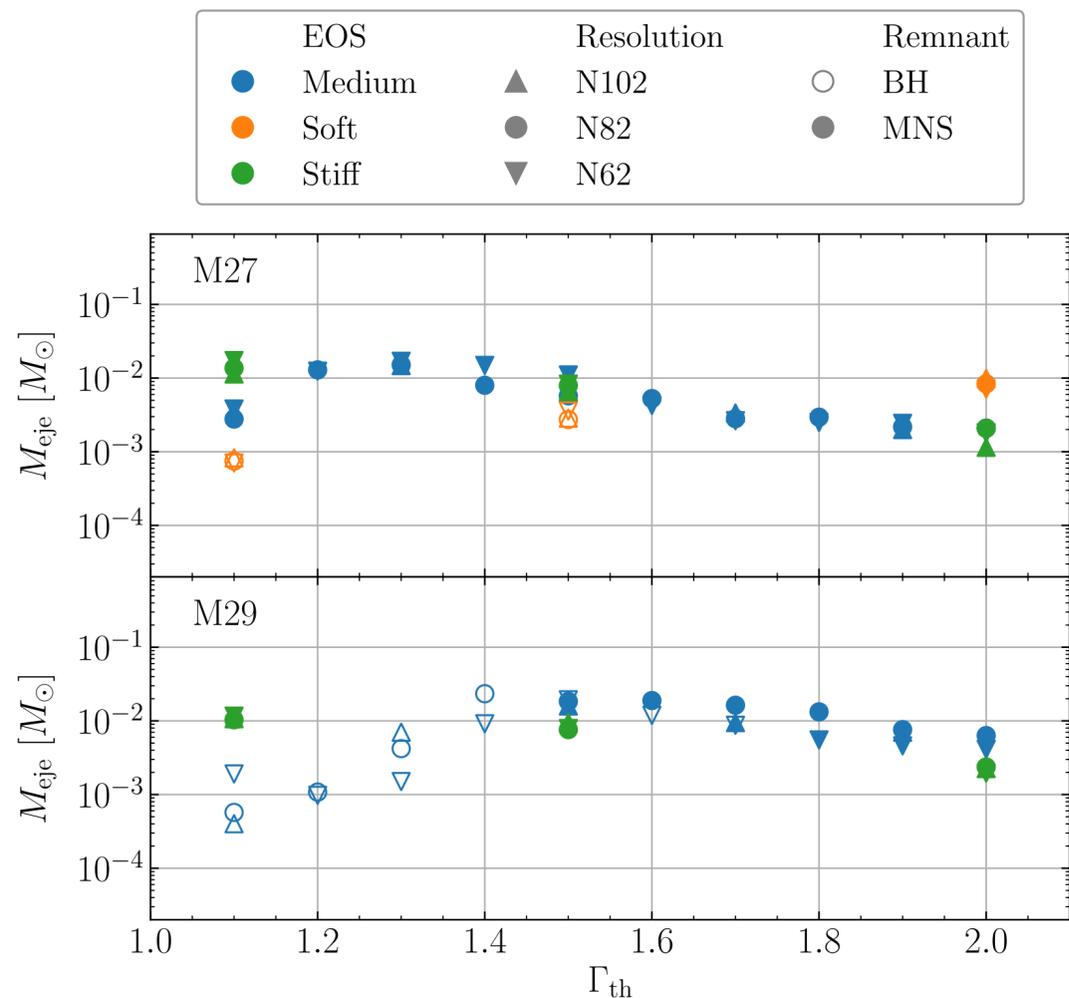
Separation ~ 45 km

eccentricity $\lesssim 0.001$

Simulation time ~ 30 ms after merger

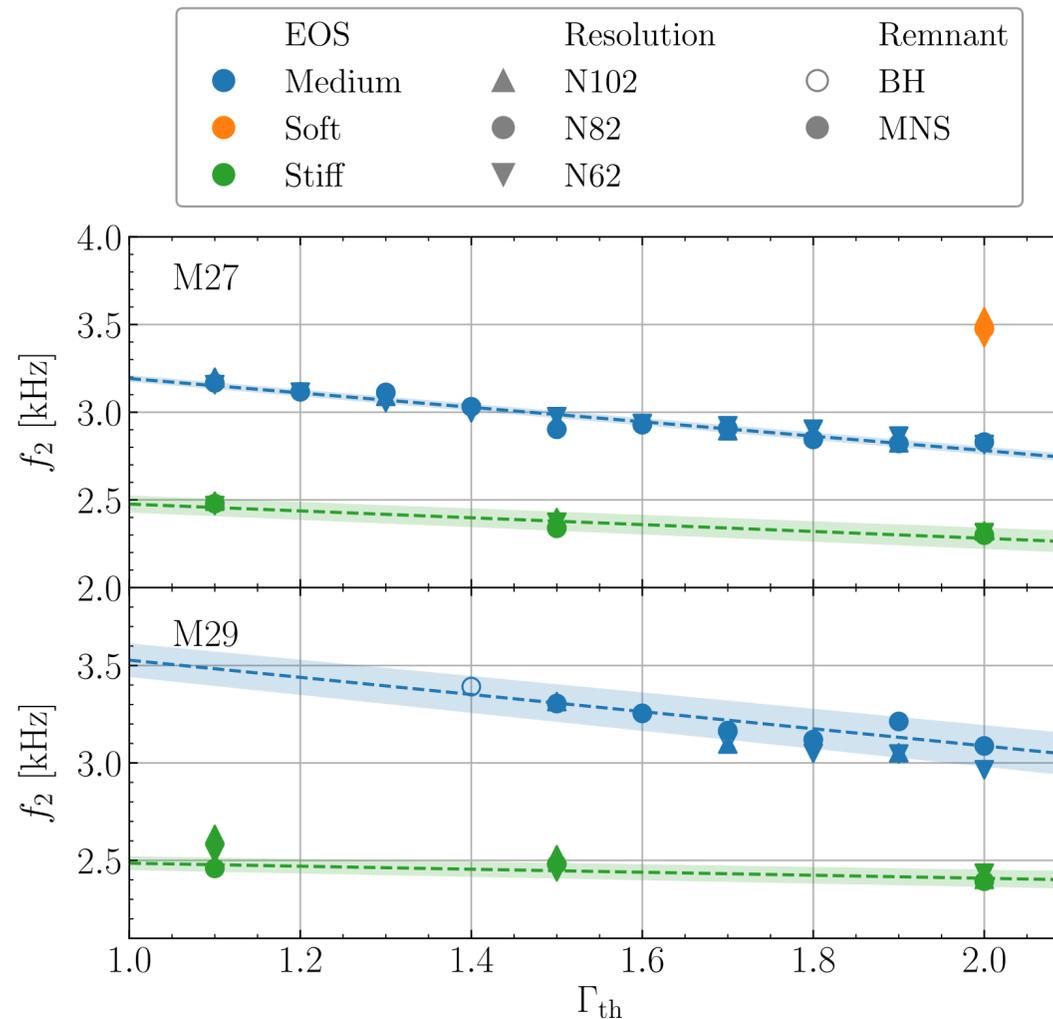
Resolution N62 ~ 150 m, N82 ~ 180 m, N102 ~ 240 m

Dynamical ejecta



Main results

f_2 peak



Fitting with:

$$f_2 = a\Gamma_{\text{th}} + b$$

Model	a [kHz]	b [kHz]	(σ_a, σ_b)
M-M27	-0.411	3.603	(0.0019, 0.0046)
M-M29	-0.440	3.967	(0.0103, 0.0301)
H-M27	-0.195	2.671	(0.0061, 0.0152)
H-M29	-0.077	2.563	(0.0043, 0.0108)

Three behaviors with different EOSs:

Soft, **monotonically** increase with Γ_{th}

Medium, **peak structure**, increase first then decrease

Stiff, **monotonically** decrease with Γ_{th}

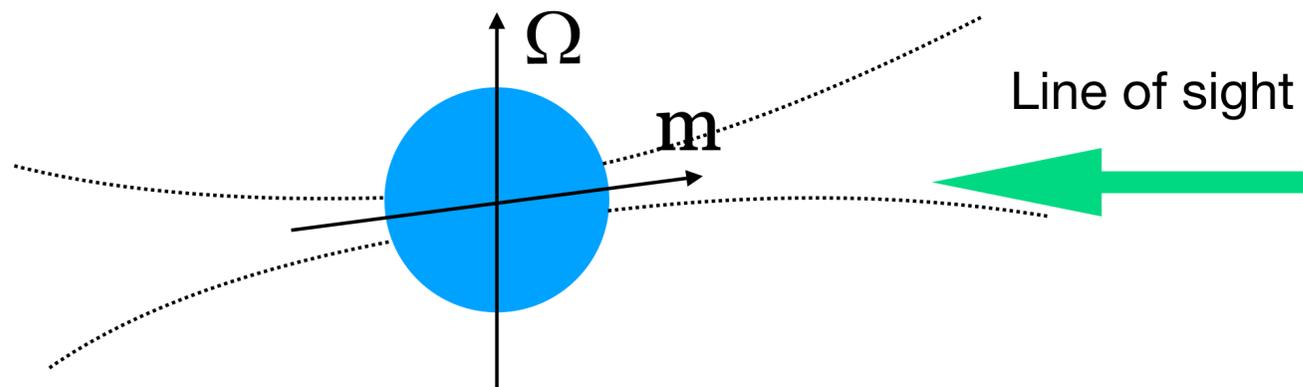
f_2 **monotonically decrease** with Γ_{th} ,

where f_2 is the dominant peak of the post-merger frequencies.

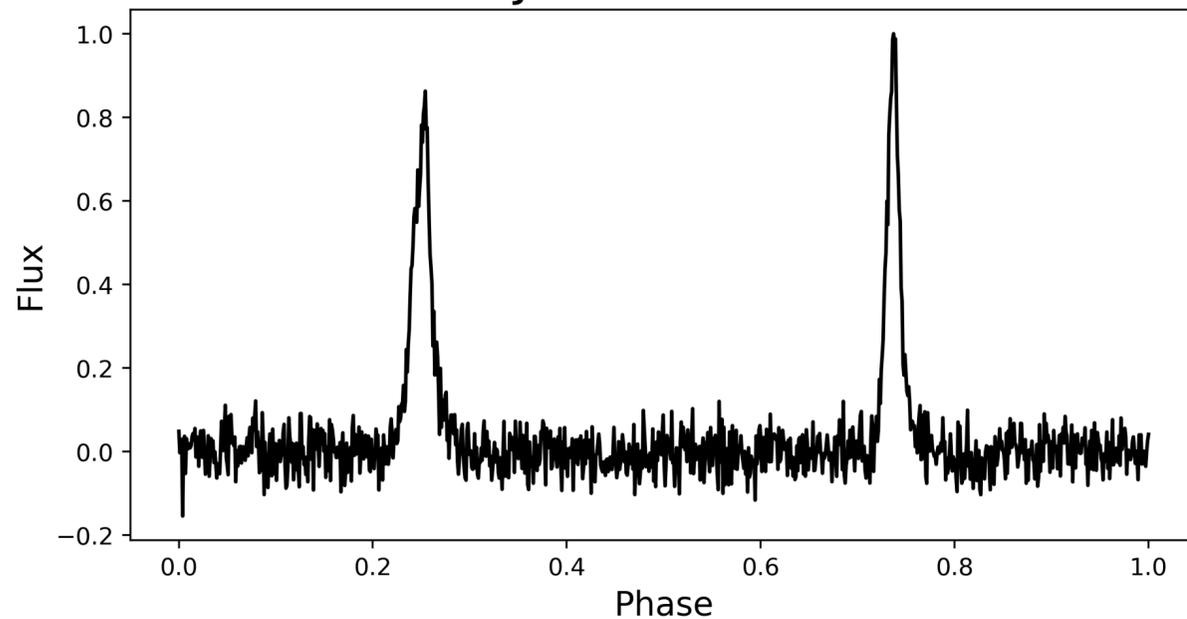
Orthogonal radio pulsars – theory and observations

Istomin A. Yu.^{1,2}, , Kniazev F. A.^{1,2}, Beskin V. S.^{1,2}

Orthogonal radio pulsars



J1909+0749



Data from Posselt et. al., MNRAS, 2023

Magnetosphere structure and «death line»

- «Classical» approach – accelerating potential is determined by Goldreich-Julian charge density
- «Modern» approach – accelerating potential is determined by the current, required for a smooth MHD-like solution

«Classical» death line

$$\dot{P} > P^{11/4} k^{7/4} \cos^{-7/4} \chi \quad k \sim 1$$

For orthogonal radiopulsars: $\cos \chi \sim \sqrt{\frac{\Omega R}{c}}$

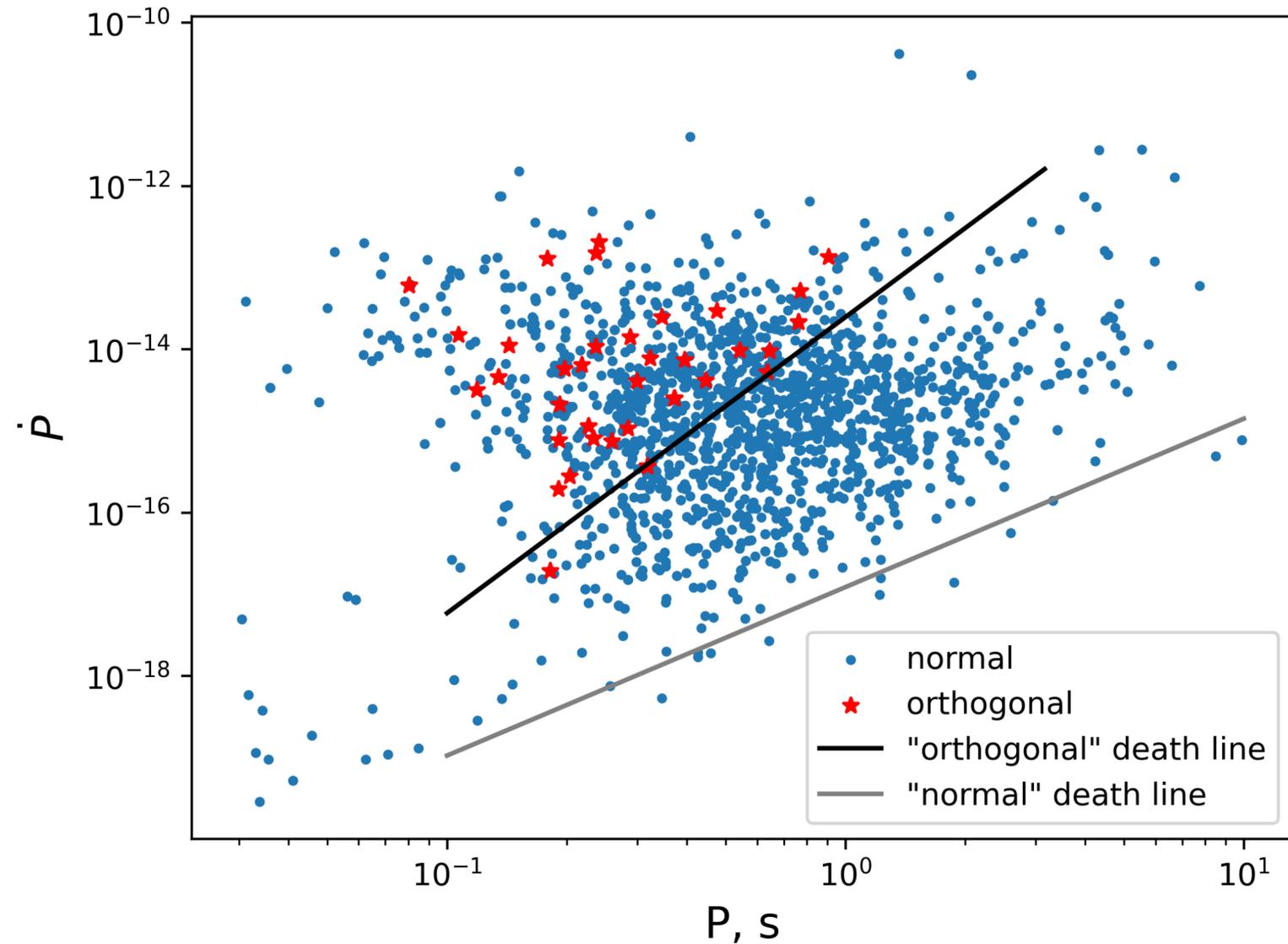
$$\dot{P} > P^{29/8} k^{7/4} \left(\frac{2\pi R}{c} \right)^{-7/4}$$

Kniazev F. A., Istomin A. Yu., Beskin V. S., Astron. Let. 2025

Orthogonal radio pulsars – theory and observations

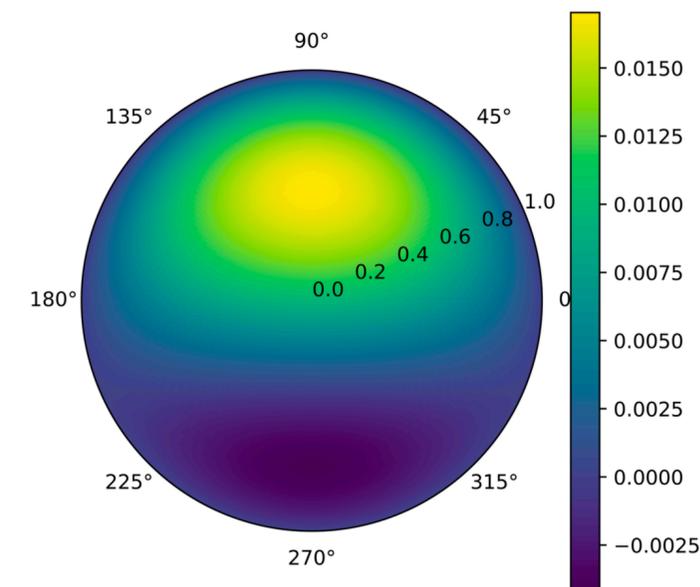
Istomin A. Yu.^{1,2}, ✉, Kniazev F. A.^{1,2}, Beskin V. S.^{1,2}

FAST and MeerKAT observations



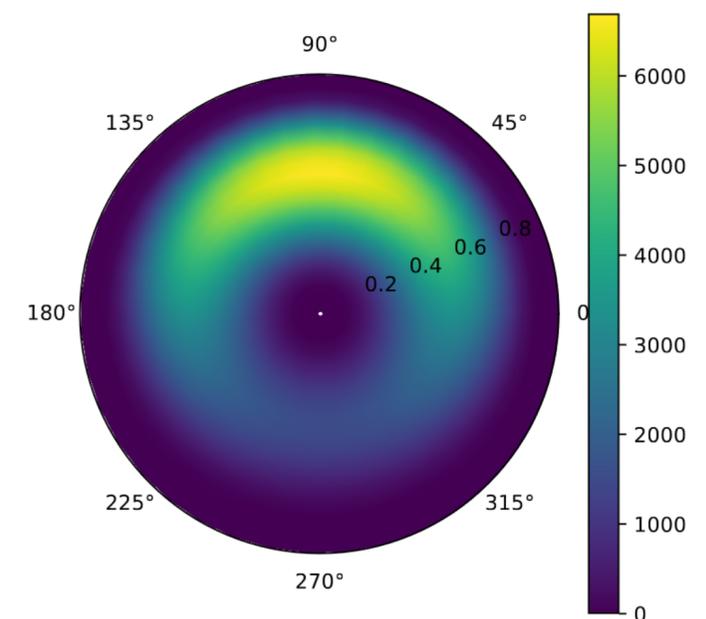
Data from Posselt et. al., MNRAS, 2023 and Wang et. al., RAA 2023

More on the theory of orthogonal pulsars



Example of accelerating potential calculation

Example of plasma number density calculation



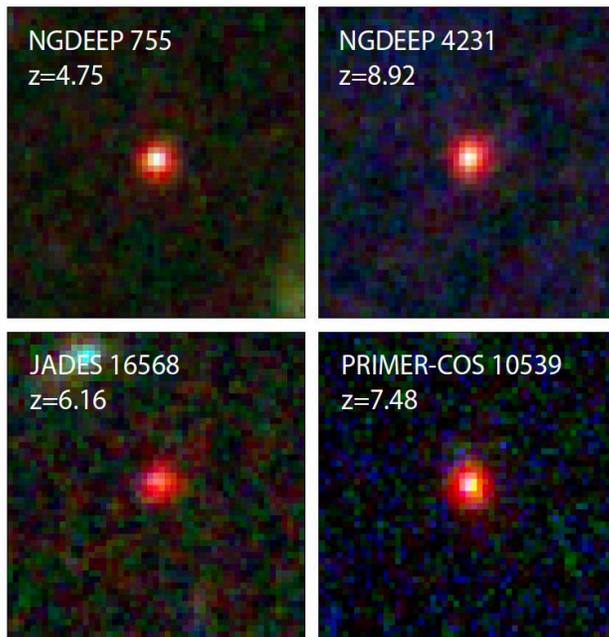
Istomin A. Yu., Kniazev F. A., Beskin V. S., Astron. Rep. 2024

Modeling of Little Red Dots and open challenges

3 main characteristics of LRDs Daisaburo Kido@RESCEU, UTokyo

Collaborators: Kunihiro Ioka, Kenta Hotokezaka, Kohei Inayoshi, Christopher Irwin

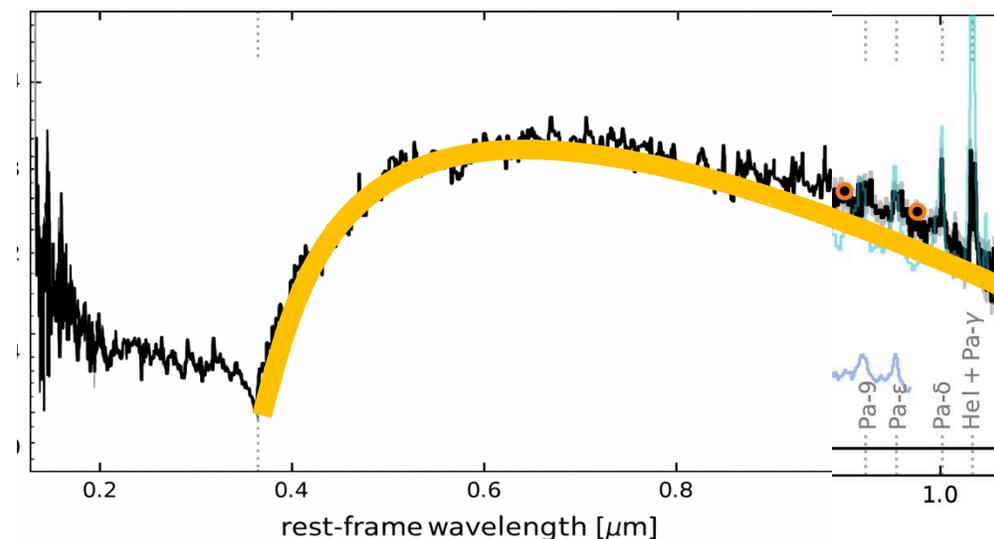
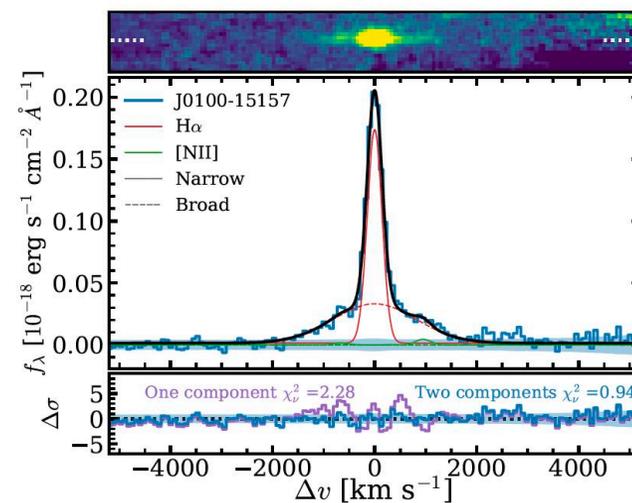
1. Very compact morphology (<100 pc)
 2. Broad hydrogen Balmer emission lines (>1000 km/s)
 3. Red in NIRC2, showing a “V-shaped” SED
- +X-ray, mid IR weak, 1~10 years variability, strong Balmer break...



Matthee+24

Kocevski+24

de Graaff +25



Our work: BH envelopes

Very gas envelope around BH

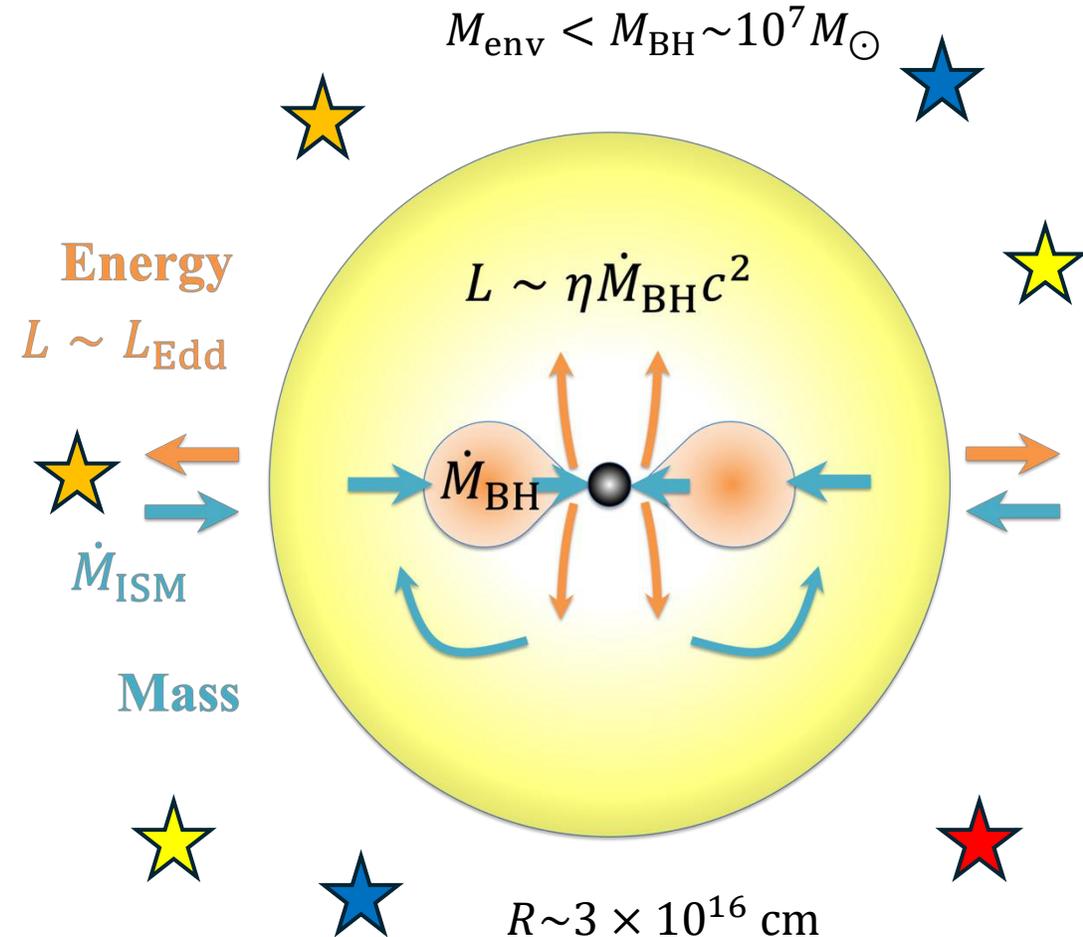
1. Hide X-ray, outflow
2. Reproduce red optical continuum
3. Dynamical timescale of 1-10 years

Reprocess outflow energy via radiation and convection

Finally, $T_{\text{ph}} \sim 5000$ K blackbody radiation is emitted from surface

$$r_{\text{ph}} = 10^{16} \left(\frac{L}{10^{45} \text{ erg s}^{-1}} \right) \left(\frac{T_{\text{ph}}}{5000 \text{ K}} \right)^{-2} \text{ cm}$$

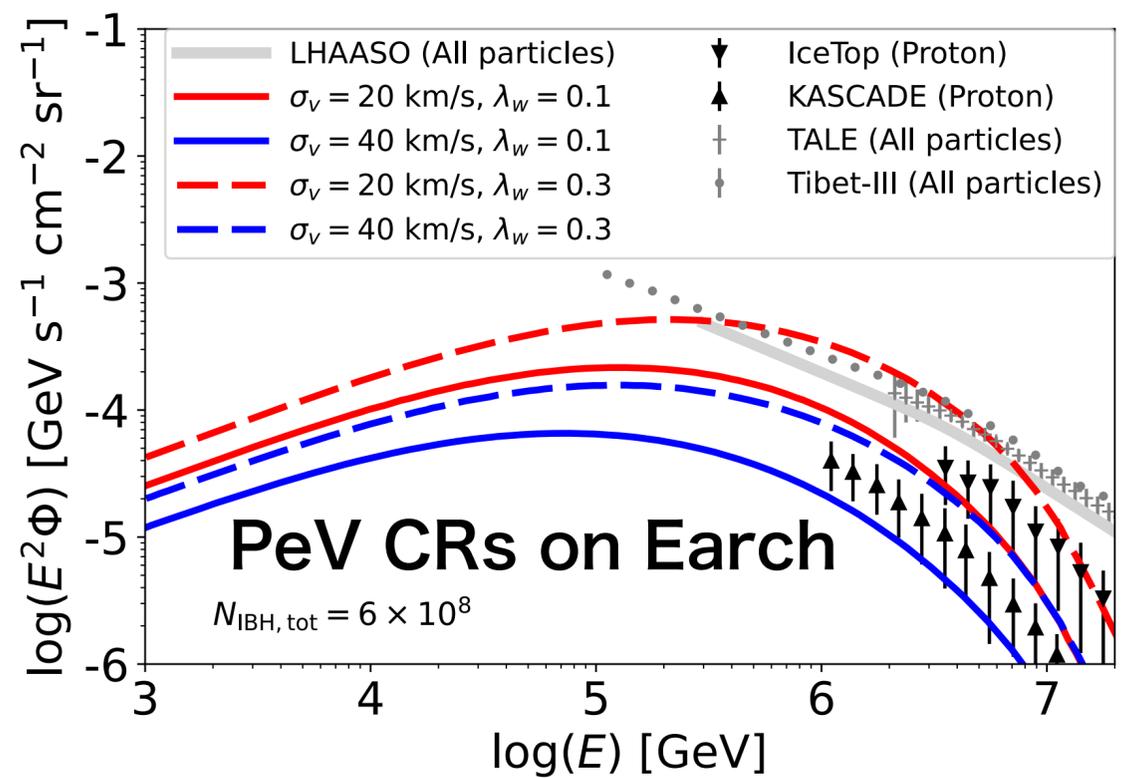
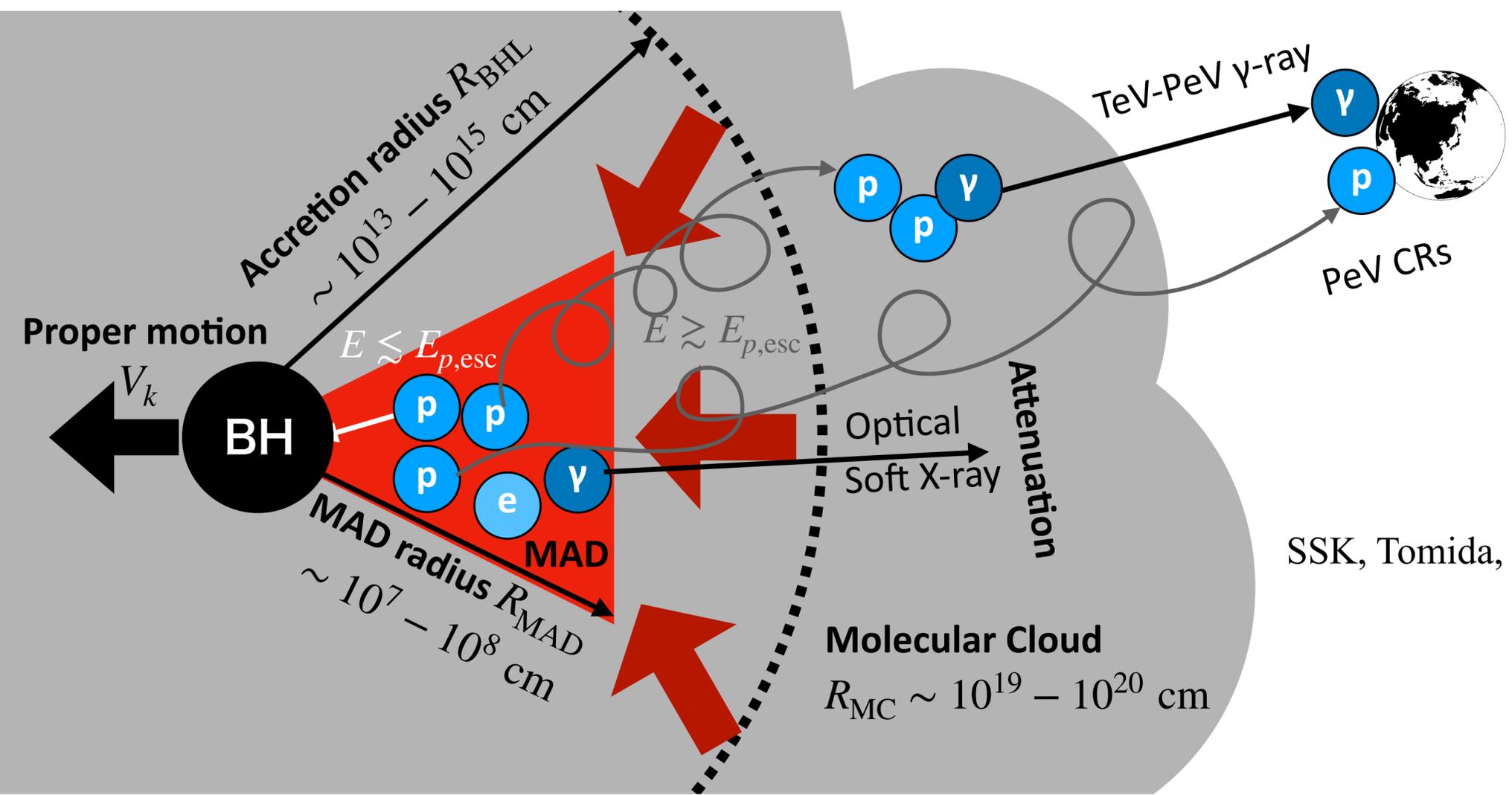
Required mass is 1~10% of BH mass



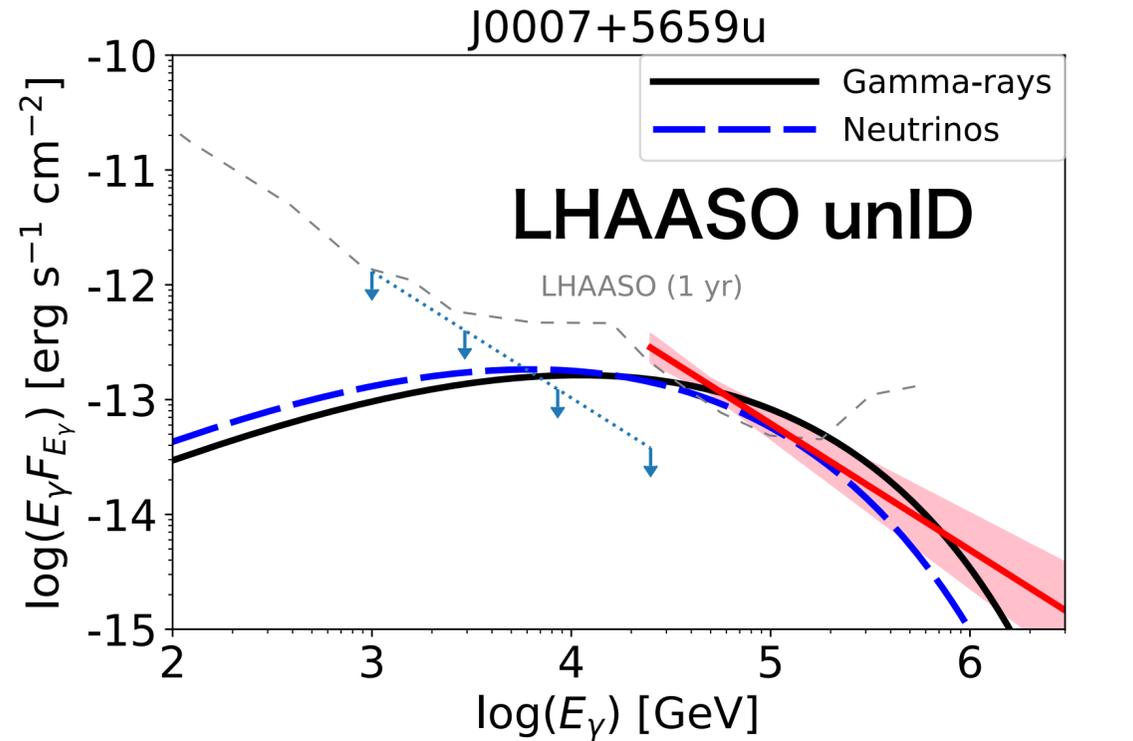
Isolated Black Holes as Potential PeVatrons

Tohoku University

Shigeo S. Kimura



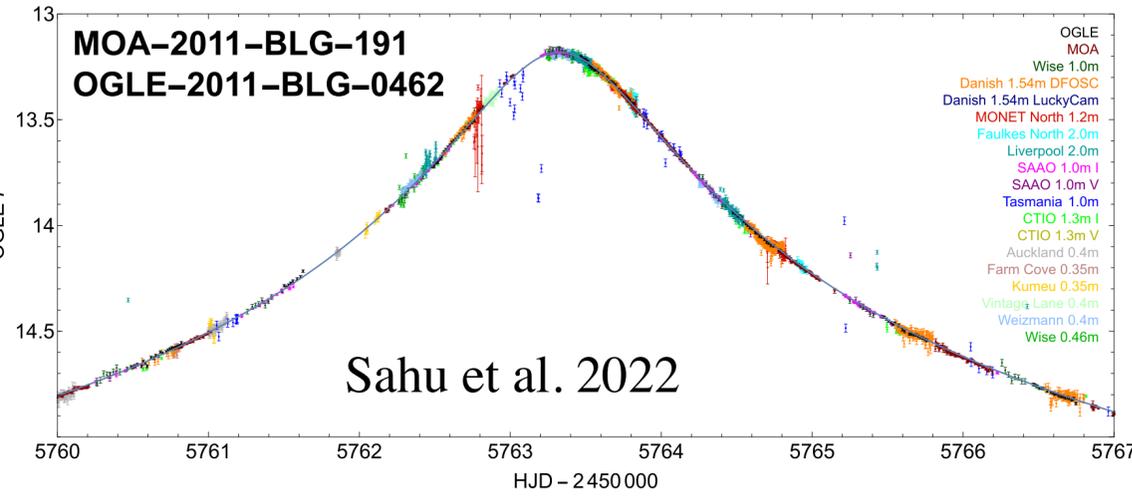
SSK, Tomida, et al. 2025 ApJL



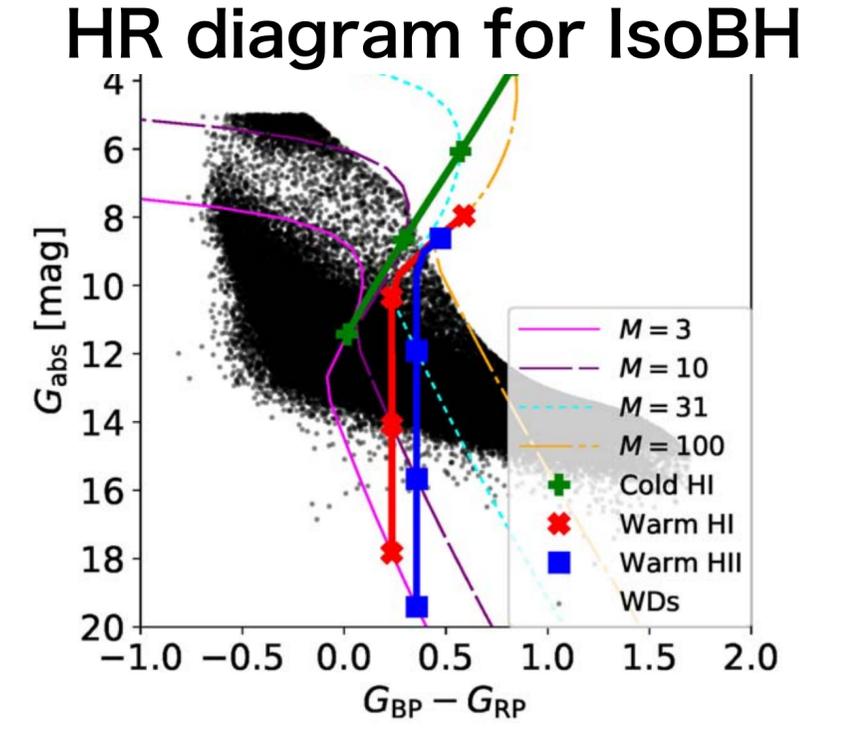
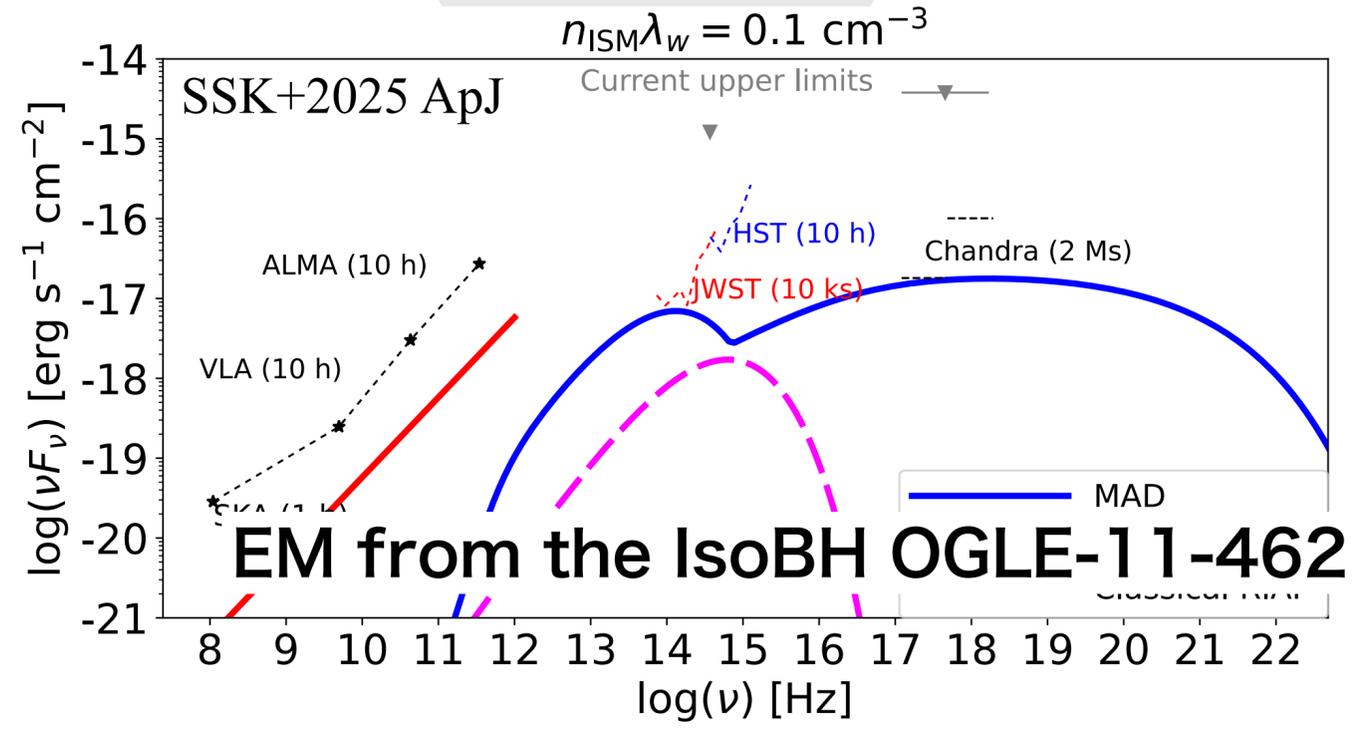
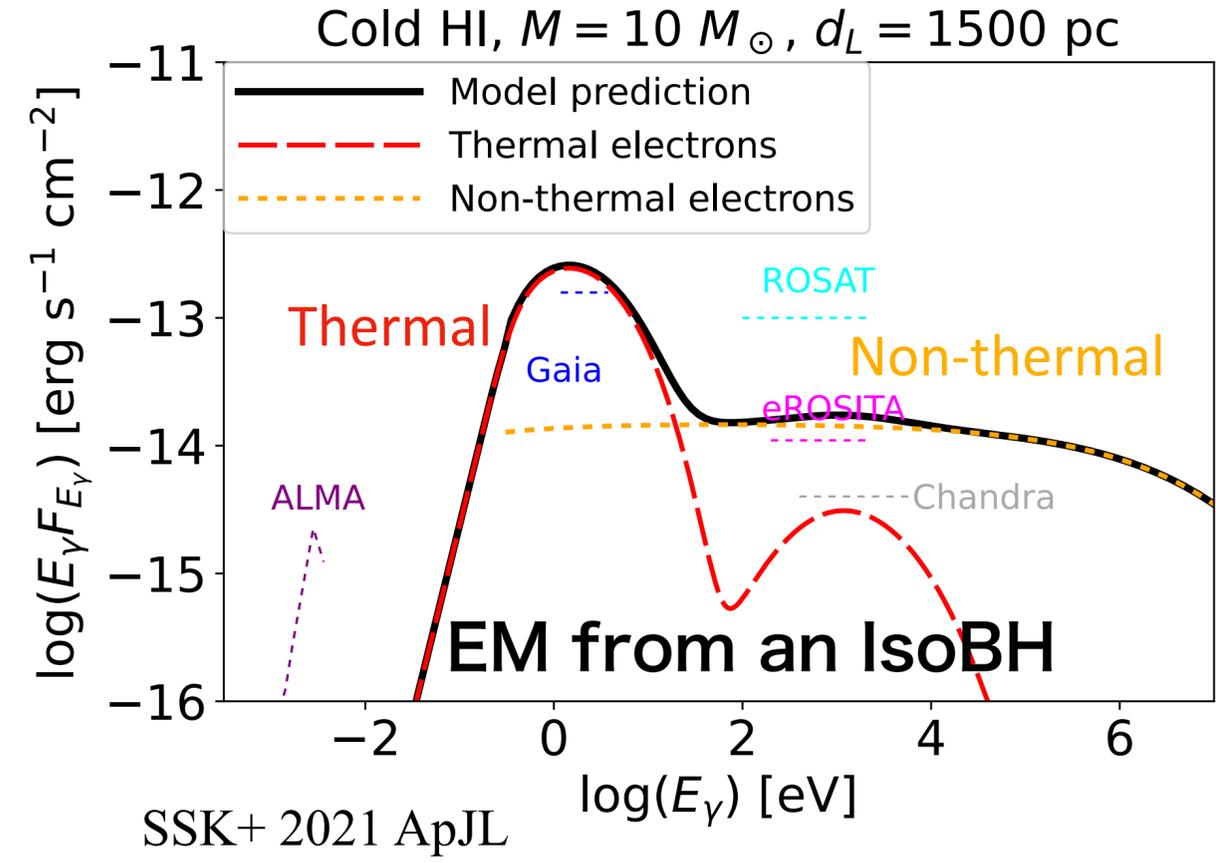
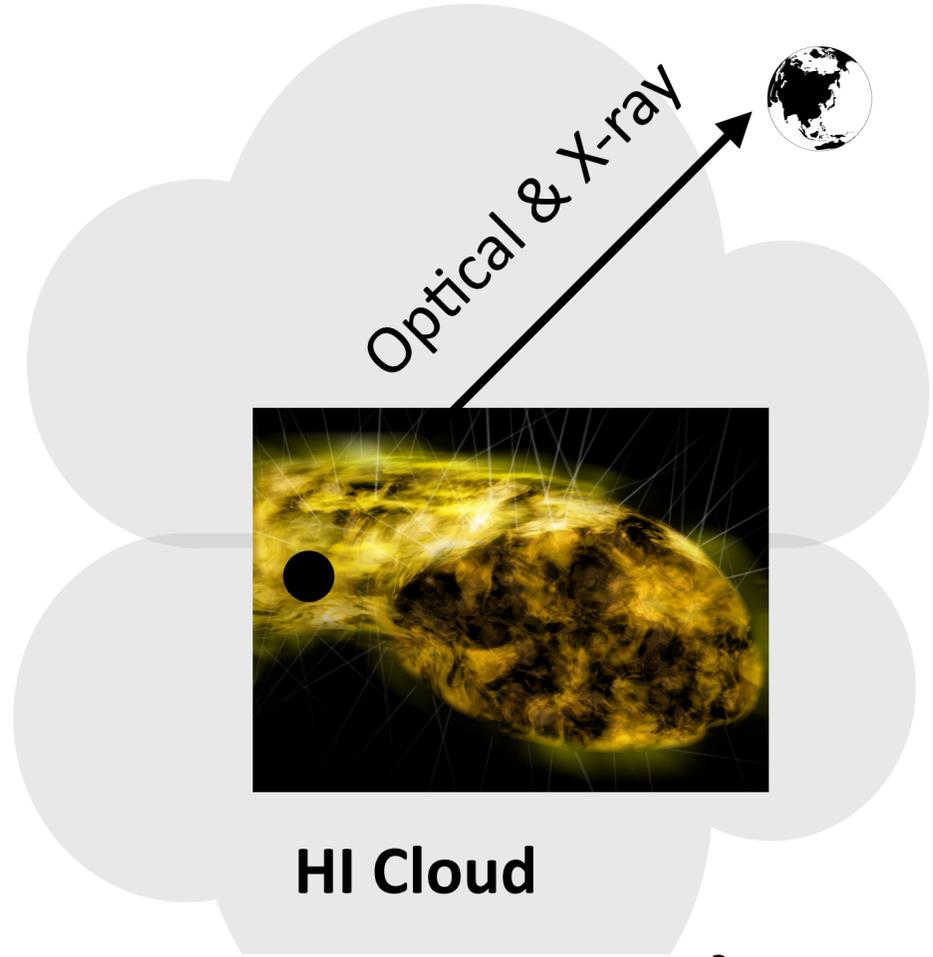
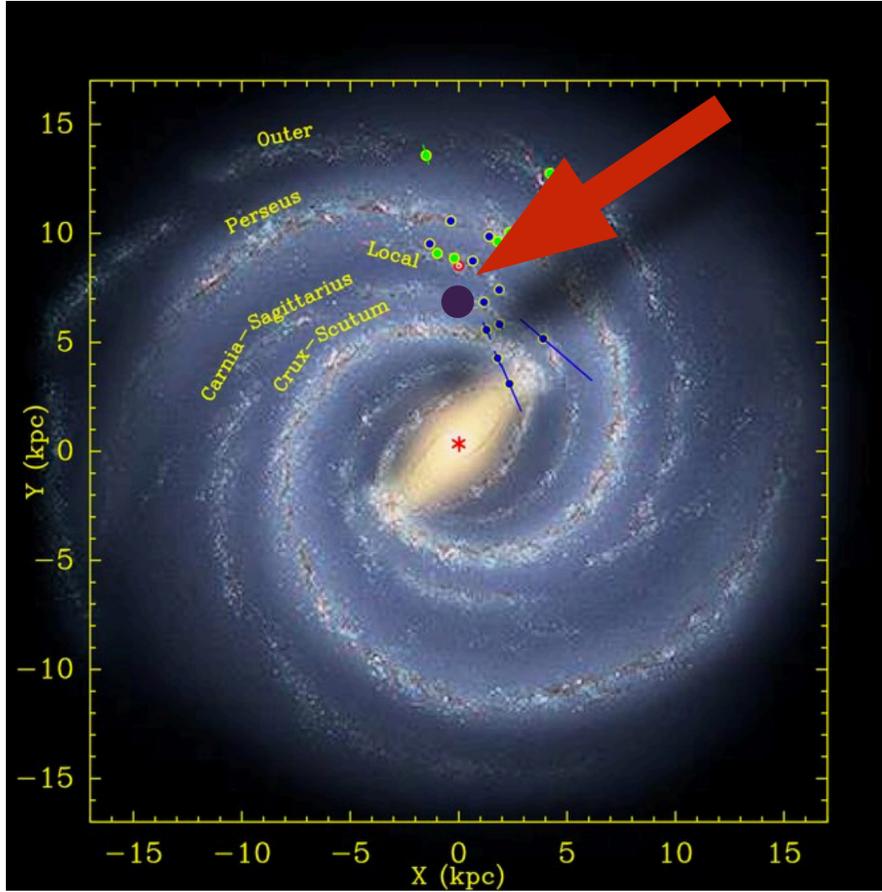
References

SSK, Kashiyama, Hotokezaka, 2021, ApJL, 922, L15
 SSK, Tomida, Kobayashi, Kin, Zhang 2025, ApJL, 981, L36
 SSK, Murchikova, Sahu, 2025, ApJ, 986, 135

EM emission from Isolated Black Holes



IsoBH OGLE-2011-BLG-0462

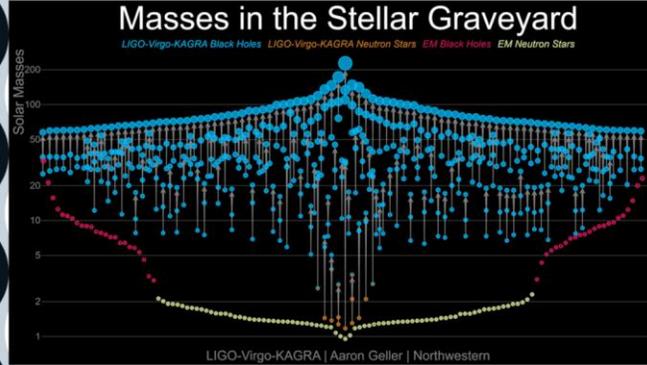


P12 First star remnants for gravitational sources

Tomoya Kinugawa (Shinshu University)

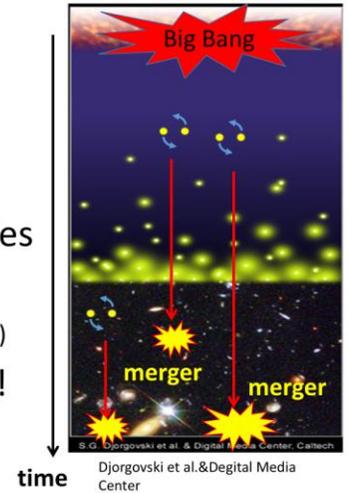
Masses of GW events

- GW events show that there are many massive BHs ($\geq 30 M_{\text{sun}}$).
- On the other hand, the typical mass of BHs in X-ray binaries is $\sim 10 M_{\text{sun}}$.



Pop III BBH remnants for gravitational wave

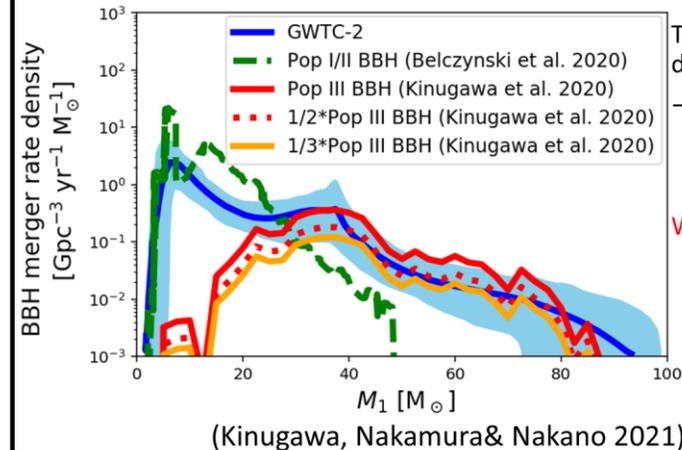
- Pop III stars were born and died at $z \gtrsim 10$.
- The typical merger time of compact binaries $\sim 10^8 - 10^{10} \text{ yr}$
 $dN/dt \propto t^{-1}$ (Kinugawa et al. 2014, 2020 Inayoshi et al. 2017)
- We can see Pop III BBH at the present day!



Wind mass loss & IMF

- Pop I (=Solar metal stars)
Typical mass is small ($IMF \propto M^{-2.35}$, $0.1 M_{\text{sun}} < M < 100 M_{\text{sun}}$)
Stars lose a lot of mass due to the strong stellar wind
- Pop II (Metal < 0.1 Solar Metal)
Typical mass is same as Pop I
But, weak wind mass loss
- Pop III (No metal)
Pop III stars are **the first stars** after the Big Bang.
Typical mass is more massive than Pop I, II
 $M_{\text{pop III}} \sim 10 - 100 M_{\text{sun}}$
No wind mass loss due to no metal.

Comparison with mass distributions of observed BBHs



The mass distribution might distinguish Pop III from Pop I/II

→ The evidence of Pop III?
Not yet

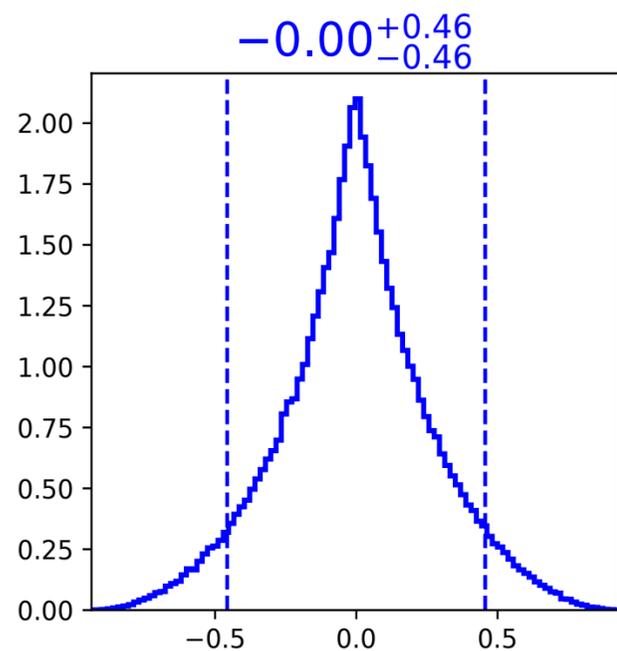
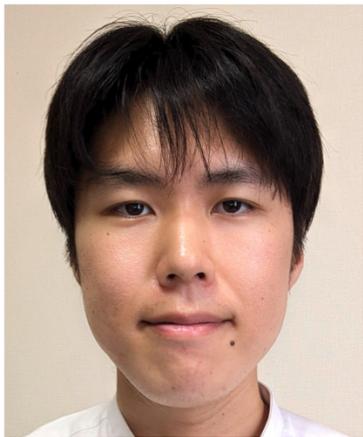
What is the smoking gun of Pop III?

→ High redshift mergers
(ET, CE, DECIGO)
Peak of BBH merger rate
Pop III: $z \sim 10$
Pop I/II: $z \sim 2-3$

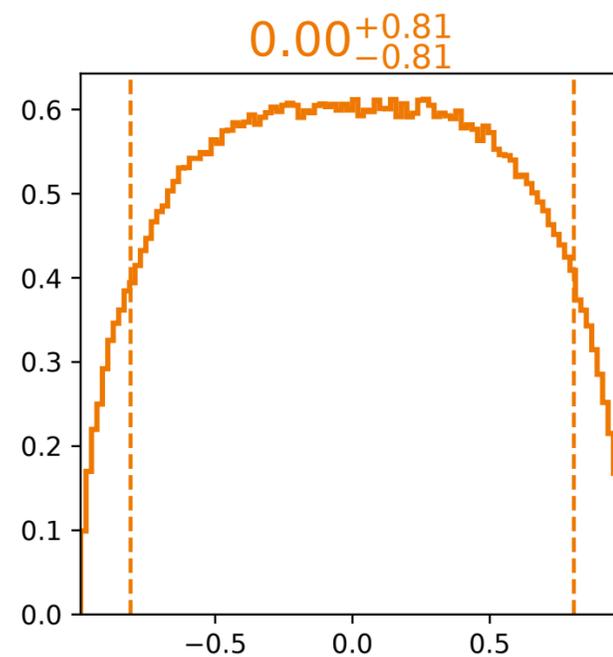
Reducing Spin-Prior Bias in Population Inference of Binary Black Holes

Kazuya Kobayashi^A, Masaki Iwaya^{B,A}, Soichiro Morisaki^A, Kenta Hotokezaka^C, Tomoya Kinugawa^{D,C}

ICRR^A, GEI Cardiff University^B, RESCEU^C, Shinshu University^D

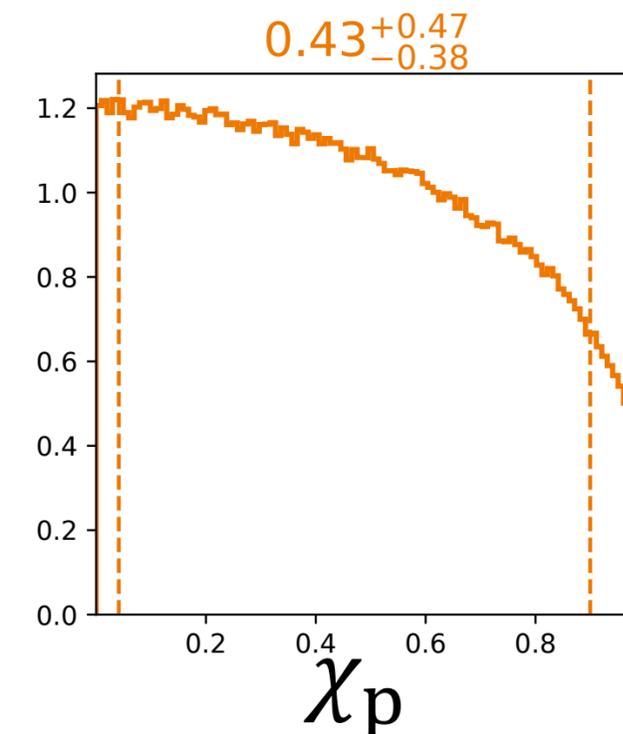
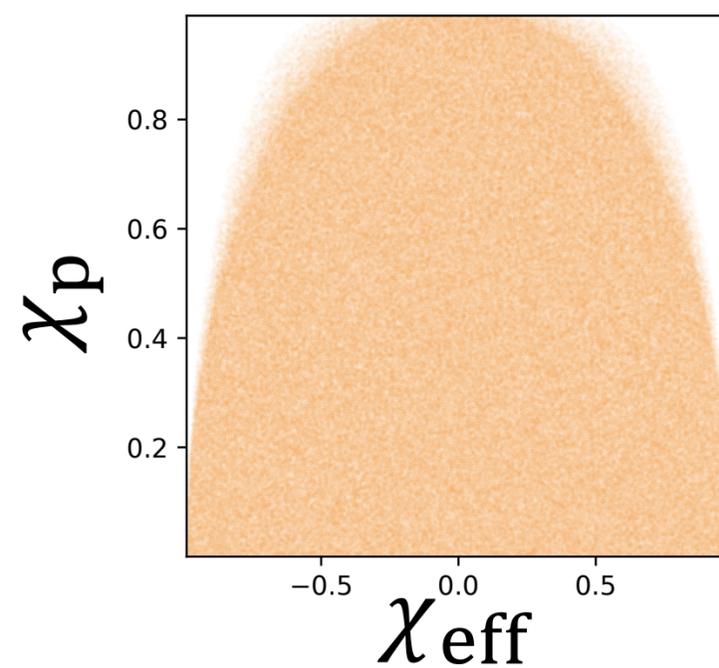
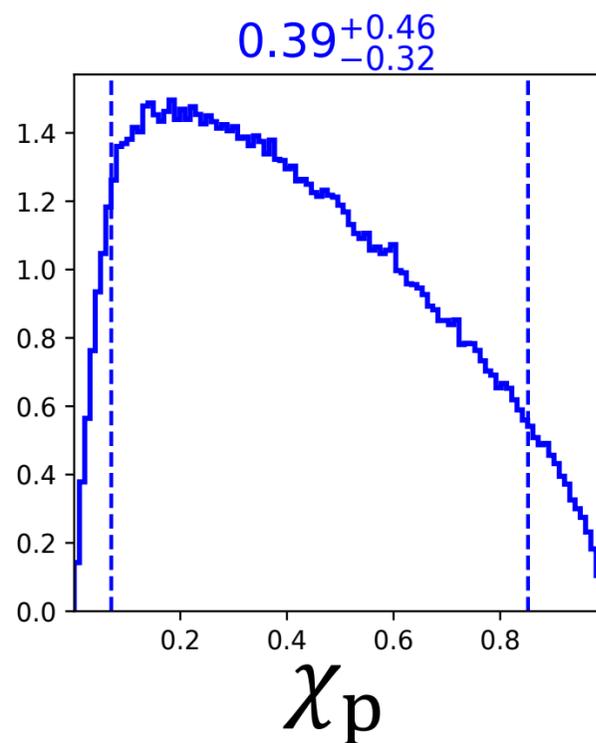
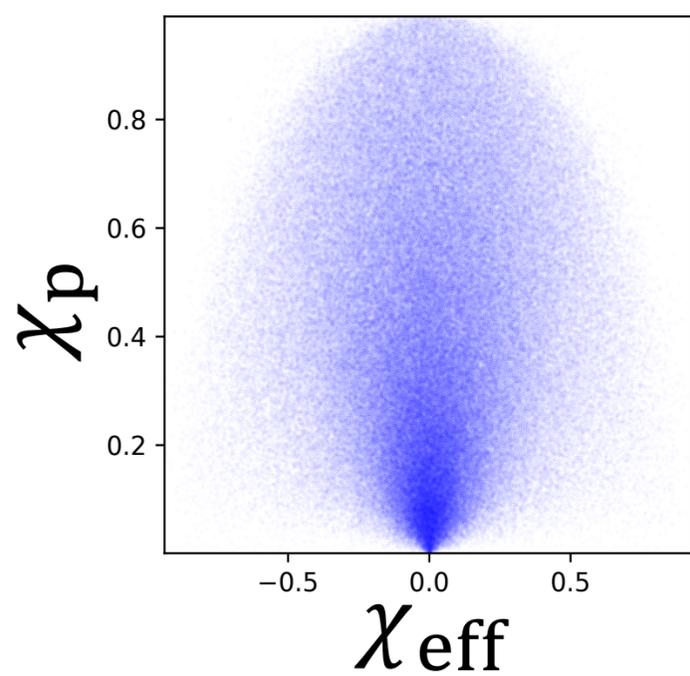
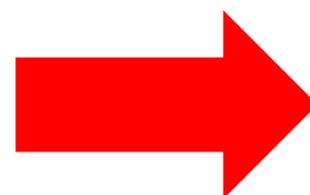


Conventional prior



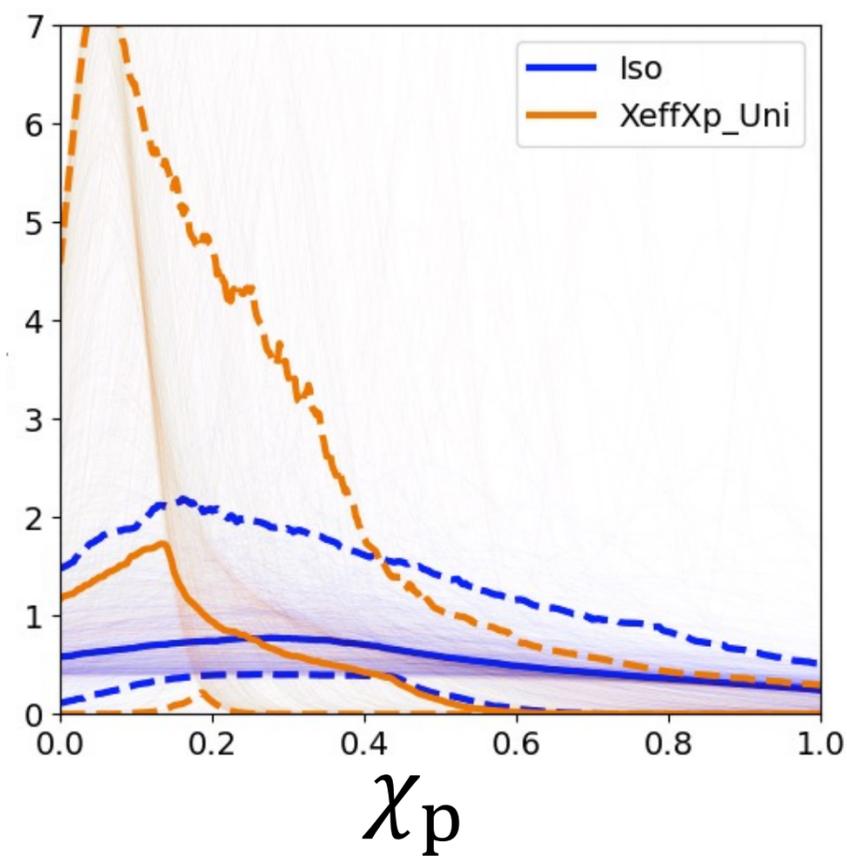
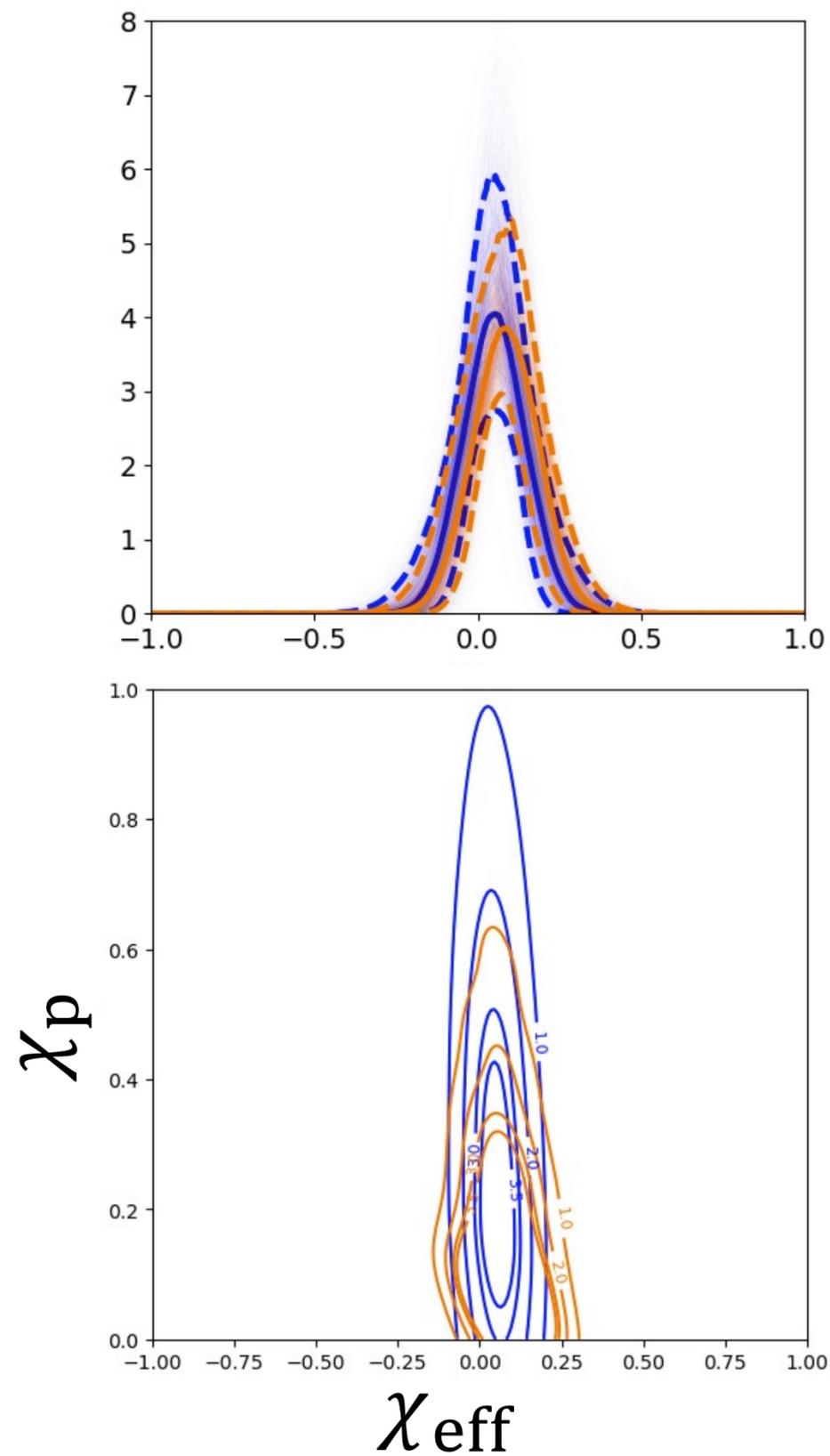
New prior

Change



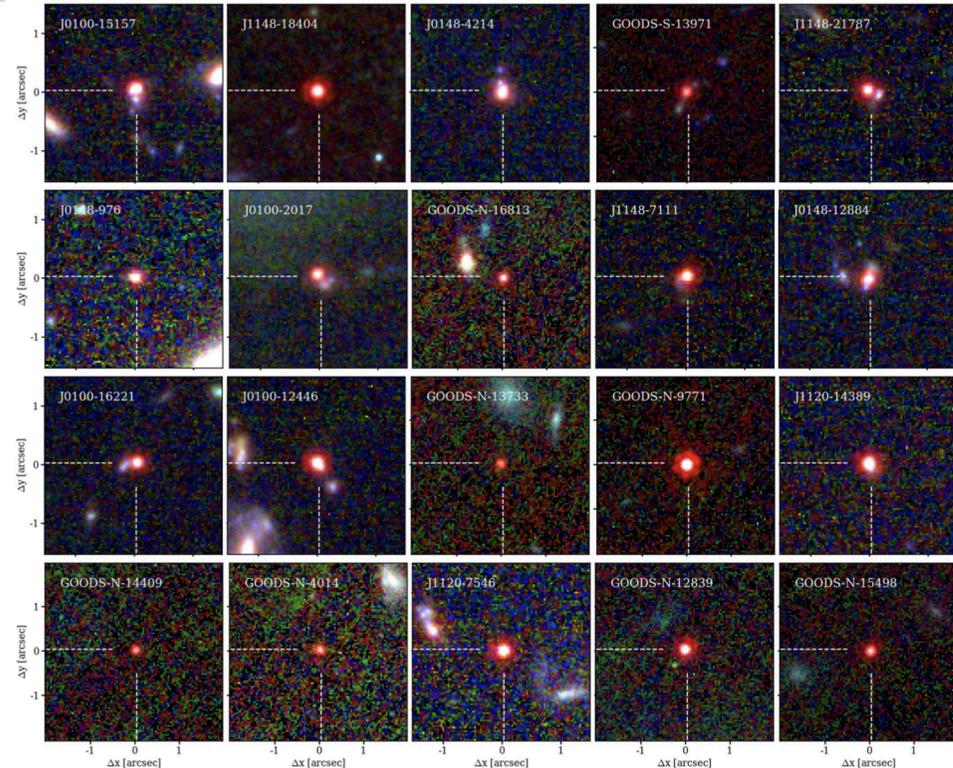
Effective spins population distribution for binary black holes observed in O3

Blue: conventional prior
Orange: new prior



Potential High-energy Neutrino Sources

Riku Kuze (YITP), Kunihiro Ioka (YITP), Kohta Murase (Penn state), Shigeo S. Kimura (Tohoku), Kohei Inayoshi (Peking)



Matthee+2024

Little Red Dots (LRDs)

- Red & Compact galaxies ($R \lesssim 100$ pc)
- Population: $\Phi(z \sim 4) \simeq 10^{-4.5} \text{ cMpc}^{-3}$
- **No counterparts in radio and X-ray bands**
- Broad emission lines: $\sigma_v \sim 1000 \text{ km s}^{-1}$

$$L_{\text{Bol}} \simeq 10^{44-45} \text{ erg s}^{-1} \rightarrow M_{\text{BH}} \sim 10^{6-7} M_{\odot}?$$

SMBH embedded with dense gaseous envelope

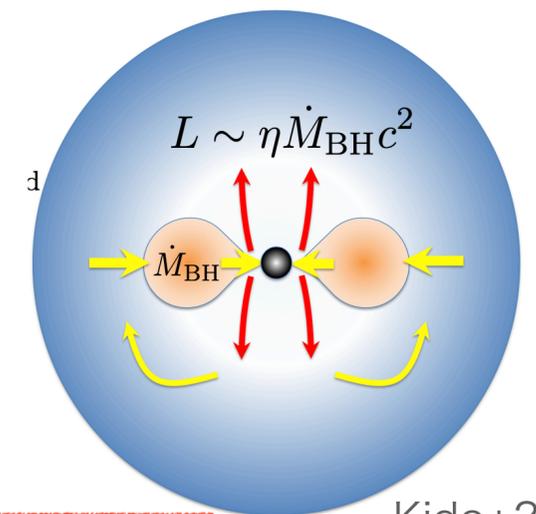
Dense envelope

- Efficient neutrino emission
- Gamma rays cannot escape

×

$$n \sim 10^{-4.5} \text{ cMpc}^{-3},$$

$$L \sim 10^{44-45} \text{ erg s}^{-1}$$



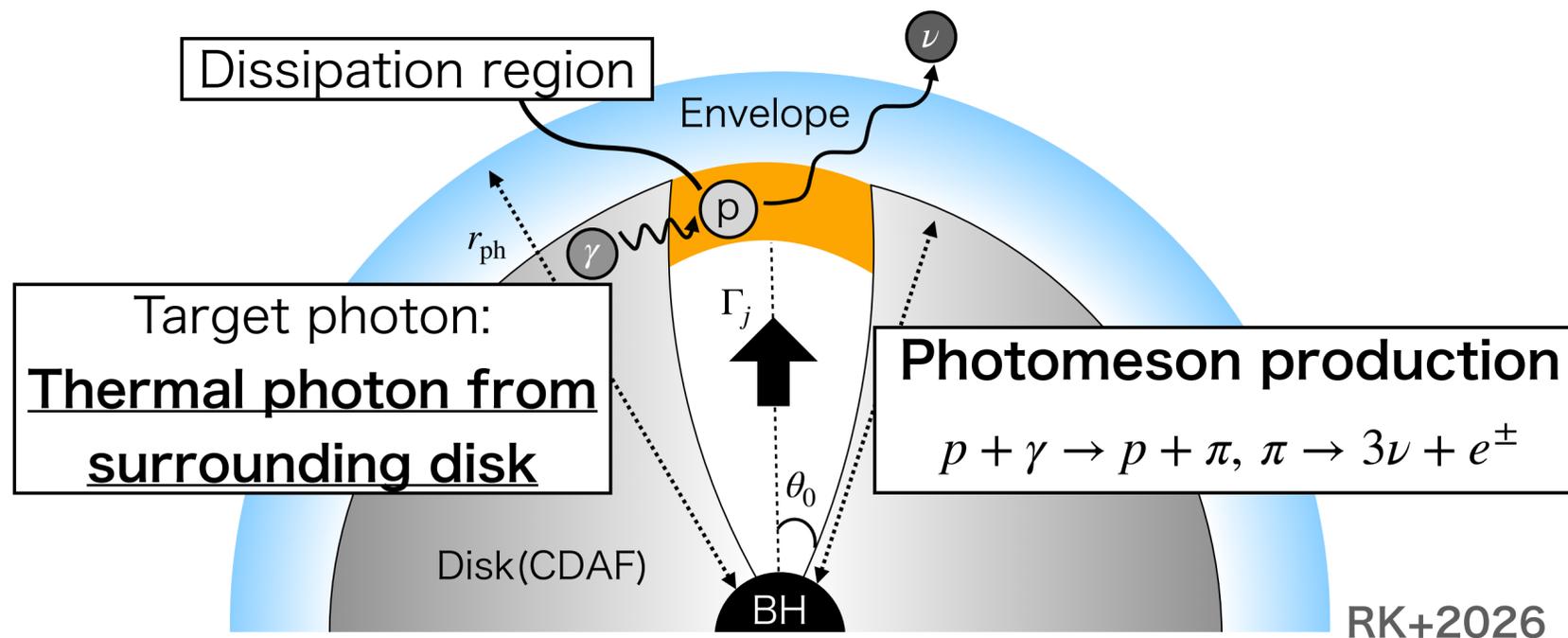
Kido+2025

Compatible with the observed diffuse neutrino background

Potential High-energy Neutrino Sources

Riku Kuze (YITP), Kunihiro Ioka (YITP), Kohta Murase (Penn state), Shigeo S. Kimura (Tohoku), Kohei Inayoshi (Peking)

Jets or outflows launched from the central BH propagate and dissipate within the BH polar funnel.



Target photon:
Thermal photon from
surrounding disk

Photomeson production
 $p + \gamma \rightarrow p + \pi, \pi \rightarrow 3\nu + e^\pm$

RK+2026

We estimate the contribution to the diffuse neutrino background.

3 types of LRD
population model

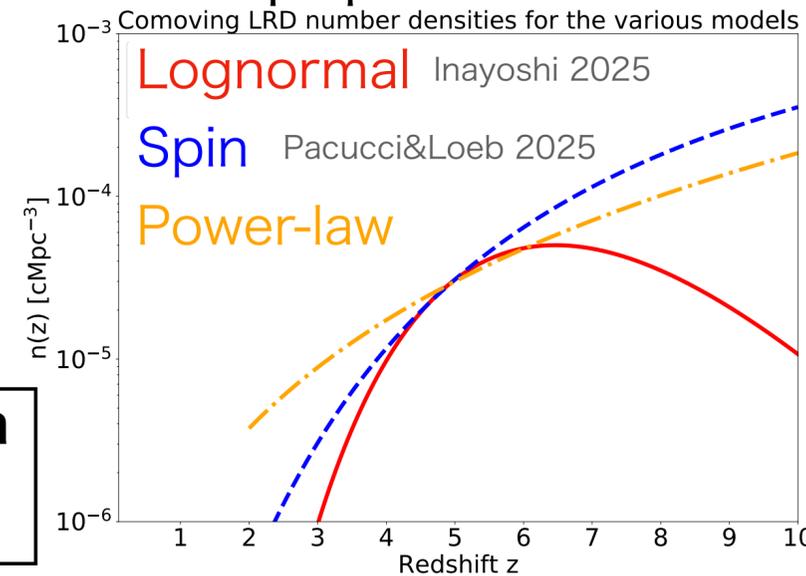
×

$$L_\nu \sim \frac{1}{8} f_{\text{bol}} \epsilon_p L_j$$

$$\sim 10^{41} \text{ erg s}^{-1} f_{\text{bol},-1} \epsilon_{p,-1} L_{j,44}$$

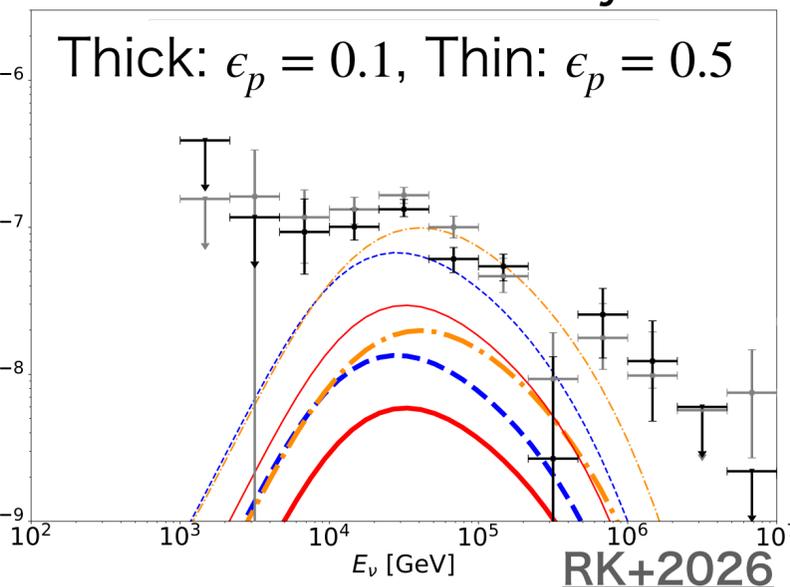
Results

3 LRD population models



Diffuse intensity

Thick: $\epsilon_p = 0.1$, Thin: $\epsilon_p = 0.5$



RK+2026

Conclusion:

- A dense BH-envelope provides the environment for efficient neutrino production and prevents gamma rays from escaping.
- LRDs can contribute to a non-negligible fraction of the diffuse neutrino background in the TeV-sub-PeV energy range.

Binary strange quark star merger in fully general-relativistic simulation

Ho Ching Luk, Lap Ming Lin

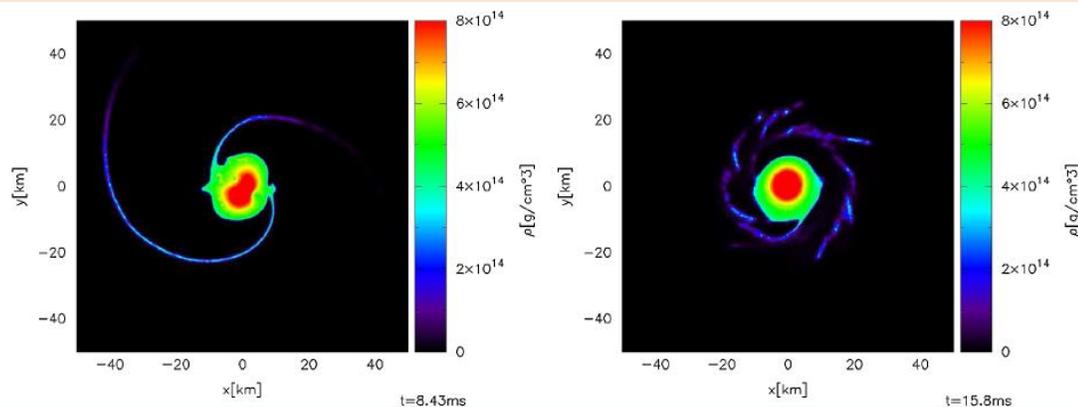
Chinese University of Hong Kong

$$\text{Quark eos: } P = B \left[\left(\frac{\rho}{\rho_s} \right)^{1+c_{ss}} - 1 \right] \longrightarrow \text{Extremely sharp surface}$$

Previous quark star merger simulation

SPH

- A. Bauswein, R. Oechslin, and H. Janka (2010)

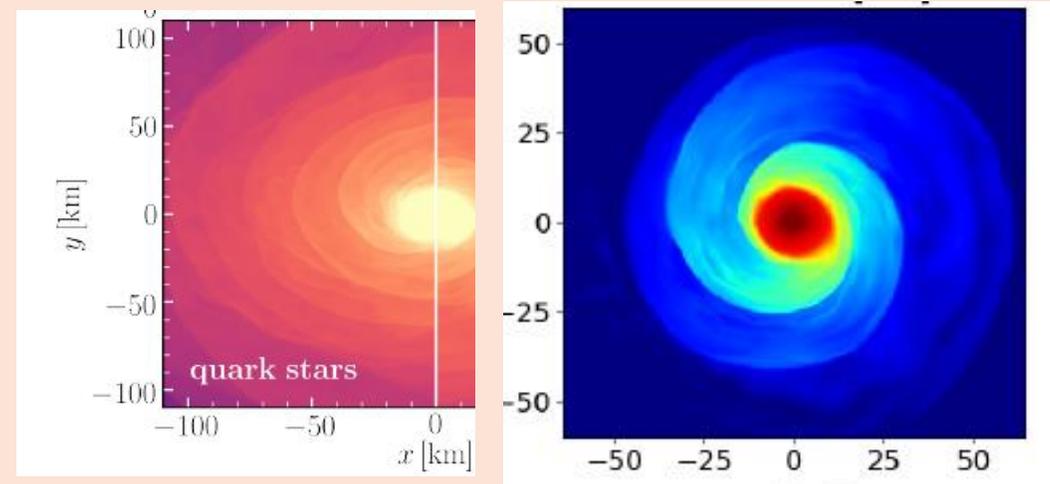


Prompt collapse: explore M_{th} but no GW extracted

- E. Zhou, K. Kiuchi, M. Shibata, A. Tsokaros, and K. Uryu (2022)

Grid base but add polytropic tail

- Z. Zhu and L. Rezzolla (2021)
- F. Grippa, A. Prakash, D. Logoteta, D. Radice, and I. Bombaci (2025)



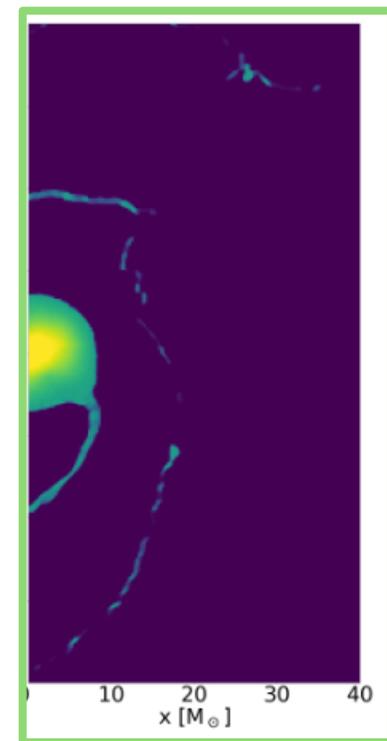
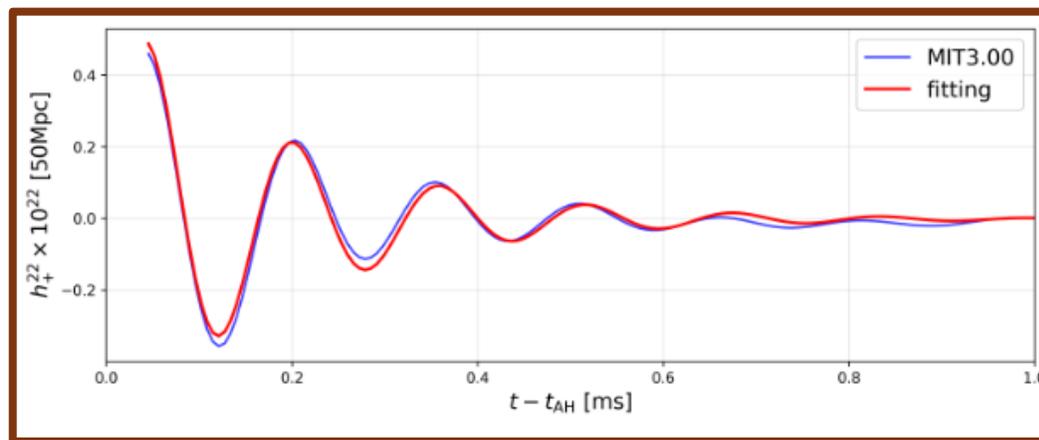
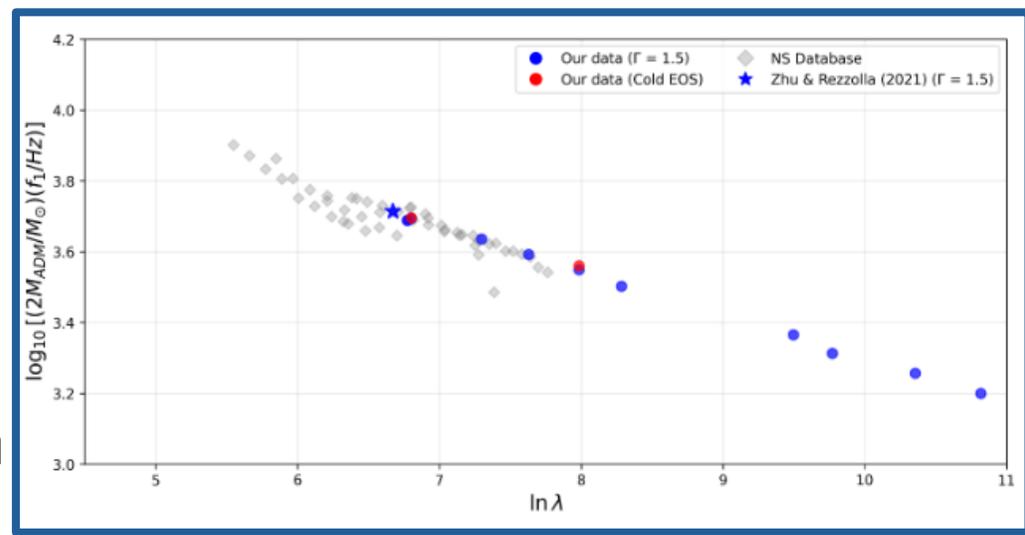
Our result

What's different?

- Positively preserving Riemann solve
- Piecewise parabolic method reconstruction
- Hydro excision inside apparent horizon
- Dust atmosphere
- Modified AMR

What's new result?

- Spiral arm and quark droplet in grid base code
- Larger mass range f_{max} , f_1 and f_2 ($0.5M_{\odot}$ to $1.5M_{\odot}$)
- Prompt collapse waveform (Compared to BH qnm)



YITP 2026

Particle-in-Cell (PIC) Simulations on Collisionless Shocks and Particle Acceleration in Black Hole Coronae

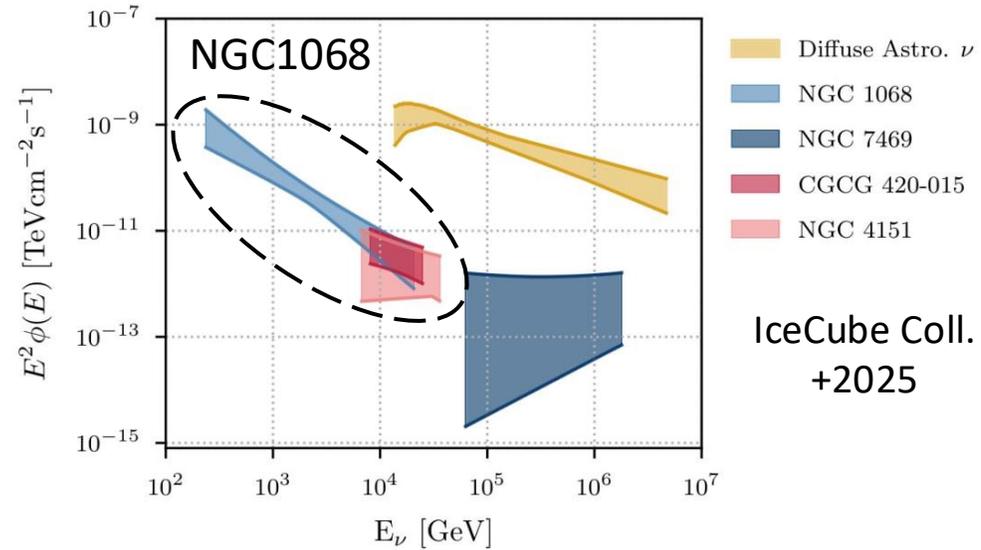
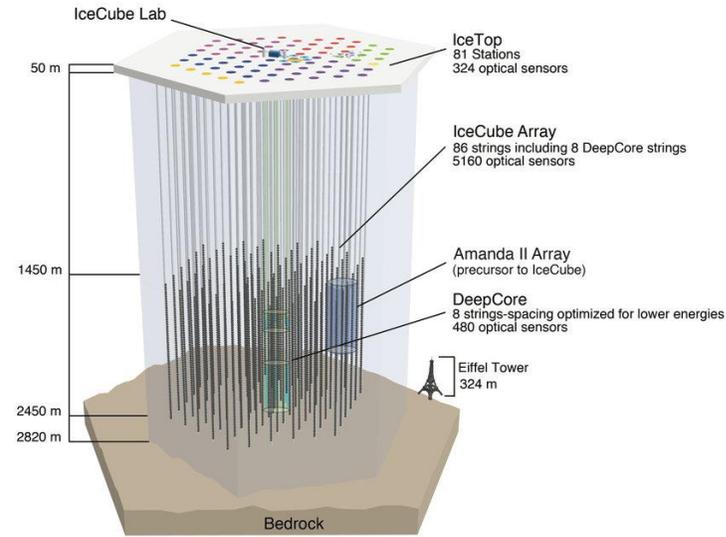
Nhat-Minh Ly¹, Takayoshi Sano^{1,2}, Yoshiyuki Inoue², Yasuhiko Sentoku¹

1. Institute of Laser Engineering, The University of Osaka
2. Department of Earth and Space Science, Graduate school of Science, The University of Osaka



(In Press) M. N. Ly, Y. Inoue, Y. Sentoku, and T. Sano, *Proton Acceleration by Collisionless Shocks in Supermassive Black Hole Coronae: Implications for High-Energy Neutrinos*, arXiv:2601.01999.

High-energy neutrinos from BH Coronae



IceCube survey of high-energy ν from Radio Quiet (RQ) AGNs

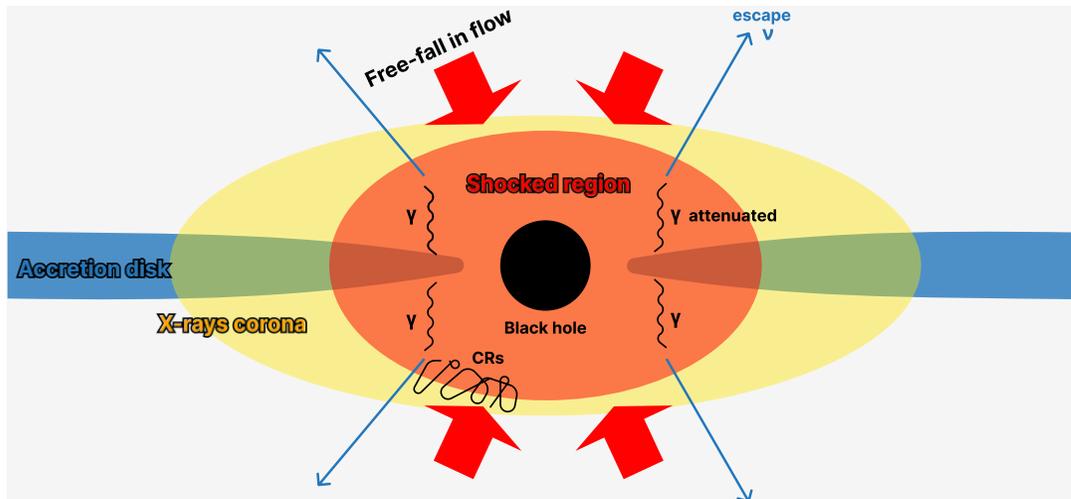
NGC1068:

Neutrino ~ 10 TeV; Proton ~ 100 TeV

No γ -rays

$\rightarrow \nu$ produce in BH Coronae

Schematic shocks picture in BH coronae



Can shocks accelerate proton to 100 TeV in BH coronae?

- **Astrophysics: could be** (Inoue et al. 2019, 2020)
- **Plasma physics: Not sure**
 - Trans-Sonic may not be efficient in accelerate protons
 - Dynamics of accelerated protons and electrons are implied by plasma physics



Coupling Electromagnetism to Torsion: Black Holes and Spin-Charge Interactions



arXiv:2507.02362 — To appear in *Phys. Rev. D*

Jorge Maggiolo Tapia

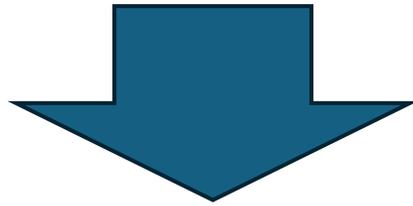
Universidad Técnica Federico Santa María, Chile

Email: jorge.maggiolo.t@gmail.com

Theory

We consider the minimal gravitational theory that consistently couples electromagnetism to torsion while preserving $U(1)$ gauge invariance. The lagrangian density is

$$\mathcal{L} = -R - k_1 F_{\mu\nu} F^{\mu\nu} + k_2 \tilde{R}^2 + k_3 F_{\mu\nu} \tilde{R}^{\mu\nu}.$$



$$ds^2 = \Psi_1(r) dt^2 - \Psi_2(r) dr^2 - r^2 (d\theta^2 + \sin^2 \theta d\phi^2)$$

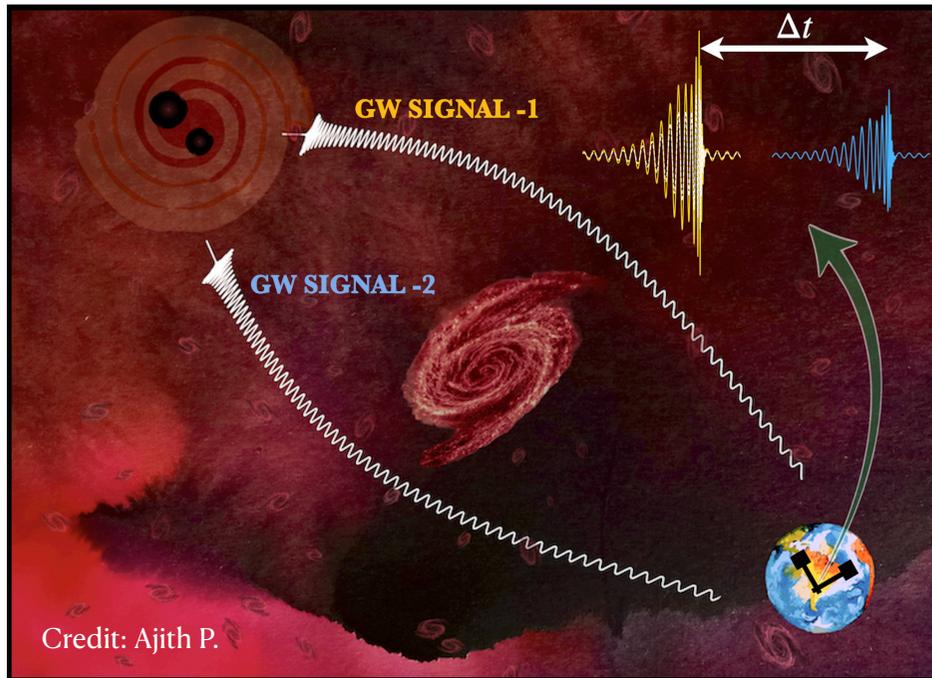

$$\Psi_1(r) = ?$$


$$\Psi_2(r) = ?$$

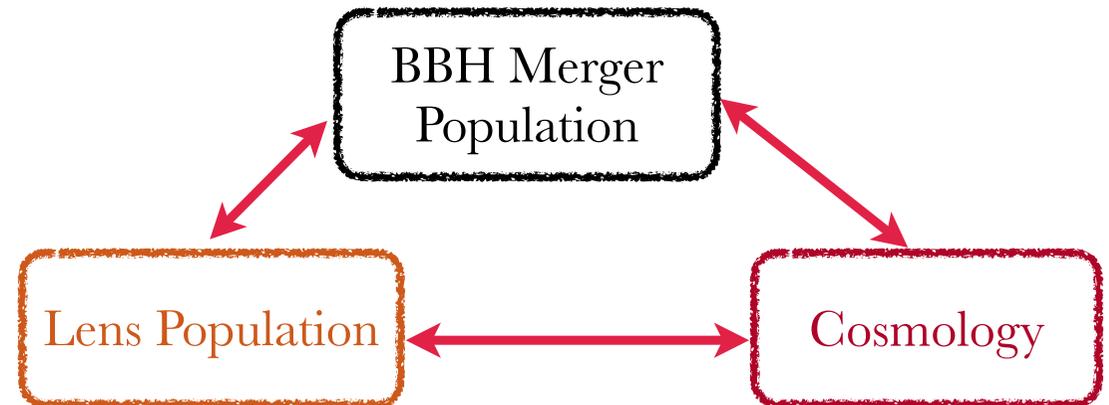
18

Inferring Cosmology and Nature of Dark Matter using Strongly Lensed Binary-Black-Hole mergers: Prospects for the Near Future

Koustav N. Maity, Souvik Jana, Tejaswi Venumadhav, Ankur Barsode & Parameswaran Ajith : arXiv:2512.15168



~ similar copies of GW signals arrive at us at different times



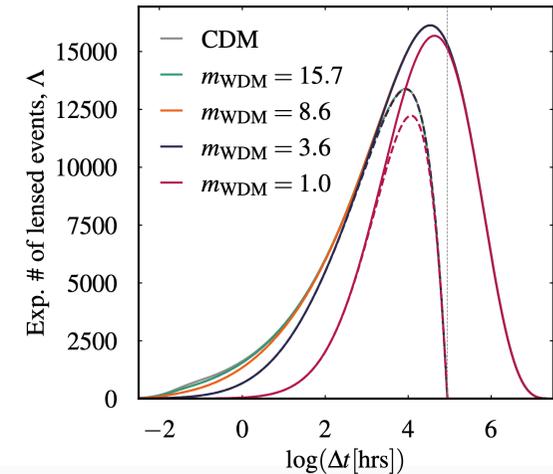
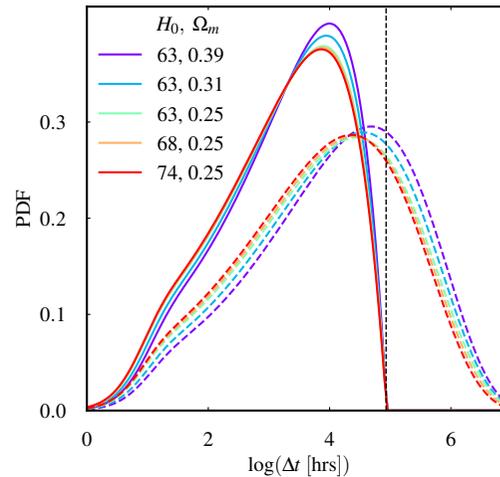
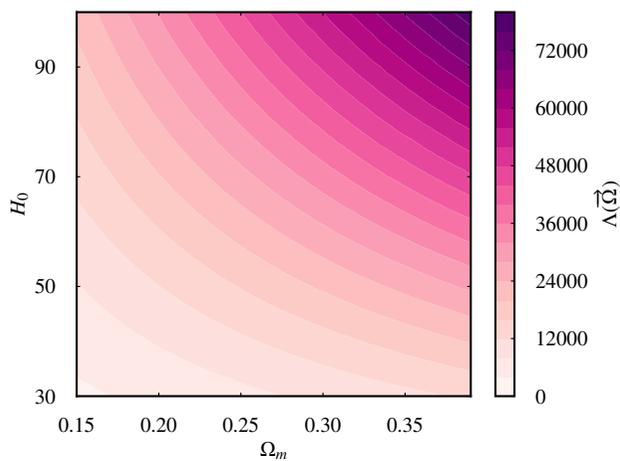
- * Changing cosmology affects structure formation and the distances.
- * The lens population changes if we change our dark matter model.

Both affect the delay-time distribution and lensing fraction.

18

Inferring Cosmology and Nature of Dark Matter using Strongly Lensed Binary-Black-Hole mergers: Prospects for the Near Future

Potential realization was already shown in XG, but **what could be done with immediate upgrades?**

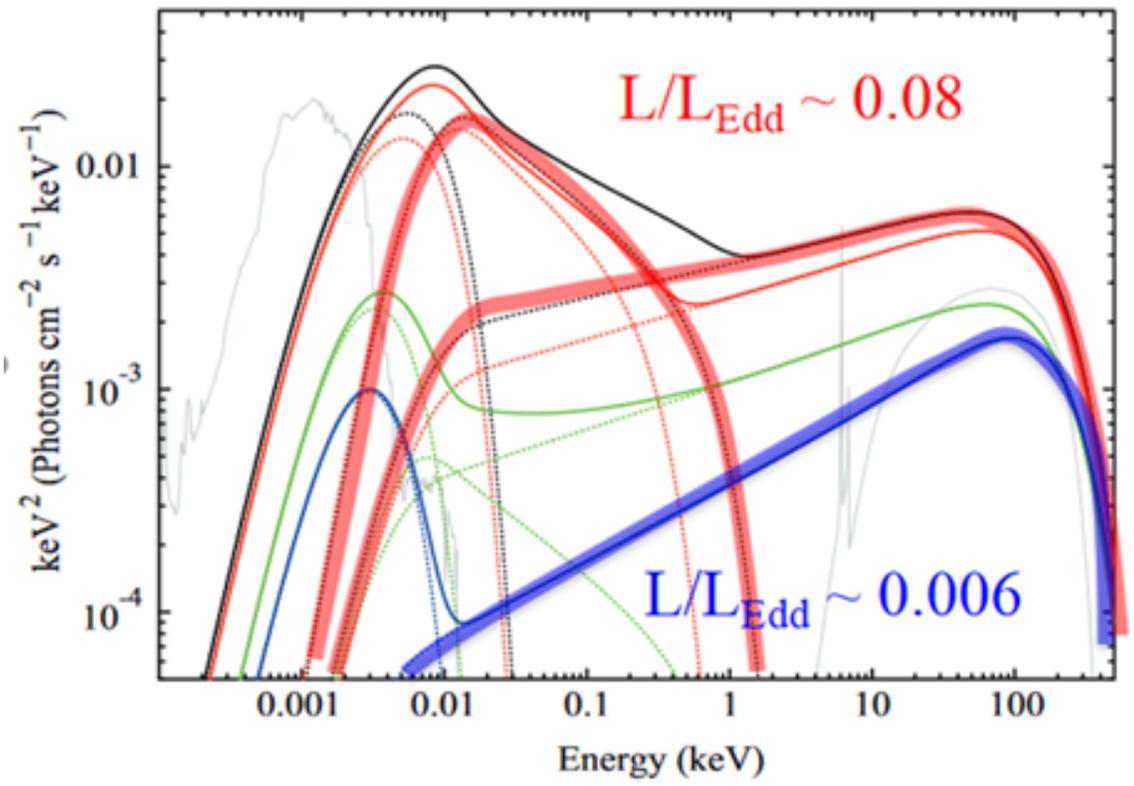


Credit: Jana. et al. (2023, 2025)

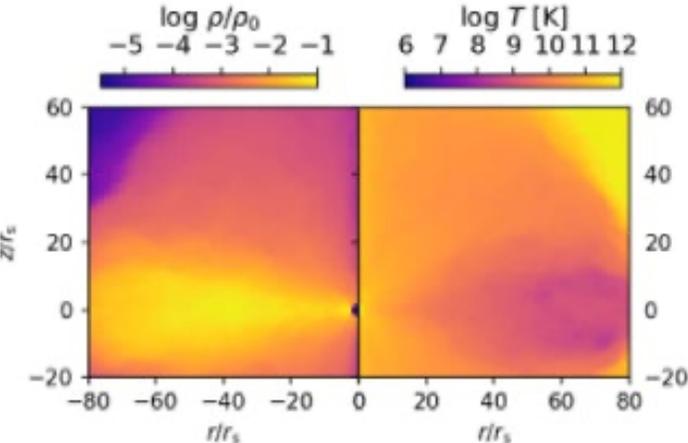
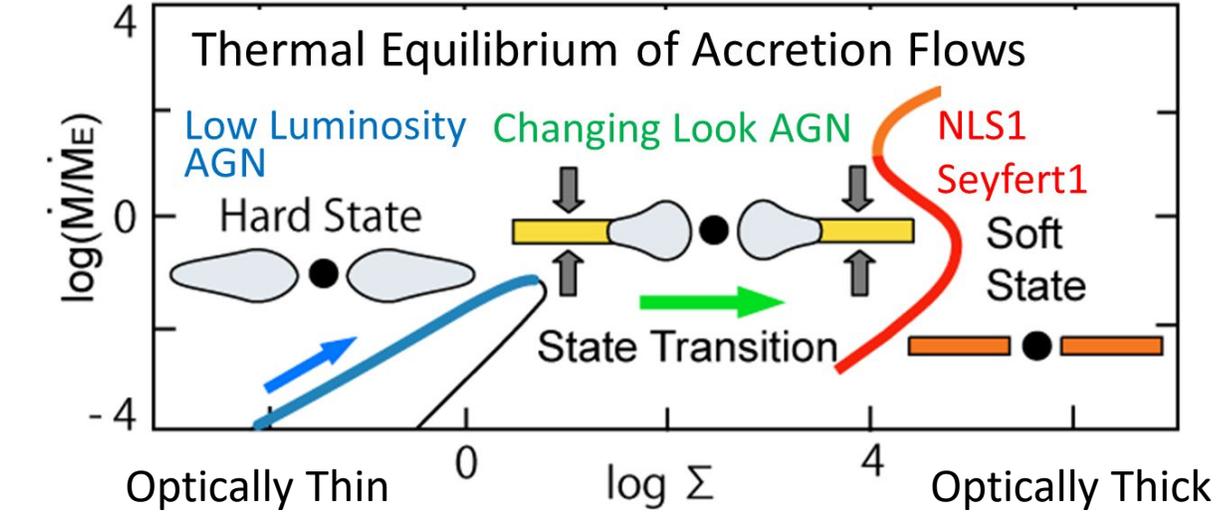
Precisely accounts for observing windows, enforces a joint SNR detectability criterion for multiple images, and incorporates time-varying detector-network sensitivity.

Global Radiation Magnetohydrodynamic Simulations of Black Hole Accretion Flows during State Transitions in AGN

Ryoji Matsumoto(Chiba Univ.), T. Igarashi(NAOJ/Rikkyo Univ.), H.R. Takahashi(Komazawa Univ.), T. Kawashima(NIT Ichinoseki), K. Ohsuga(Tsukuba Univ.), and Y. Matsumoto(Chiba Univ.)

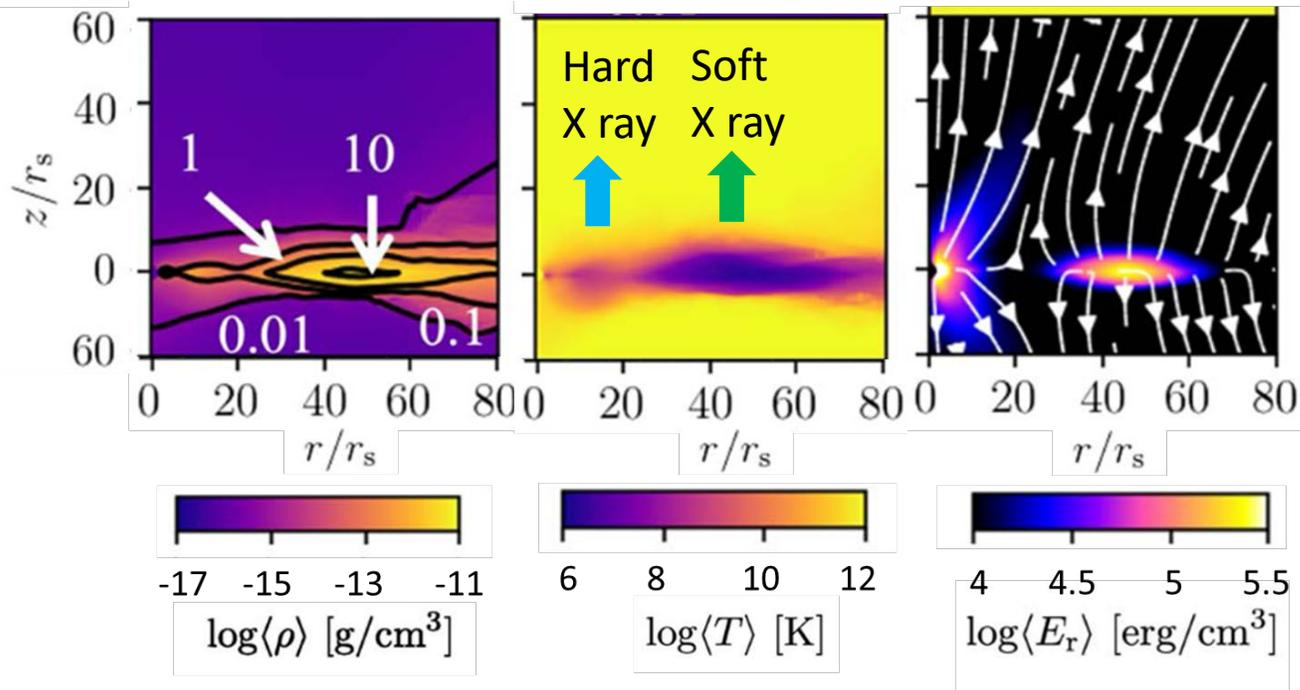


Radiation Spectra of a changing look AGN Mrk1018 (Noda and Done 2018)

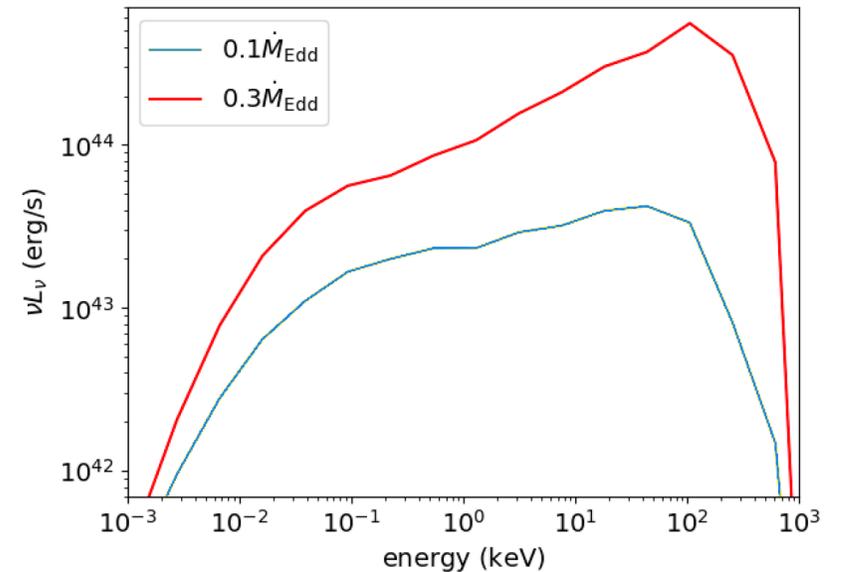
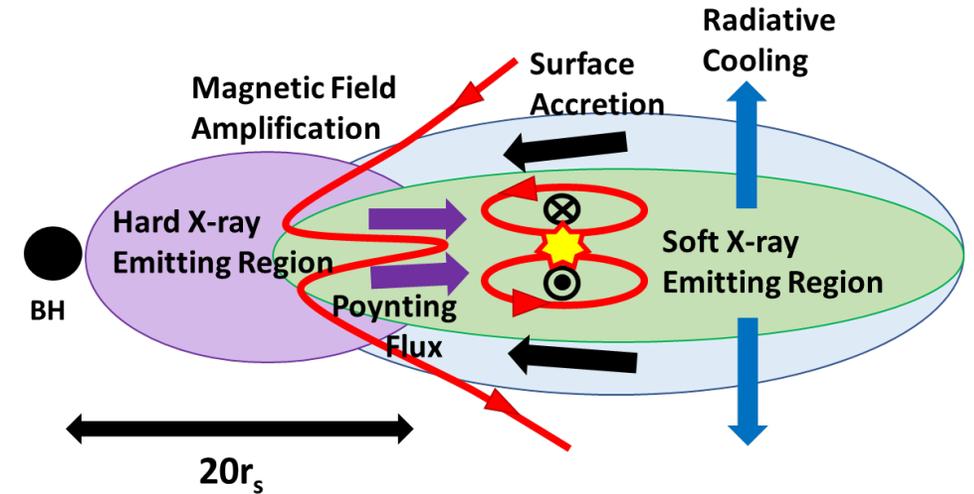


$M_{\text{BH}} = 10^7 M_{\odot}$
 Unit Length $r_s = 3 \times 10^{12} \text{cm}$
 Cooling terms are turned on after RIAF is formed.
 CANS+R (M1 code)
 $(N_r, N_{\phi}, N_z) = (464, 32, 464)$

Numerical Results and Radiation Spectra



- Results of radiation MHD simulation when the accretion rate is $0.1\dot{M}_{\text{Edd}}$ (Igarashi et al. 2024)
- Solid curves in the left panel show the contour of Thomson optical depth.
- Right panel shows the radiation energy density (color) and radiation flux (arrows)
- Soft X-ray emitting region is formed around $40r_s$



Radiation Spectra obtained by RAIKOU code

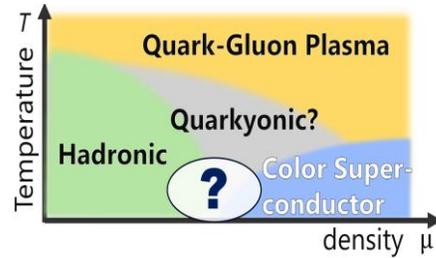
Constrain EoS from Binary Neutron Star GW

BH or NS ?

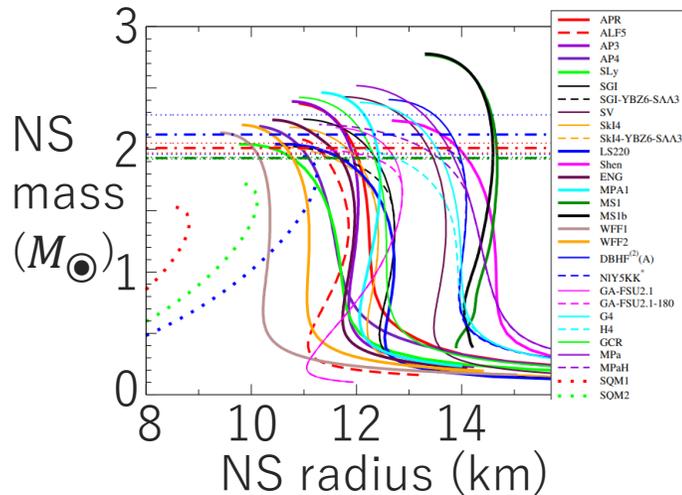
short GRB mechanism ?

QCD phase diagram ?

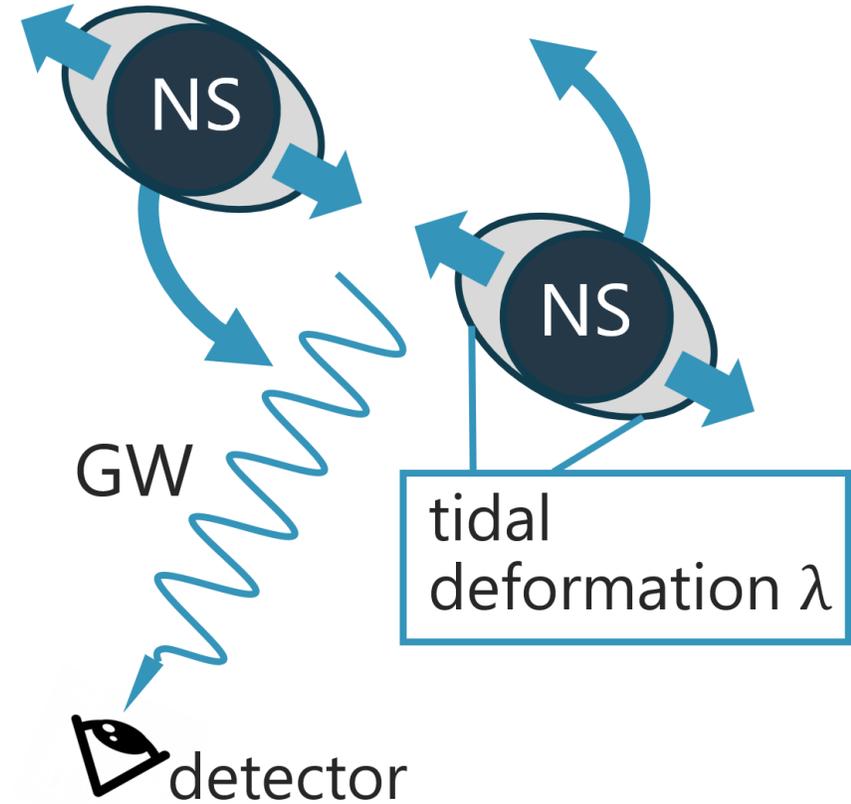
~ $2.4M_{\odot}$



Constrain Neutron Star Equation of State (EoS) !



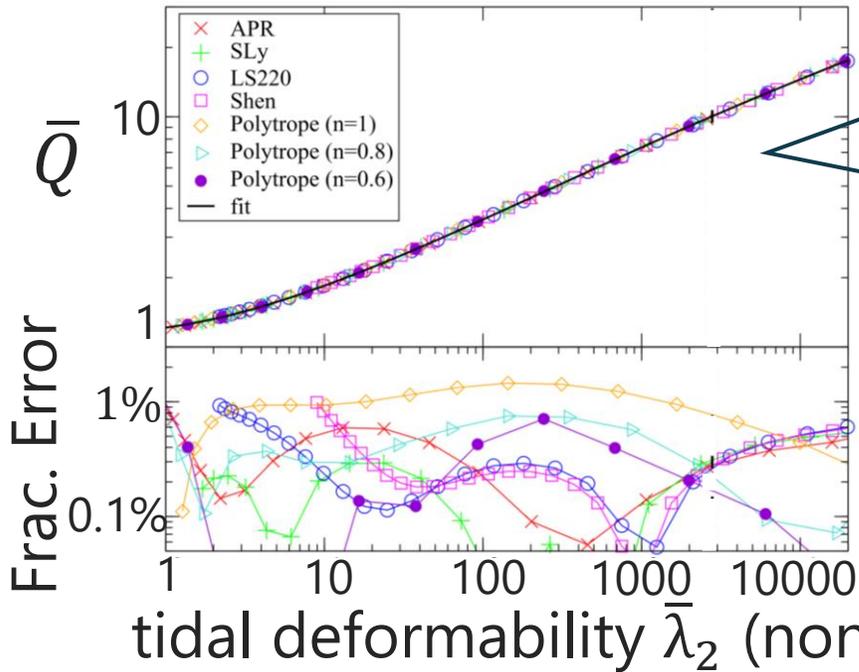
K. Yagi and N. Yunes, *Phys. Rept.*, 681, 1 (2017) (modified)



estimate params (star mass, λ , ...)
more # of params → more error

EoS-independent Universal Relations (URs)

\bar{Q} : spin-induced mass quadrupole (non-dim)



different colors
= different EoS
overlap

Max. error ~ 1 %

tidal deformability $\bar{\lambda}_2$ (non-dim) K. Yagi and N. Yunes Phys. (2013)

Reduce # of params

→ **Improved param estimation accuracy**

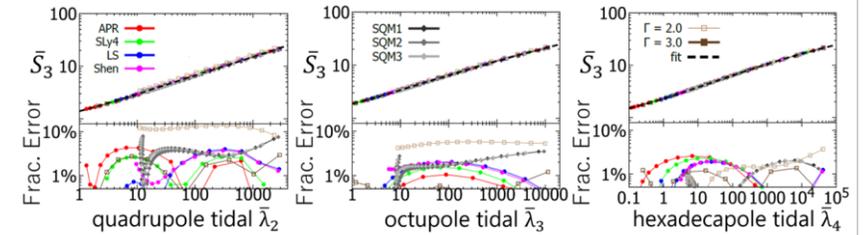
B. P. Abbott et al. *Phys. Rev. X*, 9:011001, (2019)

Reason of URs?

→ **No definitive explanation yet.**

What's New:

O(1)% URs among higher-order multipoles for future observations



← low precision high →

+

characteristic behavior of universal relations