



Heavy Elements in Late-Phase Spectra of Kilonovae

Implications for GW170817

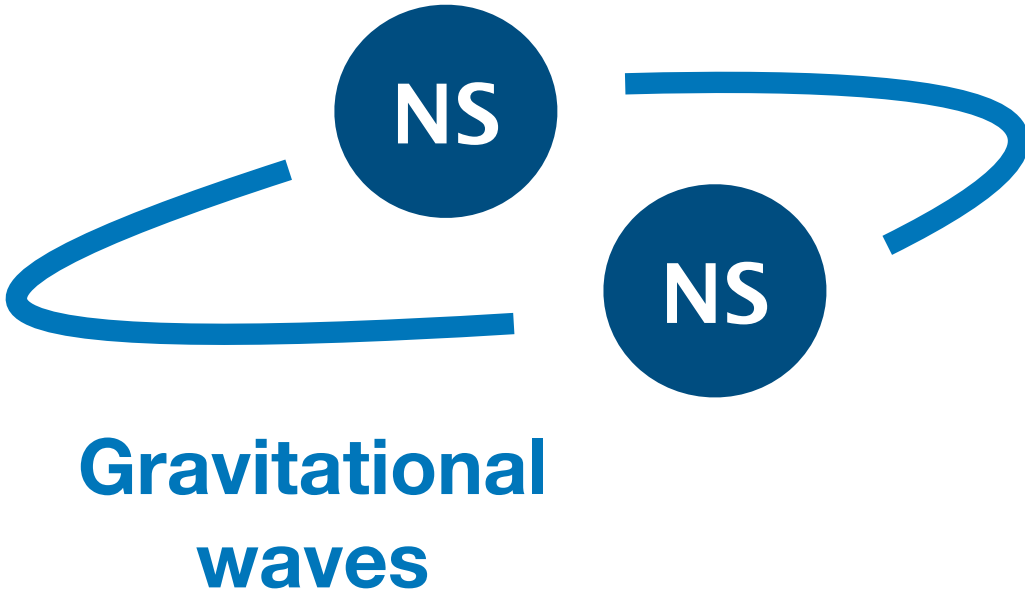
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Tohoku University

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YITP long-term workshop Feb 18th, 2026

Introduction

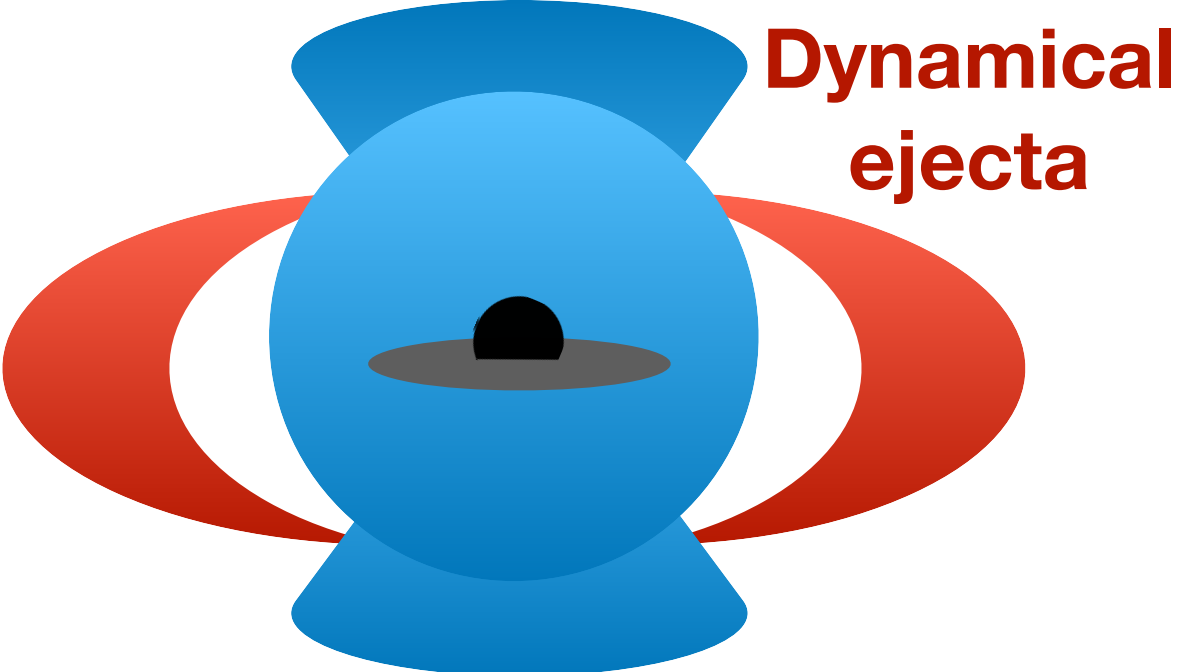
Binary Neutron Stars



Merger

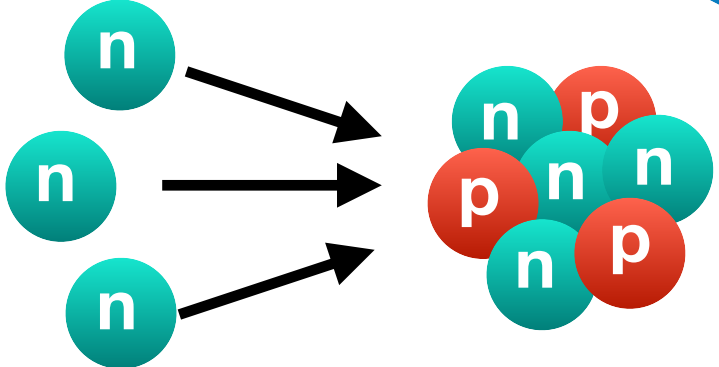
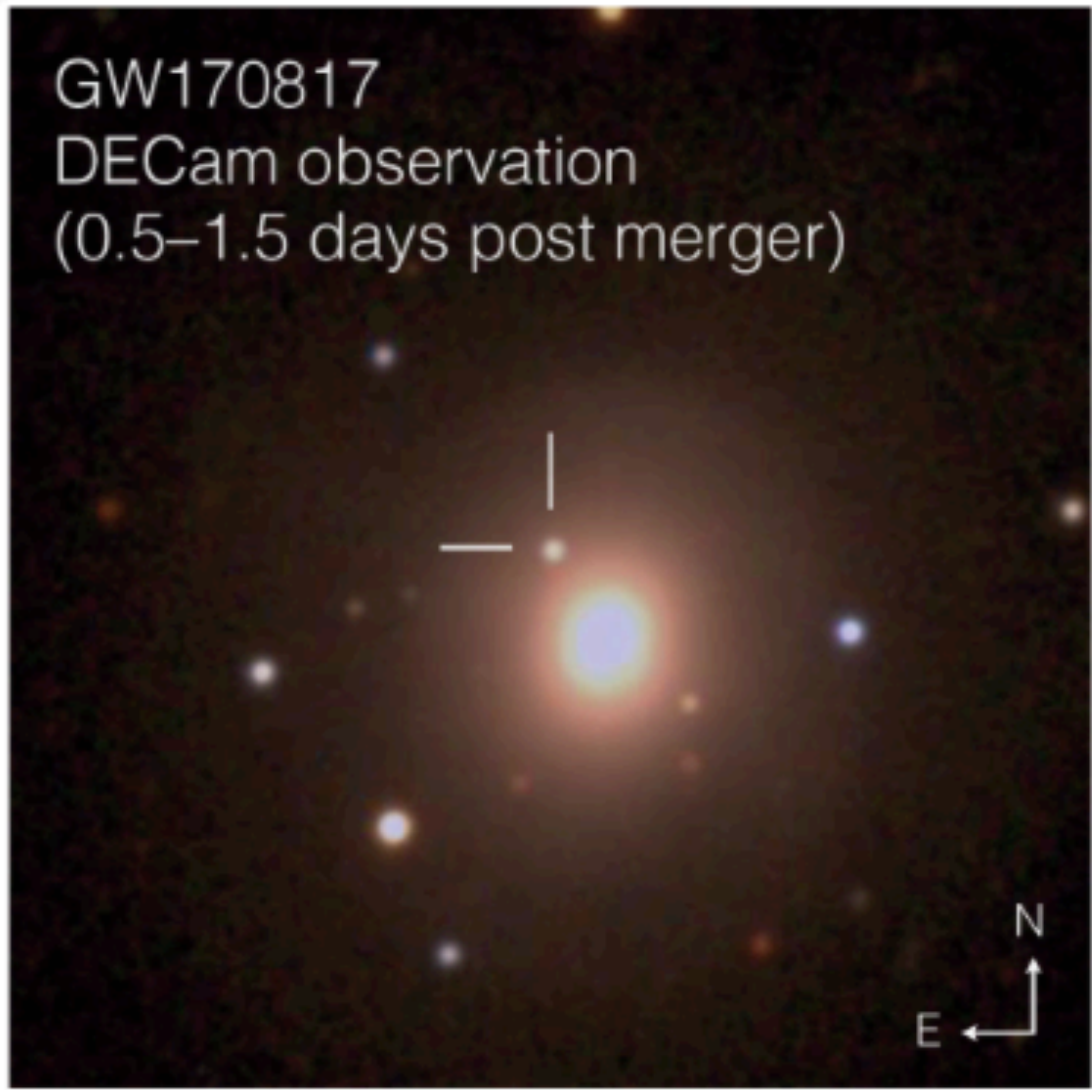


Post-merger ejecta



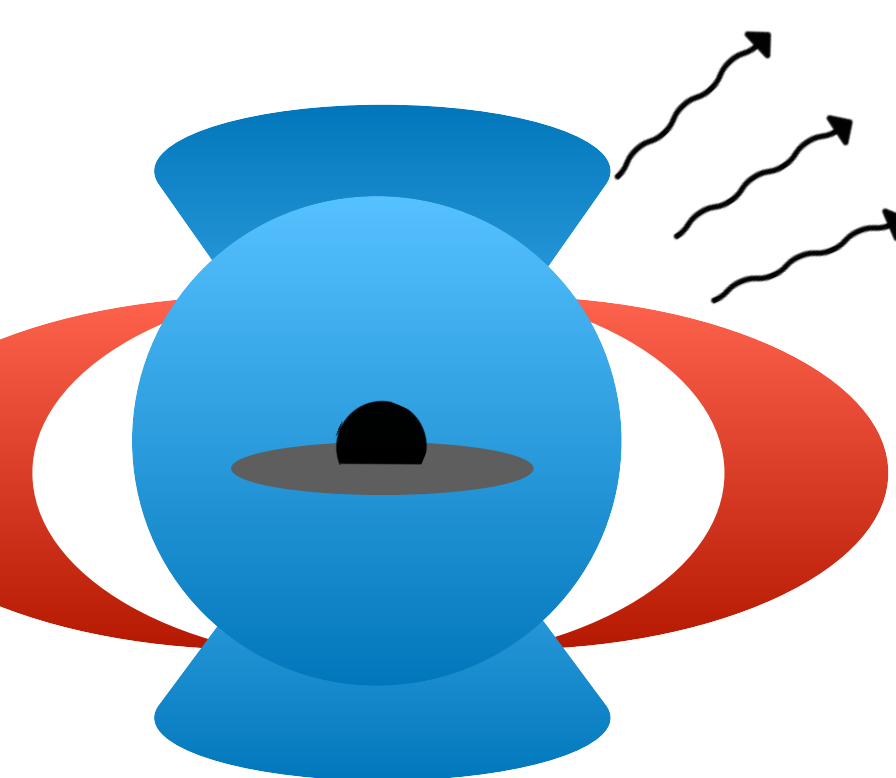
Electromagnetic Emission "Kilonova"

GW170817
DECam observation
(0.5–1.5 days post merger)



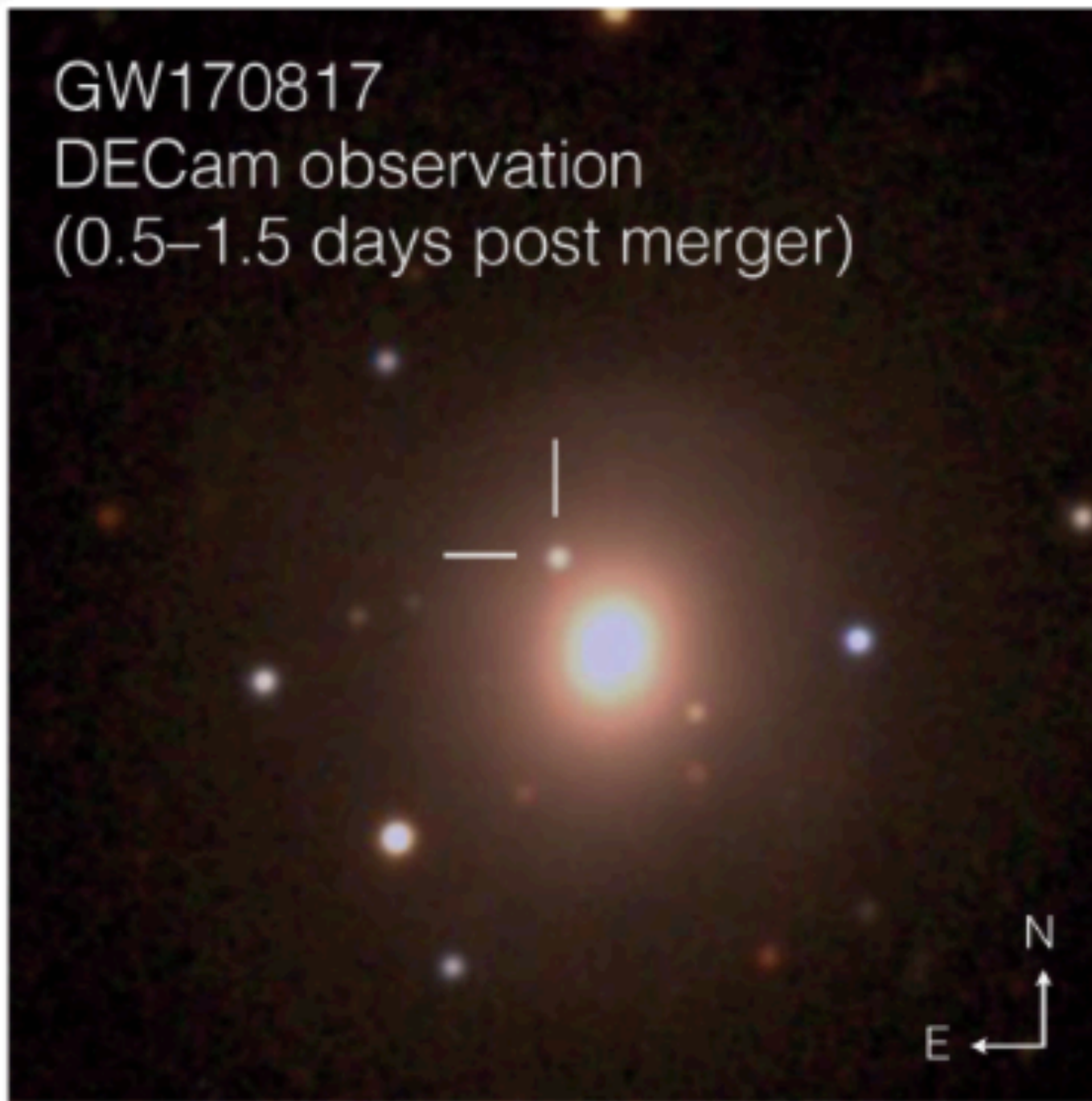
Nucleosynthesis of heavy elements by rapid neutron capture process (r-process)

Introduction



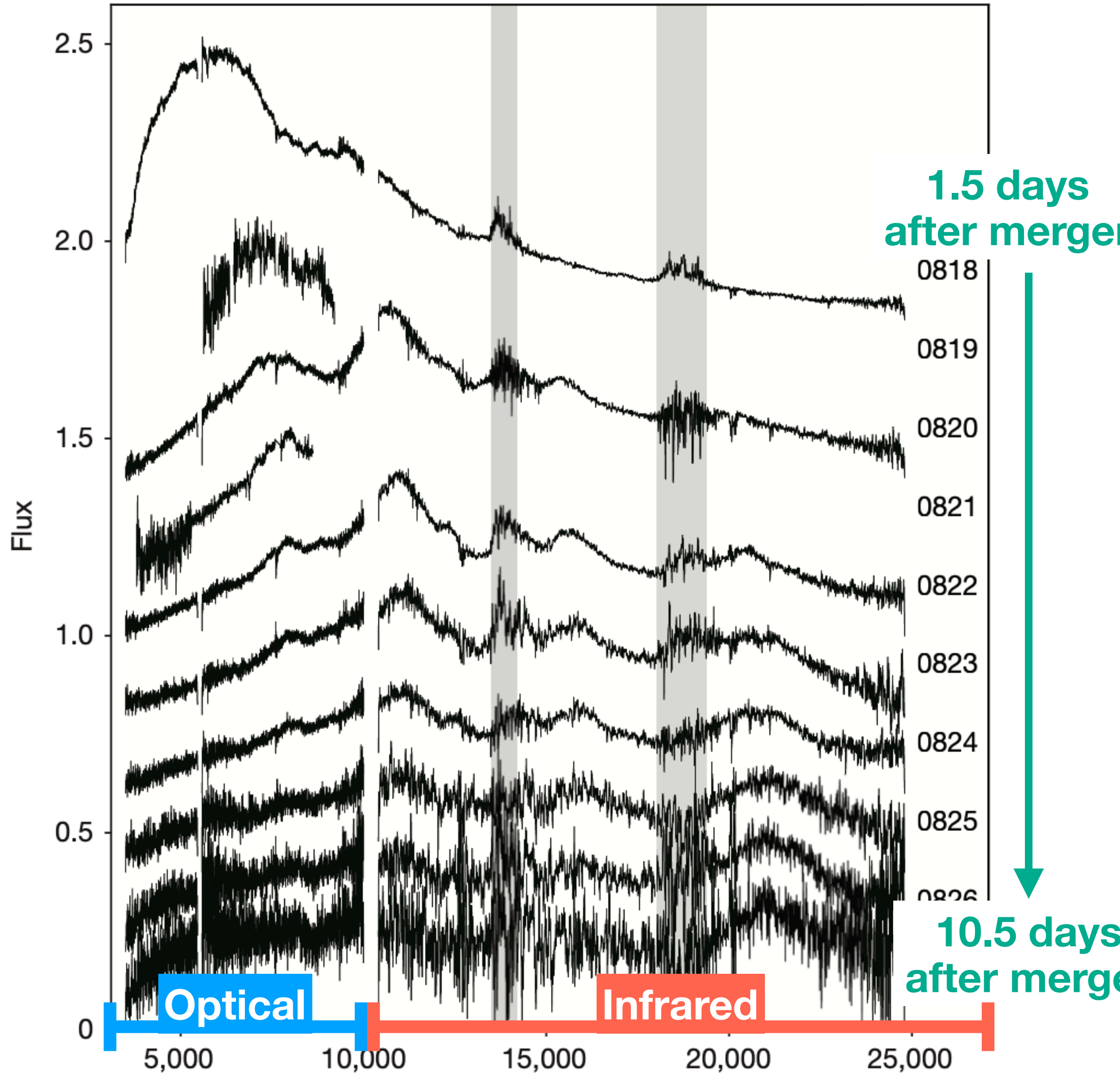
Electromagnetic emission
"Kilonova"

Kilonova of GW170817



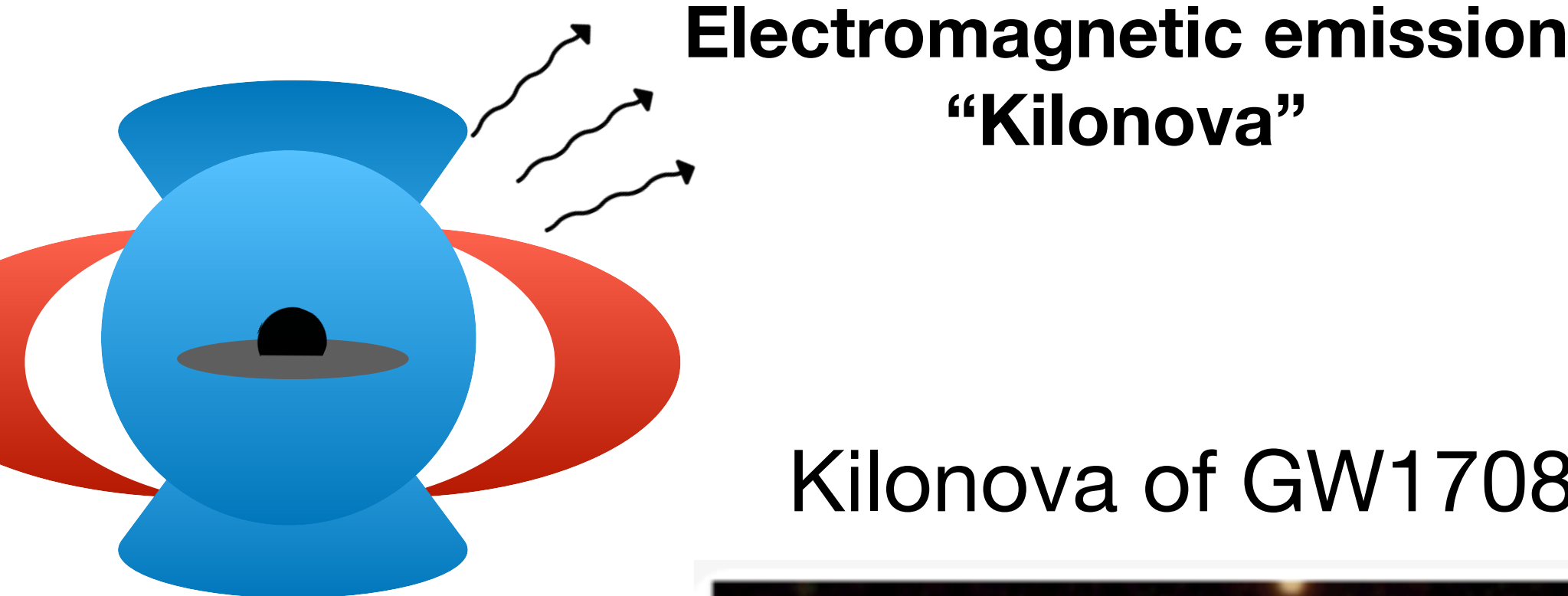
Soares-Santos et al. (2017)

GW170817 spectrum

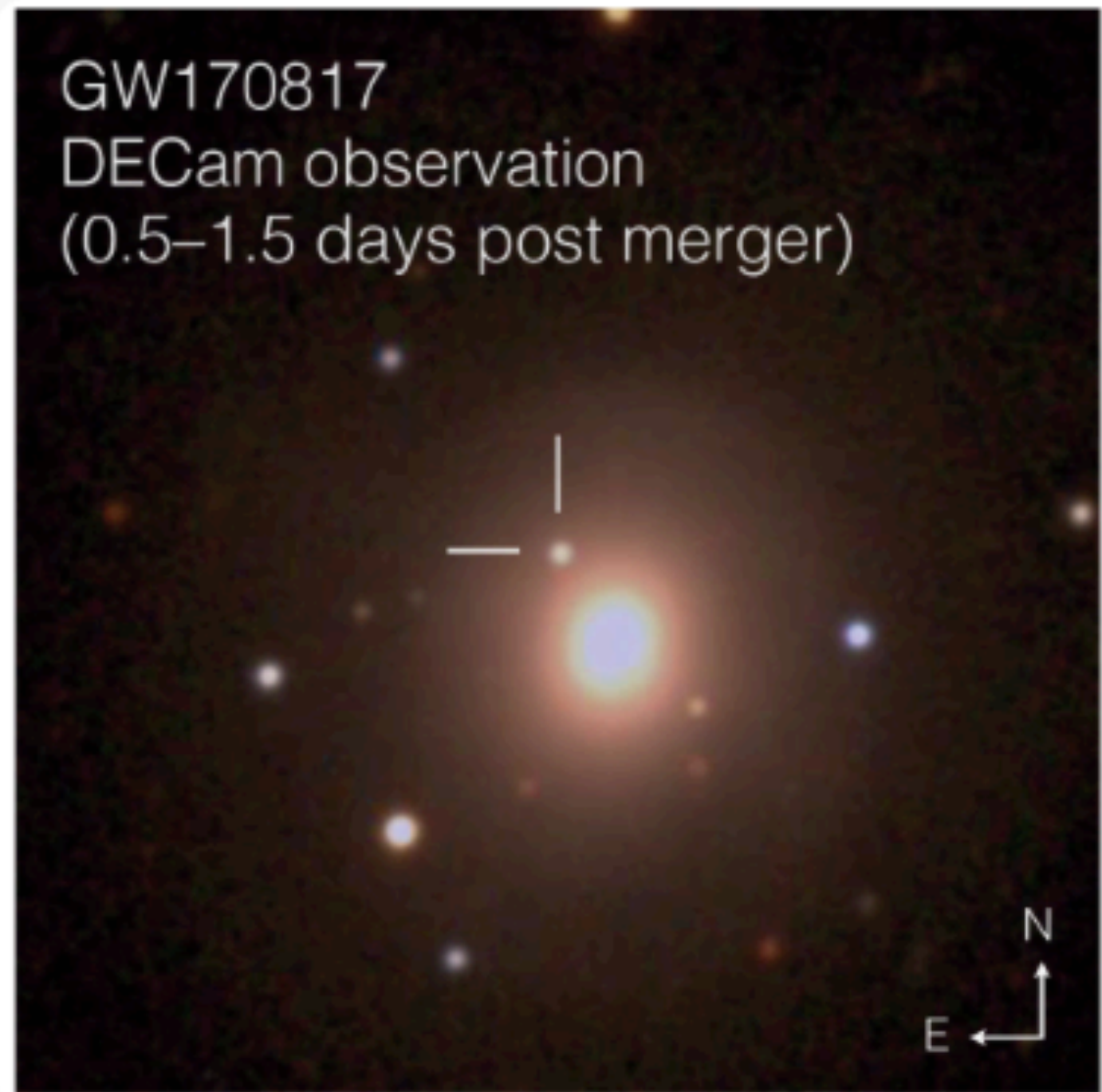


Pian et al. (2017)

Introduction

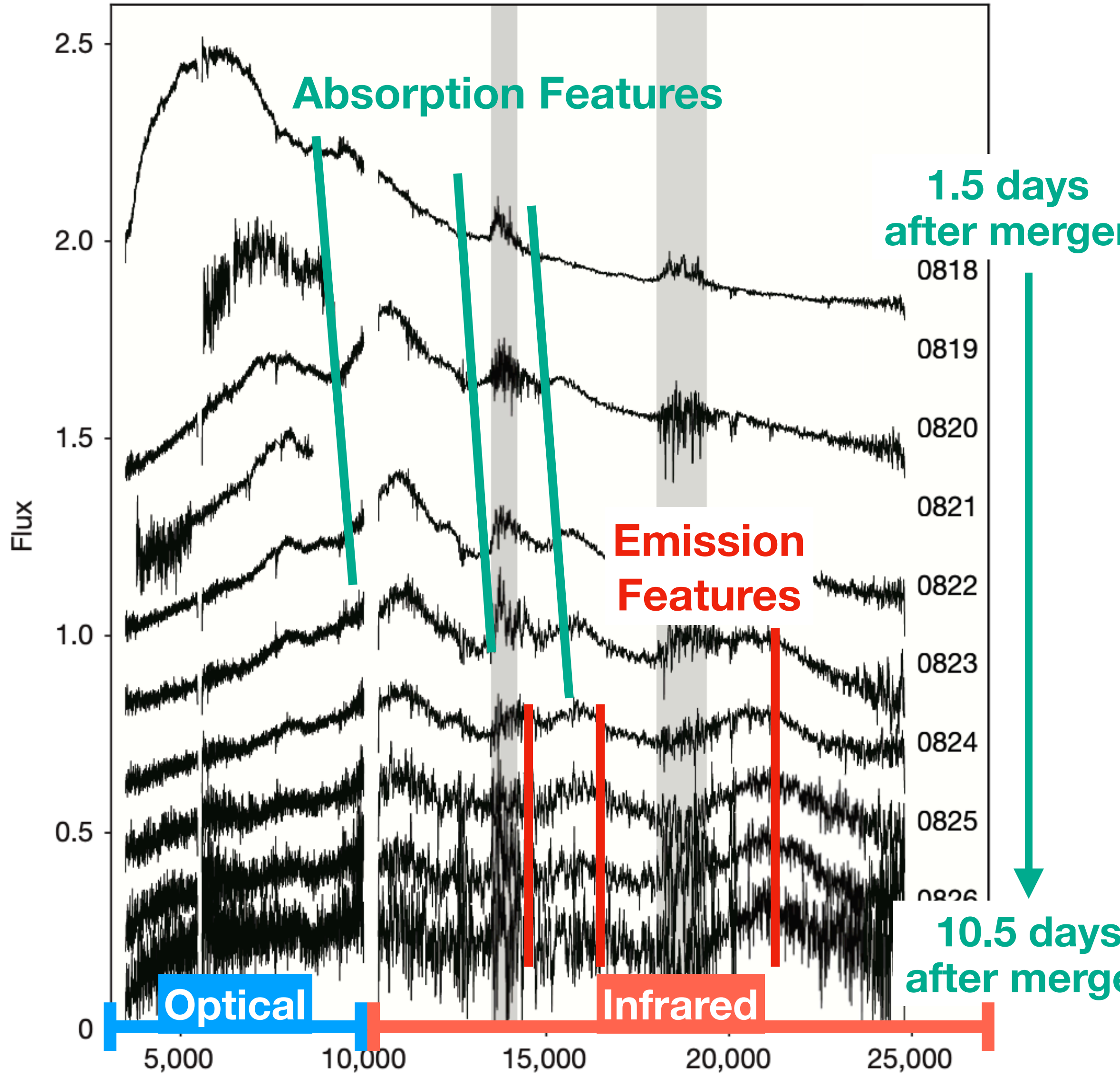


Kilonova of GW170817



Soares-Santos et al. (2017)

GW170817 spectrum



Pian et al. (2017)

Introduction

Electromagnetic emission
"Kilonova"

Goal

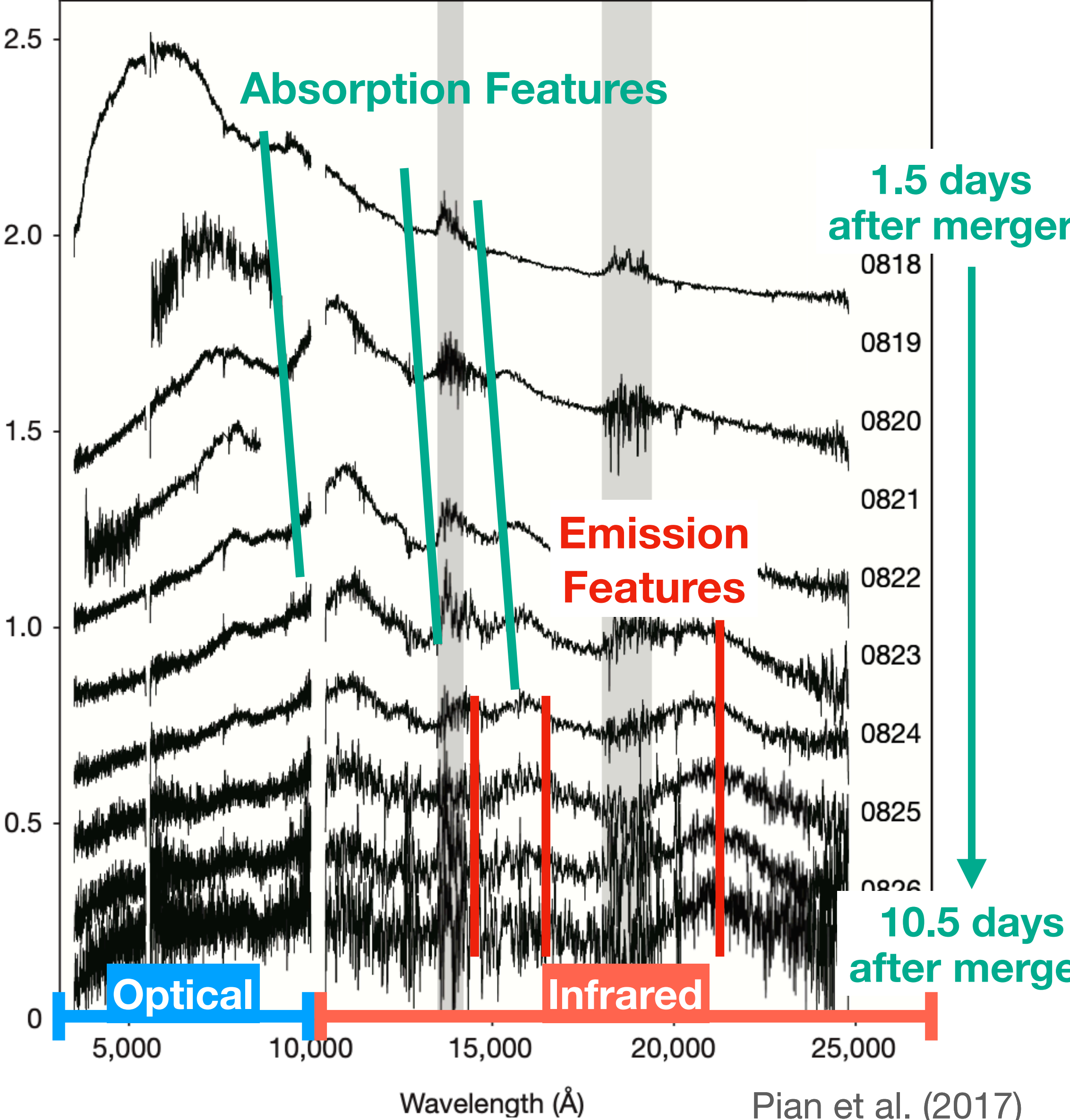
Quantify the mass fraction of heavy elements in NSMs to ultimately understand nucleosynthesis of heavy elements

First step: Spectral Investigation



Soares-Santos et al. (2017)

GW170817 spectrum

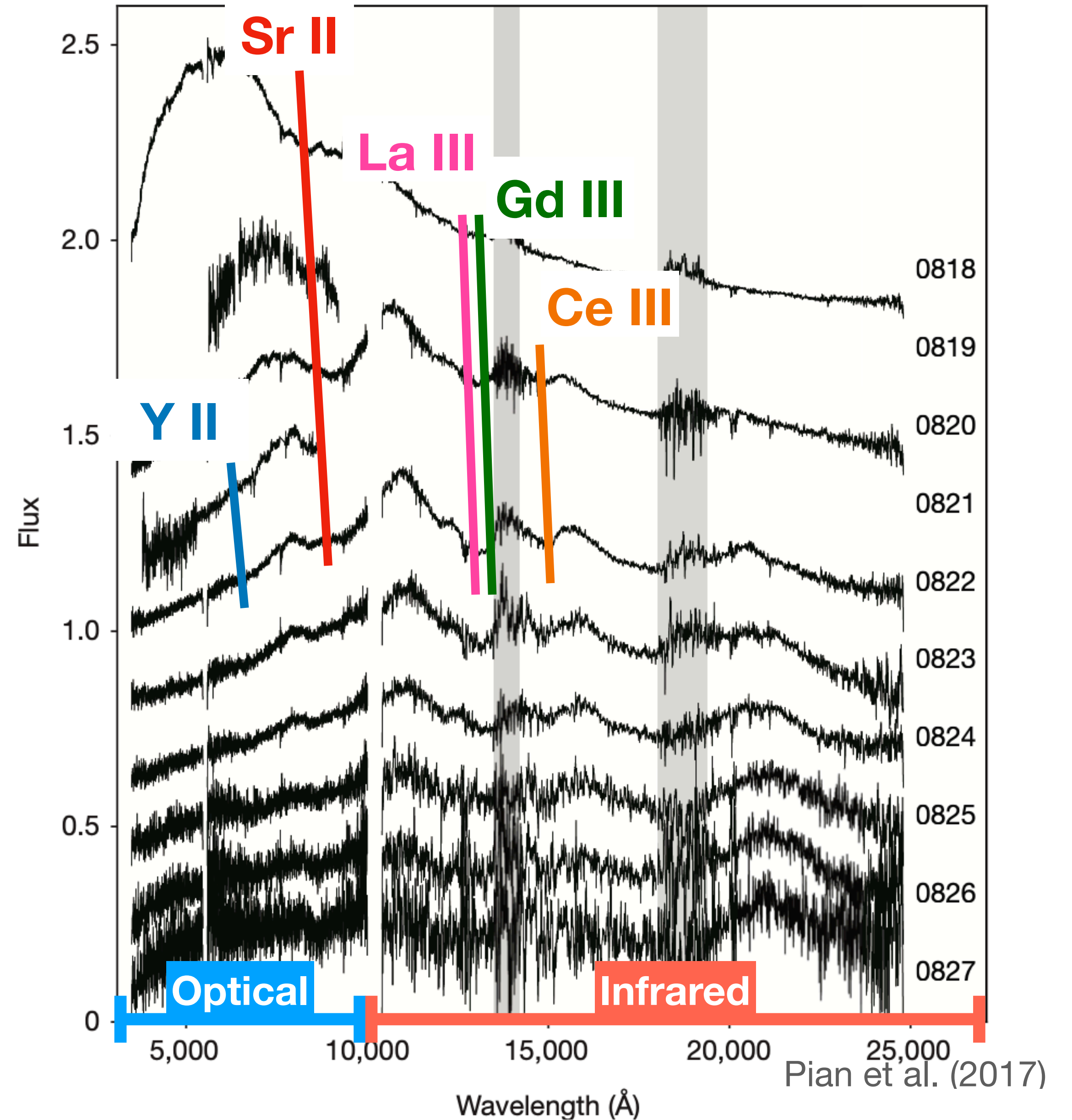


Early-Phase Spectral Investigation

Identification of Most Features

(Watson+19, Domoto+22, Sneppen+23, SR+25)

1 H																	2 He
3 Li	4 Be											5 B	6 C	7 N	8 O	9 F	10 Ne
11 Na	12 Mg											13 Al	14 Si	15 P	16 S	17 Cl	18 Ar
19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	35 Br	36 Kr
37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	53 I	54 Xe
55 Cs	56 Ba	72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	84 Po	85 At	86 Rn	
87 Fr	88 Ra	104 Rf	105 Db	106 Sg	107 Bh	108 Hs	109 Mt	110 Ds	111 Rg	112 Cn	113 Nh	114 Fl	115 Mc	116 Lc	117 Ts	118 Og	
57 La	58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb	71 Lu			
89 Ac	90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No	103 Lr			

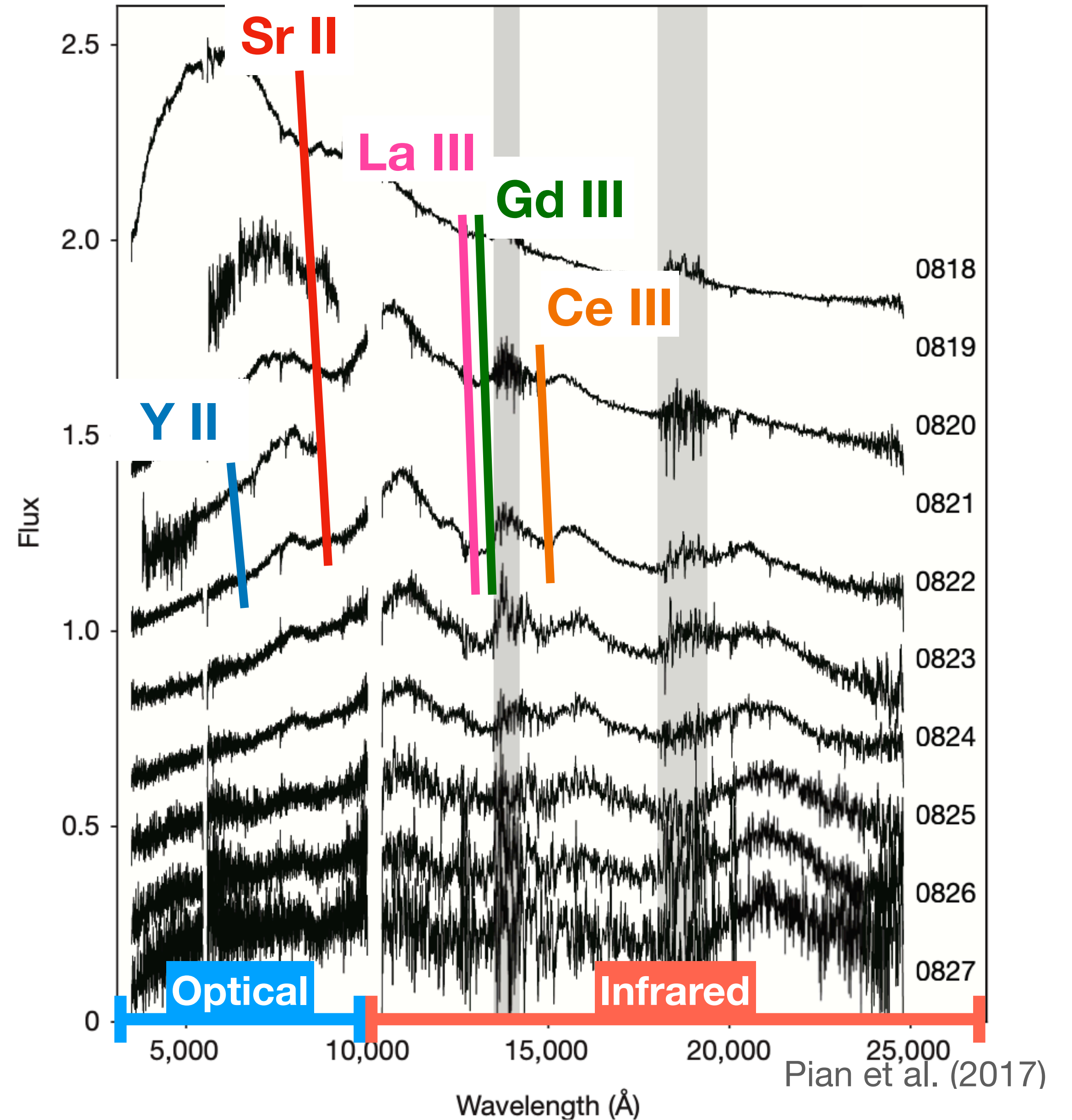
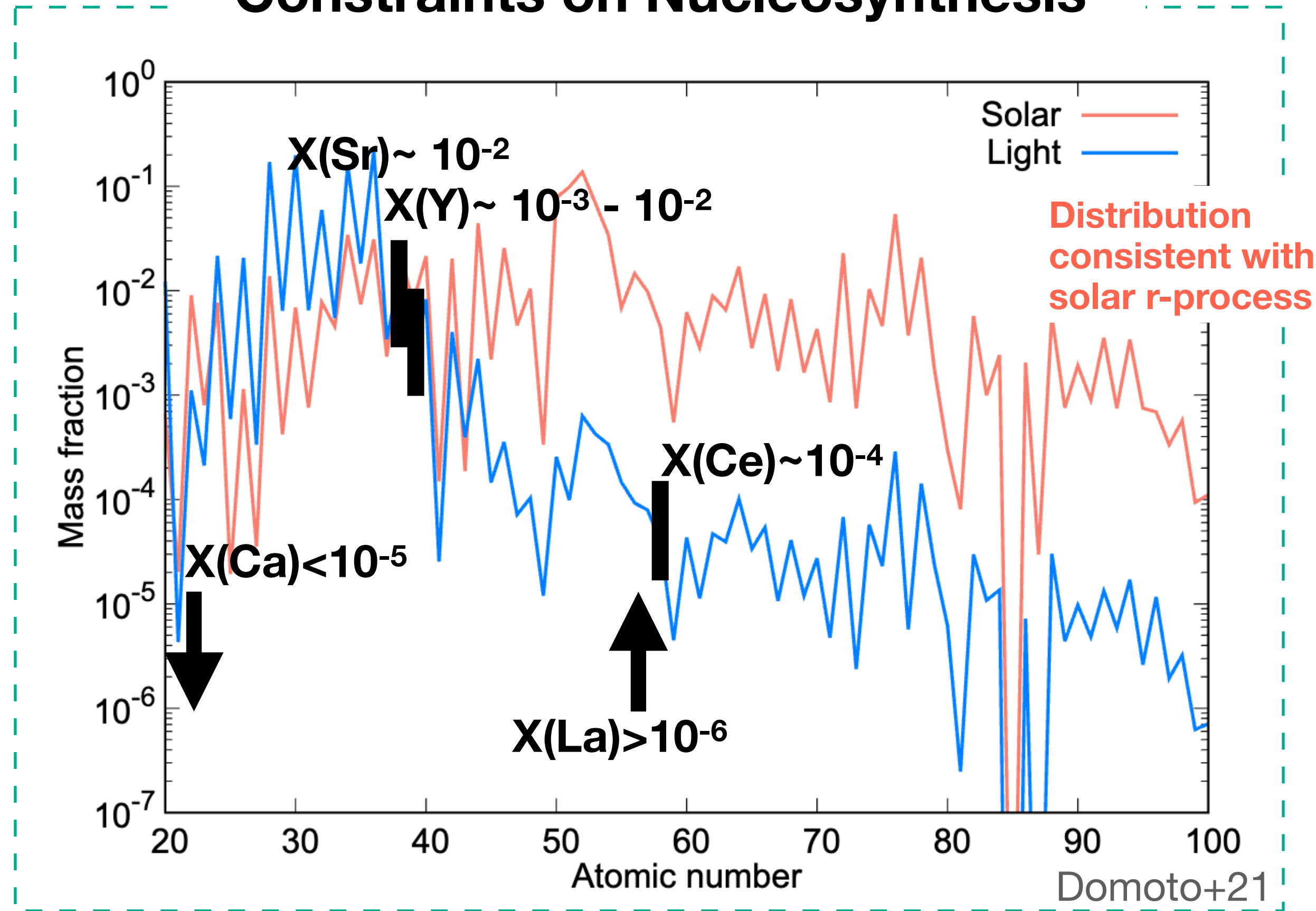


Early-Phase Spectral Investigation

Identification of Most Features

(Watson+19, Domoto+22, Sneppen+23, SR+25)

Constraints on Nucleosynthesis

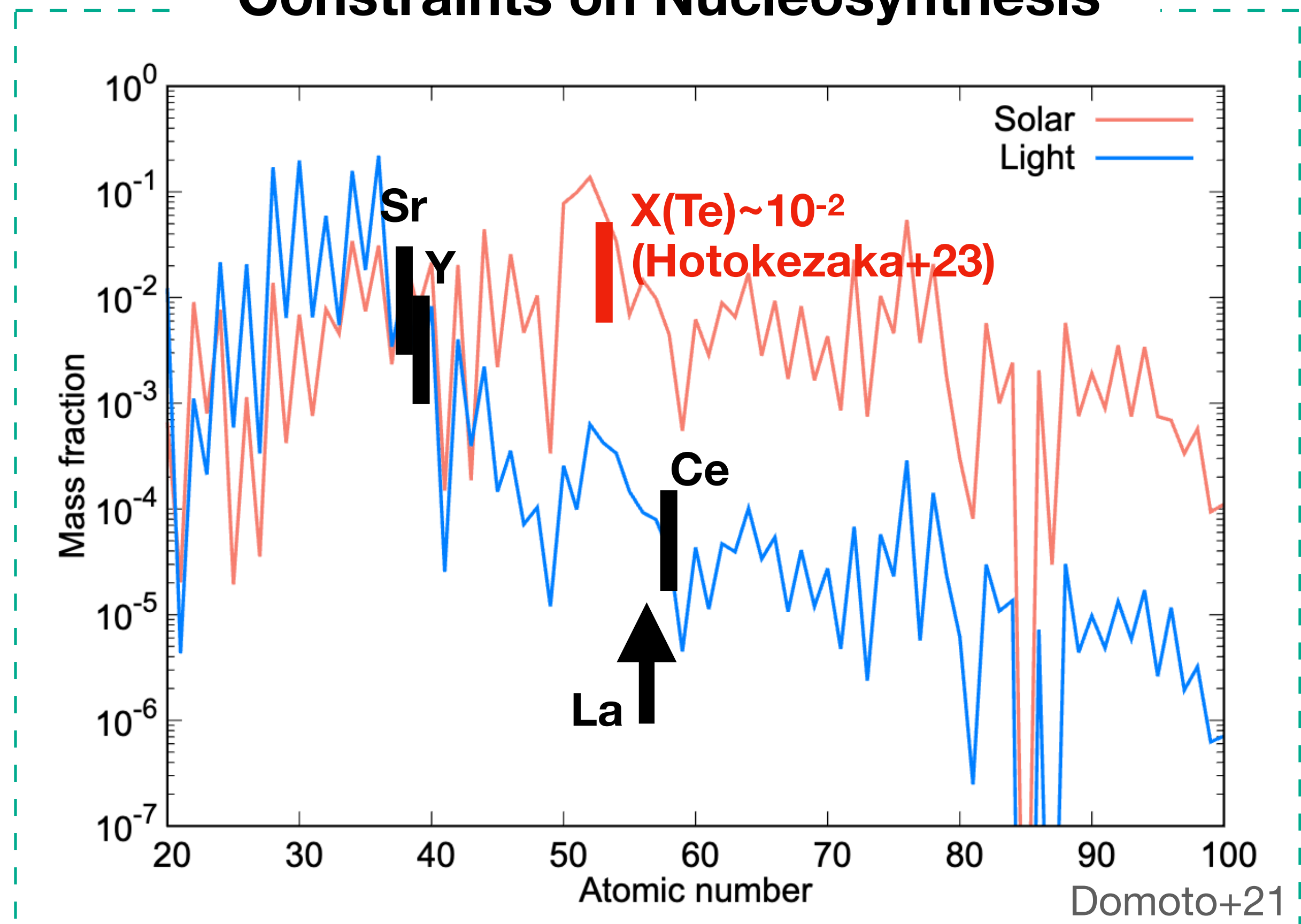


Late-Phase Spectral Investigation

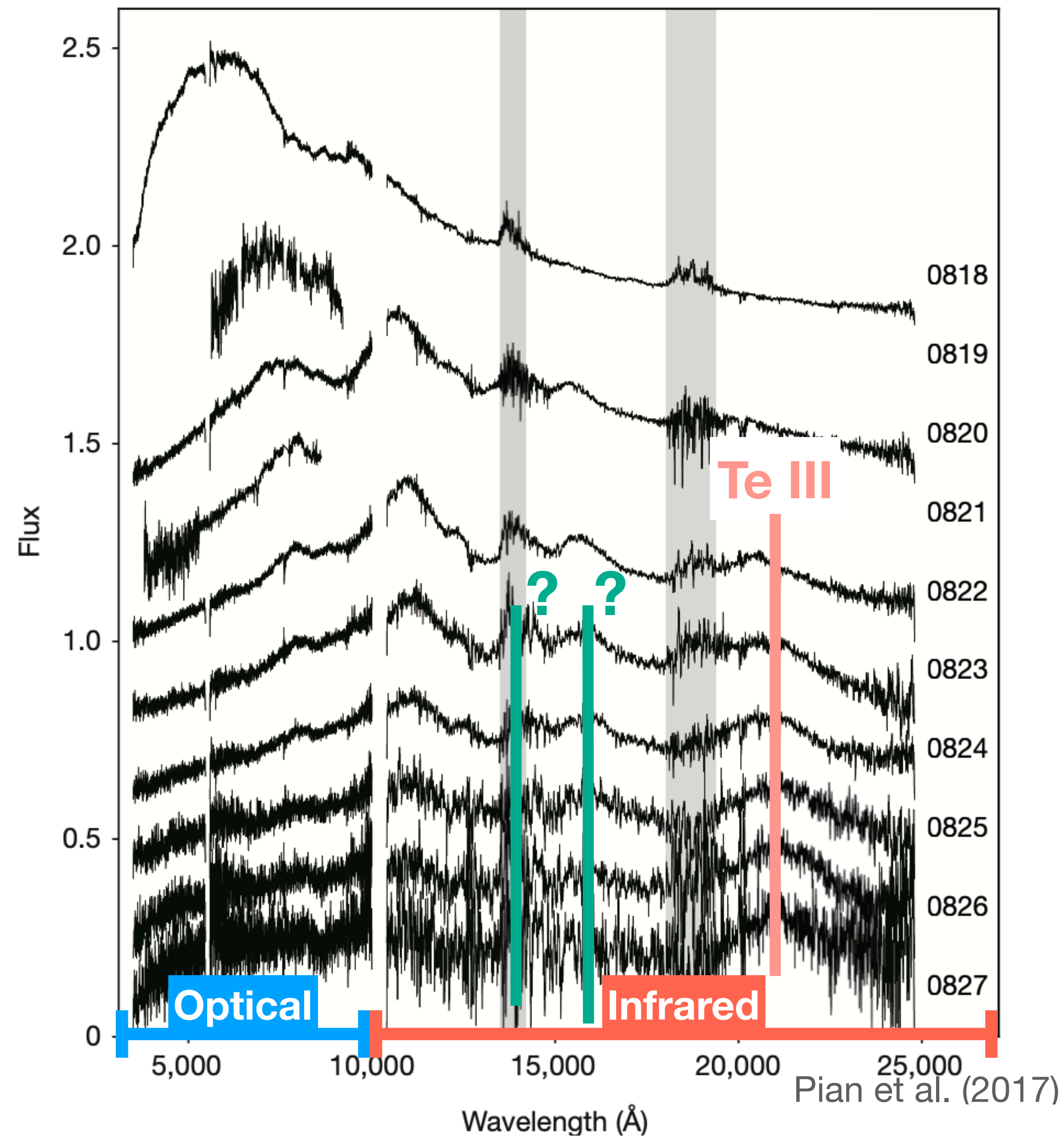
Identification of a Single Feature

(Hotokezaka+23)

Constraints on Nucleosynthesis



Inconsistency?



Late-Phase Spectral Investigation

Notable Previous Works

Gillanders+23: Several candidates for each feature

Pognan+23, 25: non-LTE radiative transfer simulations for few selected elements

Jerkstrand+25: non-LTE simulations for weak r-process elements

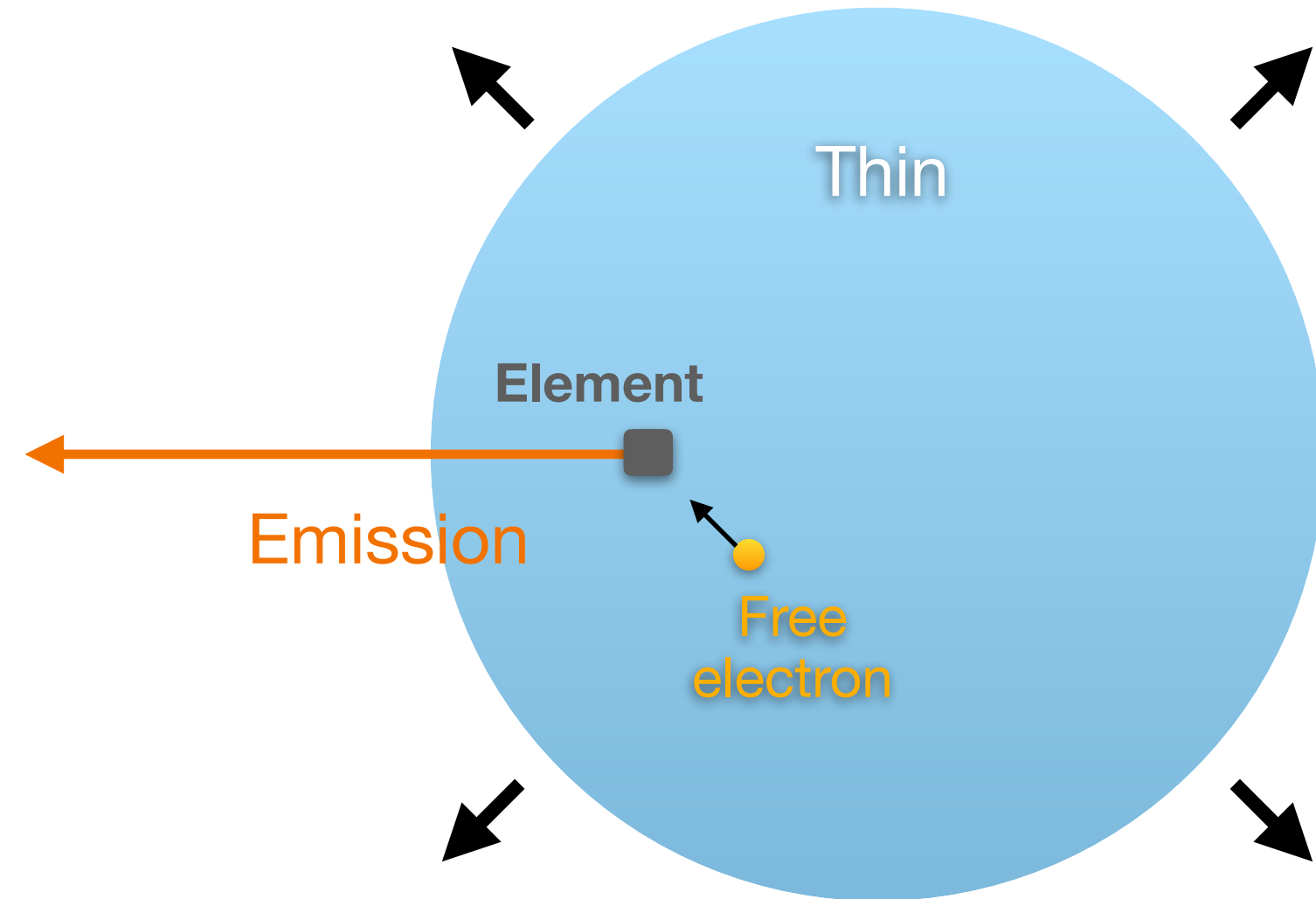
No robust Identification of all late-phase features has been performed so far

This work

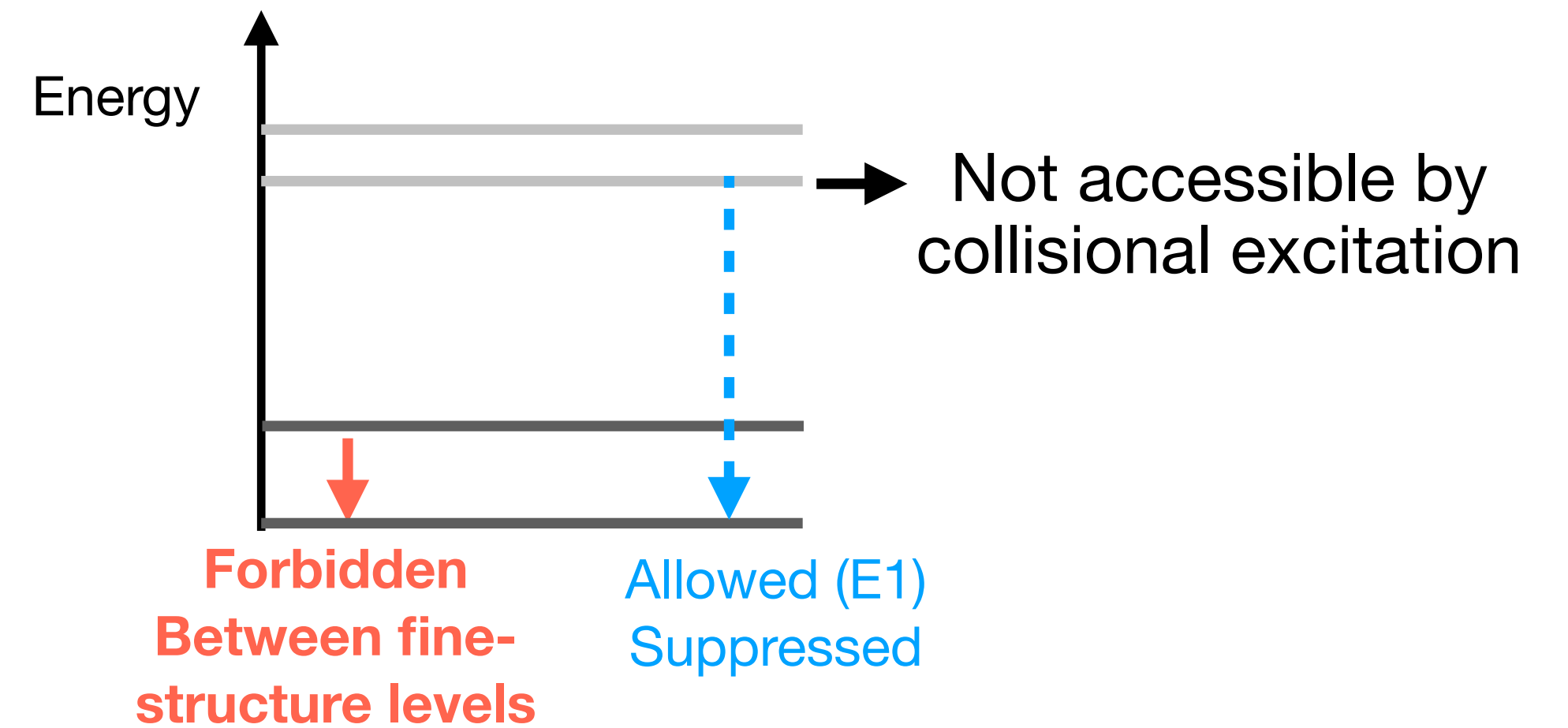
Comprehensive analysis using a simple spectral model including all elements to explain observed features

Spectral Model

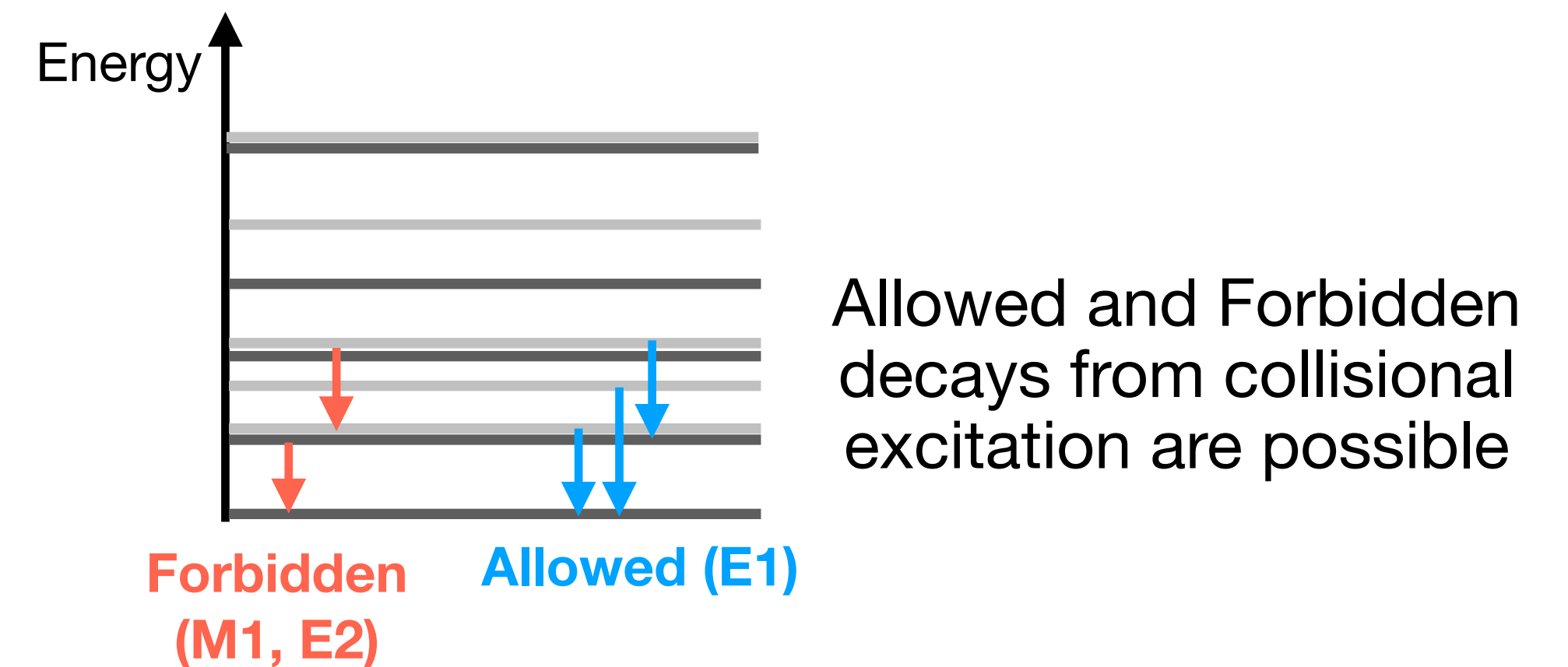
Late-phase features are caused by radiative decay of collisionally excited elements



Light Elements: no mixing between configurations



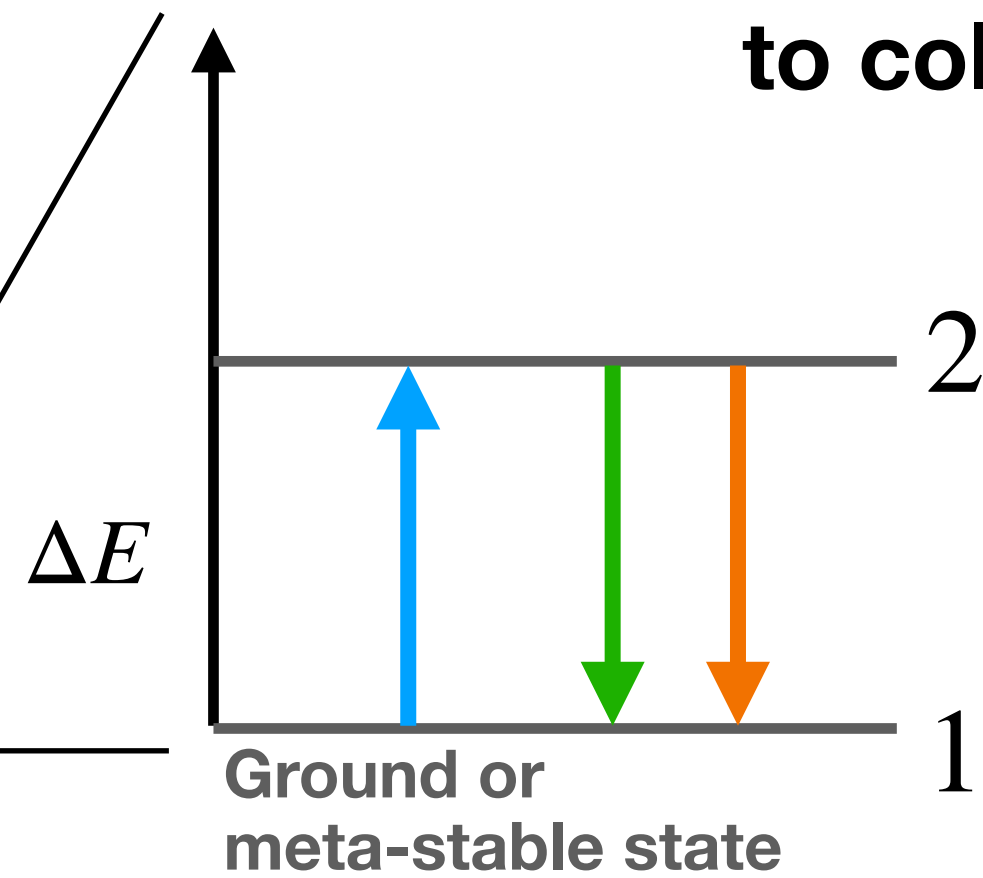
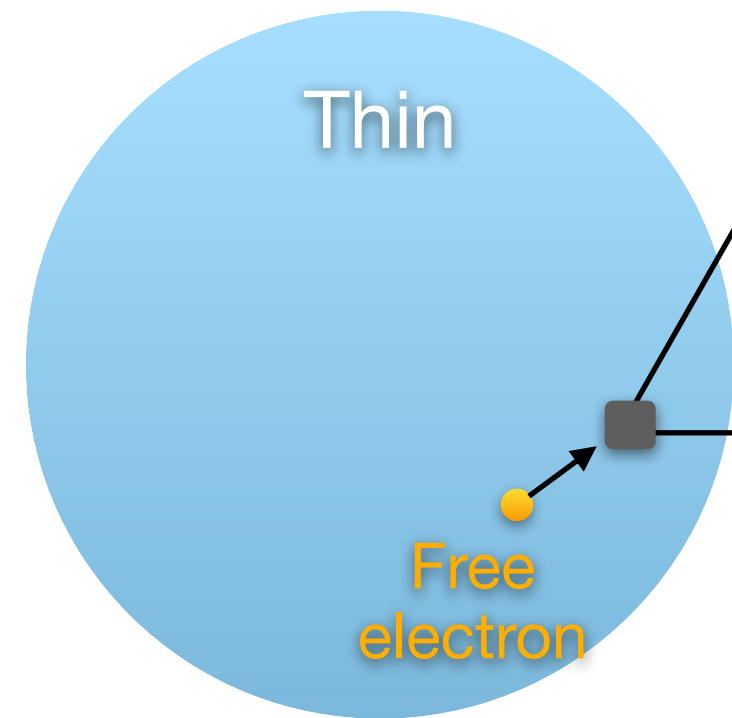
Heavy Elements: large mixing between configurations



Spectral Model

Emission observed at late-phase is due to collisionally excited elements

Late-phase ejecta



Radiation rate: $n_2 A_{2 \rightarrow 1}$

Collisional deexcitation rate: $n_2 n_e q_{2 \rightarrow 1}$

Collisional excitation rate: $n_1 n_e q_{1 \rightarrow 2}$

$$n_2 A_{2 \rightarrow 1} + n_2 n_e q_{2 \rightarrow 1} = n_1 n_e q_{1 \rightarrow 2}$$

Luminosity (per volume): $L = \Delta E n_2 A_{2 \rightarrow 1}$

Rate coefficient: $q_{1 \rightarrow 2} = \frac{8.629 \times 10^{-6}}{T^{1/2}} \frac{\Upsilon_{1 \rightarrow 2}}{g_1} \exp(-\Delta E/kT_e)$ Atomic property

Electron density $n_e \sim 10^7 \text{ cm}^{-3}$

Electron temperature $T_e \sim 2000 \text{ K}$

Einstein coefficient:

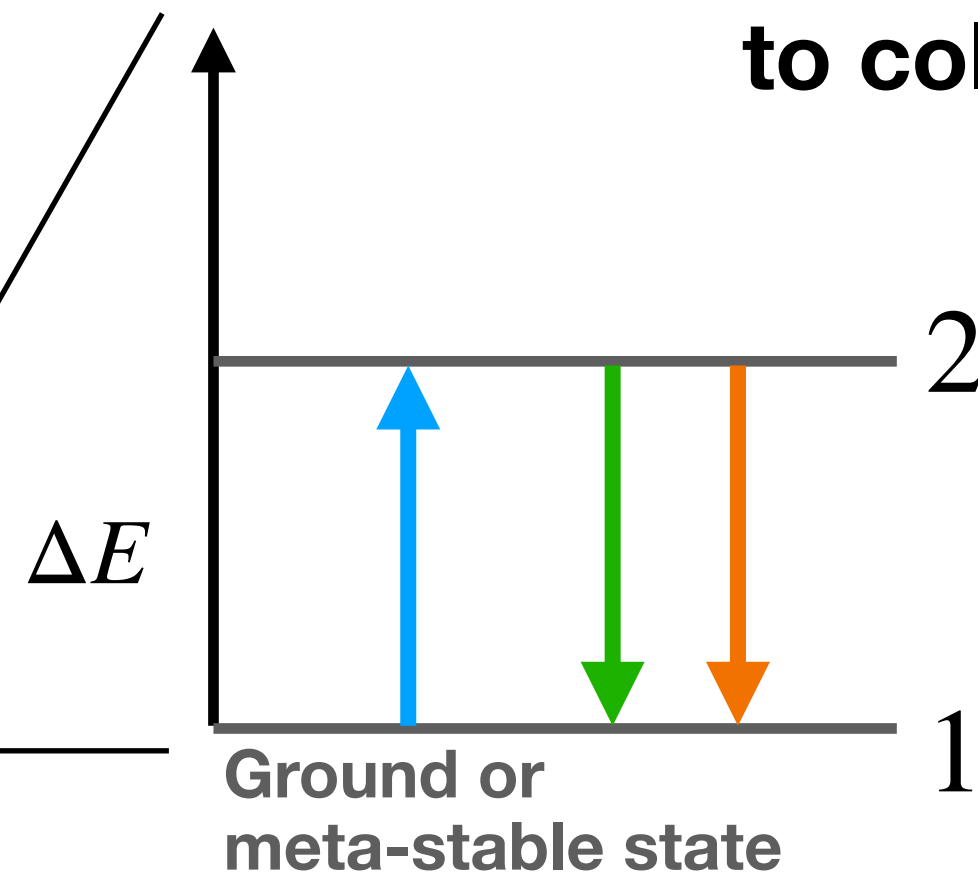
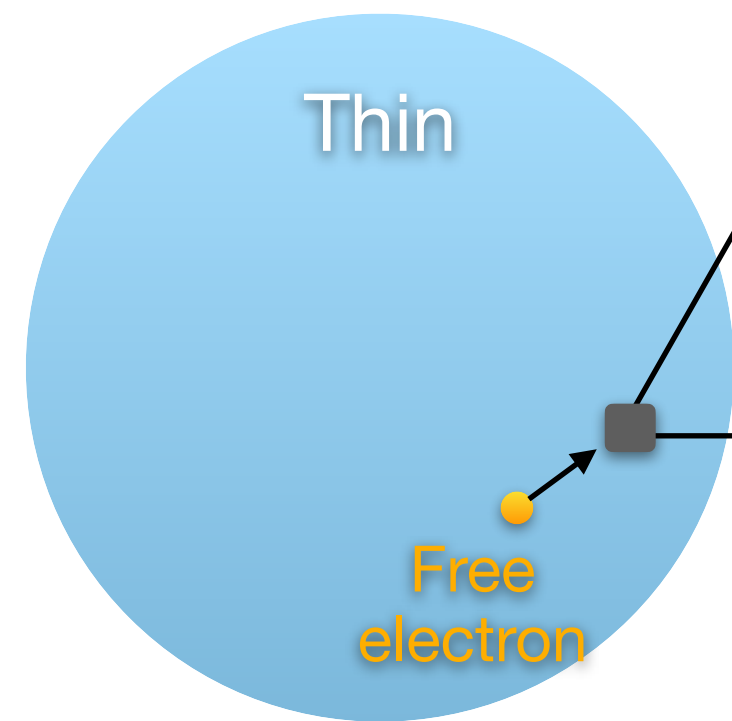
Allowed: $A_{2 \rightarrow 1} \sim 10^6 \text{ s}^{-1}$

Forbidden: $A_{2 \rightarrow 1} \sim 1 \text{ s}^{-1}$

Spectral Model

Emission observed at late-phase is due to collisionally excited elements

Late-phase ejecta



Radiation rate: $n_2 A_{2 \rightarrow 1}$

Collisional deexcitation rate: $n_2 n_e q_{2 \rightarrow 1}$

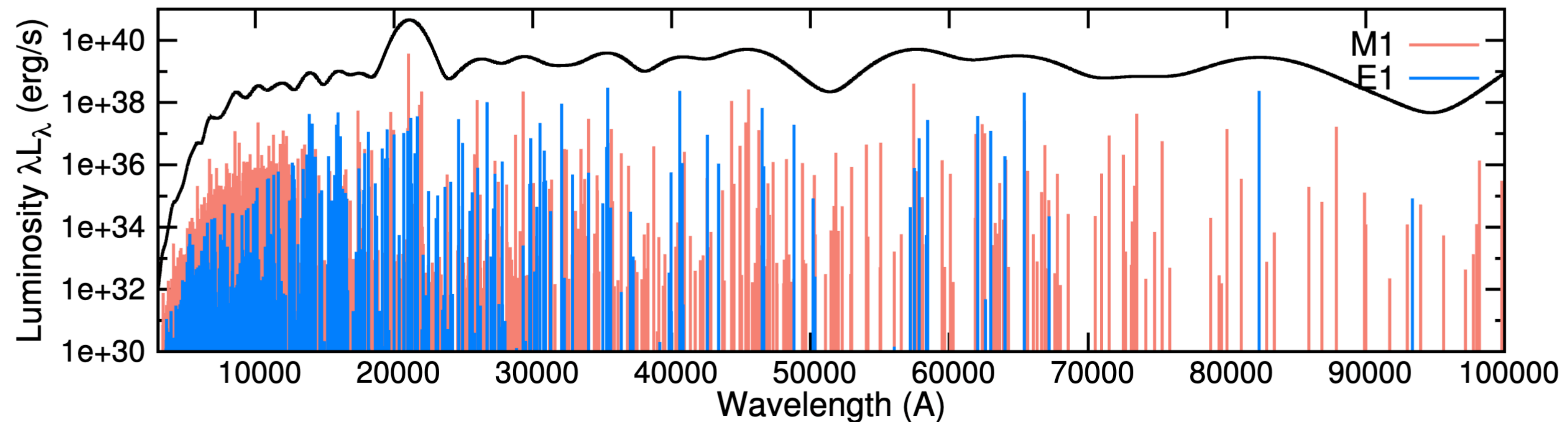
Collisional excitation rate: $n_1 n_e q_{1 \rightarrow 2}$

Luminosity (per volume): $L = \Delta E n_2 A_{2 \rightarrow 1}$

Resulting Spectrum:

- Selection of forbidden and allowed transitions of all elements
- Doppler broadening of $0.07c$

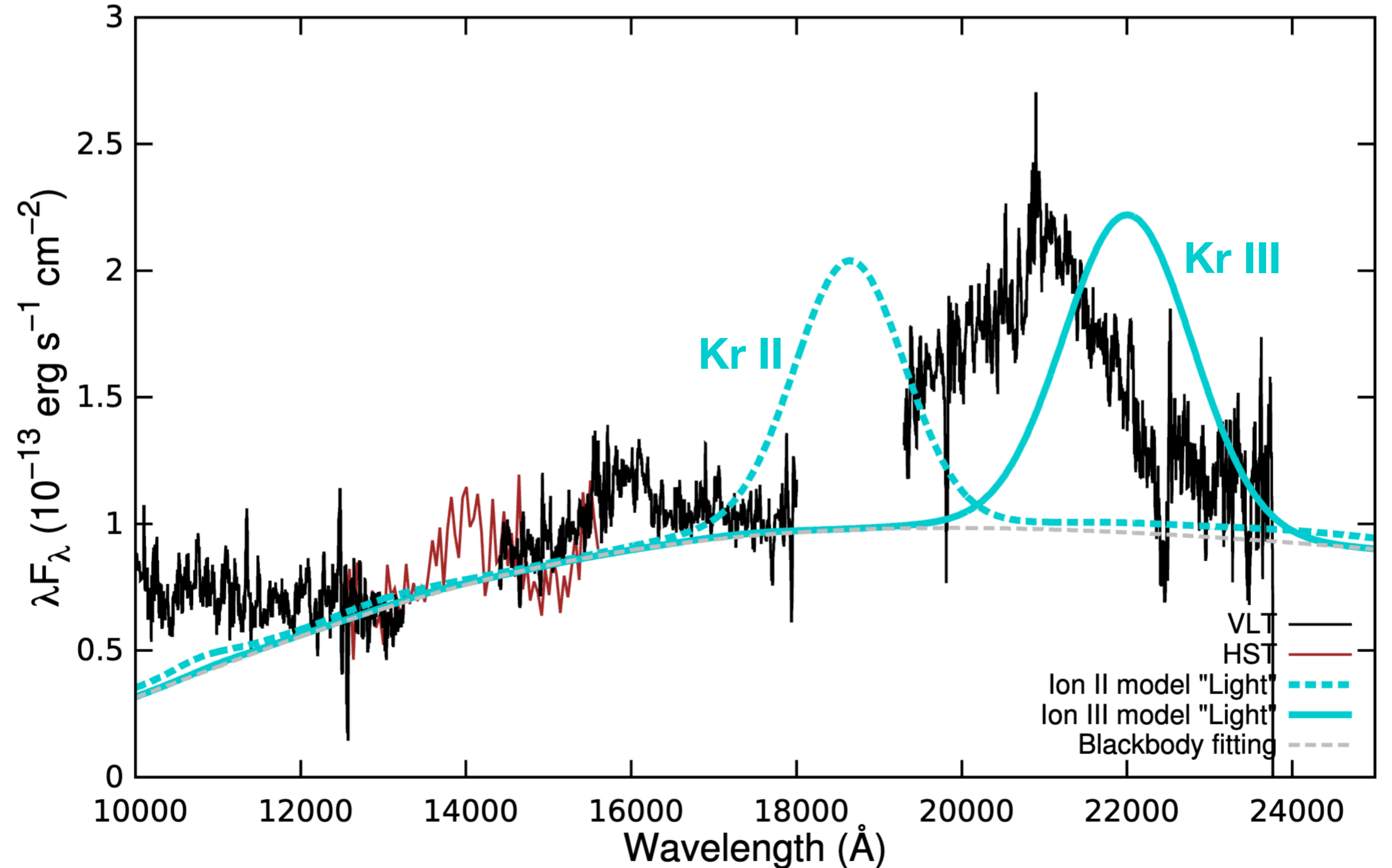
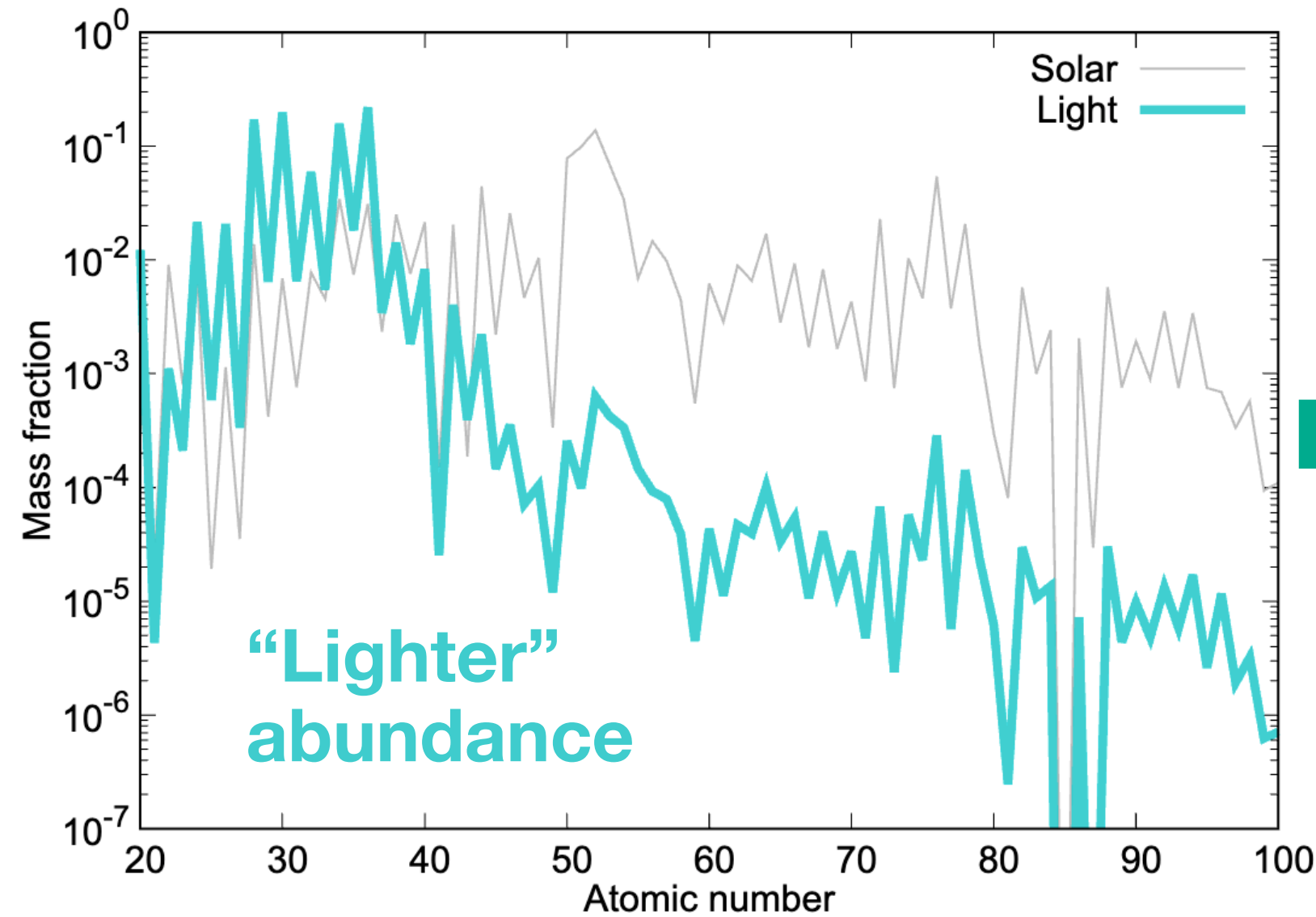
Example of doubly ionized state with solar abundance



Comparison with GW170817

Model spectrum + Blackbody continuum compared with GW170817 observed spectrum at 9.4 days

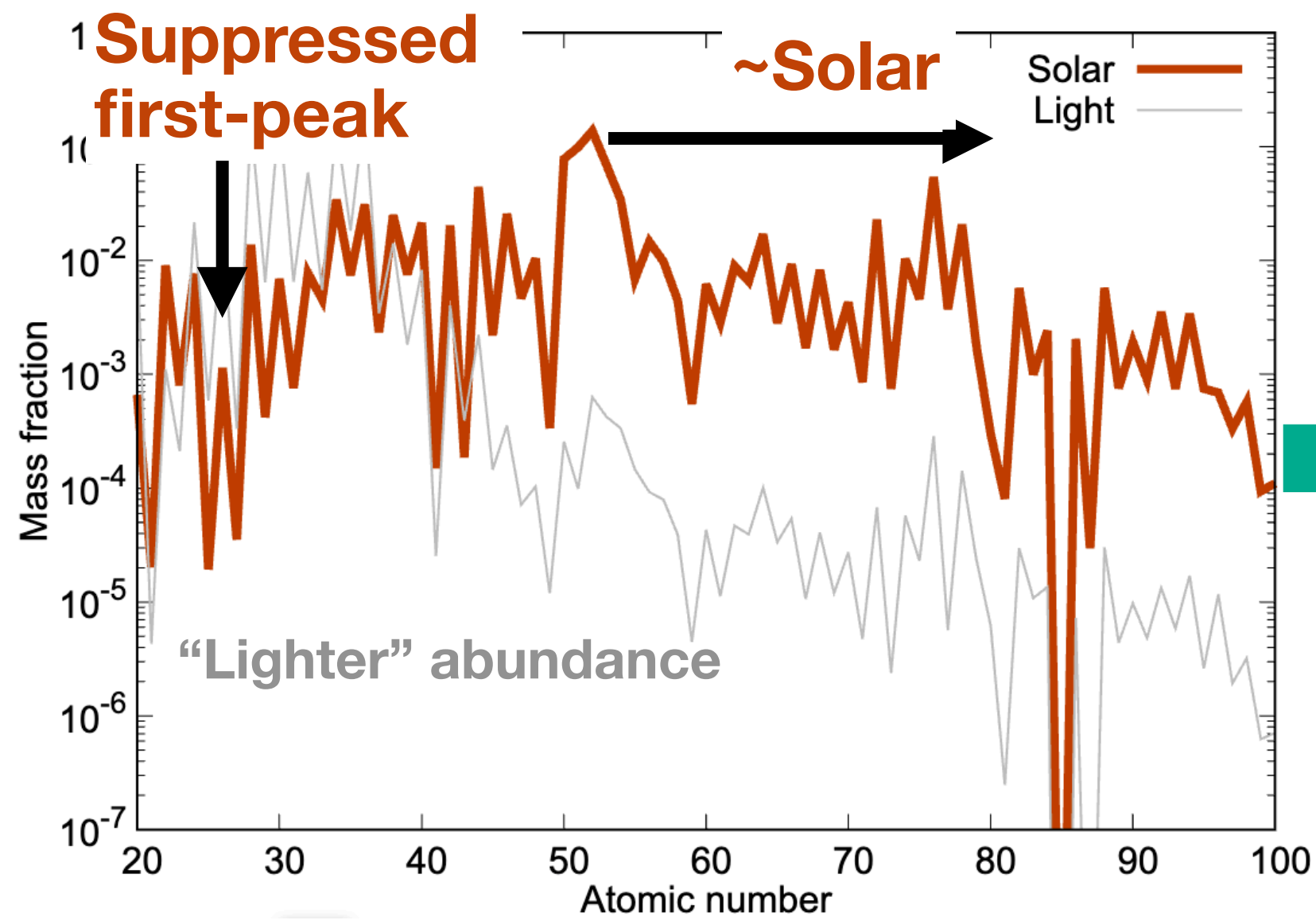
Assuming same abundance as in early phase



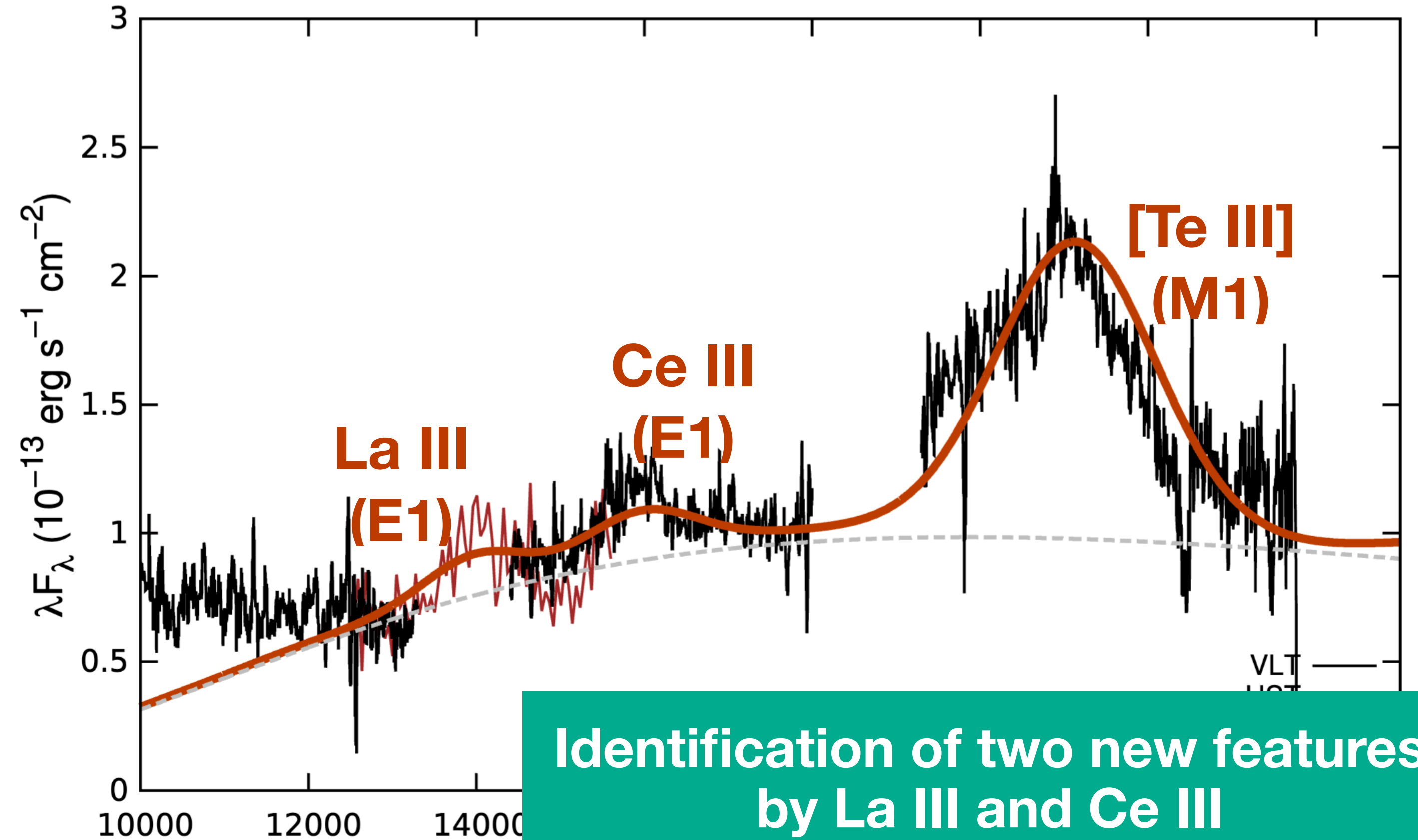
We cannot reproduce any of GW170817 late-phase features with light abundance assumption

Comparison with GW170817

Assuming abundance consistent with solar r-process



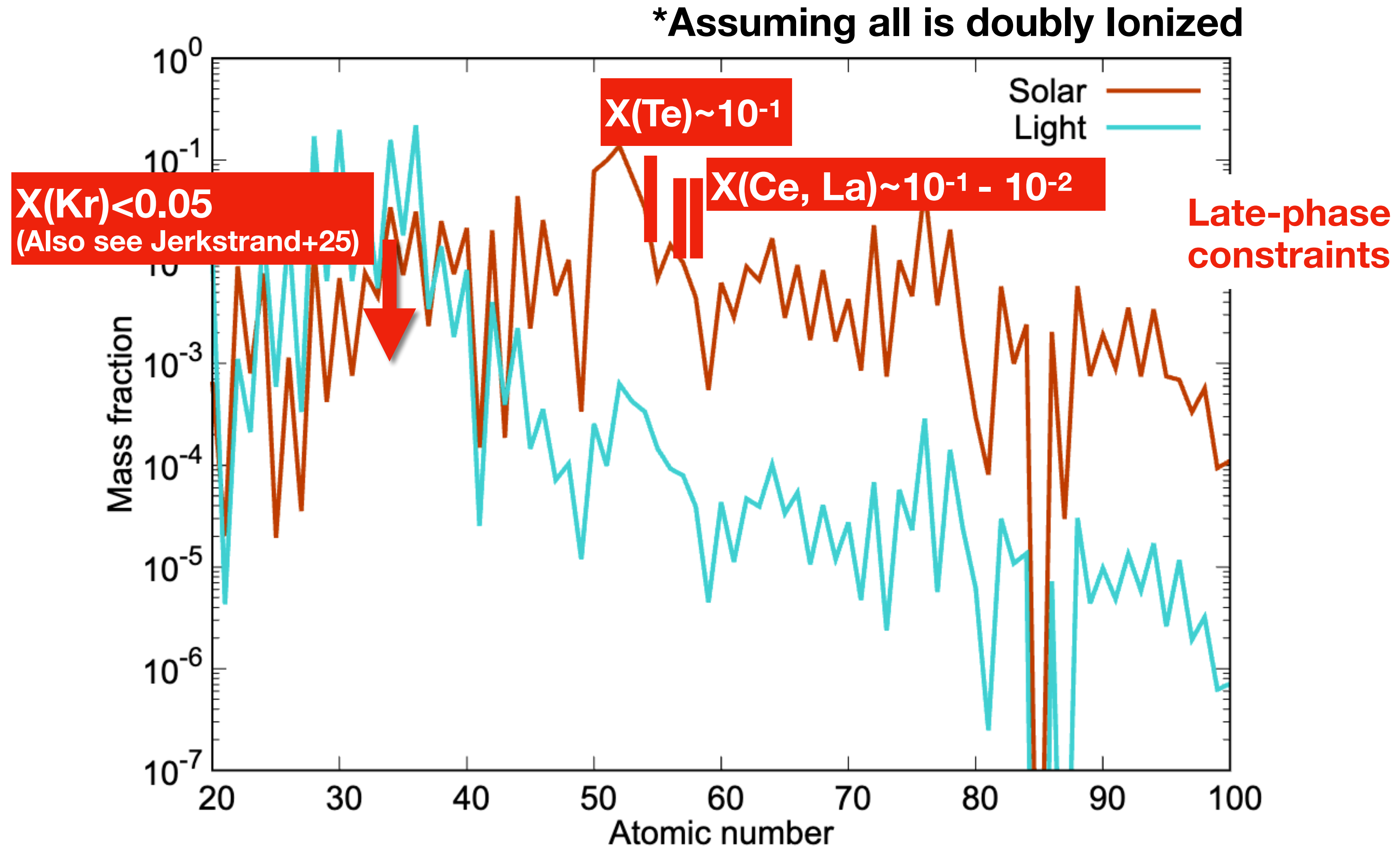
Best-fit model spectrum + Blackbody continuum with GW170817 observed spectrum at 9.4 days



Identification of two new features by La III and Ce III + confirmation of Te III line

*Effect of other lines at each feature is not significant

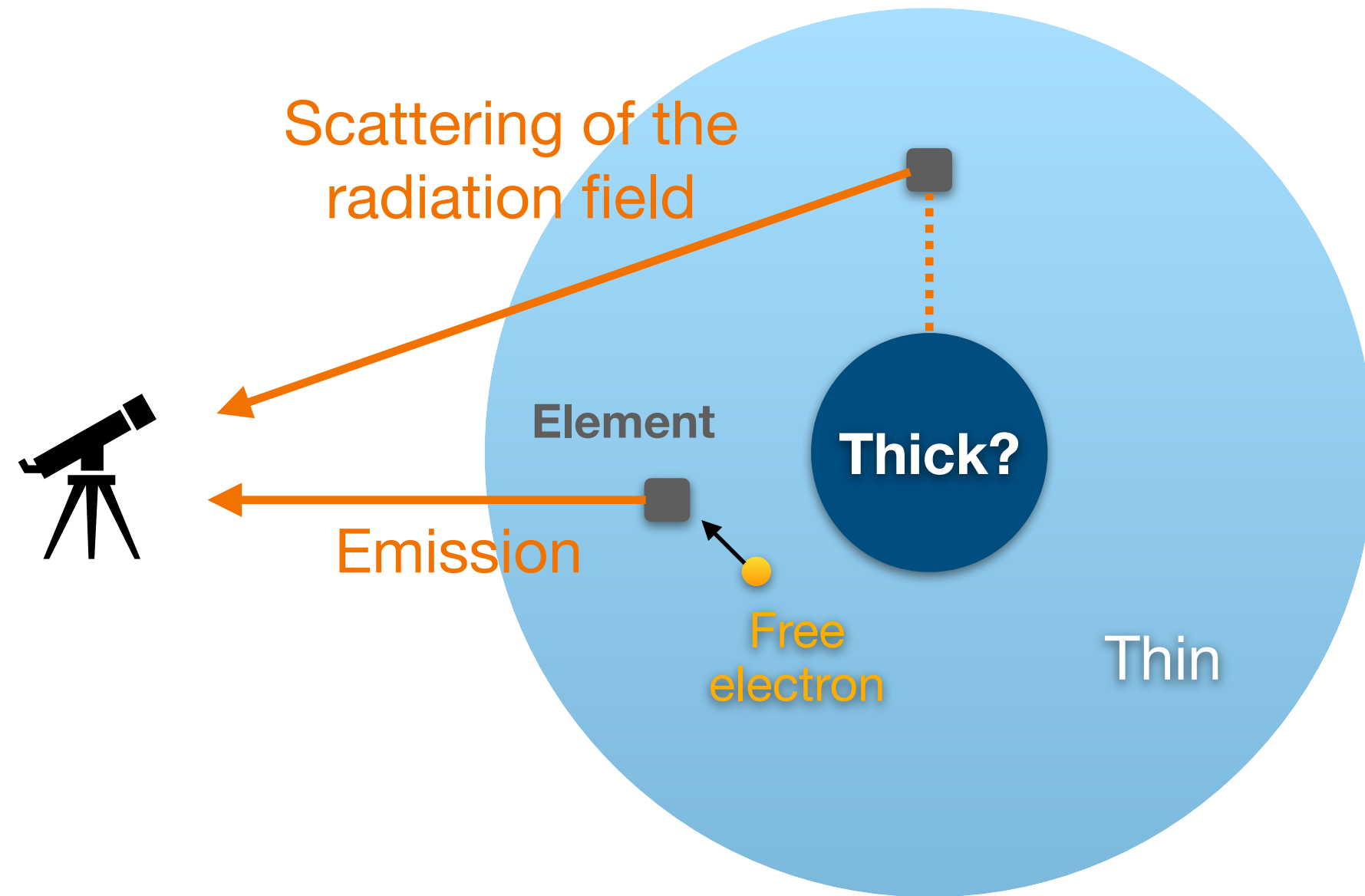
Nucleosynthesis Constraints



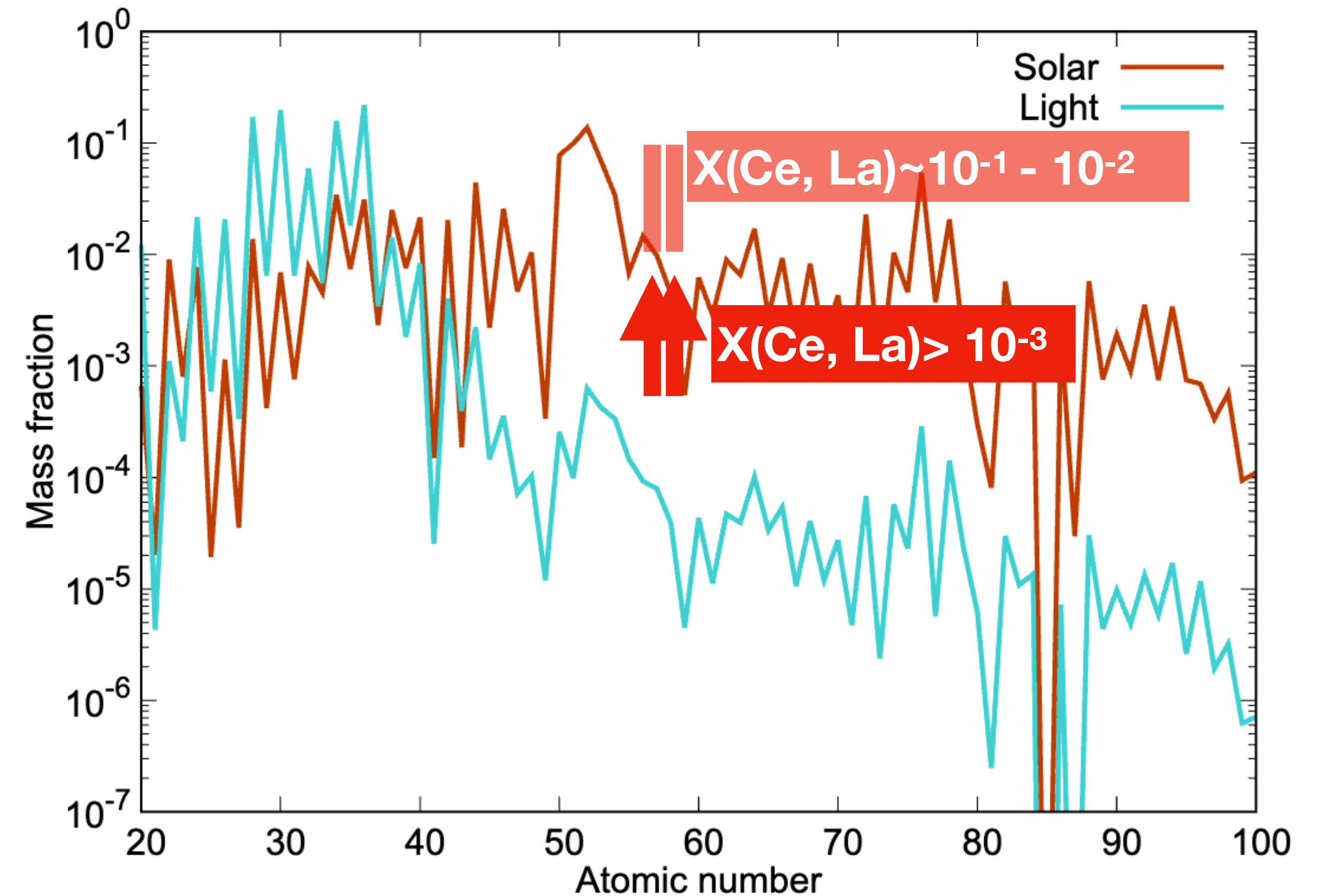
Nucleosynthesis Constraints

Good, but what about the continuum?

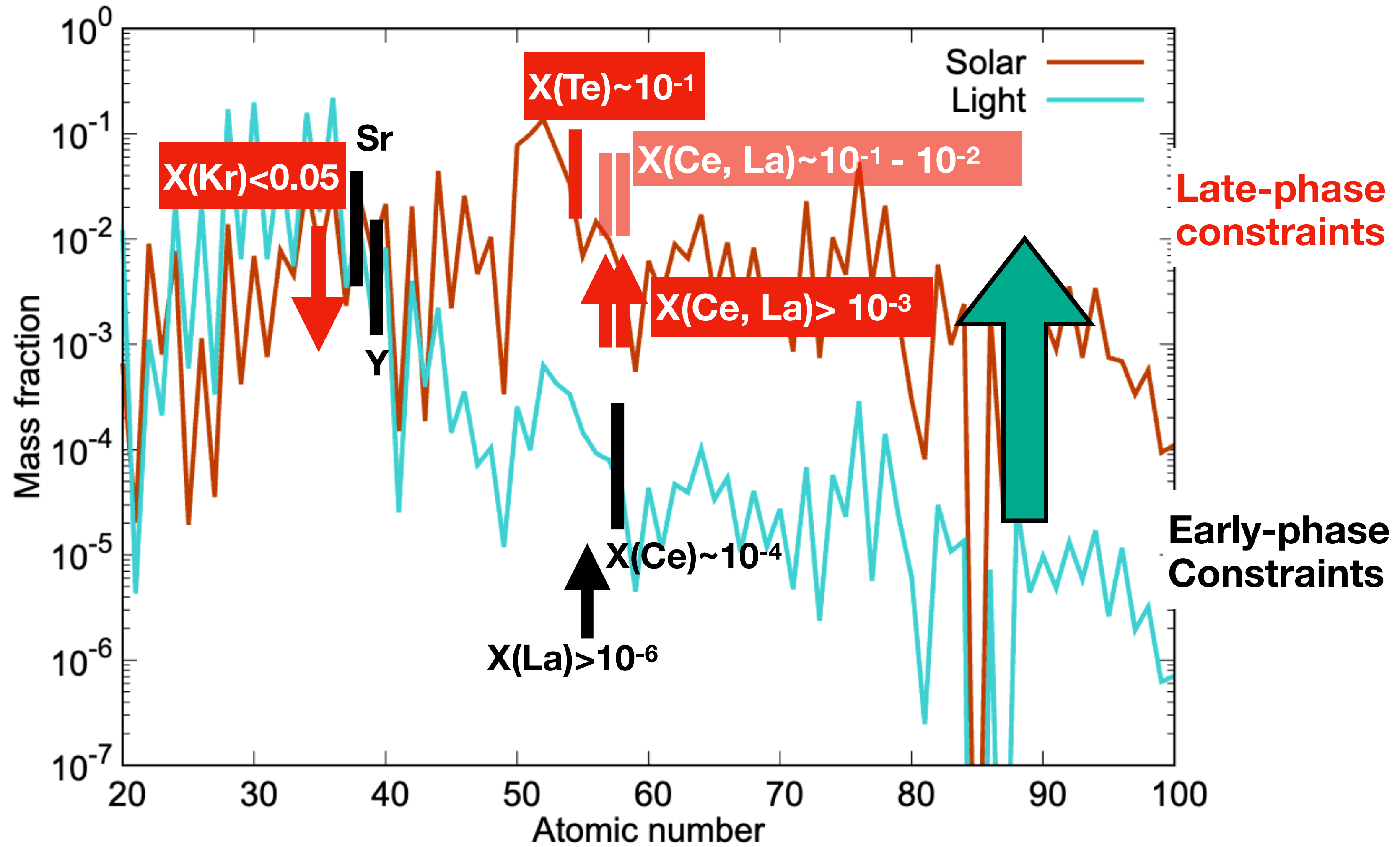
If the radiation field is coming from an inner optically thick region, allowed (E1) transitions can be caused by two mechanisms



→ La III and Ce III become optically thick at $\sim 10^{-3}$



Nucleosynthesis Constraints

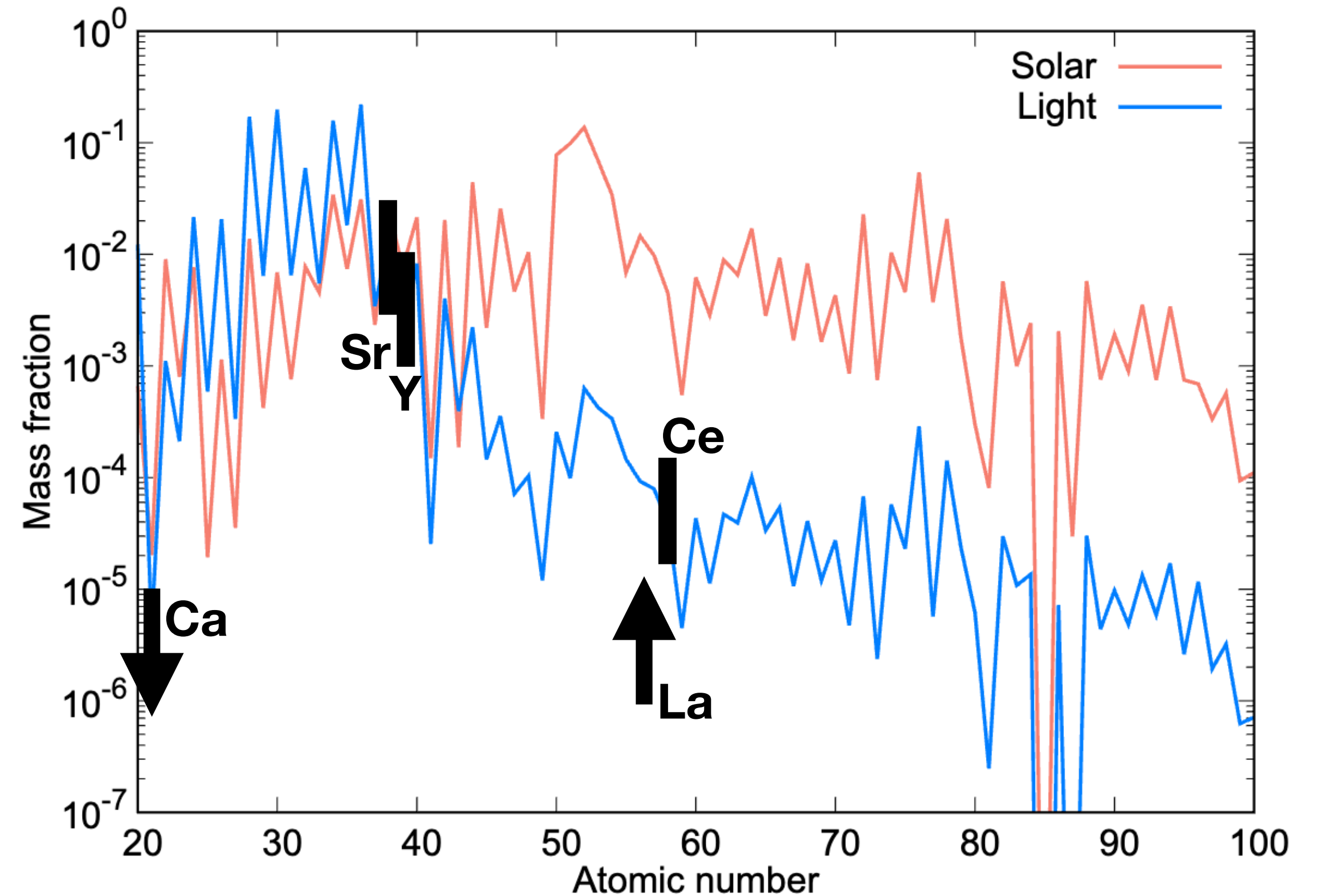
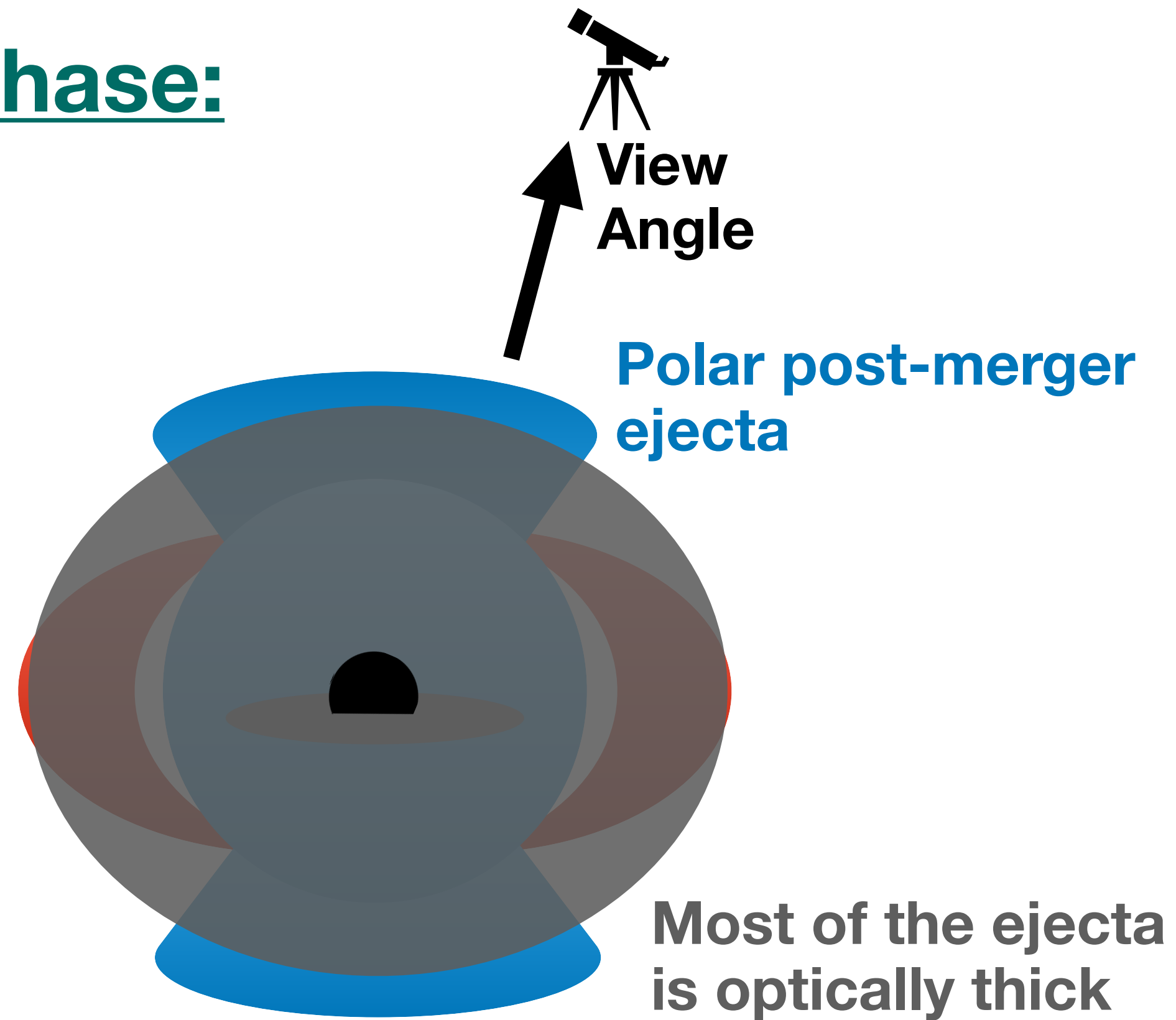


Mass fraction inferred for heavy elements is higher at late time

Discussion

Why different abundance distributions are needed to explain different epochs?

Early phase:

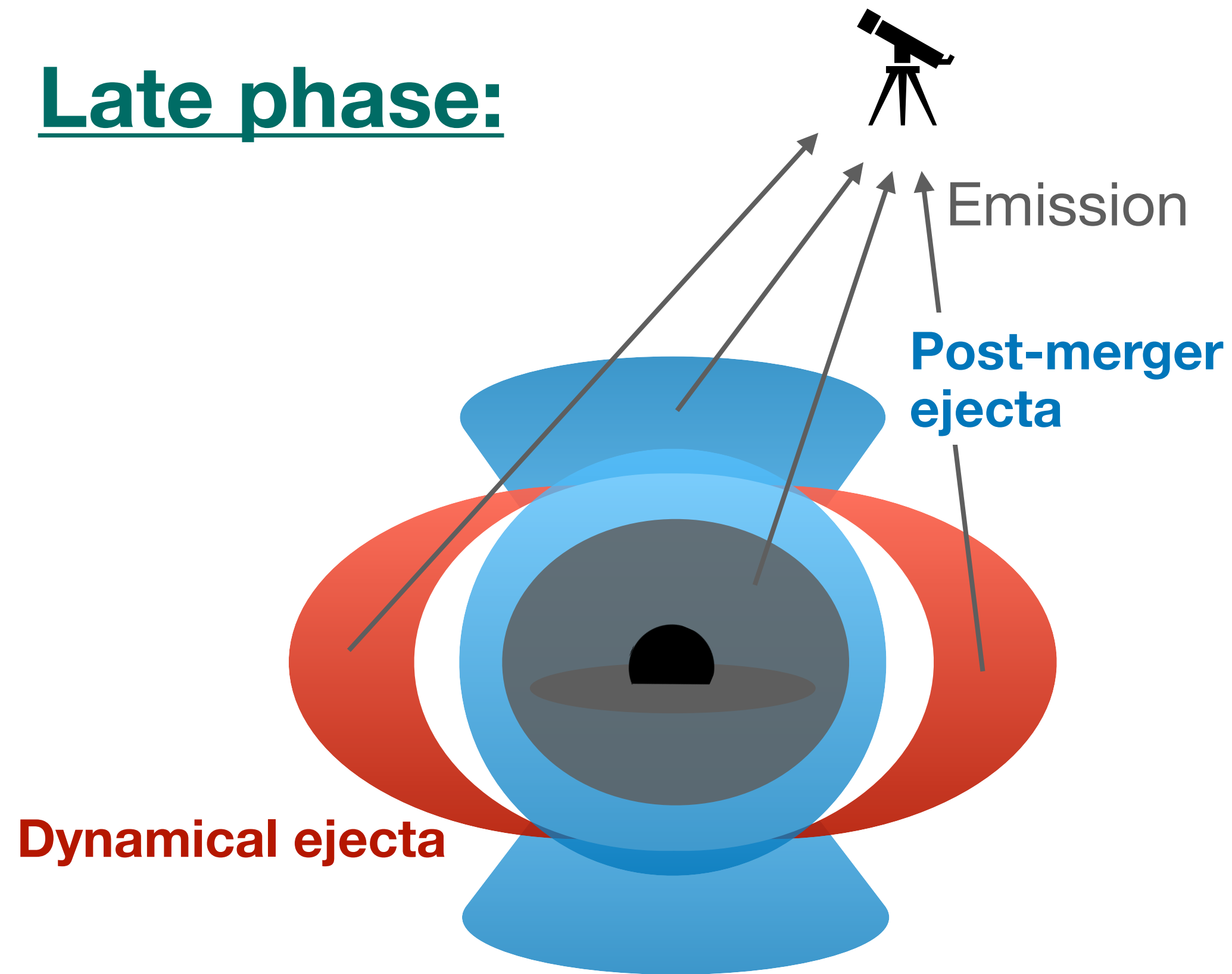


Absorption features appear due to elements along line of sight in the polar region
→ **Lower inferred mass fraction**

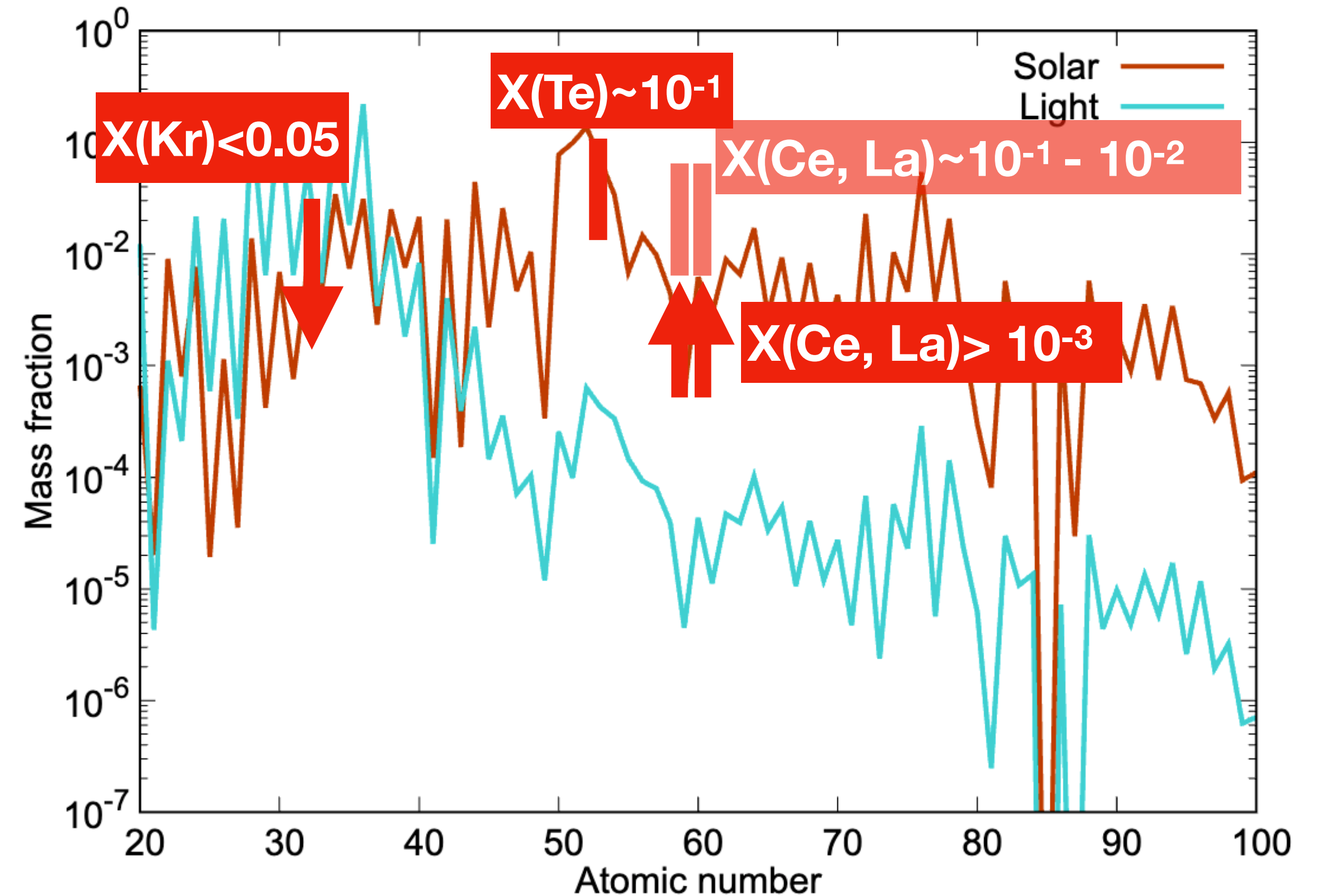
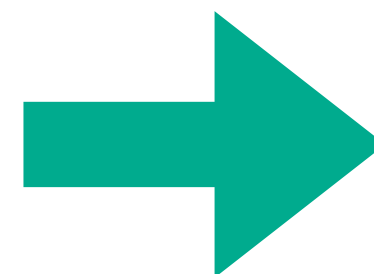
Discussion

Why different abundance distributions are needed to explain different epochs?

Late phase:



Most of the ejecta is optically thin and emission from different parts is observed



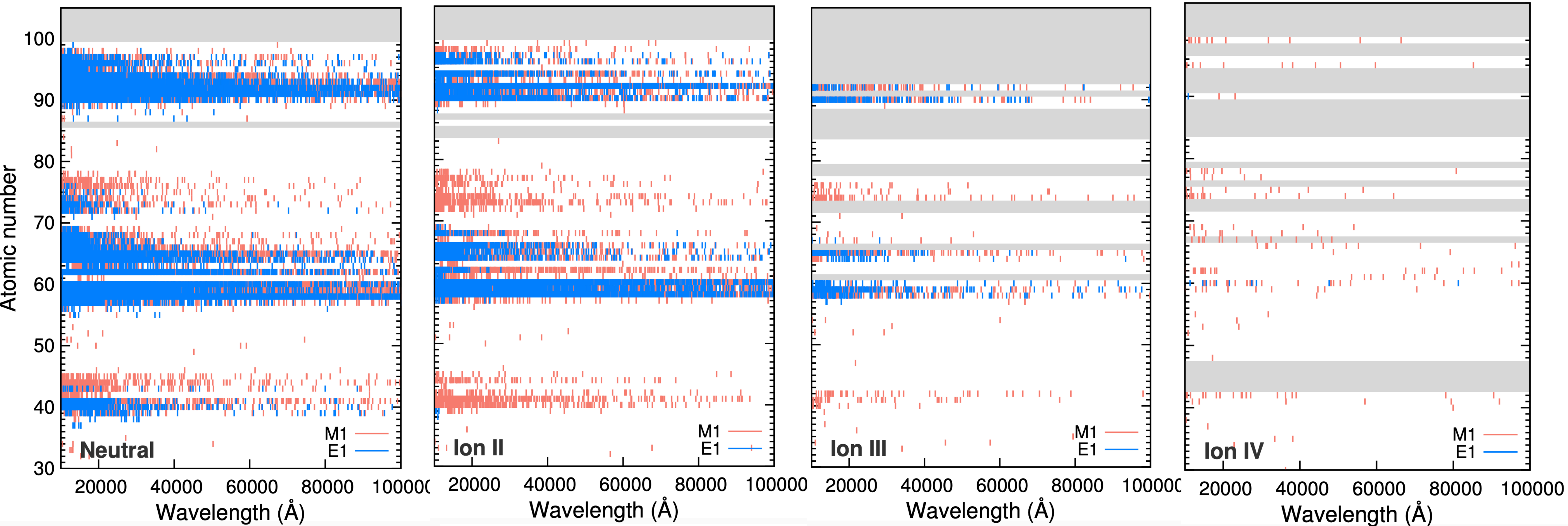
Mass fraction inferred reflects the overall synthesized matter in the merger

GW170817 nucleosynthesis is more consistent with solar r-process!

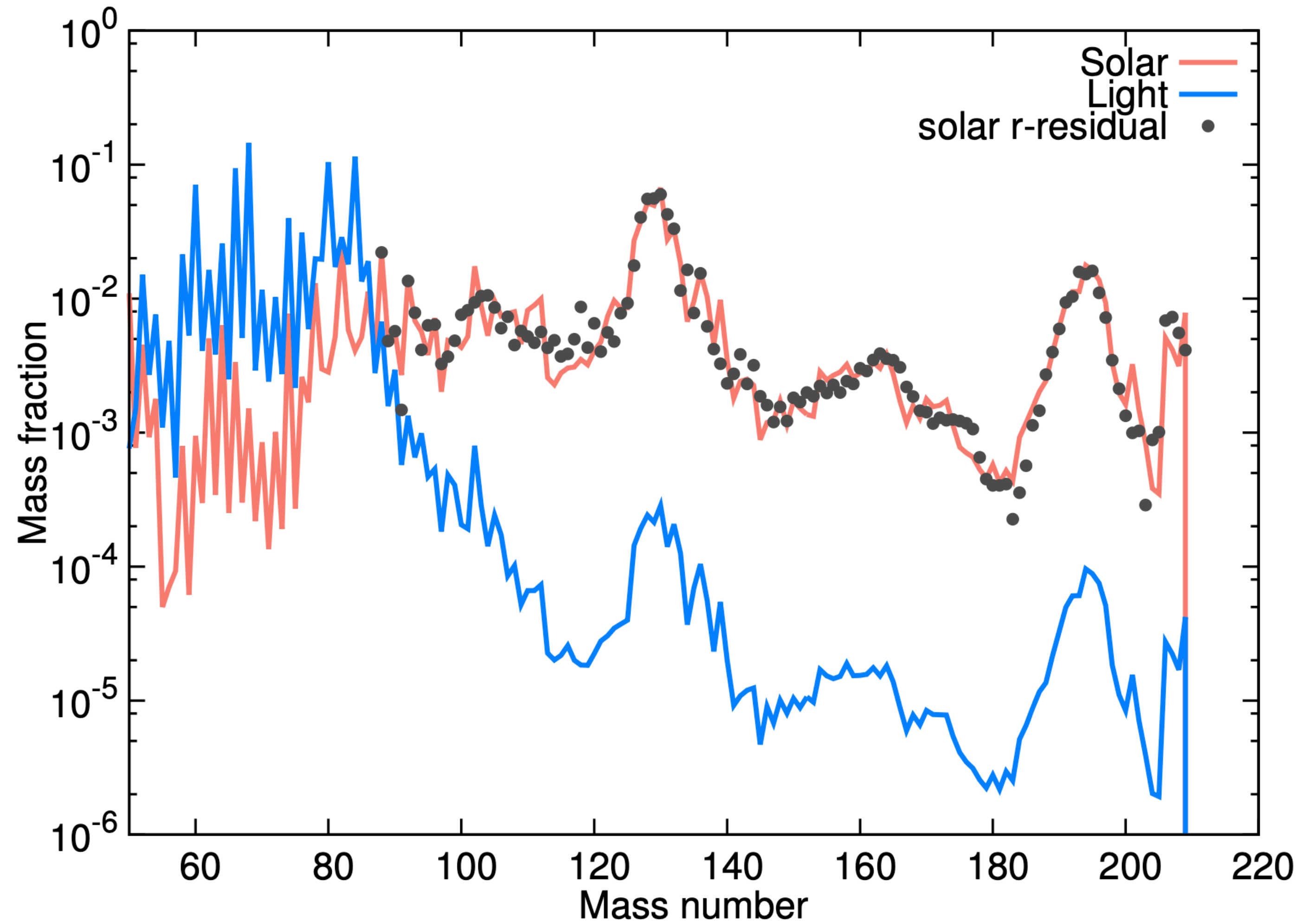
Summary

- **La III, Ce III, and Te III** as main contributors to three late-phase features
- Late phase estimates reflect the overall ejecta abundance
- Nucleosynthesis of heavy elements in GW170817 is **consistent with solar r-process abundance**

Selected Transitions



Abundance distributions used compared with the solar r-process residuals



Atomic properties estimates

$n_{cr} < n_e$ (M1 transitions)

$$L = \Delta E n_X \frac{g_2}{\sum_k g_k e^{-\Delta E_k/kT}} e^{-E_2/kT} A_{2 \rightarrow 1}$$

For fine-structure:

$$A_{2 \rightarrow 1} = 1.3 \text{ s}^{-1} \left(\frac{\lambda}{4 \mu\text{m}} \right)^{-3} f(J, L, S)$$

Where

$$f(L, S, J) = \frac{(J^2 - (L - S)^2)((L + S + 1)^2 - J^2)}{12J(2J + 1)} \quad \text{if } J_u = J_l + 1$$

$$f(L, S, J) = \frac{((J + 1)^2 - (L - S)^2)((L + S + 1)^2 - (J + 1)^2)}{12(J + 1)(2J + 1)} \quad \text{if } J_u = J_l + 1$$

Pasternack (1940)

Other M1 transitions: $A_{2 \rightarrow 1} = 1$

$n_{cr} > n_e$ (E1 transitions and a few M1)

$$L = \Delta E n_X \frac{g_1}{\sum_k g_k e^{-\Delta E_k/kT}} e^{-E_1/kT} q_{2 \rightarrow 1}$$

$$q_{1 \rightarrow 2} = \frac{8.629 \times 10^{-6}}{T^{1/2}} \frac{\Upsilon_{1 \rightarrow 2}}{g_1} \exp(-\Delta E/kT)$$

Where $\Upsilon_{2 \rightarrow 1} = \int_0^\infty \Omega_{2 \rightarrow 1} e^{-E/kT} d\left(\frac{E}{kT}\right)$

Calibrated theoretical calculations by Flörs+25

For E1:

$$\Omega_{2 \rightarrow 1} = 2.388 \bar{G} \left(\frac{\lambda}{1 \mu\text{m}} \right)^3 \left(\frac{g A_{2 \rightarrow 1}}{10^6 \text{ s}^{-1}} \right)$$

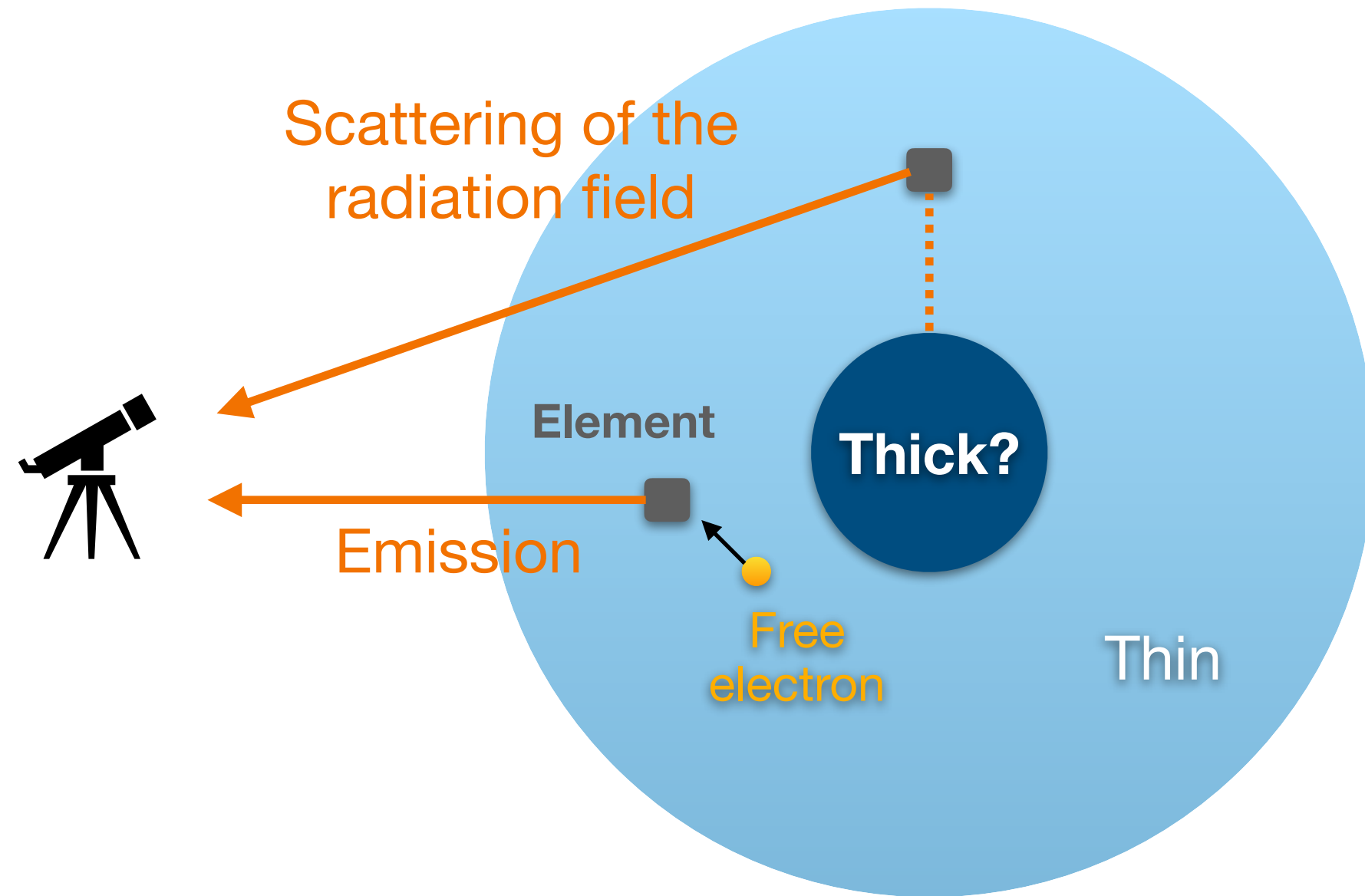
Van Regemorter (1962)

For M1, we assume a typical $\Omega_{2 \rightarrow 1} = 1$

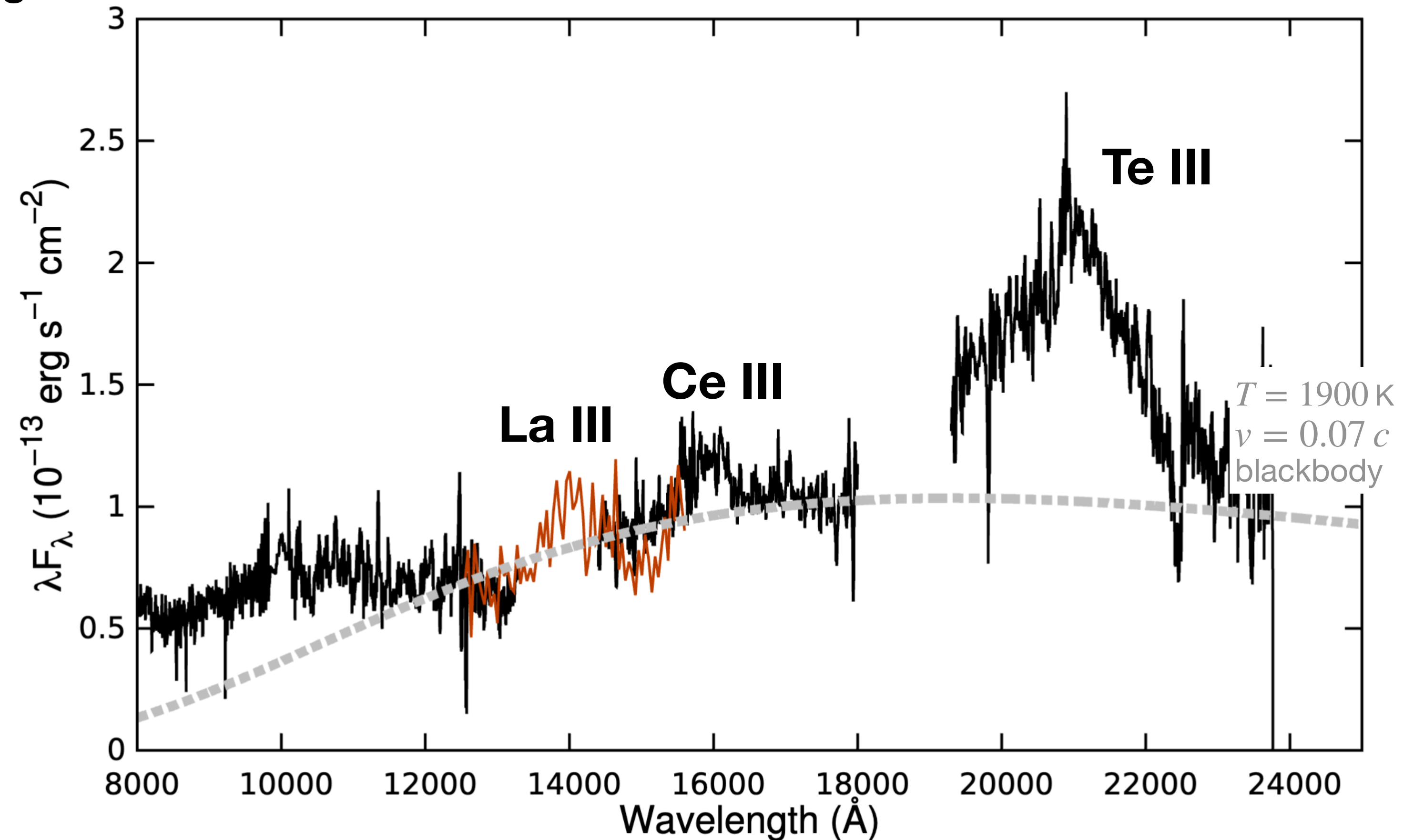
Nucleosynthesis Constraints

Good, but what about the radiation field?

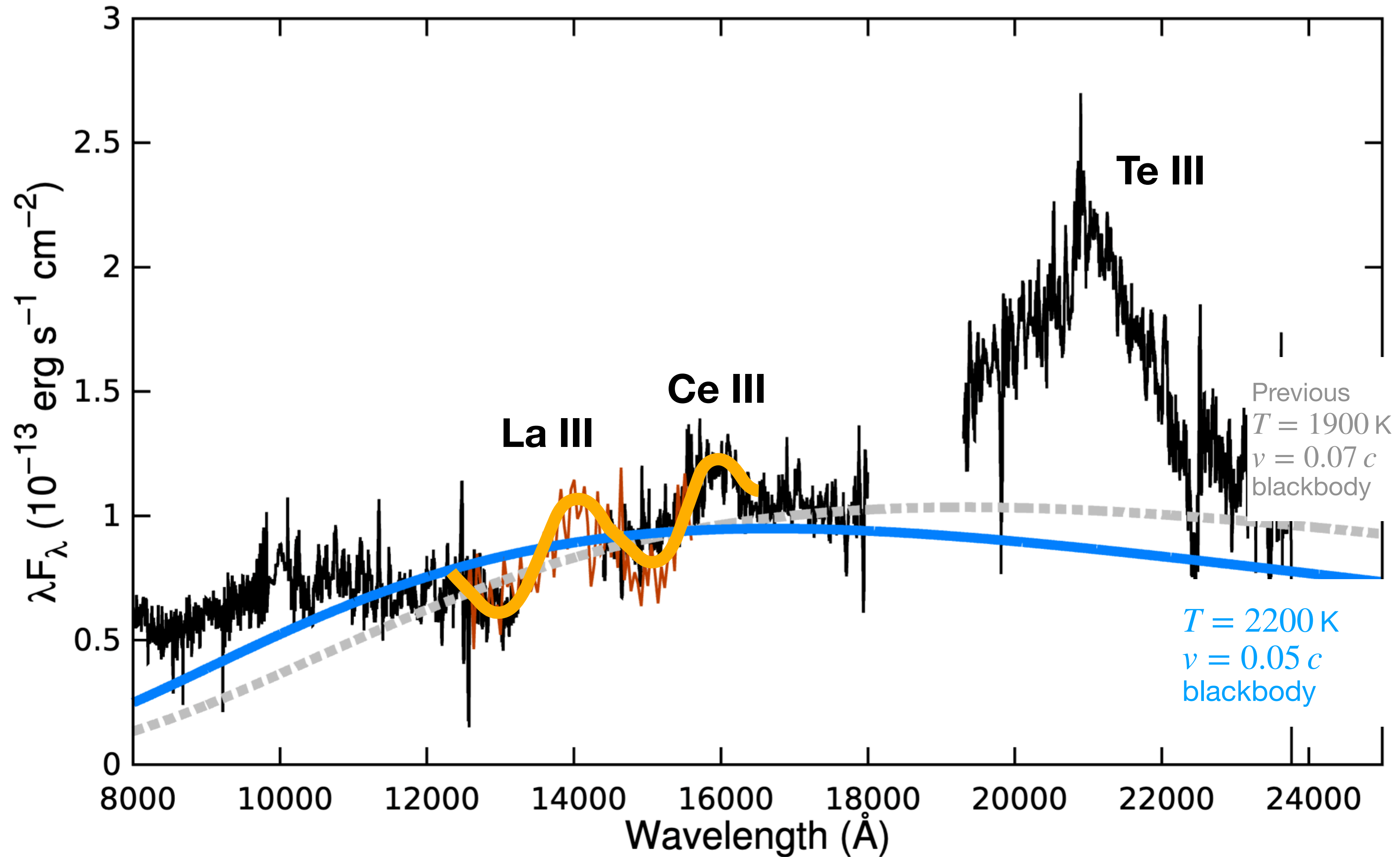
If the radiation field is coming from an inner optically thick region, allowed (E1) transitions can be caused by two mechanisms



→ La III and Ce III should show P-Cygni features if optically thick

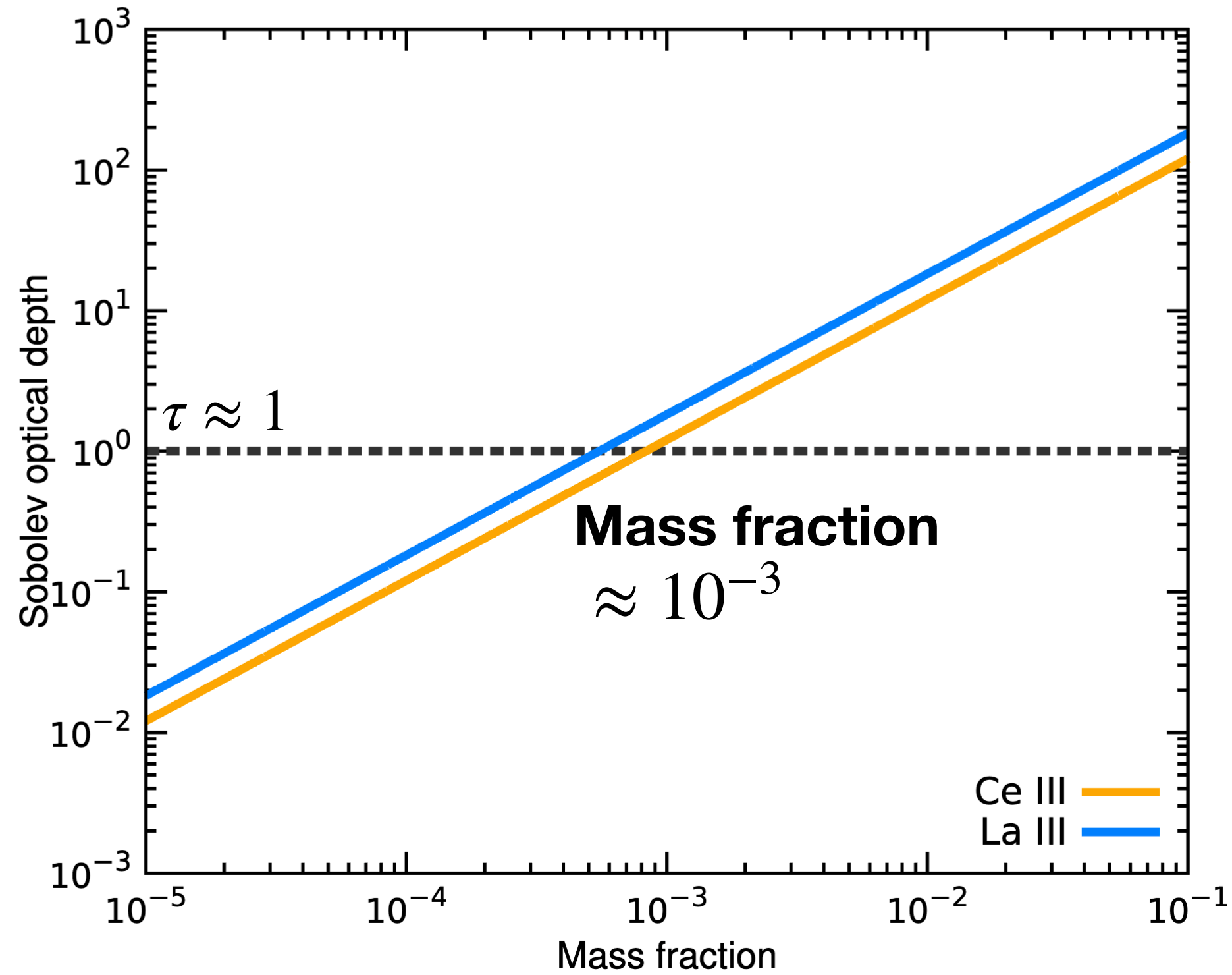


Are La III and Ce III optically thick?

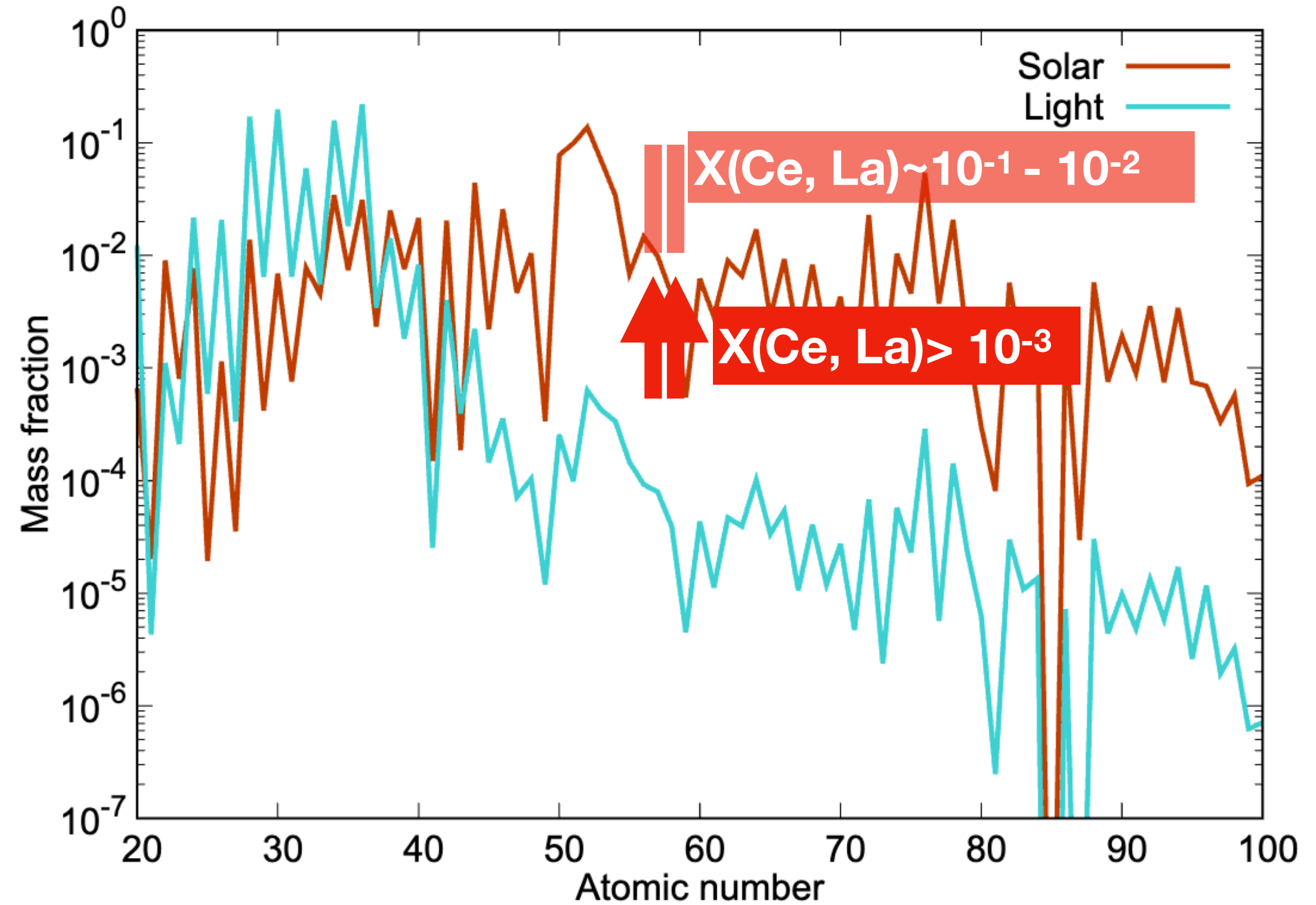
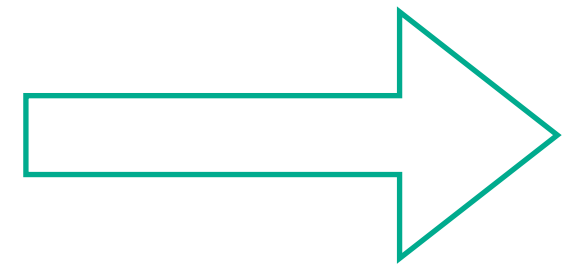


Nucleosynthesis Constraints

Sobolev optical depth of Ce III and La III as a function of their mass fractions



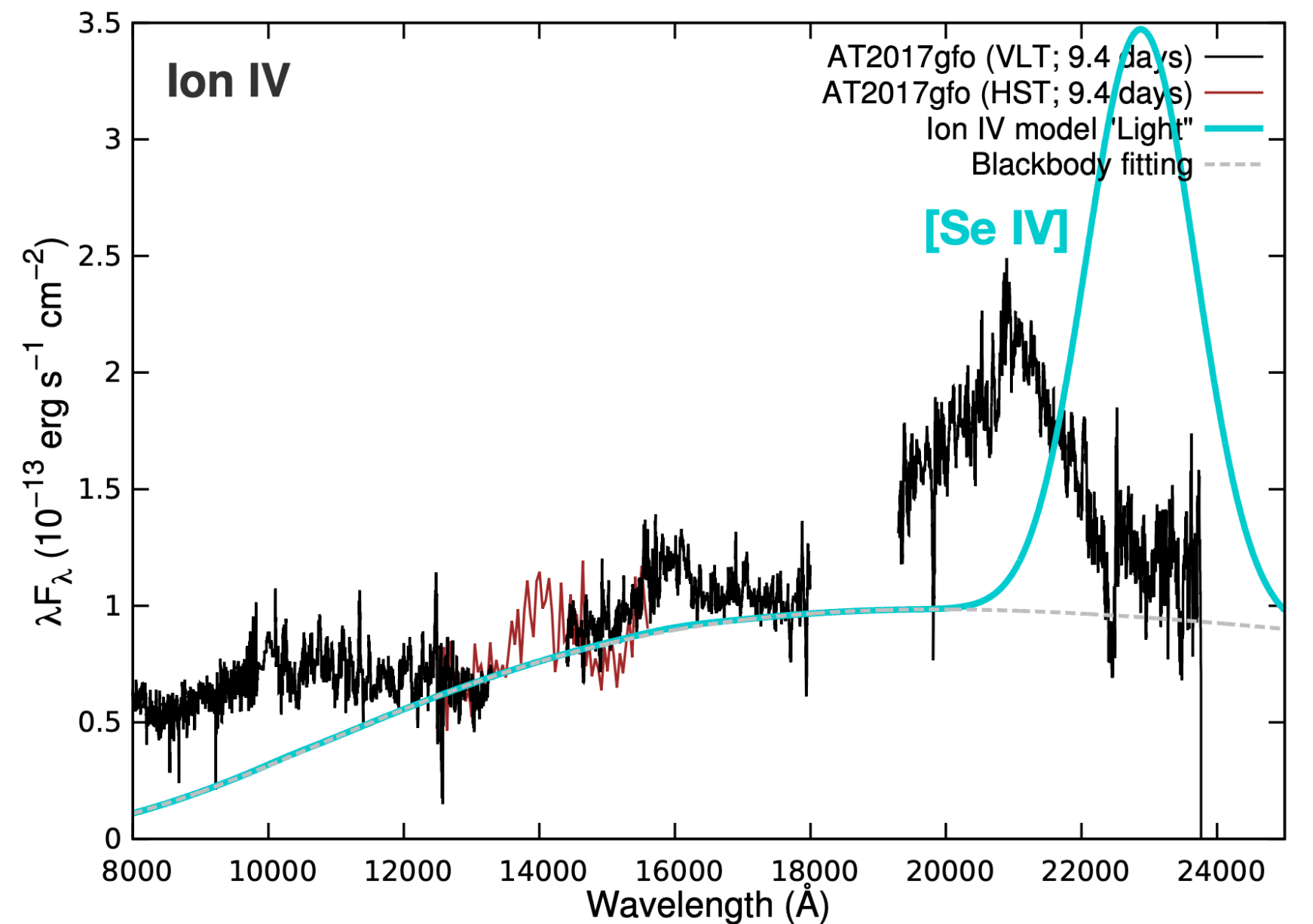
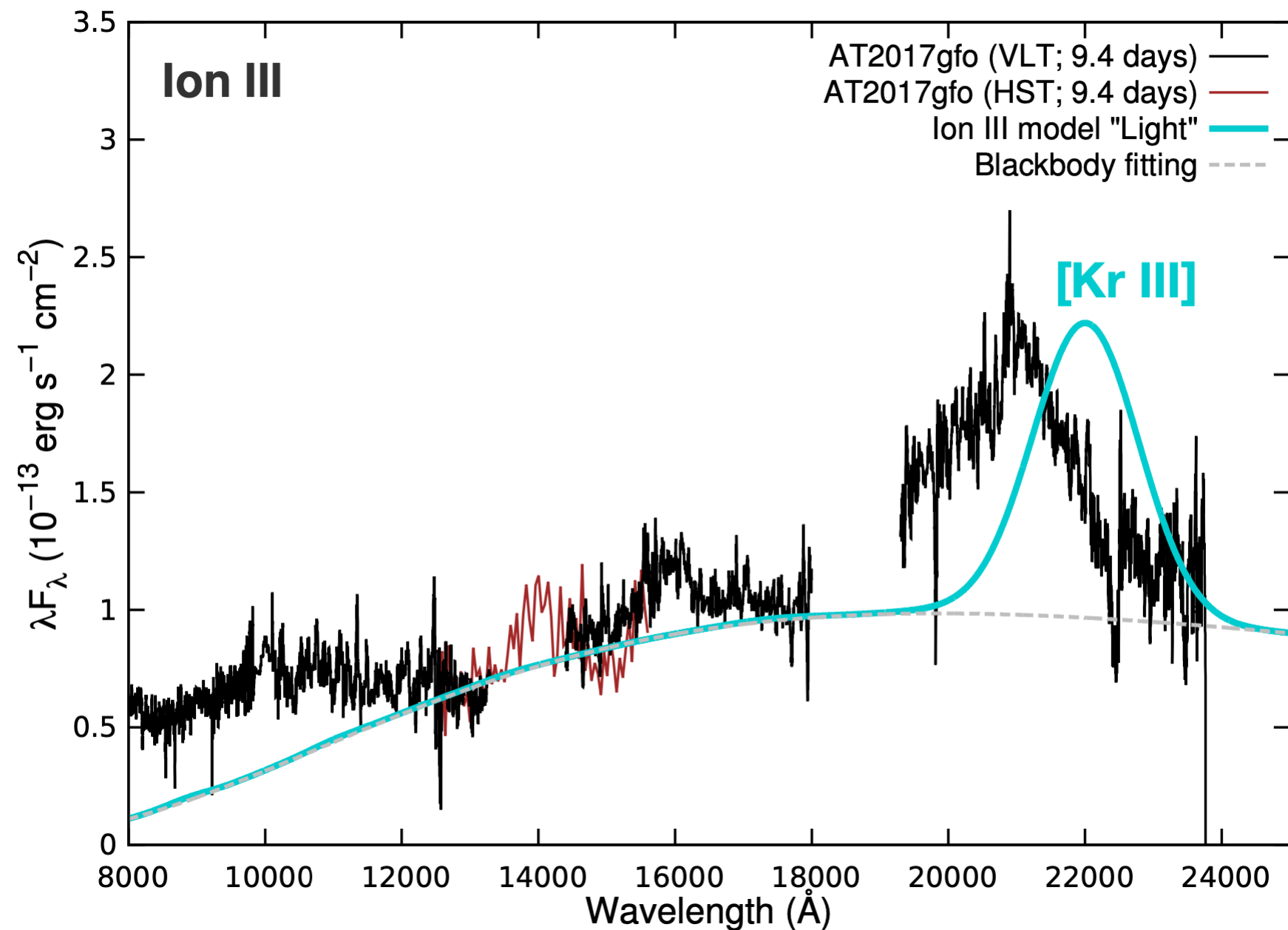
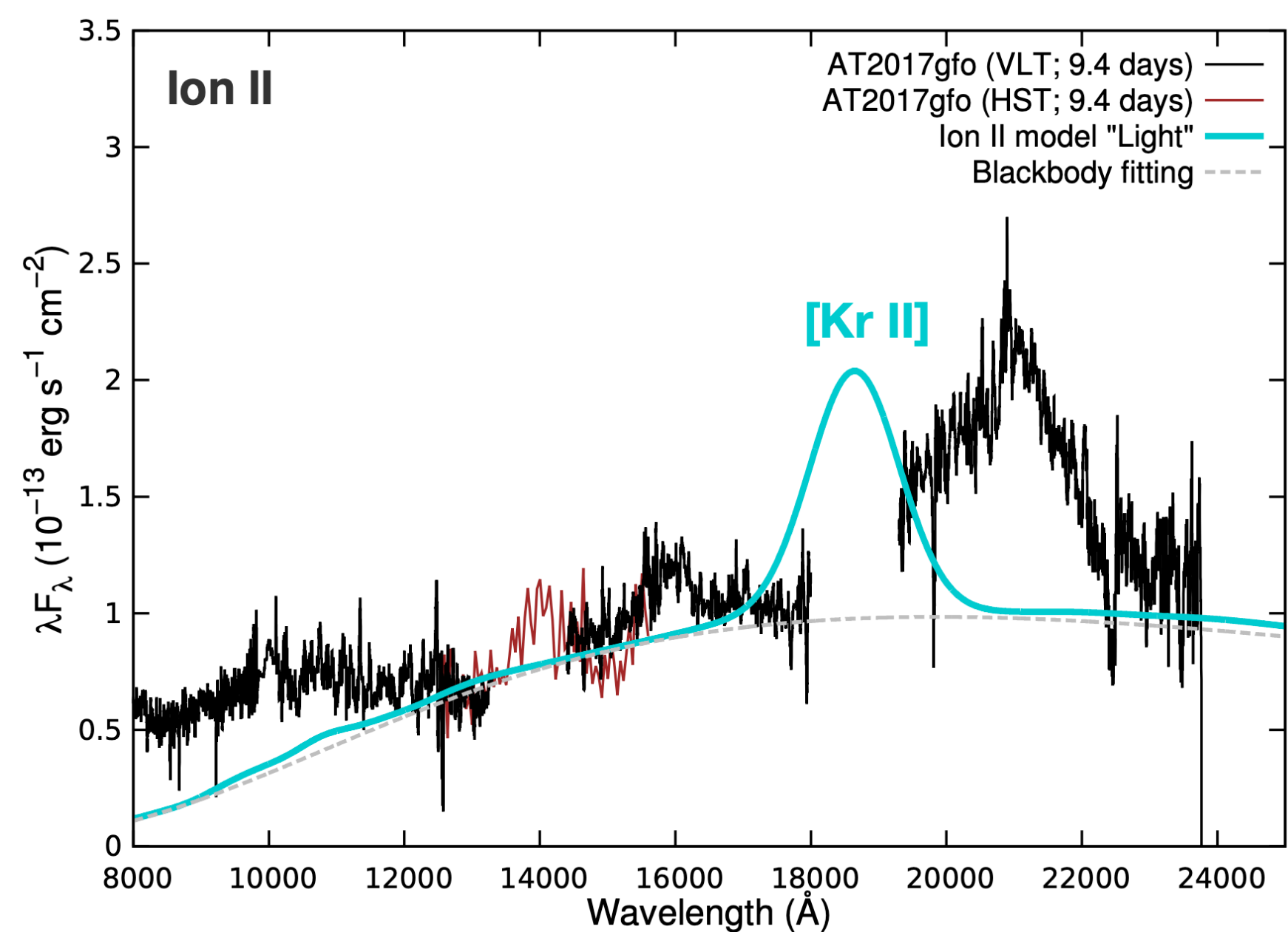
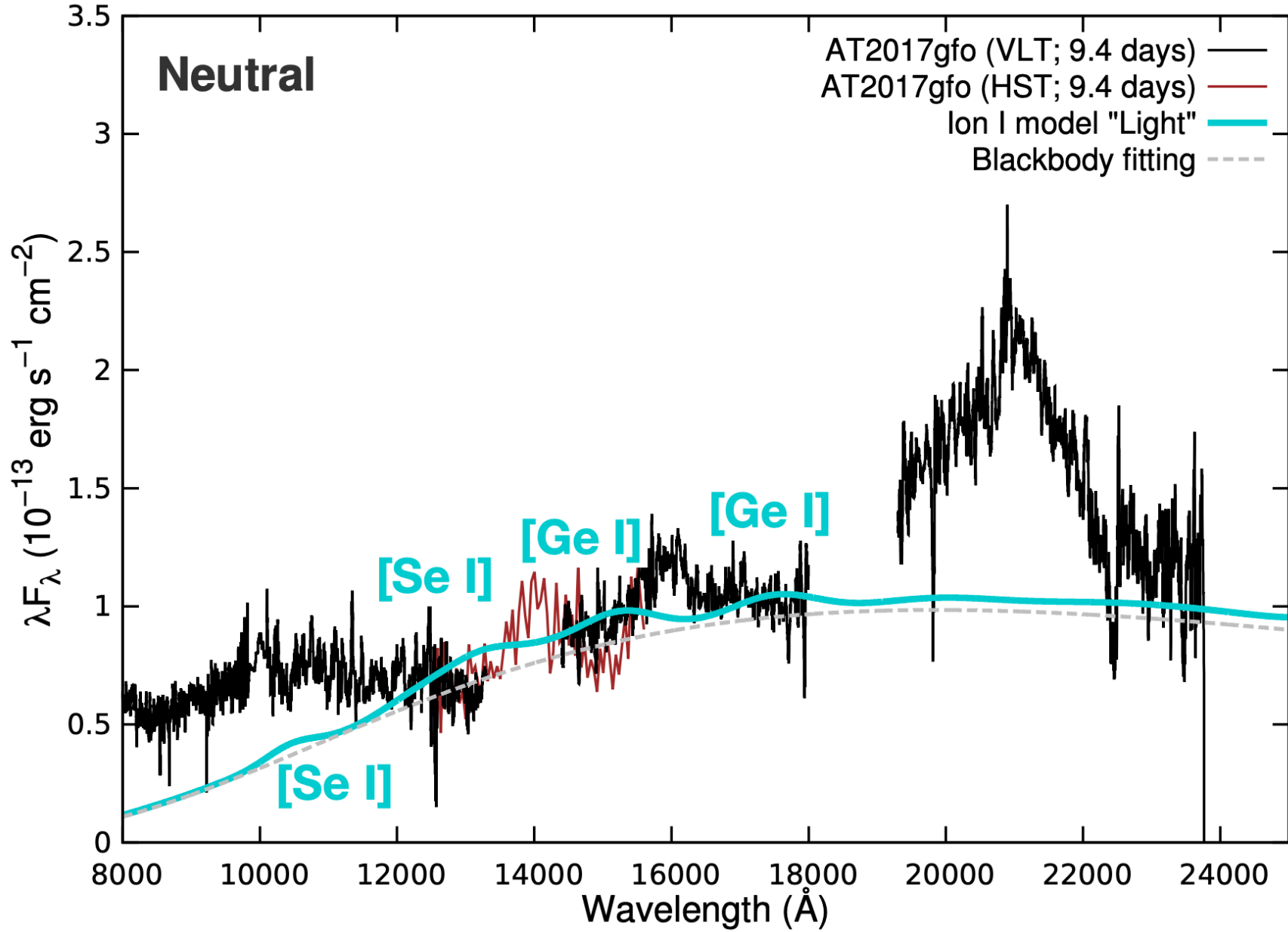
Transition is optically thick if $\tau \approx 1$



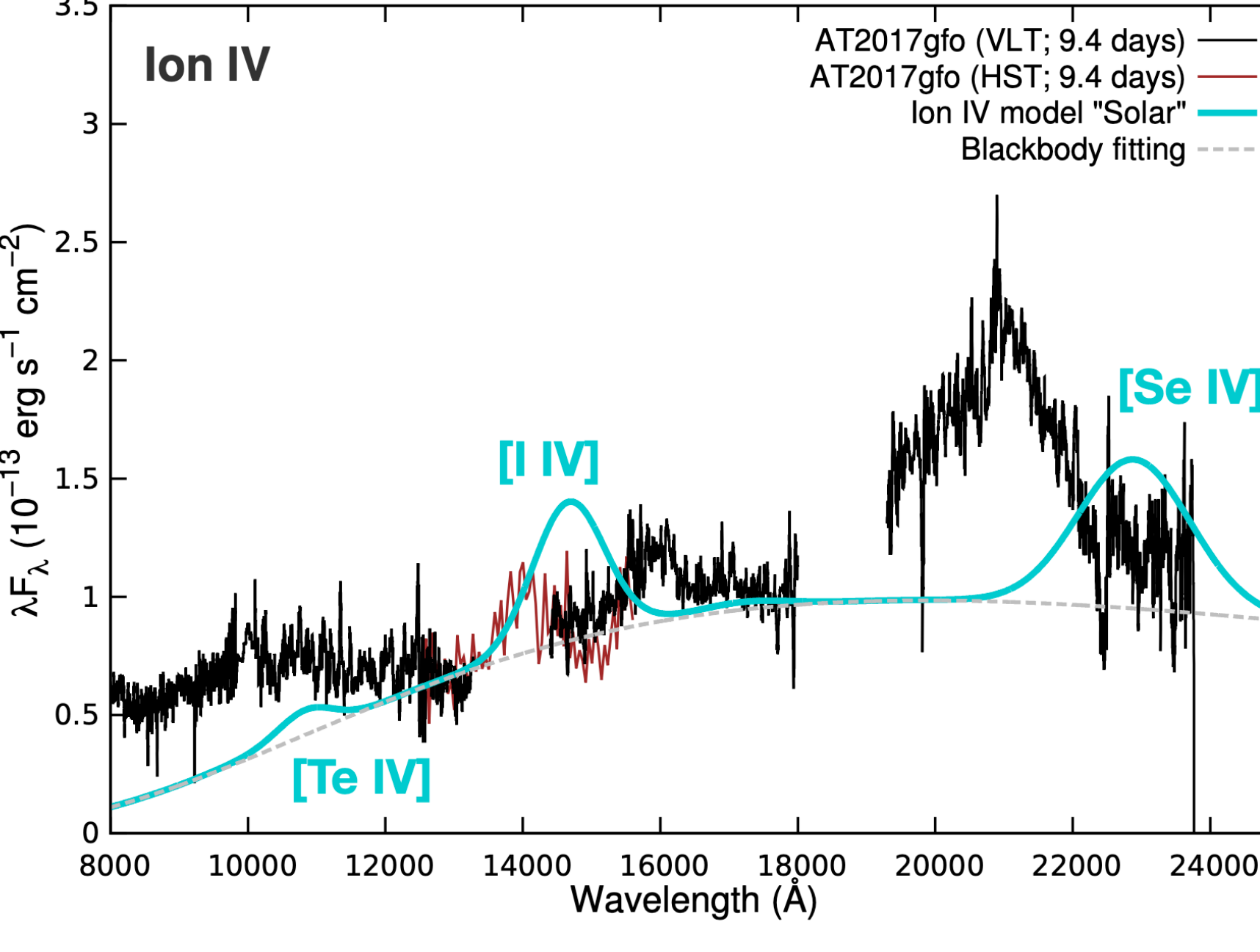
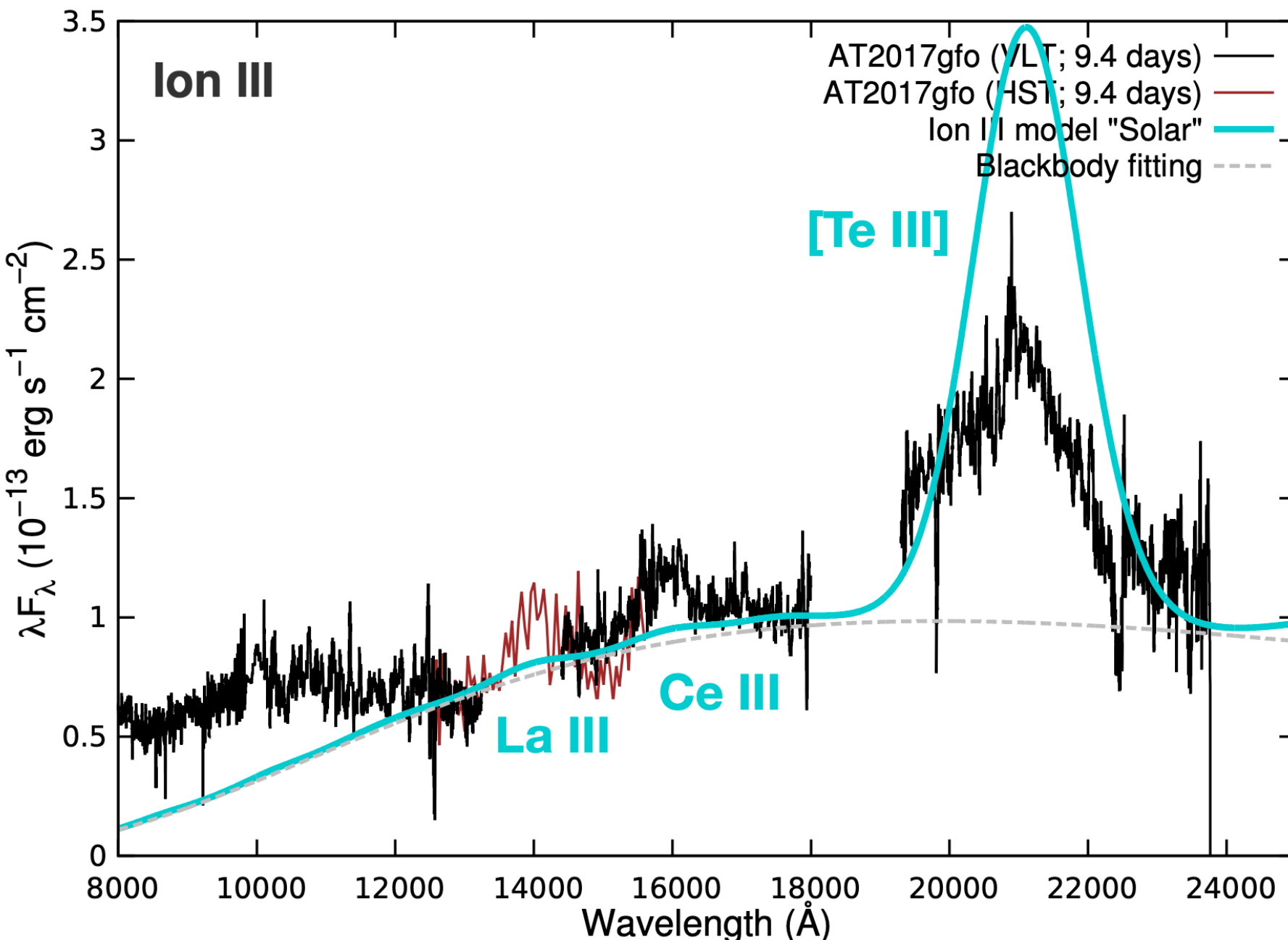
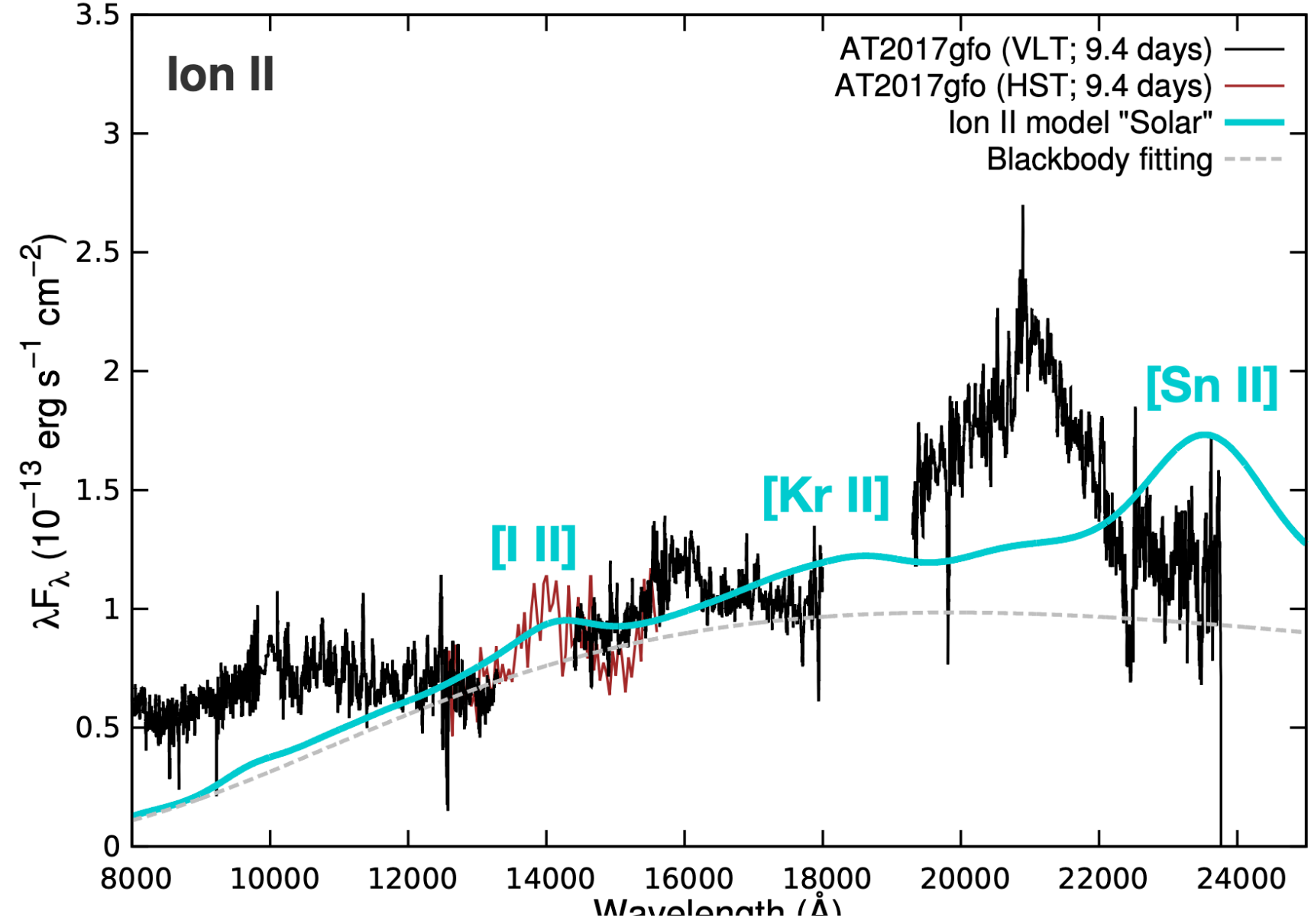
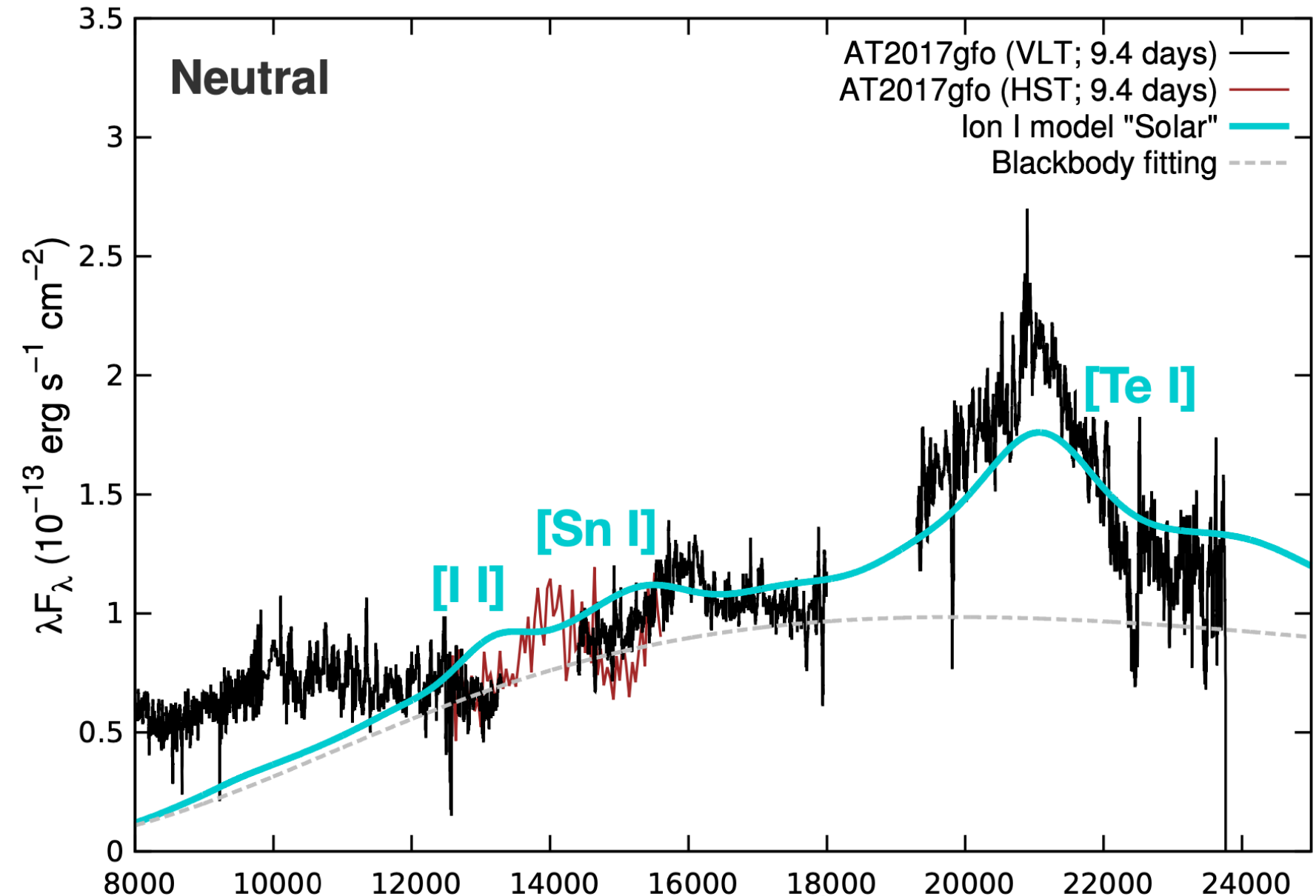
c.f. Sobolev optical depth

$$\tau = \frac{\pi e^2}{m_e c} n_{i,j} t \lambda g f \frac{1}{\Sigma} e^{-\frac{E_l}{kT}}$$

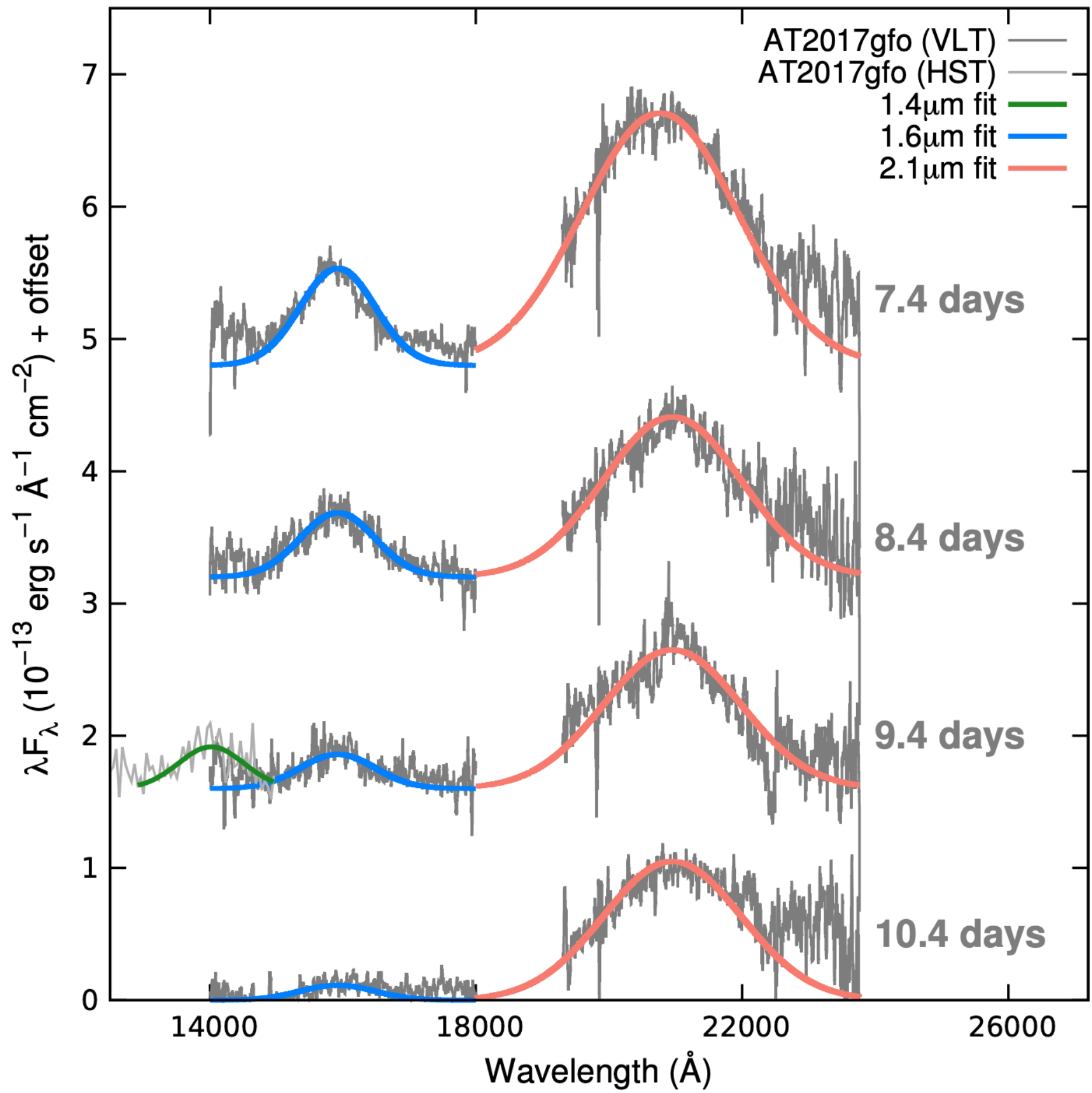
Models and Spectrum Comparison: light abundance



Models and Spectrum Comparison: solar abundance



Gaussian Fitting best-fit parameters



Feature	λ_{peak} (Å)	v_{FWHM}	Luminosity (erg s $^{-1}$)
7.4 days			
1.6 μm	15900	0.12 c	1.26×10^{39}
2.1 μm	20700	0.24 c	4.89×10^{39}
8.4 days			
1.6 μm	15900	0.12 c	8.11×10^{38}
2.1 μm	20900	0.16 c	2.87×10^{39}
9.4 days			
1.4 μm	14000	0.12 c	5.40×10^{38}
1.6 μm	15900	0.12 c	4.44×10^{38}
2.1 μm	20900	0.16 c	2.50×10^{39}
10.4 days			
1.6 μm	15900	0.12 c	1.93×10^{38}
2.1 μm	20900	0.16 c	2.4×10^{39}

Required mass fraction fro Ce, La, and Te

Ce III feature
is due to
three lines

$$L = \Delta E_{1.59} n_e \frac{g_l}{\sum_k g_k e^{-E_k/kT}} \frac{X_{Ce} M_{ej}}{A_{Ce} m_u} \frac{8.6 \times 10^{-6}}{\sqrt{T}} \frac{\Upsilon_{1.59}}{g_l} e^{-\Delta E_{1.59}/kT}$$

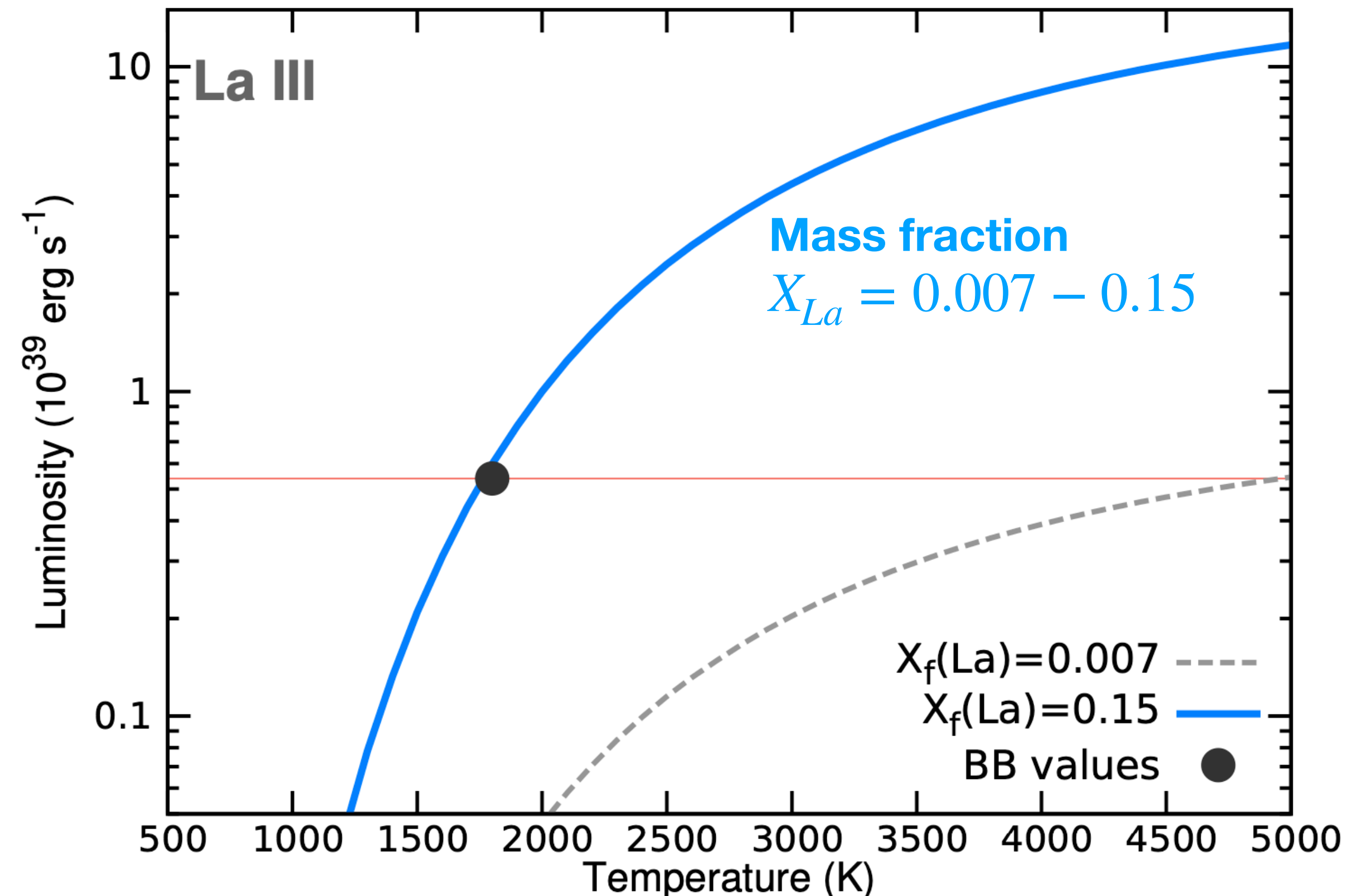
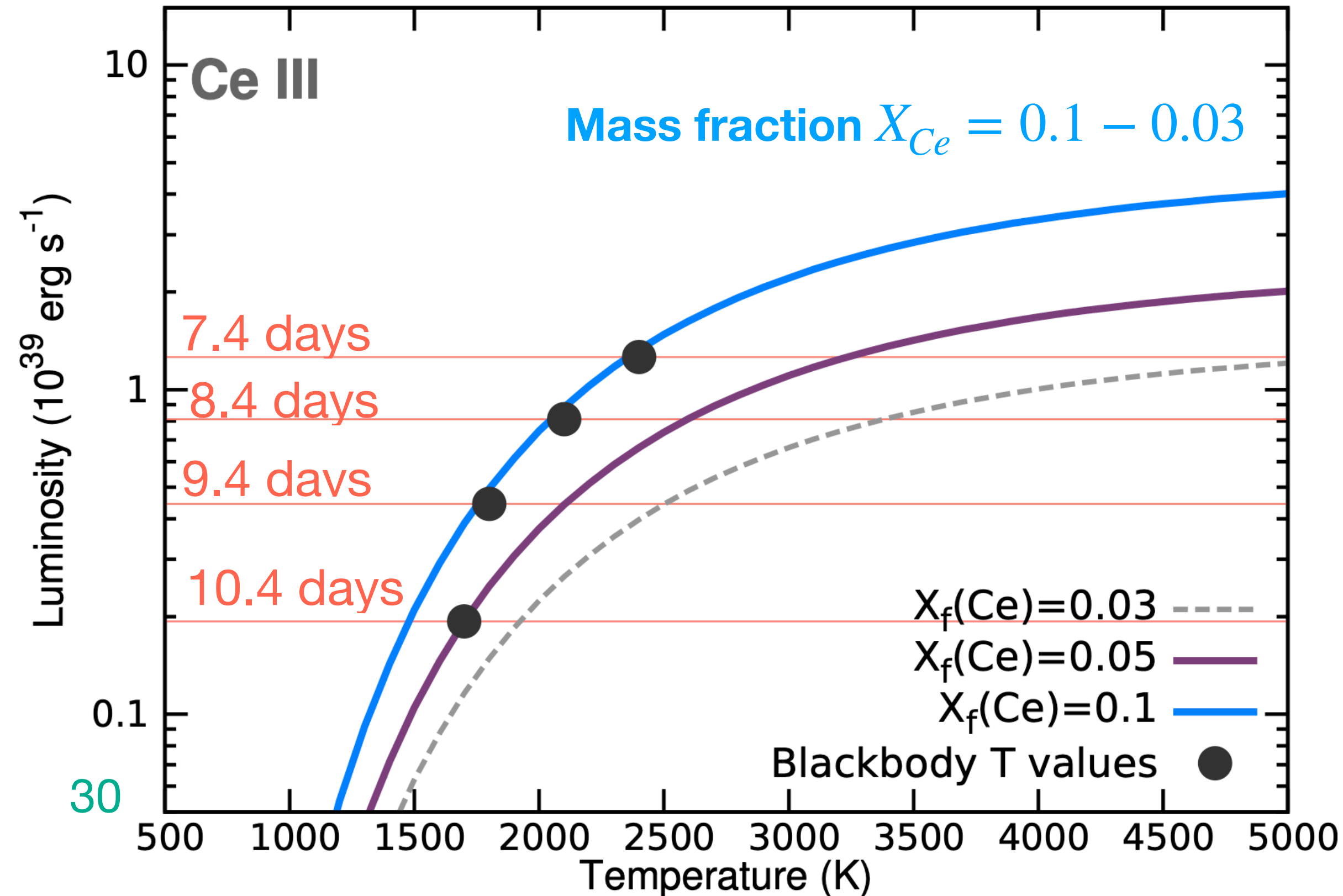
$$+ \Delta E_{1.58} n_e \frac{g_l e^{-E_l/kT}}{\sum_k g_k e^{-E_k/kT}} \frac{X_{Ce} M_{ej}}{A_{Ce} m_u} \frac{8.6 \times 10^{-6}}{\sqrt{T}} \frac{\Upsilon_{1.58}}{g_l} e^{-\Delta E_{1.58}/kT}$$

$$+ \Delta E_{1.61} n_e \frac{g_l e^{-E_l/kT}}{\sum_k g_k e^{-E_k/kT}} \frac{X_{Ce} M_{ej}}{A_{Ce} m_u} \frac{8.6 \times 10^{-6}}{\sqrt{T}} \frac{\Upsilon_{1.61}}{g_l} e^{-\Delta E_{1.61}/kT}$$

La III feature
is due to two
line

$$L = \Delta E_{1.38} n_e \frac{g_l}{\sum_k g_k e^{-E_k/kT}} \frac{X_{La} M_{ej}}{A_{La} m_u} \frac{8.6 \times 10^{-6}}{\sqrt{T}} \frac{\Upsilon_{1.38}}{g_l} e^{-\Delta E_{1.38}/kT}$$

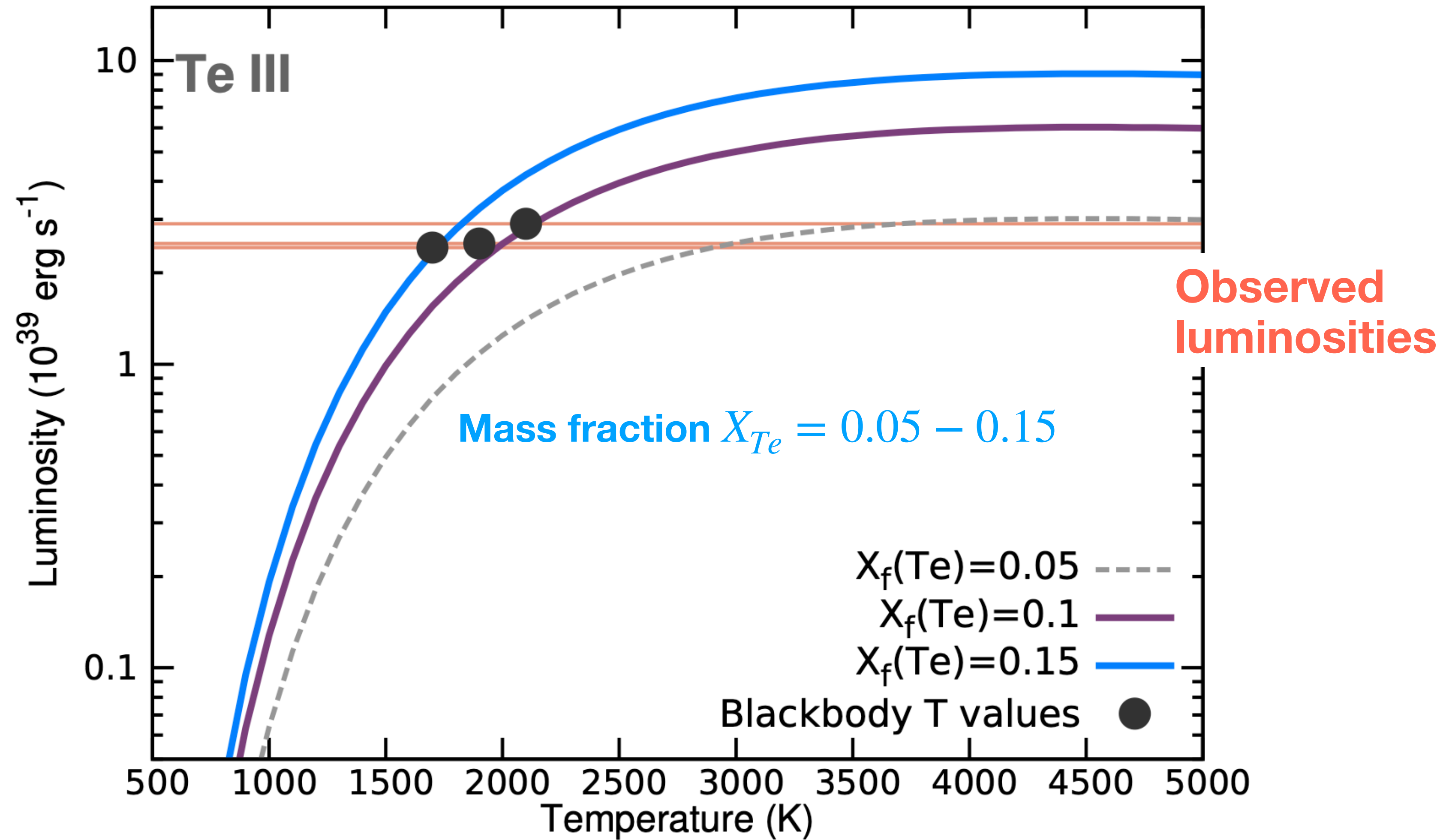
$$+ \Delta E_{1.41} n_e \frac{g_l}{\sum_k g_k e^{-E_k/kT}} \frac{X_{La} M_{ej}}{A_{La} m_u} \frac{8.6 \times 10^{-6}}{\sqrt{T}} \frac{\Upsilon_{1.41}}{g_l} e^{-\Delta E_{1.41}/kT}$$



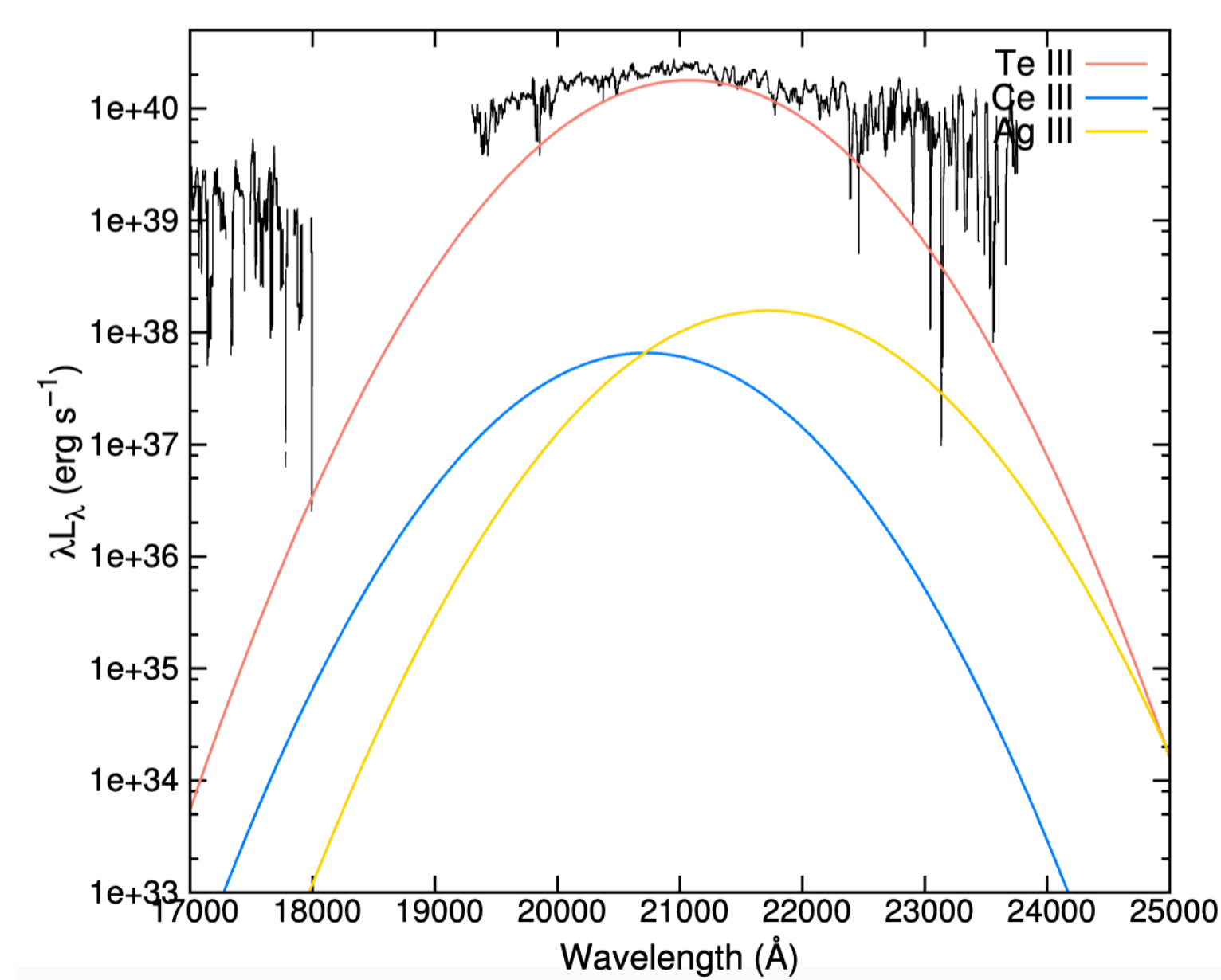
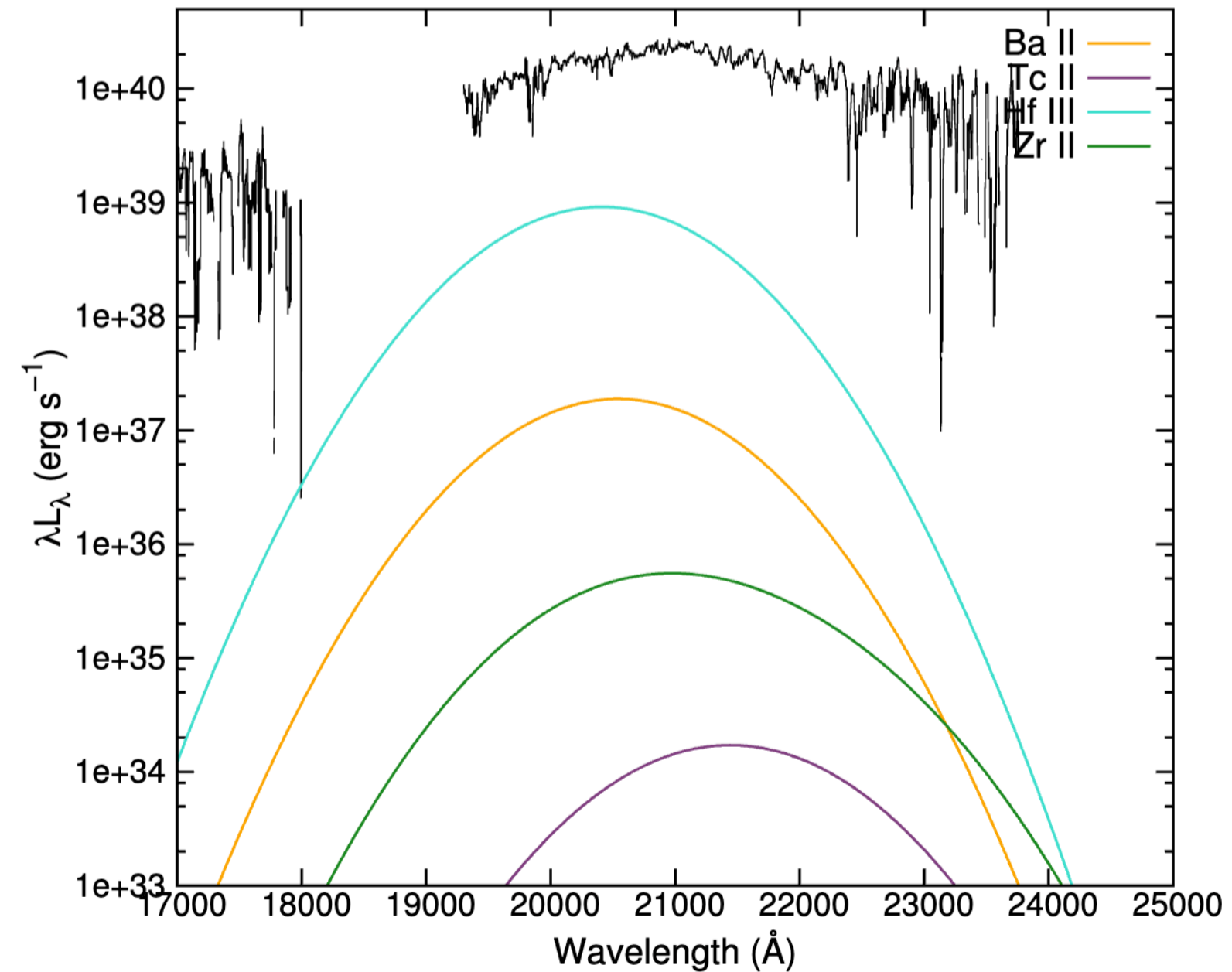
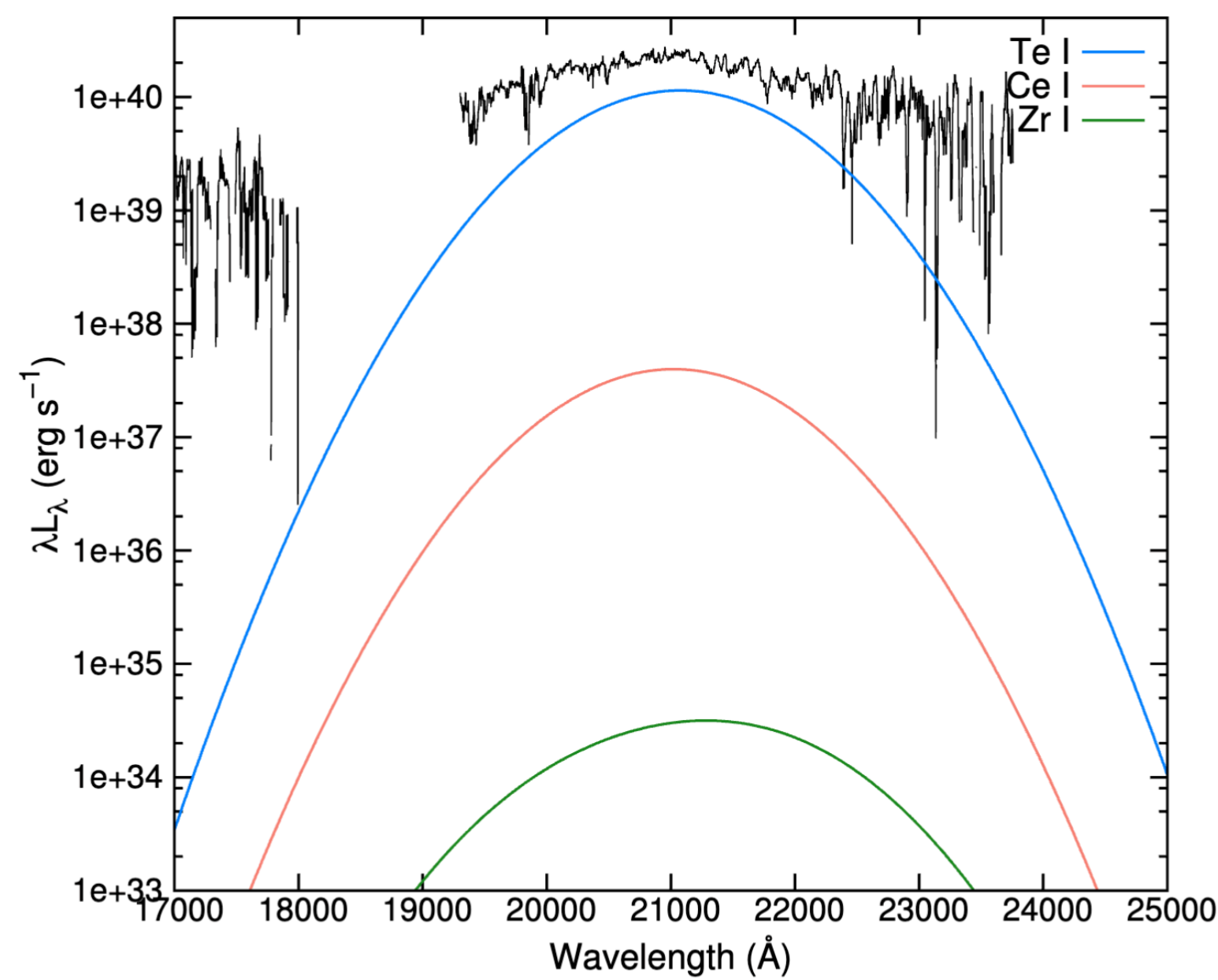
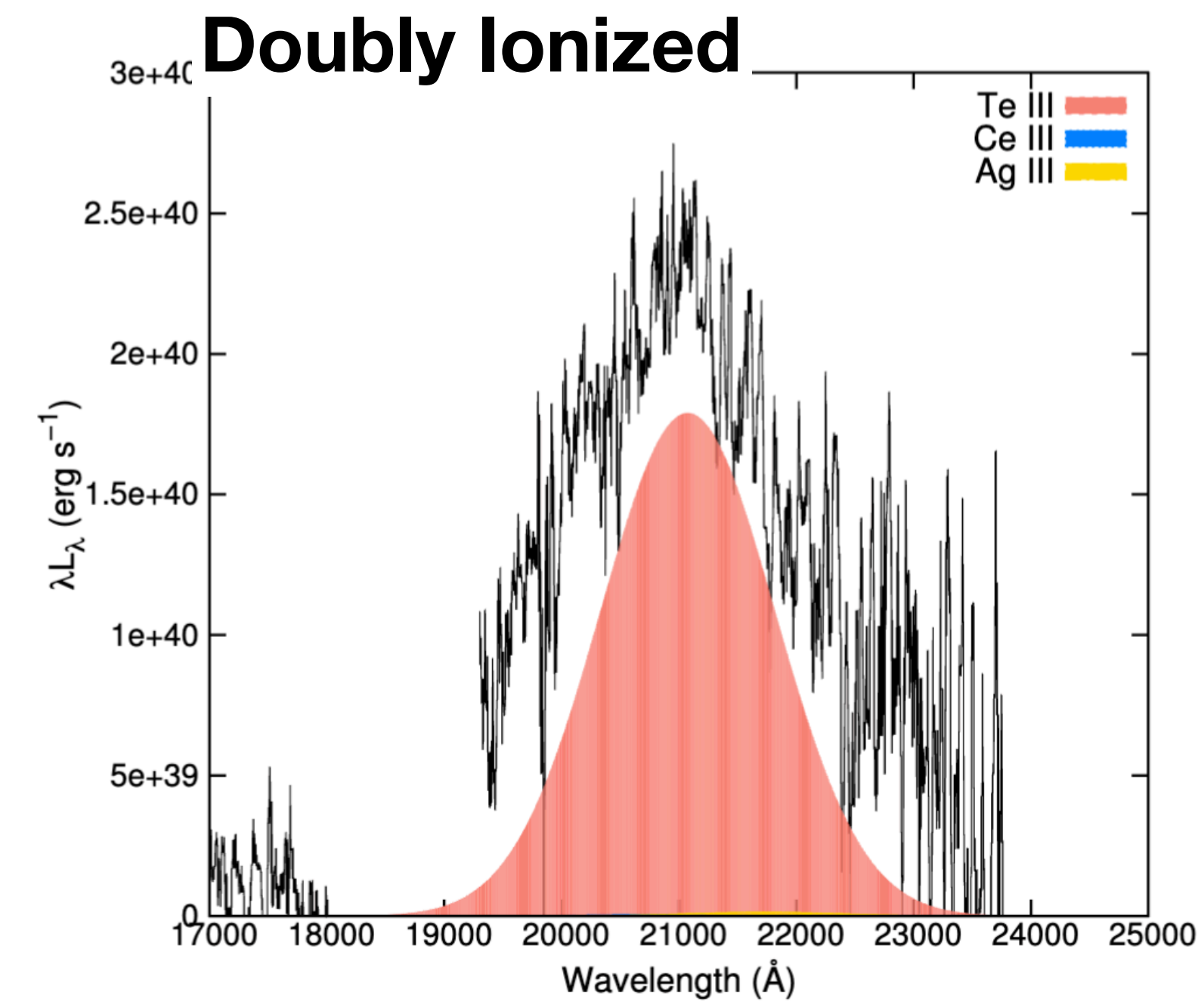
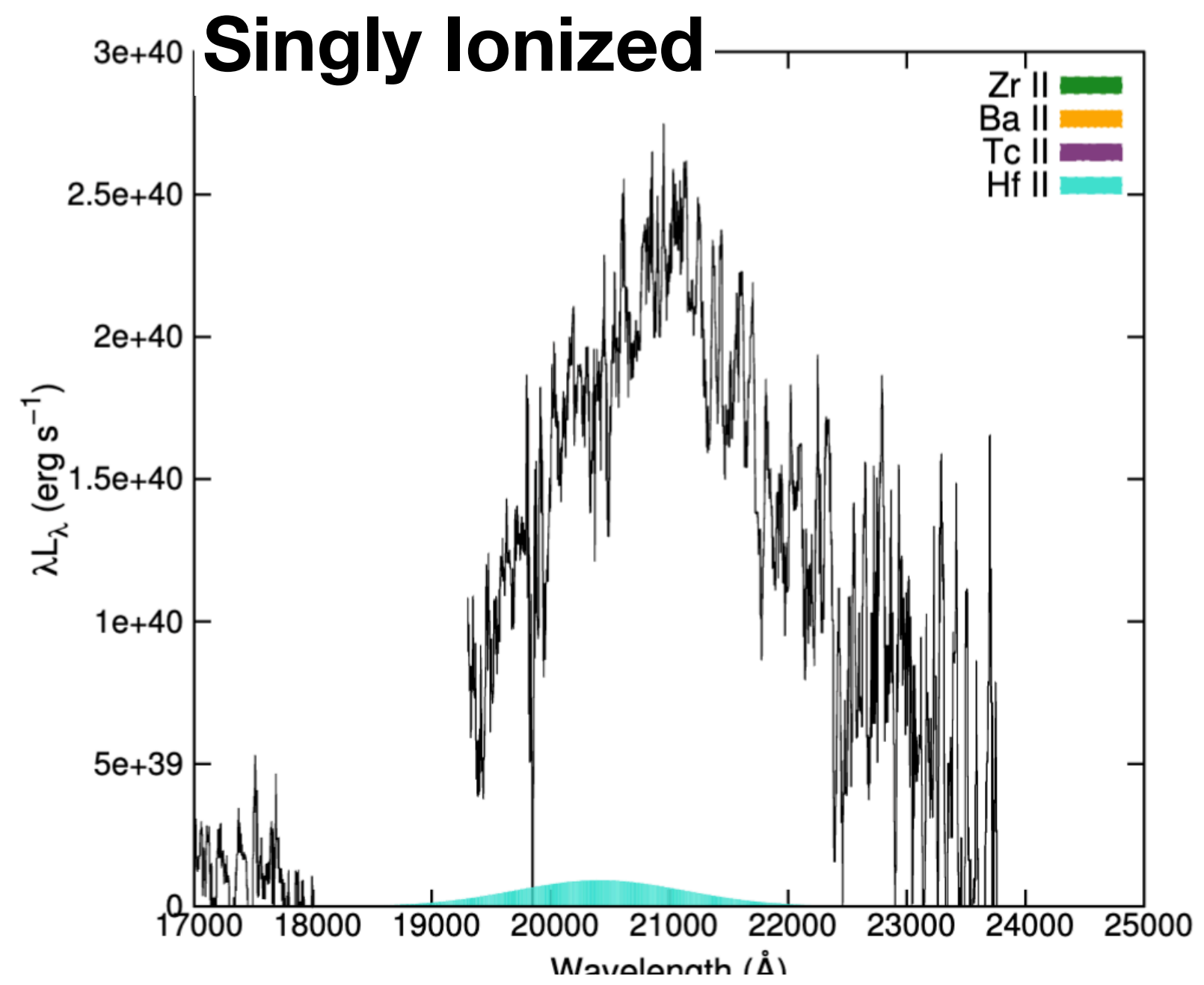
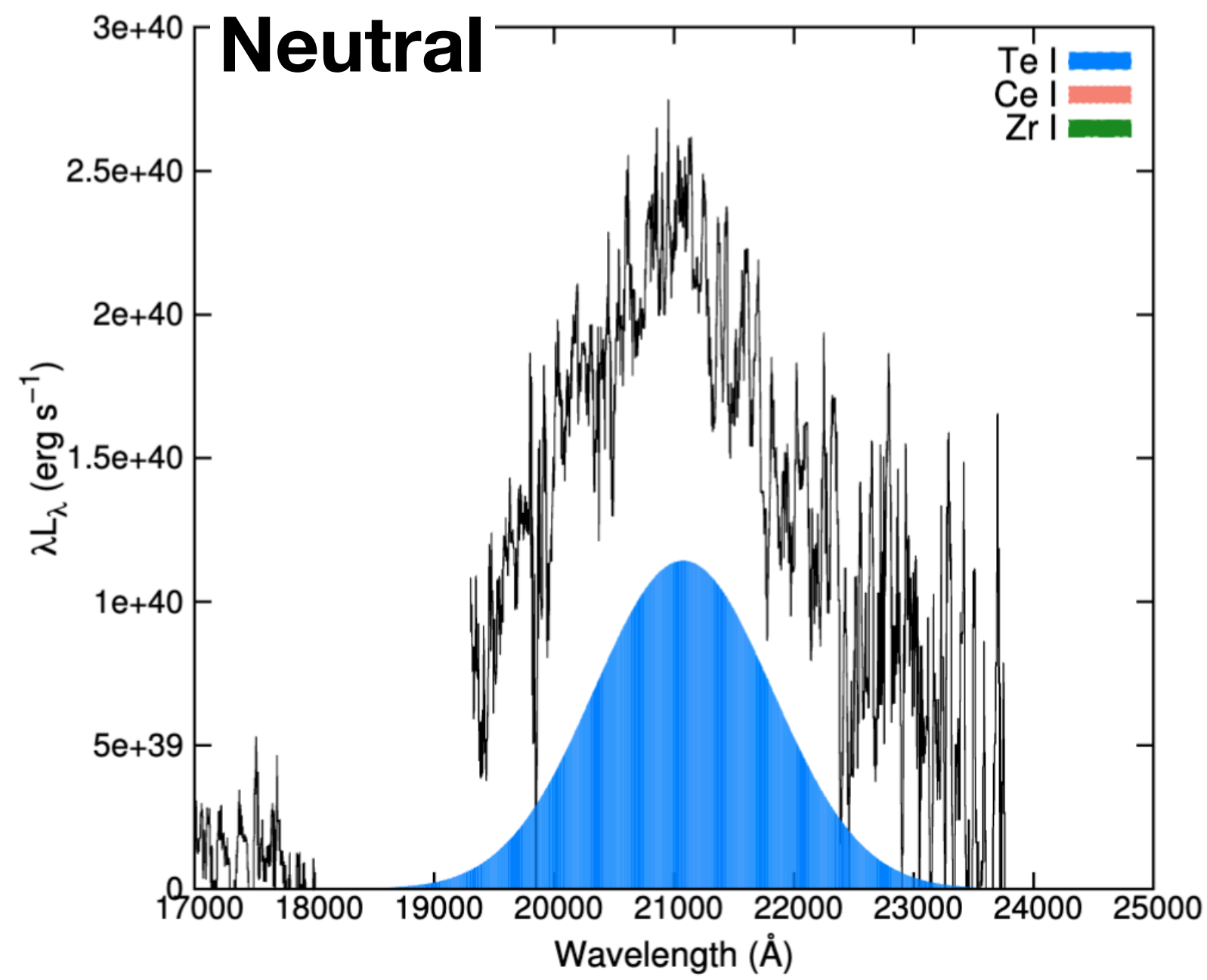
Required mass fraction fro Ce, La, and Te

Te III feature
is due to one
line

$$L = \Delta E_{2.1} n_e \frac{g_l}{\sum_k g_k e^{-E_k/kT}} \frac{X_{Te} M_{ej}}{A_{Te} m_u} \frac{8.6 \times 10^{-6}}{\sqrt{T}} \frac{\Upsilon_{2.1}}{g_l} e^{-\Delta E_{2.1}/kT}$$



Important Candidates (2.1 feature)



Important Candidates for 2.1 feature (Atomic data)

Element	Wavelength	Lower energy	Computed luminosity (erg/s)	Energy-average collision strength	A-value
Te I	21048	0.0	1.0E+39	1.16	6.00
Ce I	20691	0.0	3.5E+36	13.05	5.58E+08
Zr I	20531	0.0	9.6E+32	0.45	5.2E-05
	21373	0.0	7.4E+32	1.64	1.8E-05
	21514	0.07	1.6E+33	2.89	1.4E-04
Ba II	20517	0.0	1.7E+37	2.7	0.79
Zr II	20657	0.35	2.4E+34	1.17	1.16E-02
	21014	0.41	1.9E+34	0.16	1.14E-02
	21656	0.35	1.3E+34	0.107	6.1E-03
Tc II	21416	0.0	1.5E+33	0.16	3.98E-03
Te III	21048	0.0	1.6E+39	1.5	5.8
Ag III	21695	0.0	1.4E+37	0.4	7.4
Ce III	20691	0.19	5.7E+36	1.34	4.7E+04